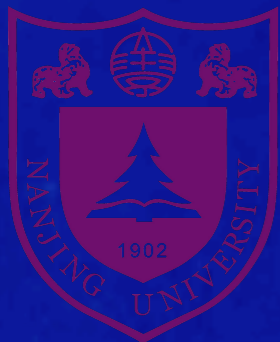


Initiation and early kinematic evolution of solar eruptions

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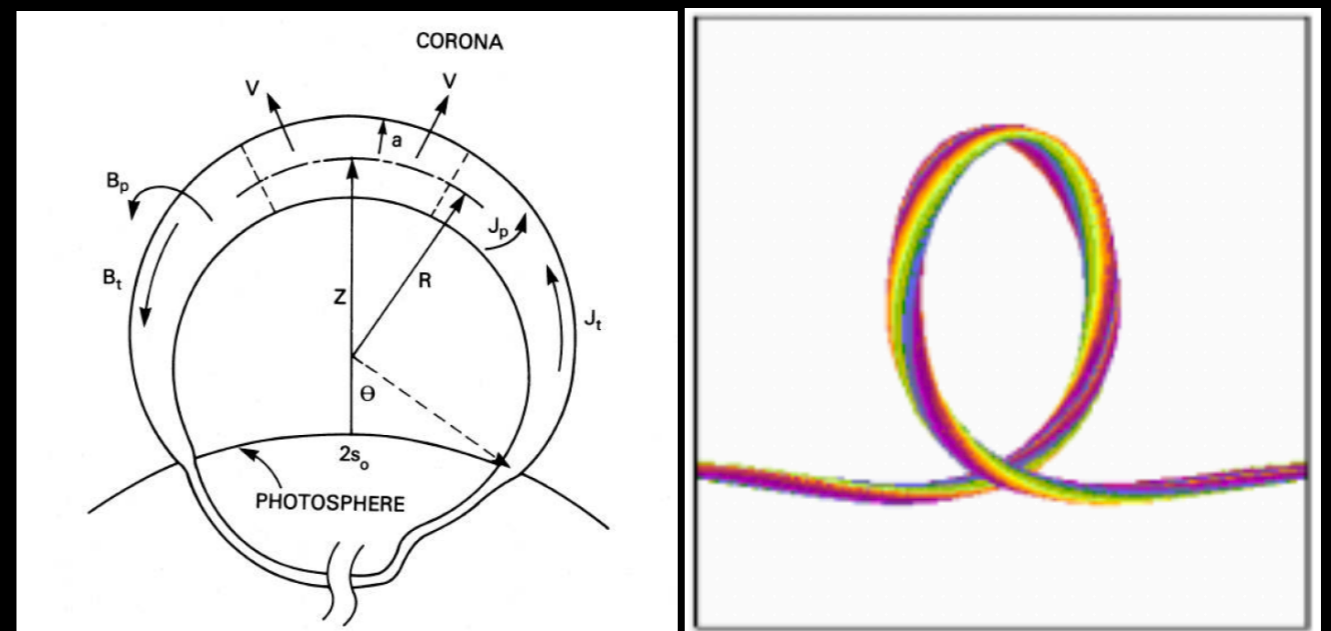
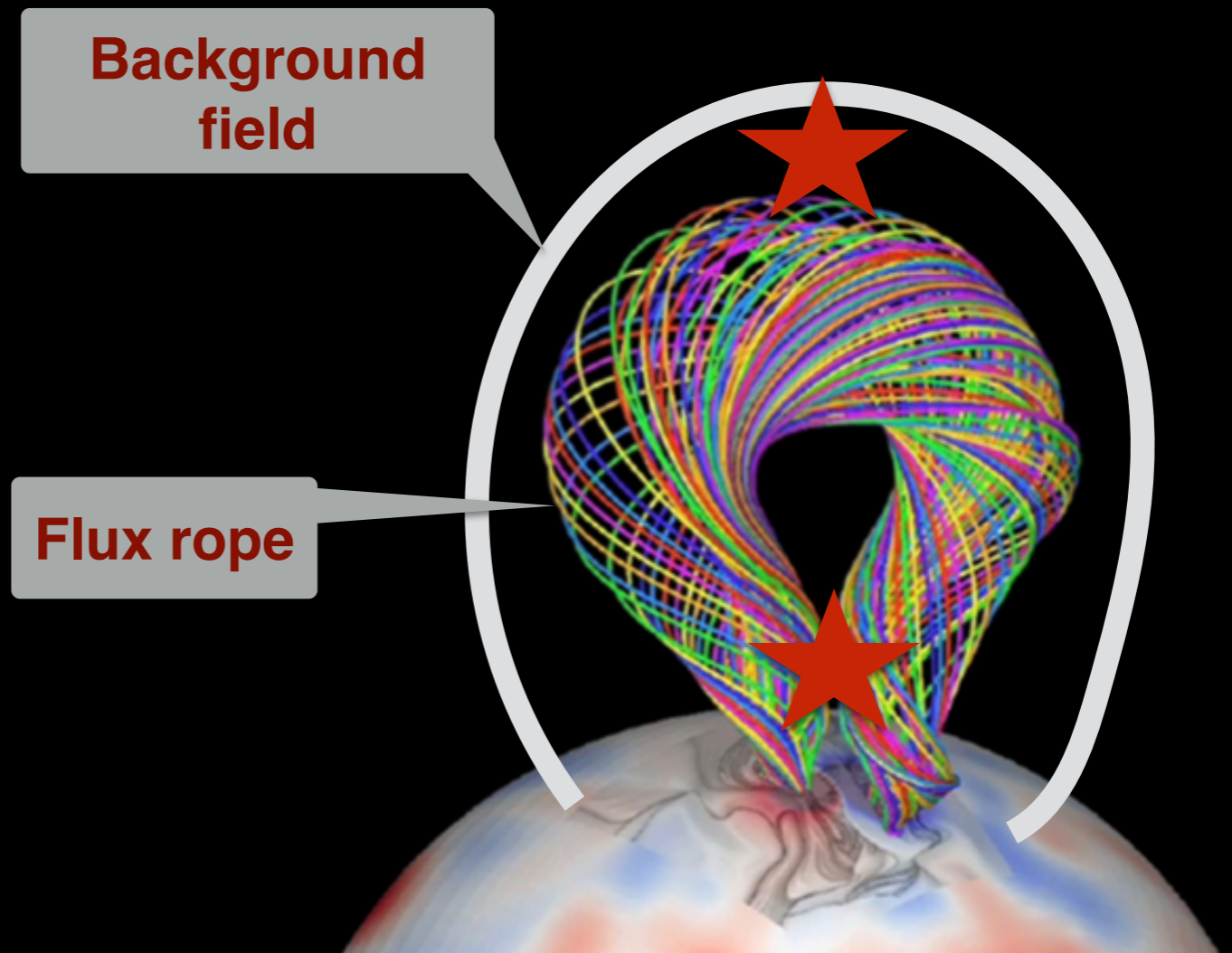
Initiation models

Reconnection type

- **Tether-cutting** (Moore et al. 2001)
- **Breakout** (Antiochos et al. 1999)
- **Flux emergence** (Chen et al. 1999)

Ideal instability type

- **Torus instability** (Kliem & Torok 2006)
- **Kink instability** (Torok & Kliem 2004)
- **Loss of equilibrium** (Lin et al. 2001)

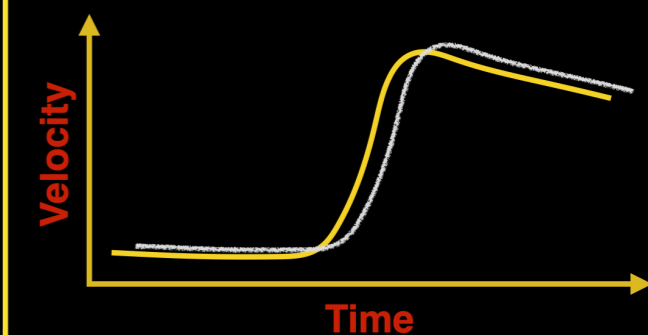
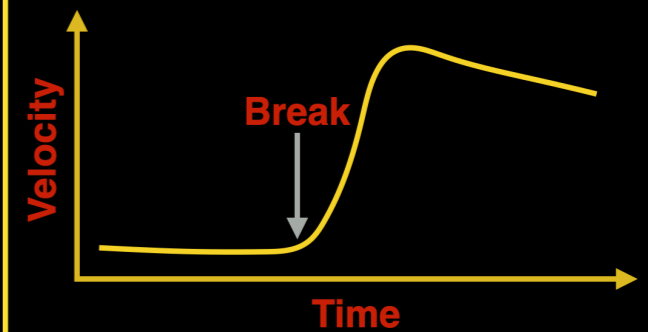


Motivation:

To distinguish initiation models based on kinematics

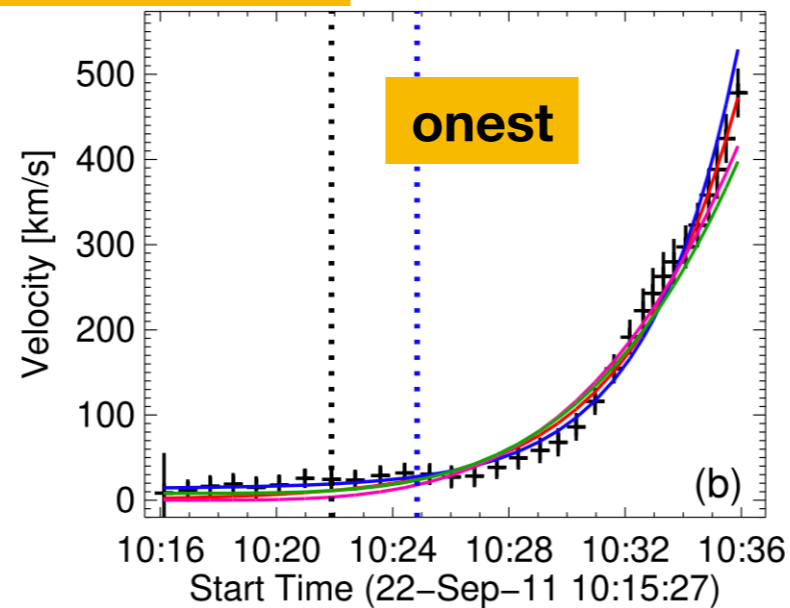
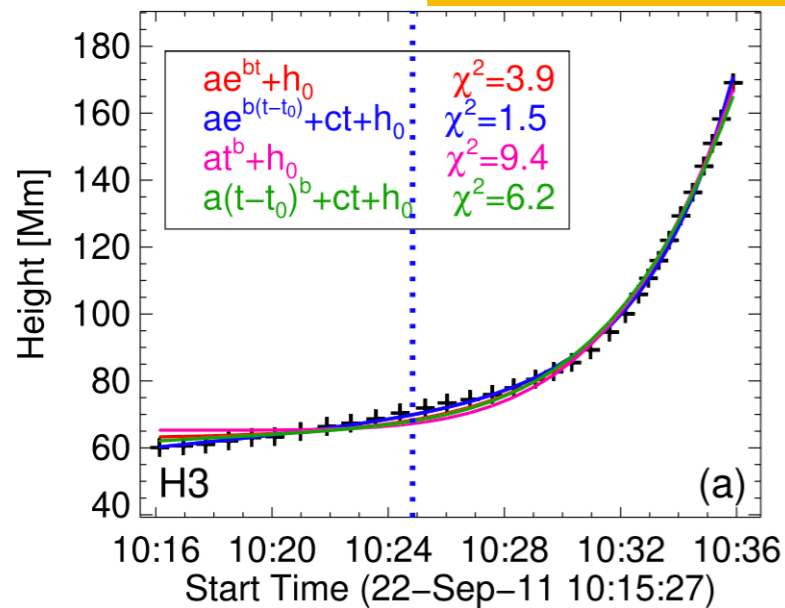
Some valuable characteristics for diagnosing

- A **break** exists in the early H-T profiles? if yes, against the tether-cutting model (Moore et al. 2001) as only one process — “runaway” reconnection — drives the eruptions
- Temporal **offset** between H-T profile and flare light (Fsxr) curve, favors ideal MHD (reconnection) models if H-T profile (Fsxr) precedes.
- **Correlation** between the onset of eruptions, either the slow-rise phase or main-acceleration phase, and the threshold of torus or kink instability favors ideal MHD models, as it is not required in reconnection models.

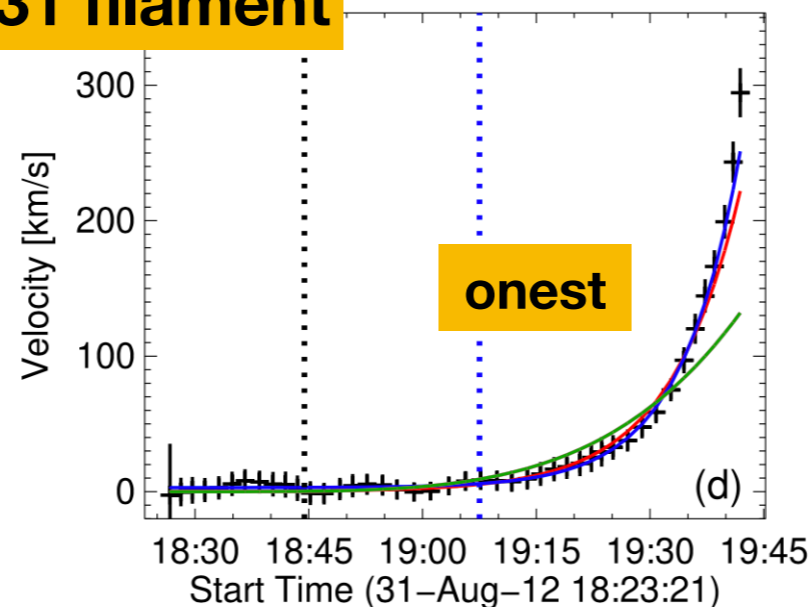
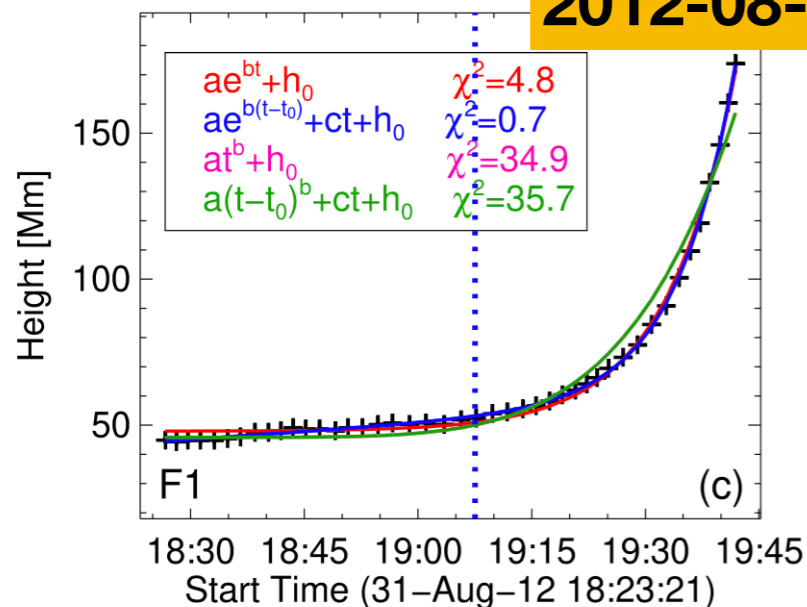


H-T profiles and fitting of CMEs

2011-09-22 hot channel



2012-08-31 filament



➤ H-T presents two phases: slow rise+impulsive acceleration

➤ Fitting with functions:

$$h_1(t) = a_1 \exp(b_1 t) + d_1,$$

$$h_2(t) = a_2 \exp[b_2(t - t_0)] + c_2 t + d_2,$$

$$h_3(t) = a_3 t^{b_3} + d_3,$$

$$h_4(t) = a_4(t - t_0)^{b_4} + c_4 t + d_4,$$

➤ Defining the **onset of the acceleration (nonlinear) phase**, i.e., taking over the linear one.

Table 2. Metrics for fitting goodness of different functions.

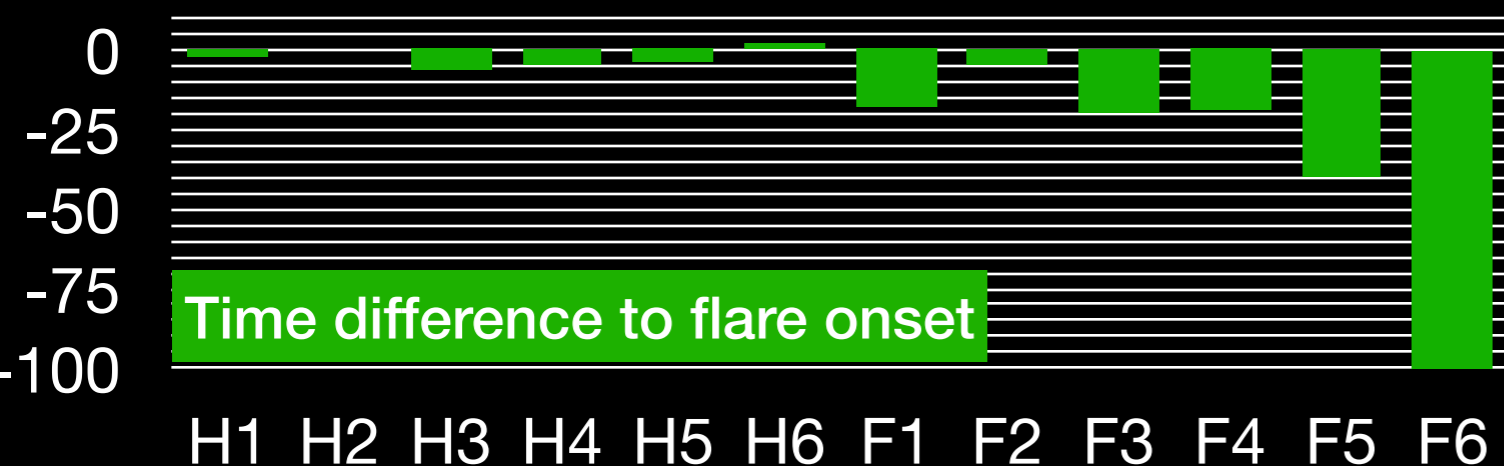
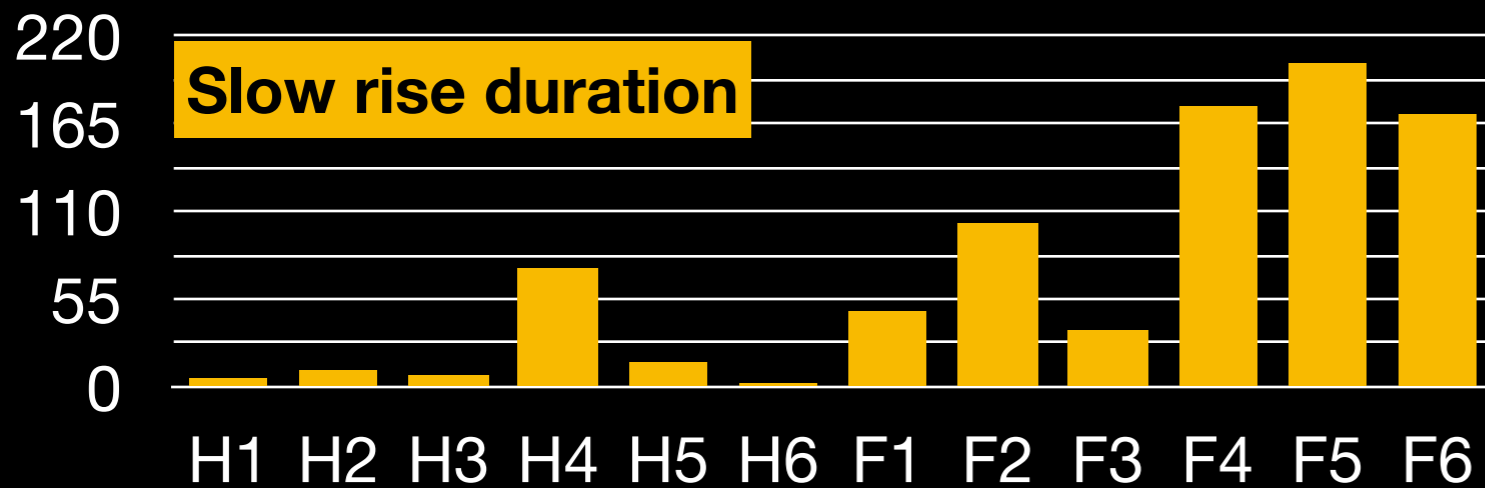
Events	$a_1 e^{b_1 t} + h_0$ $\chi^2_{\nu 1}$	$a_2 e^{b_2(t-t_0)} + c_2 t + h_0$ $\chi^2_{\nu 2}$	$a_3 t^{b_3} + h_0$ $\chi^2_{\nu 3}$	$a_4 (t - t_0)^{b_4} + c_4 t + h_0$ $\chi^2_{\nu 4}$
H1	3.7	0.5	12.7	1.0
H2	10.6	0.9	16.1	5.4
H3	3.9	1.5	9.4	6.2
H4	37.0	1.7	62.6	17.3
H5	2.6	2.8	9.5	1.8
H6	3.5	3.8	1.7	1.5
F1	4.8	0.7	34.9	35.7
F2	4.5	1.4	7.6	9.1
F3	3.7	0.5	6.0	3.7
F4	14.5	1.6	29.2	33.1
F5	42.2	1.5	926.3	949.8
F6	7.8	0.9	18.8	21.7

➤ The best one is the function consisted of **the linear + exponential (9) or power-law (3),**

➤ Duration of the slow-rise phase for the hot channels (several mins) is mostly **shorter** than that of the filaments (>35 mins)

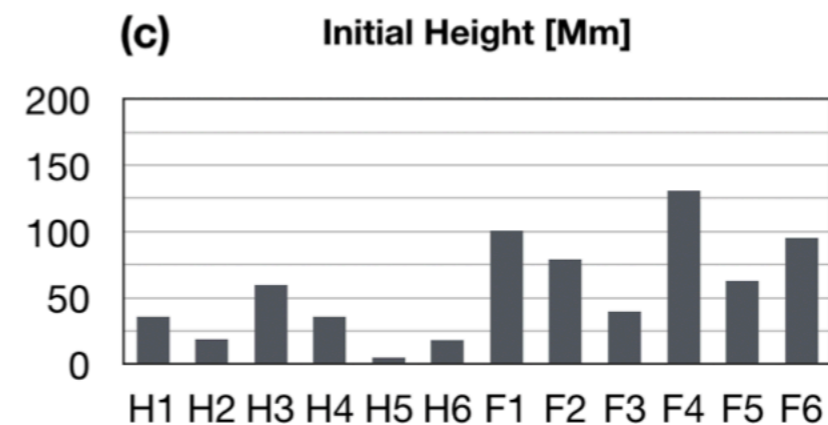
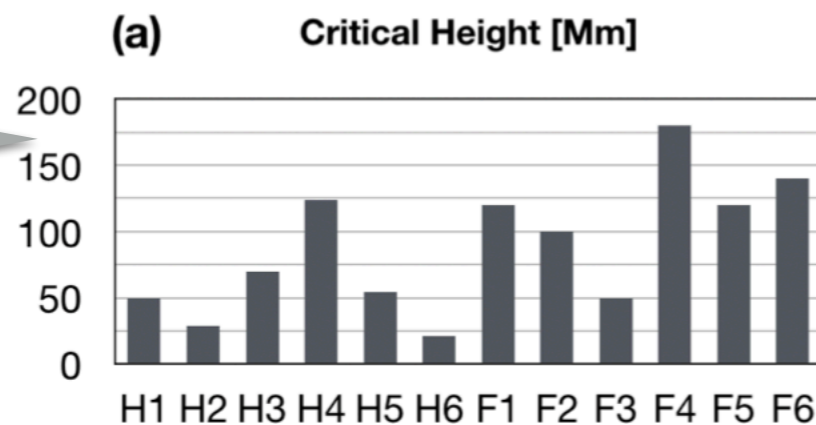
➤ For hot channels, **the time difference is very small**, the speed evolution is completely synchronised with the variation of the flare emission,

➤ For quiescent filaments, **the time difference is large**, the main-acceleration starts earlier than that of flares.



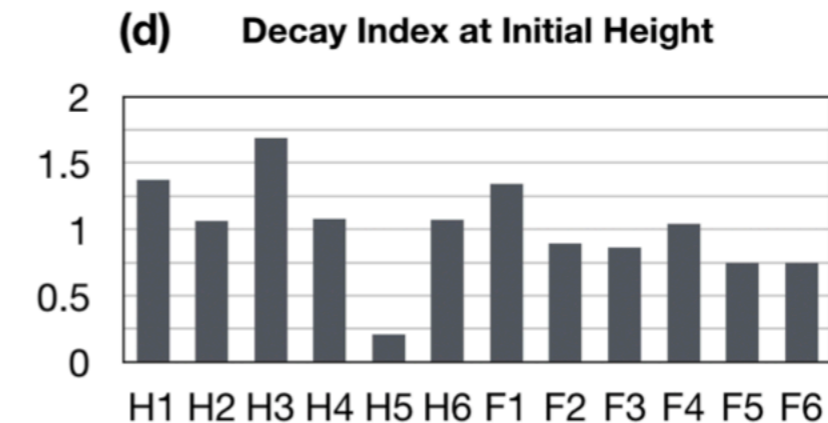
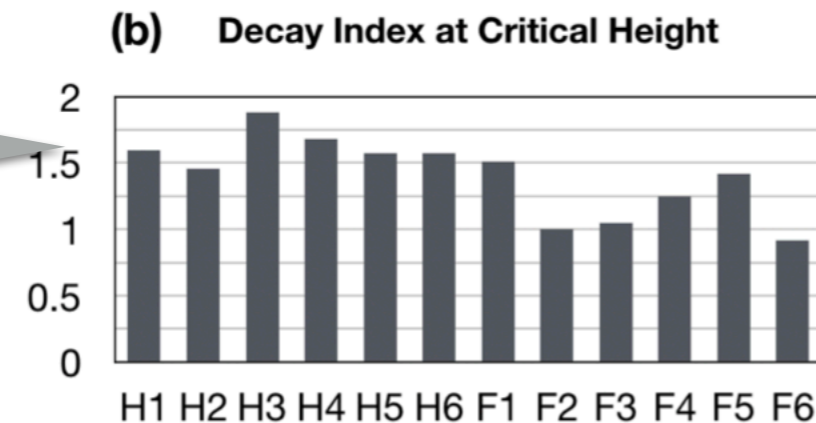
Decay index at critical and initial heights

Critical height



Initial height

Critical decay index



Initial decay index

- The average of decay indices are close to the threshold of torus instability.
- The critical decay indices have a range: 1.46 ± 0.08 – 1.88 ± 0.03 for the former; 0.92 ± 0.11 – 1.51 ± 0.24 for the latter.
- The decay indices at initial heights are too small to initiate torus instability.

Conclusions:

1. The early evolution of all events consists of **a slow-rise phase followed by a main-acceleration phase**, the height-time profiles of which **differ markedly** and can be best fit, respectively, by **a linear and an exponential function**. This indicates that **different physical processes dominate in these phases, which is at variance with models that involve a single process**.
2. The kinematic evolution of the eruptions tends **to be synchronized with the flare light curve in both phases**. The synchronization is often but not always close. **A delayed onset** of the impulsive flare phase is found in **the majority of the filament** eruptions (5 out of 6). This delay, and its trend to be larger for slower eruptions, **favor ideal MHD instability models**.
3. The average decay index at the onset heights of the main acceleration is **close to the threshold of the torus instability** for both groups of events, suggesting that this **instability initiates and possibly drives the main acceleration**.

Reference: Cheng, X., Zhang, J., Kliem, B., Torok, T., Xing, C., Zhou, Z. J., Inhester, B. & Ding, M. D., Initiation and Early Kinematic Evolution of Solar Eruptions, 2020, **ApJ**, 894, 85