

THE TIME MACHINE FACTORY CONFERENCE 2019 UNIVERSITÁ DI TORINO, 22 SEPTEMBER 2019

# BLACK HOLES ARE TIME MACHINES

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Illustration by Chris Gash







# TIME DILATATION





#### The lifetime of a black hole

#### Francesca Vidotto

# SCHWARZSCHILD IS BOTH BLACK AND WHITE



- White Holes and Black Holes shares the same Schwarzschild spacetime
- Only checking if matter is outgoing or ingoing they can be distinguished
- This could be practically impossible! A long time may be needed!





How does a Black Hole die?







Haggard, Rovelli 1407.0989



# WHITE HOLES AS REMNANTS

Bianchi, Cristodoulou, D'Ambrosio, Haggard, Rovelli 1802.04264



 $(\tau, n)^2$ ], and the coefficients have  $1 \sigma$  across the shell. The metric is of the leading terms are determined

$$h_2(\tau)n^2 + h_3(\tau)n^3 + \cdots,$$
 (2.8)

$$g_{2}(\tau)n^{2}+g_{3}(\tau)n^{3}+\cdots$$
 (2.9)

equation for the stress tensor comidary of the collapsing shell there is is for u and  $\sigma$ , and the stress tensor we have to describe this system now

the shell is of the form it the classical Lagrangian for the

$$(\partial A)^2 - m^2 A^2 + \cdots ].$$
 (2.10)

ing shell, the matching equation sive mode of the string. In the rest the collapsing shell couples to the

 $np^{00}$ 

$$-\frac{Q}{2R}\left[\left|\dot{R}^{2}+\left(1-\frac{Q}{R}\right)^{2}\right|^{1/2}\right]$$

$$u(\tau, 0).$$
 (2.1)

$$\frac{\mathbf{x}}{2R} \left\| \vec{R}^{2} + \left( 1 - \frac{\mathbf{x}}{R} \right) \right\|$$
$$(\tau, 0). \tag{2}$$

$$u(\tau, 0).$$
 (2.1

$$\frac{\mathcal{L}}{2R} \left\| \dot{R}^{2} + \left[ 1 - \frac{\mathcal{L}}{R} \right] \right\|$$

$$u(\tau, 0). \qquad (2.1)$$

$$u(\tau, 0).$$
 (2.)

$$\frac{\mathcal{L}}{2R} \left\| \dot{R}^{2} + \left| 1 - \frac{\mathcal{L}}{R} \right| \right\|$$

$$\frac{2\mathbf{K}}{u(\tau,0)} \begin{bmatrix} \mathbf{K} \\ \mathbf{K} \end{bmatrix}$$

$$\frac{1}{2R} \left\| R^{2} + \left[ 1 - \frac{\varepsilon}{R} \right] \right\|$$
$$u(\tau, 0). \tag{2.}$$

$$u(\tau, 0).$$
 (2.

$$u(\tau, 0).$$
 (2.

$$\frac{\mathcal{L}}{2R} \left\| \dot{R}^{2} + \left| 1 - \frac{\mathcal{L}}{R} \right| \right\|$$
$$u(\tau, 0). \tag{2}$$

$$u(\tau,0).$$
 (2

$$u(\tau,0). \tag{2}$$

ymmetric nonsingular vacuum solu-

FIG. 2. Inst

g cornucopion.

shell.

In Einstein's theory, there is

lerivation are in Appendix B

<sup>9</sup>Appendix C.

WHI

tails.

$$u(\tau, 0).$$
 (2.

$$-\frac{Q}{2R} \left\| \dot{R}^{2} + \left[ 1 - \frac{Q}{R} \right] \right\|$$

$$u(\tau, 0).$$
 (2.)

$$\frac{2R}{2R} \left[ \frac{K^{2} + \left[ 1 - \frac{2}{R} \right]}{R} \right]$$

$$u(\tau, 0). \qquad (2.1)$$

$$(\tau,0). \tag{2}$$

$$u(\tau, 0)$$
. (2.1

$$2R \qquad [ \qquad R \qquad ] \qquad [ \qquad R \qquad ] \qquad [ \qquad (\tau, 0). \qquad (2.1)$$

$$u(\tau, 0).$$
 (2.

$$\frac{\mathcal{L}}{2R} \left\| \dot{R}^{2} + \left| 1 - \frac{\mathcal{L}}{R} \right| \right\|$$

$$\iota(\tau, 0). \qquad (2.1)$$

$$\begin{bmatrix} 2\mathbf{K} \\ \mathbf{K} \end{bmatrix} \begin{bmatrix} \mathbf{K} \\ \mathbf{K} \end{bmatrix}$$

$$(\tau, 0).$$
(2)

$$\frac{2\mathbf{K}}{(\tau,0)} \begin{bmatrix} \mathbf{K} \\ \mathbf{K} \end{bmatrix}$$

$$\frac{\varepsilon}{2R} \left\| R^{2} + \left( 1 - \frac{\varepsilon}{R} \right) \right\|$$

$$i(\tau, 0). \qquad (2.12)$$

$$R^{2} + \left[1 - \frac{z}{R}\right]$$

$$(2.12)$$

$$u(\tau,0)$$
. (2.12)  
ist be more specific about the fields

$$\left|\dot{R}^{2}+\left(1-\frac{Q}{R}\right)^{2}\right|^{1/2}$$

$$+\left(1-\frac{Q}{R}\right)^2\Big]^{1/2}$$

$$\left[-\frac{Q}{2}\right]^{2}|^{1/2}$$
 (2.11) or th

shell we have made the fairly
$$f_1(\tau) = \frac{a}{R(\tau)}.$$

(2.7)

tion. smooth functions of This gives us a single first equation for  $R(\tau)$ .<sup>9</sup> The so guarantee the existence that the coefficients in the ex solution to check that the of like  $R(\tau) \approx Q + e^{-\gamma \tau}$ , as  $\mathcal{T}$ We can continue this proced leading order, are, well behav The so

enable us to demonstrate smooth perturbation expansion around the shell, does 1 such a solution exists. smooth collapsing solution, We continue to seare Of courser under my and and an≷ever∕ a sensible ansatz that v where smooth so

 $r_s = \overline{0}$ 











## HOLES TIMES



## HOLE TIMES

**TRANSITION TIME**  $\sim M^0$  (Planckian)

 $\blacksquare \mathsf{INTERNAL} \mathsf{BOUNCE} \mathsf{TIME} \sim \mathbf{M}$ 

## EXTERNAL BOUNCE TIME

- Death by evaporation  $\sim M^3$
- Earlier quantum instability  $\sim M^2$

#### **REMNANT LIFETIME** $\sim M^4$

Ingredients:

1. General Relativity

\* No modifications ( $\Lambda$ + included) but relax conditions on manifold

## 2. Quantum Mechanics

\* Violation of Einstein eq.s in a finite region  $\Rightarrow$  NO central BH singularity

3. Non-perturbative Methods

\* Effects that are not captured by standard QFT

Example: TUNNELLING

perturbasival corrections non-perturbative corrections

## TIME DILATATION

Vidotto, Rovelli 1401.6562



The lifetime of a black hole

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# LOOP QUANTUM GRAVITY



# BOUNDARY STATE



- $\blacksquare Boundary: B_3 \ U \ B_3 \quad (joined \ on \ a \ S_2)$
- Each B<sub>3</sub> can be triangulated by 4 isosceles tetrahedra
- The bulk can be approximated to first order by two 4-simplices joined by a tetrahedron





## HOLE TIMES

**TRANSITION TIME** ~  $M^0$  (Planckian)

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#### **Black hole explosions?**

QUANTUM gravitational effects are usually ignored in calculations of the formation and evolution of black holes. The justification for this is that the radius of curvature of spacetime outside the event horizon is very large compared to the Planck length  $(G\hbar/c^3)^{1/2} \approx 10^{-33}$  cm, the length scale on which quantum fluctuations of the metric are expected to be of order unity. This means that the energy density of particles created by the gravitational field is small compared to the space-time curvature. Even though quantum effects may be small locally, they may still, however, add up to produce a significant effect over the lifetime of the Universe  $\approx 10^{17}$  s which is very long compared to the Planck time  $\approx 10^{-43}$  s. The  $\beta_{ij}$  will not be zero because the time dependence of the metric during the collapse will cause a certain amount of mixing of positive and negative frequencies. Equating the two expressions for  $\phi$ , one finds that the  $b_i$ , which are the annihilation operators for outgoing scalar particles, can be expressed as a linear combination of the ingoing annihilation and creation operators  $a_i$  and  $a_i^+$ 

$$b_i = \sum_i \{\bar{\alpha}_{ij}a_j - \bar{\beta}_{ij}a_j^+\}$$

Thus when there are no incoming particles the expectation value of the number operator  $b_i^{\dagger}b_i$  of the *i*th outgoing state is

$$< 0_{-} |b_{i}^{+}b_{i}| 0_{-} > = \sum_{i} |\beta_{ii}|^{2}$$

The number of particles created and emitted to infinity in a gravitational collapse can therefore be determined by calculating the coefficients  $\beta_{ij}$ . Consider a simple example in which

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Hawking evaporation: **m<sup>3</sup>~10**<sup>50</sup> Hubble time

#### Nature Vol. 248 March 1 1974

the collapse is spherically symmetric. The angular dependence of the solution of the wave equation can then be expressed in terms of the spherical harmonics  $Y_{im}$  and the dependence on retarded or advanced time u, v can be taken to have the form  $\omega^{-1/2} exp$  ( $i\omega u$ ) (here the continuum normalisation is used). Outgoing solutions  $p_{im\omega}$  will now be expressed as an integral over incoming fields with the same l and m:

$$p_{\omega} = \int \left\{ \alpha_{\omega \omega'} f_{\omega'} + \beta_{\omega \omega'} \bar{f}_{\omega'} \right\} d\omega'$$

(The lm suffixes have been dropped.) To calculate  $\alpha_{\omega\omega'}$  and  $\beta_{\omega\omega'}$  consider a wave which has a positive frequency  $\omega$  on  $I^+$  propagating backwards through spacetime with nothing crossing the event horizon. Part of this wave will be scattered by the curvature of the static Schwarzschild solution outside the black hole and will end up on  $I^-$  with the same frequency  $\omega$ . This will give a  $\delta(\omega - \omega')$  behaviour in  $\alpha_{\omega\omega'}$ . Another part of the wave will propagate backwards into the star, through the origin and out again onto  $I^-$ . These waves will have a

Beckenstein<sup>6</sup> suggested on thermodynamic grounds that some multiple of  $\kappa$  should be regarded as the temperature of a black hole. He did not, however, suggest that a black hole could emit particles as well as absorb them. For this reason Bardeen, Carter and I considered that the thermodynamical similarity between  $\kappa$  and temperature was only an analogy. The present result seems to indicate, however, that there may be more to it than this. Of course this calculation ignores the back reaction of the particles on the metric, and quantum fluctuations on the metric. These might alter the picture.

Further details of this work will be published elsewhere. The author is very grateful to G. W. Gibbons for discussions and help. S. W. HAWKING

Department of Applied Mathematics and Theoretical Physics and Institute of Astronomy University of Cambridge For something quantum to happens, semiclassical approximation must fail. Typically in quantum gravity: high curvature Curvature ~ (L<sub>P</sub>)-<sup>2</sup> Small effects can pile up: small probability per time unit gives a probable effect on a long time! Typically in quantum tunneling: Curvature × (time) ~ (L<sub>P</sub>)-<sup>1</sup>

$$\frac{1}{m^2} T_b \sim 1$$

 $\implies$  the hole lifetime must be longer or of the order of  $\sim m^2$ 

Haggard, Rovelli 1407.0989

Quantum Break Time Dvali, Gomez 1112.3359

## Black-to-White Tunnelling

- In the quantum world, things happen as soon as they can!
- Indications from a full LQG computations. Chistodoulou, Rovelli, Speziale, Vilensky 1605.05268

#### LARGE EXTRA DIMENSIONS

1st order topological phase transition from black string to black hole occurring because of the Gregory-Laflamme metrical instability

 $M_{\rm BH}^{(3+2n)/(1+n)}$ 

#### **BRANES**

Large black holes localized on infinite Randall-Sundrum branes: period of rapid decay via Hawking radiation of CFT modes Emparana, Garcia-Bellido, Kaloper 2003

Quantum effects shorten the lifetime of black holes!

Casadio and Harms 2000/01 Gubser 2002, Kol 2002 Gregory and Laflamme 2002

## Assumptions:

Almheiri, Marolf, Polchinski, Sully 1207.3123

- General Relativity: Equivalence Principle
- Quantum Mechanics: Unitary Evolution
- QFT in Curved Spacetime (fixed smooth background)

Firewall argument See also Rovelli: "The Subtle Unphysical Hypothesis of the Firewall Theorem" after the Page time i.e. when about half of black-hole mass has evaporated particles emitted needs to break entanglement releasing an enormous energy

## Black Hole Lifetime

Quantum Gravity effects should manifest before the Page time  $\implies$  the hole lifetime must be shorter or of the order of  $\sim m^3$ 

# See also Quantum Break Time

Dvali, Gomez 1112.3359

Vidotto, Rovelli 1401.6562

# Scenario 3 FAST EXPLOSION

# PRIMORDIAL BLACK HOLES

- PBHs are the least exotic beast in the dark universe zoo of theories
- PBHs are a viable DARK MATTER candidate

\* careful with old constraints in the literature!

PBHs are interesting even if they are not all DARK MATTER
 \* PBHs can be used to test QUANTUM GRAVITY

# (QUANTUM) PBH DARK MATTER

• Today, black holes smaller than  $m(t)|_{t=t_H}$  have already exploded.

It decreases with time. (but for later accretion/merging)





Effects on late cosmology

Galaxy clusters surveys



Raccanelli, Vidotto, Verde 1708.02588

## EFFECT ON GALAXY CLUSTERS

Raccanelli, Vidotto, Verde 1708.02588

$$C_{\ell}^{XY}(z_i, z_j) = \left\langle a_{\ell m}^X(z_i) \ a_{\ell m}^{Y^*}(z_j) \right\rangle$$

 angular positions and redshifts perturbed by peculiar velocities, gravitational lensing and potentials

Choice of redshift distribution:





How does a Black Hole die?

Characterisation of the signal

# PBH EXPLOSIONS

- fast process (few milliseconds?)
- the source disappears with the burst
- very compact object: big flux  $E = mc^2 \sim 1.7 \times 10^{47} \text{ erg}$

• exploding today: 
$$m = \sqrt{\frac{t_H}{4k}} \lesssim 1.2 \times 10^{23} \text{ kg}$$
  $R = \frac{2Gm}{c^2} \lesssim .02 \text{ cm}$ 

- **HIGH ENERGY:** energy of the particle liberated  $\approx Tev$
- SYNCHROTRON EMISSION (REES' MECHANISM)

• LOW ENERGY: size of the source  $\approx$  wavelength  $\lambda_{predicted} \gtrsim .2$  cm

GRAVITATIONAL WAVES !!!

**FRB** (?)

Barrau, Bolliet, Vidotto, Weimer 1507.1198

■ shorter lifetime — smaller wavelength

#### Low energy channel

#### High energy channel



- detection of arbitrarily far signals
- better single-event detection



- PBH: mass temperature relation
- different scaling

# THE SMOKING GUN: DISTANCE/ENERGY RELATION



distant signals originated in younger, smaller&hotter sources

How does a Black Hole die?

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# Remnants

# QUANTUM EFFECTS MAKE REMNANTS STABLE

#### ■ LARGE INTERNAL VOLUME ~ M<sub>o</sub><sup>4</sup>

It depends only on the original mass  $M_o$  at the BH formation  $\gamma \gamma$ 

#### Christodoulou, Rovelli 1411.2854 Christodoulou, De Lorenzo1604.07222

#### **REMNANT LIFETIME** ~ $M_0^4$

Bianchi, Cristodoulou, D'Ambrosio, Haggard, Rovelli 1802.04264

Time for information to leak out from such a large volume trough the small WH surface.

#### PROCESSES

- 1. BH volume increase & WH volume decrease
- 2. White to black instability
- 3. Hawking evaporation
- 4. Black to white tunnelling

From the outside, at a finite time, no distinction between black and white holes

#### STABILITY

The minimal area yields a minimal mass!

$$|R\rangle = \frac{\sqrt{\frac{a}{b}}|B,\mu\rangle - |W,\mu\rangle}{\sqrt{1 + \frac{a}{b}}}$$

oscillation between black and white hole states

Rovelli, Vidotto 1805.03872

How does a Black Hole die?

# DARK MATTER: REMNANTS FORMED AFTER INFLATION

Rovelli, Vidotto 1805.03872

■ PBHs form at the reheating, evaporates and evolve in a long-living remnant

## ■ REMNANT LIFETIME COMPATIBLE WITH FORMATION AT REHEATING $\mathbf{M}_{o}^{4} \ge \mathbf{t}_{\text{Hubble}}$ $\mathbf{M}_{o}^{3} < \mathbf{t}_{\text{Hubble}}$ $\implies 10^{10} gr \le \mathbf{M}_{o} < 10^{15} gr$ $\implies 10^{-18} cm \le R_{o} < 10^{-13} cm$

#### NUCLEOSYNTHESIS

BH evaporation should not modify D/H, Li6/Li7, and He3/D ratio

# DARK MATTER: REMNANTS FROM BEFORE THE BOUNCE

Rovelli, Vidotto 1805.03224

Quintin, Brandenberger 1609.02556

Carr, Clifton, Coley 1704.02919

#### BOUNCING BLACK HOLES IN A BOUNCING UNIVERSE

Planckian PBH remands from a previous eon (Penrose's **EREBONS**) Planck size particles can pass trough the bounce.

#### PAST LOW ENTROPY

Matter near thermal equilibrium: geometry has low entropy A volume of the universe outside BH as low as only  $1/T_H^2 \sim 10^{-120}$  of the total could have been outside the remnants at the bounce!

#### **DARK MATTER**

We want  $\mathbf{M}_{o^4} \ge \mathbf{t}_{\text{Hubble}}$  for them to survive till today.

Inflation dilutes PBH: 
$$\frac{1}{T_H^2} \sim \left(\frac{\dot{a}}{a}\right)^2 \sim \rho_M \qquad \rho_b \sim \rho T_H^3 \sim T_H \qquad V_{int} = \rho_b V_{WH} > T_H^2$$

■ MATTER BOUNCE: PBH as pressureless component

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# to conclude

**TRANSITION TIME** ~  $M^0$  (Planckian)

### **INTERNAL BOUNCE TIME** ~ M

## **EXTERNAL BOUNCE TIME**

- Death by evaporation  $\sim M^3$
- Earlier quantum instability  $\sim M^2$

### REMNANT LIFETIME $\sim M^4$

## QG allows for Black-to-White tunnelling

\* Mystery solved: this is how a BH dies!
\* Instability possibly before Hawking
evaporation time: new phenomenology

**TRANSITION TIME**  $\sim M^0$  (Planckian)

 $\blacksquare$  INTERNAL BOUNCE TIME  $\sim \mathbf{M}$ 

## EXTERNAL BOUNCE TIME

- Death by evaporation  $\sim M^3$
- Earlier quantum instability  $\sim M^2$

#### **REMNANT LIFETIME** ~ M<sup>4</sup>

## **REMNANTS AS DARK MATTER**

\* compatible with PBH formation at reheating\* stability via minimal area/mass

**BOUNCE<sup>2</sup>**: Bouncing BH in a Bouncing Universe