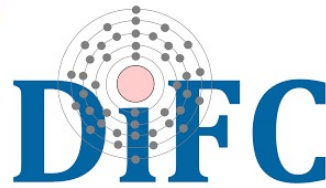




Università
degli Studi
di Palermo



DIAGNOSTICS AND FINE STRUCTURING OF CORONAL HEATING FROM 3D MHD MODELING OF TWISTED/BRAIDED CORONAL LOOPS

F. Reale, G. Peres – University of Palermo, Italy

S. Orlando, A. Petralia – INAF Osservatorio Astronomico di Palermo, Italy

M. Guarrasi – CINECA/Italy

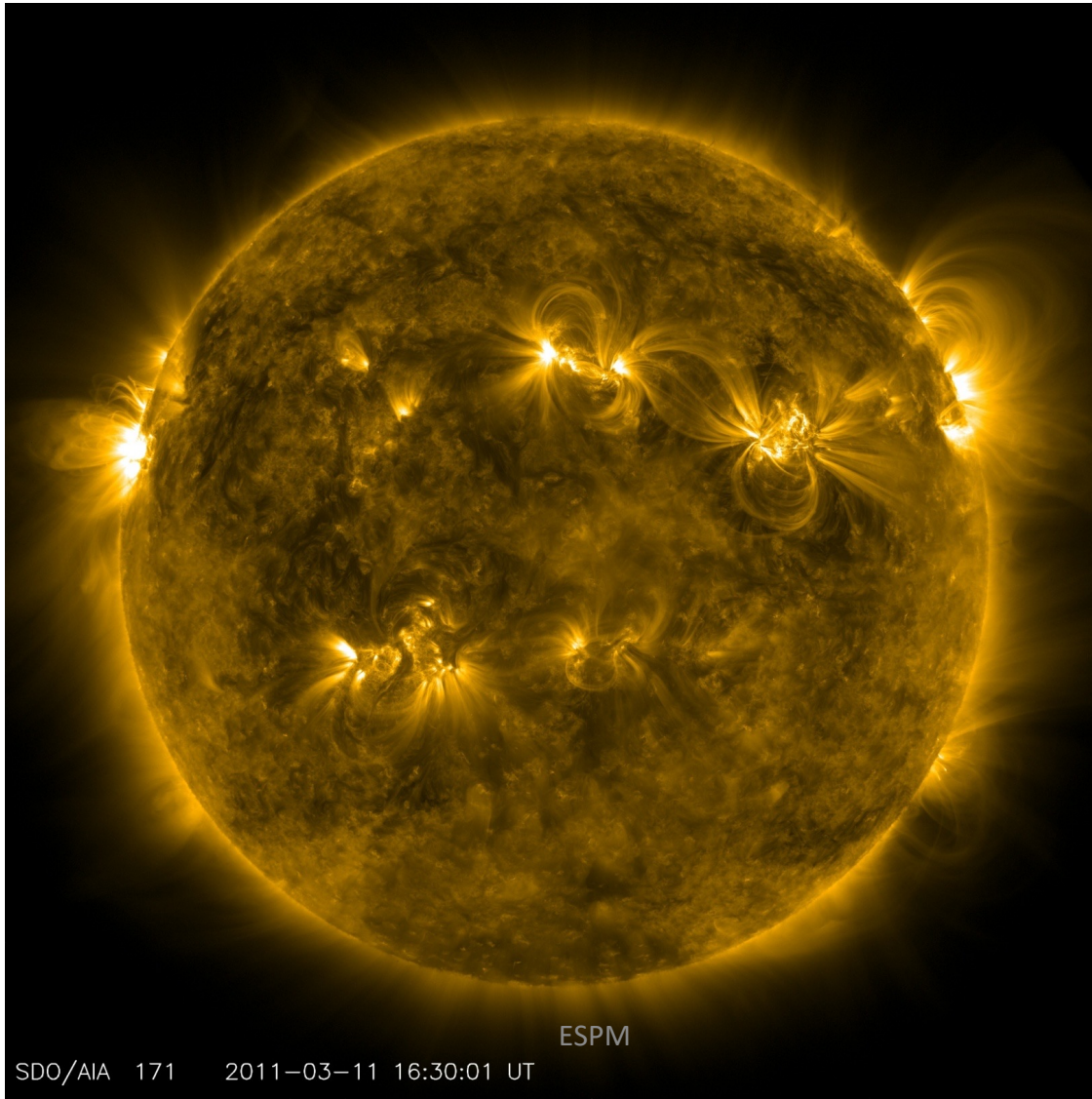
A. Mignone – University of Turin, Italy

A. Hood, E. R. Priest – University of St. Andrews, UK

B. DePontieu – LMSAL, USA

P. Testa – Harvard-Smithsonian CfA, USA

The corona

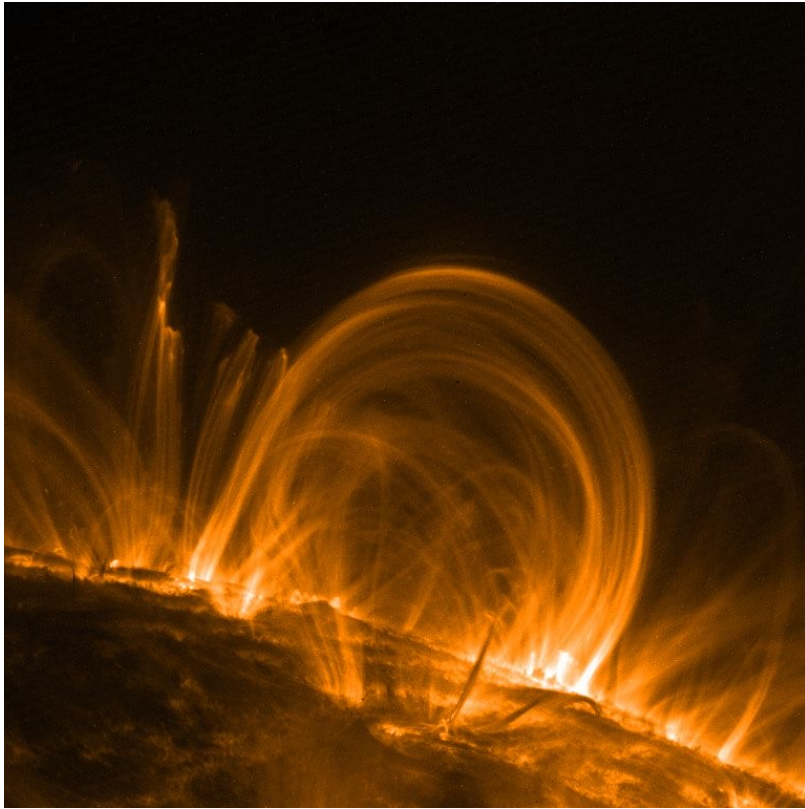


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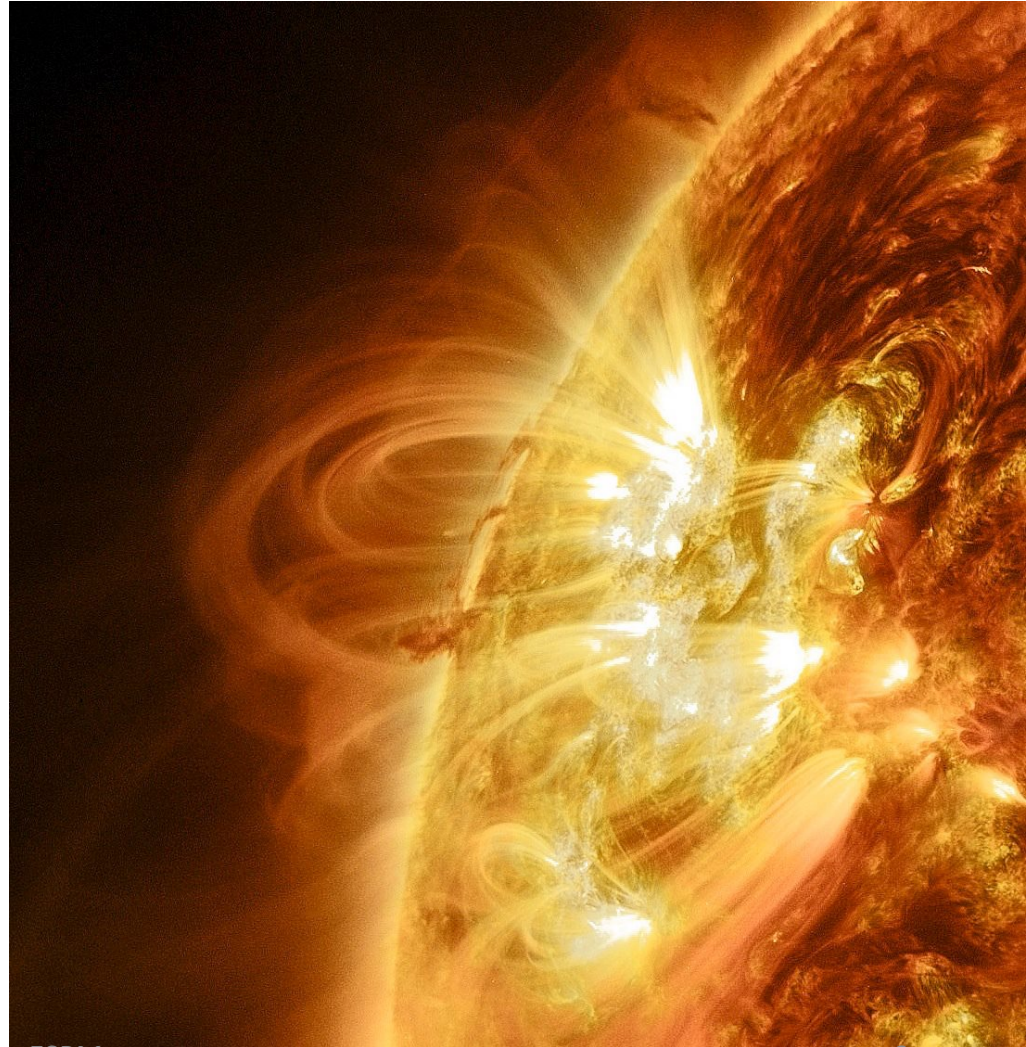
ESPM

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The confined corona: coronal loops

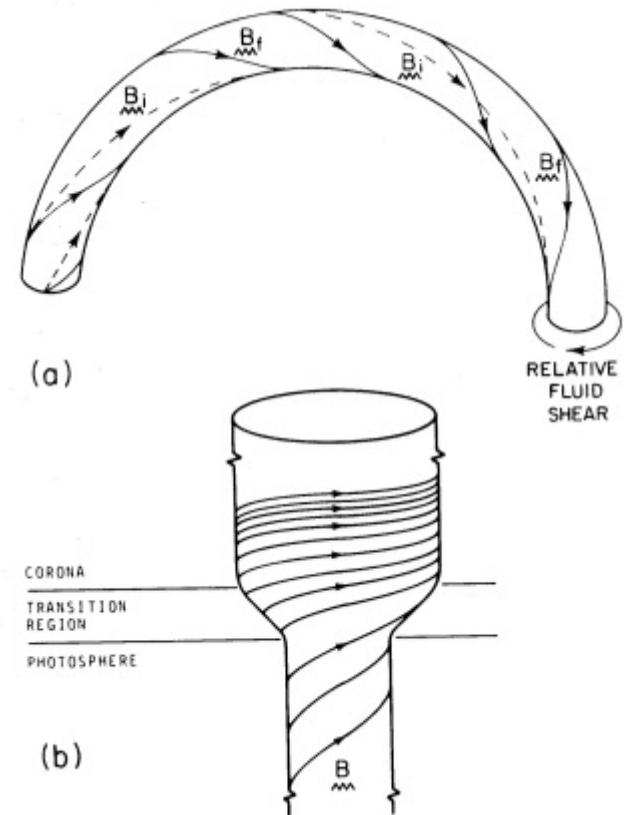
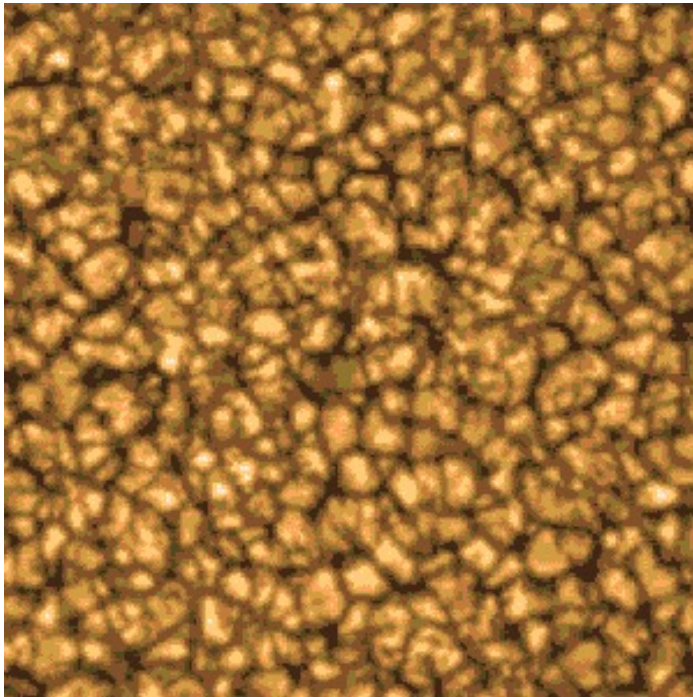


06/09/21



ESPM

MHD modeling of twisted coronal loops



Rosner, Golub, Coppi, Vaiana 1978

3D MHD MODELING OF TWISTED CORONAL LOOPS

F. REALE^{1,2}, S. ORLANDO², M. GUARRASI³, A. MIGNONE⁴, G. PERES^{1,2}, A. W. HOOD⁵, AND E. R. PRIEST⁵

¹ Dipartimento di Fisica & Chimica, Università di Palermo, Piazza del Parlamento 1, I-90134 Palermo, Italy; fabio.reale@unipa.it

² INAF-Osservatorio Astronomico di Palermo, Piazza del Parlamento 1, I-90134 Palermo, Italy

³ CINECA—Interuniversity consortium, via Magnanelli 6/3, I-40033, Casalecchio di Reno, Bologna, Italy

⁴ Dipartimento di Fisica Generale, Università di Torino, via Pietro Giuria 1, I-10125, Torino, Italy

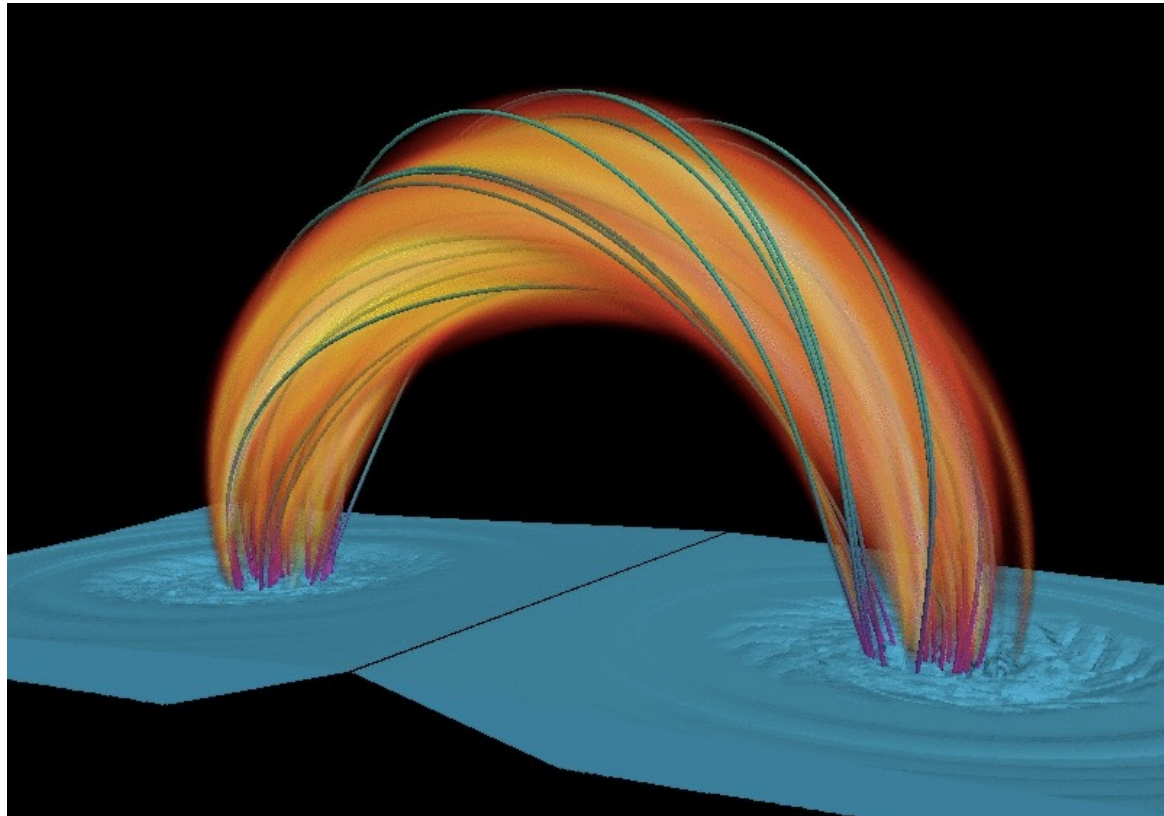
⁵ School of Mathematics and Statistics, University of St. Andrews, St. Andrews, KY16 9SS, UK

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The 3D MHD model

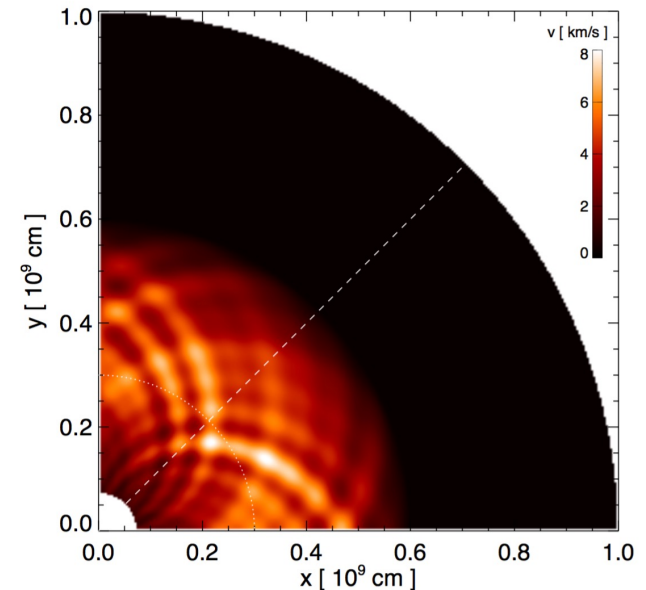
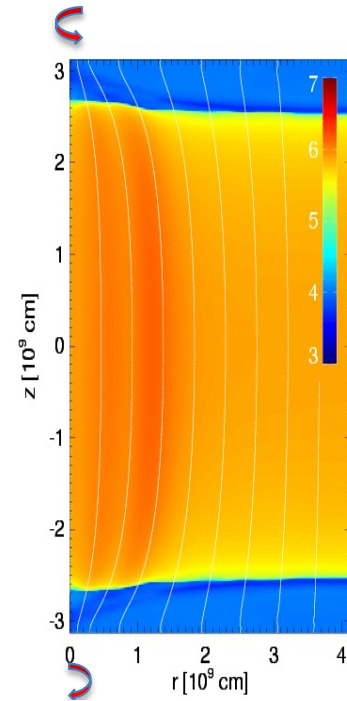
(Reale et al. 2016)

- Twisting is driven by rotation of the footpoints rooted in the photosphere
- We assume Joule heating above a current threshold (Hood+ 2009)



The basic loop model

- The loop field: tapering down to the chromosphere
- The twisting/braiding
 - Footpoint rotation
 - Basic constant angular speed ω
 - Maximum: 5 km/s (both footpoints)
 - Radius: $r = 3000$ km
 - **Braiding:** RANDOM COMPONENT ADDED TO ROTATION VELOCITY AT THE FOOTPOINTS (see pattern aside)



The 3D MHD model

(PLUTO 4 - Mignone+ 2007)

Three-dimensional cylindrical coordinates: r, φ, z

One quarter domain: $0 < \varphi < \pi/2, r_0 = 0.07 \cdot 10^9 \text{ cm}$

The equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0 ,$$

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \mathbf{u} - B \mathbf{B} + I P_t) = \rho \mathbf{g} ,$$

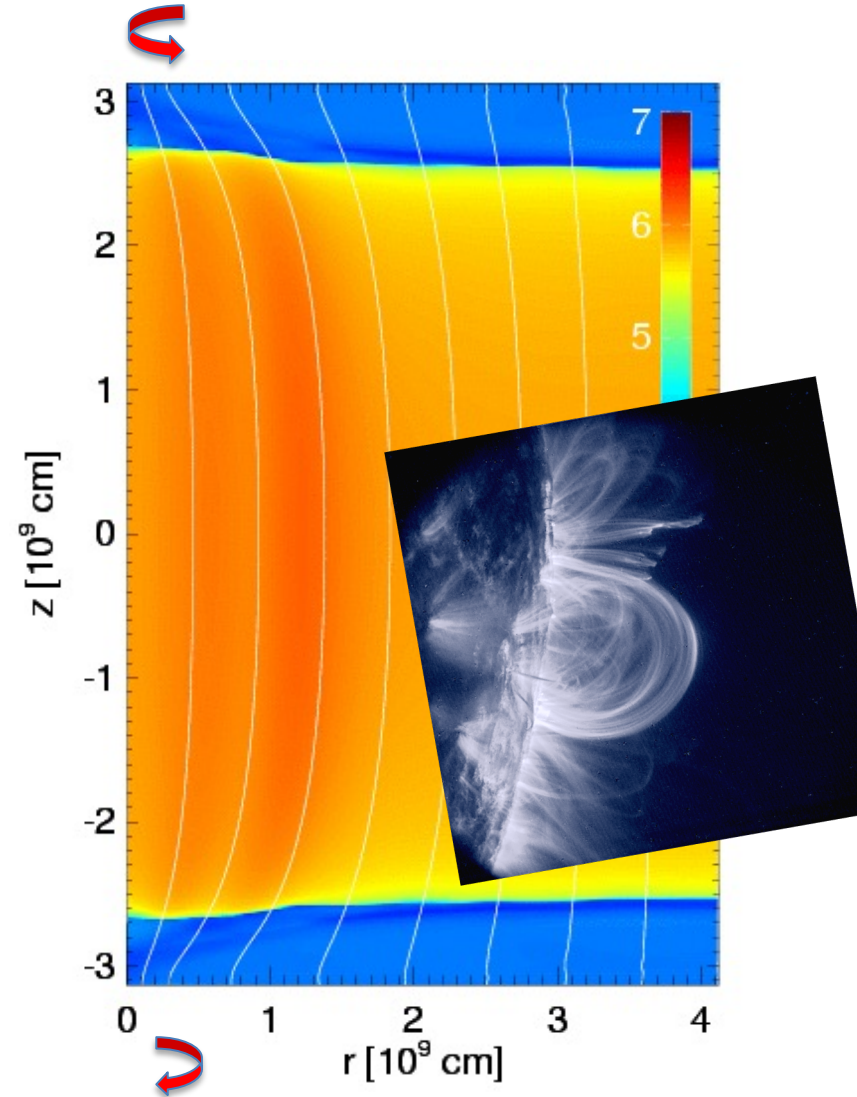
$$\begin{aligned} \frac{\partial \rho E}{\partial t} + \nabla \cdot [\mathbf{u}(\rho E + P_t) - B(\mathbf{u} \cdot \mathbf{B})] = \\ - \nabla \cdot [(\boldsymbol{\eta} \cdot \mathbf{J}) \times \mathbf{B}] + \rho \mathbf{u} \cdot \mathbf{g} - \nabla \cdot \mathbf{F}_c - n_e n_H \Lambda(T) \end{aligned}$$

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{u} \mathbf{B} - \mathbf{B} \mathbf{u}) = -\nabla \times (\boldsymbol{\eta} \cdot \mathbf{J}) ,$$

where

$$P_t = P + \frac{\mathbf{B} \cdot \mathbf{B}}{2} , \quad E = \epsilon + \frac{\mathbf{u} \cdot \mathbf{u}}{2} + \frac{\mathbf{B} \cdot \mathbf{B}}{2\rho} ,$$

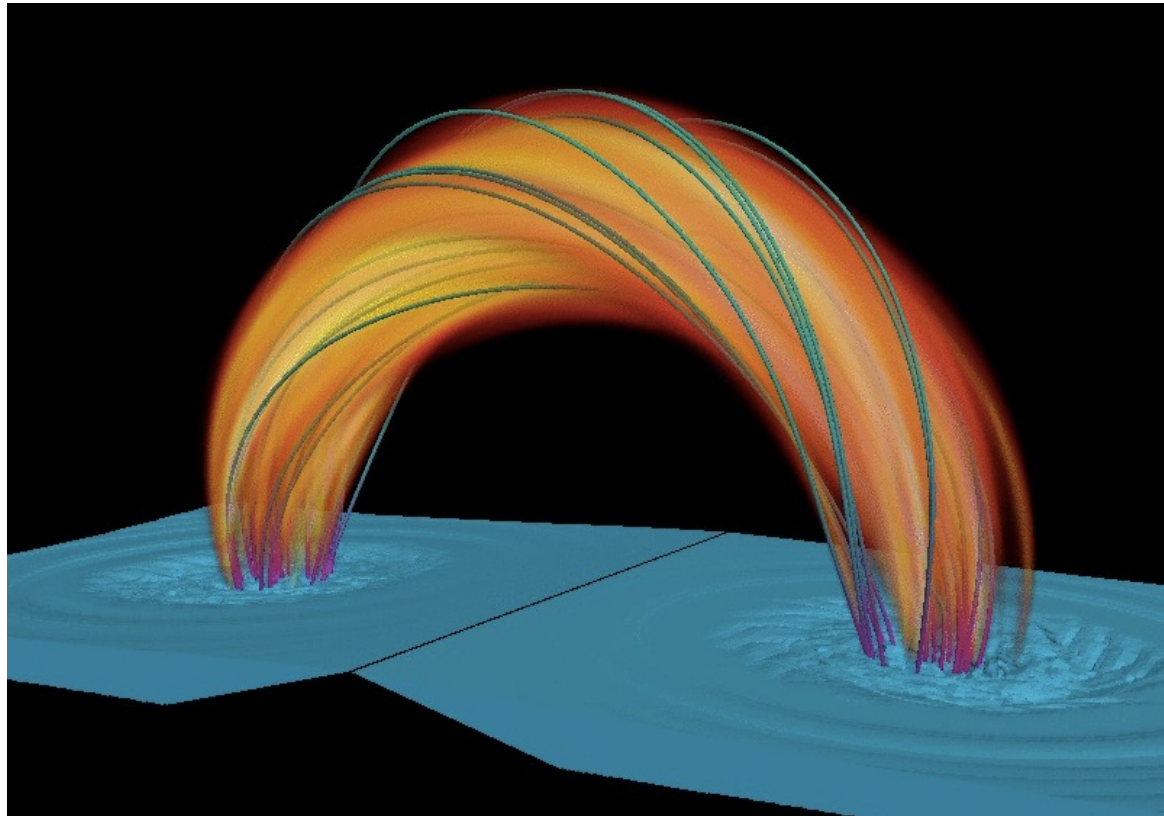
$$\mathbf{F}_c = \frac{F_{sat}}{F_{sat} + |\mathbf{F}_{class}|} \mathbf{F}_{class}$$



The 3D MHD model

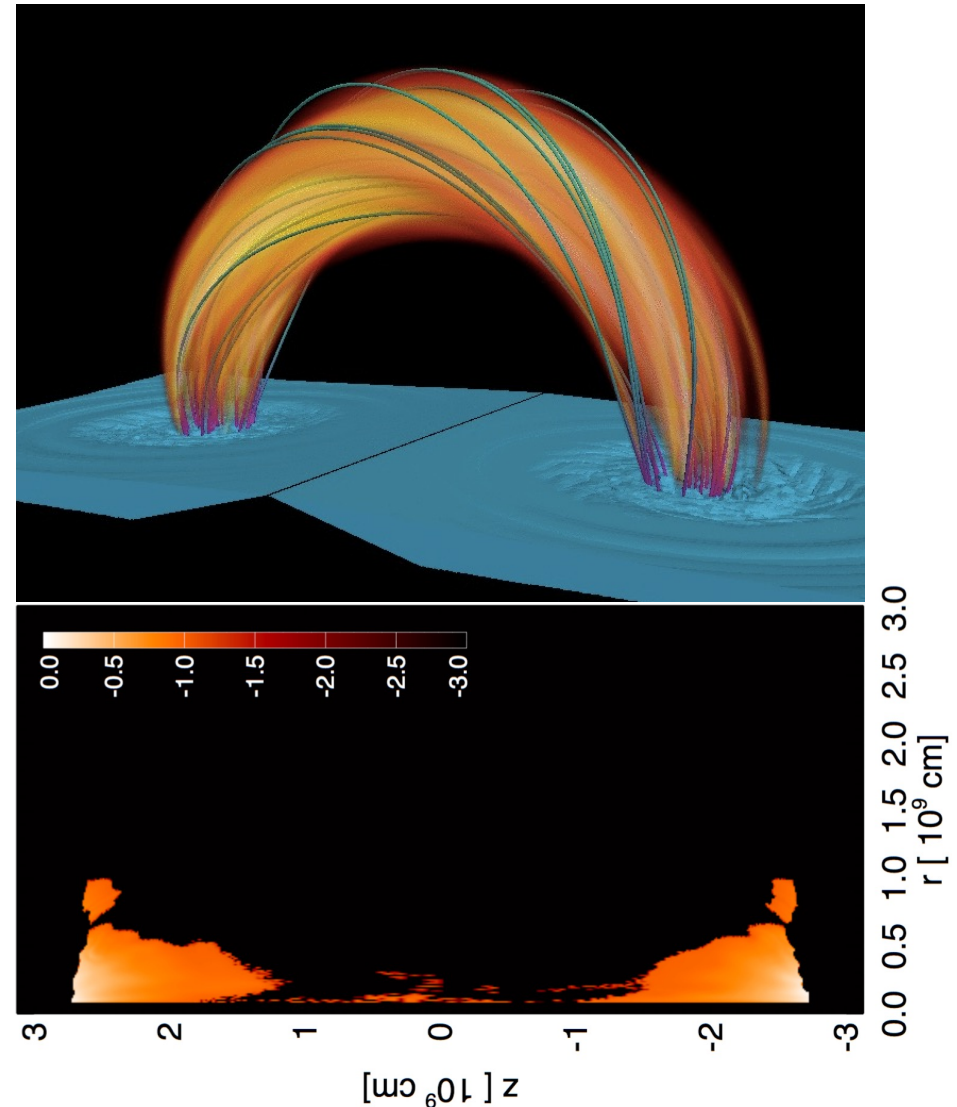
(Reale et al. 2016)

- HPC:
 - [384x256x768] cells
 - 16000/32000cores
 - ~10 Mhrs on CINECA/FERMI, Italy

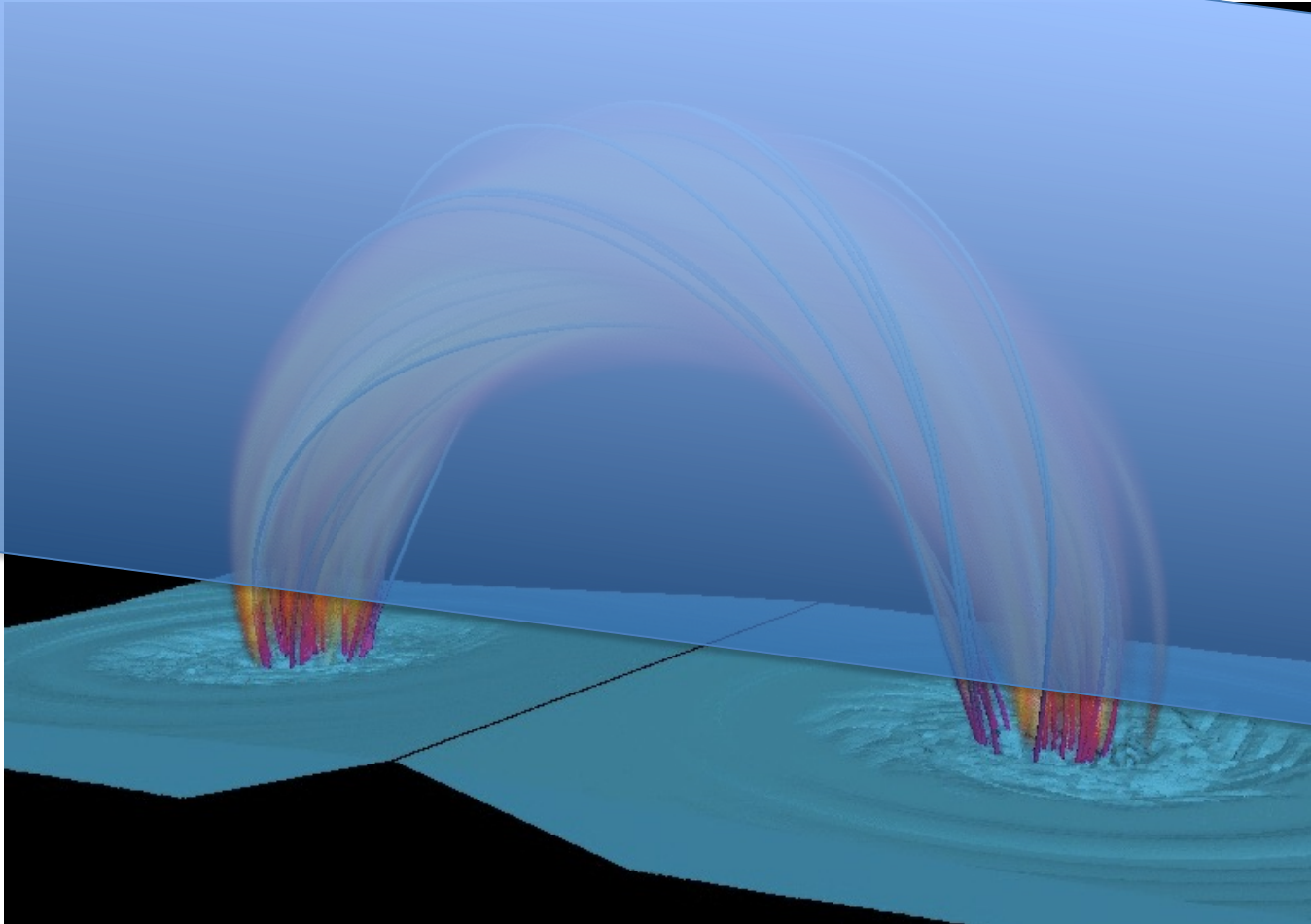


Final heating distribution

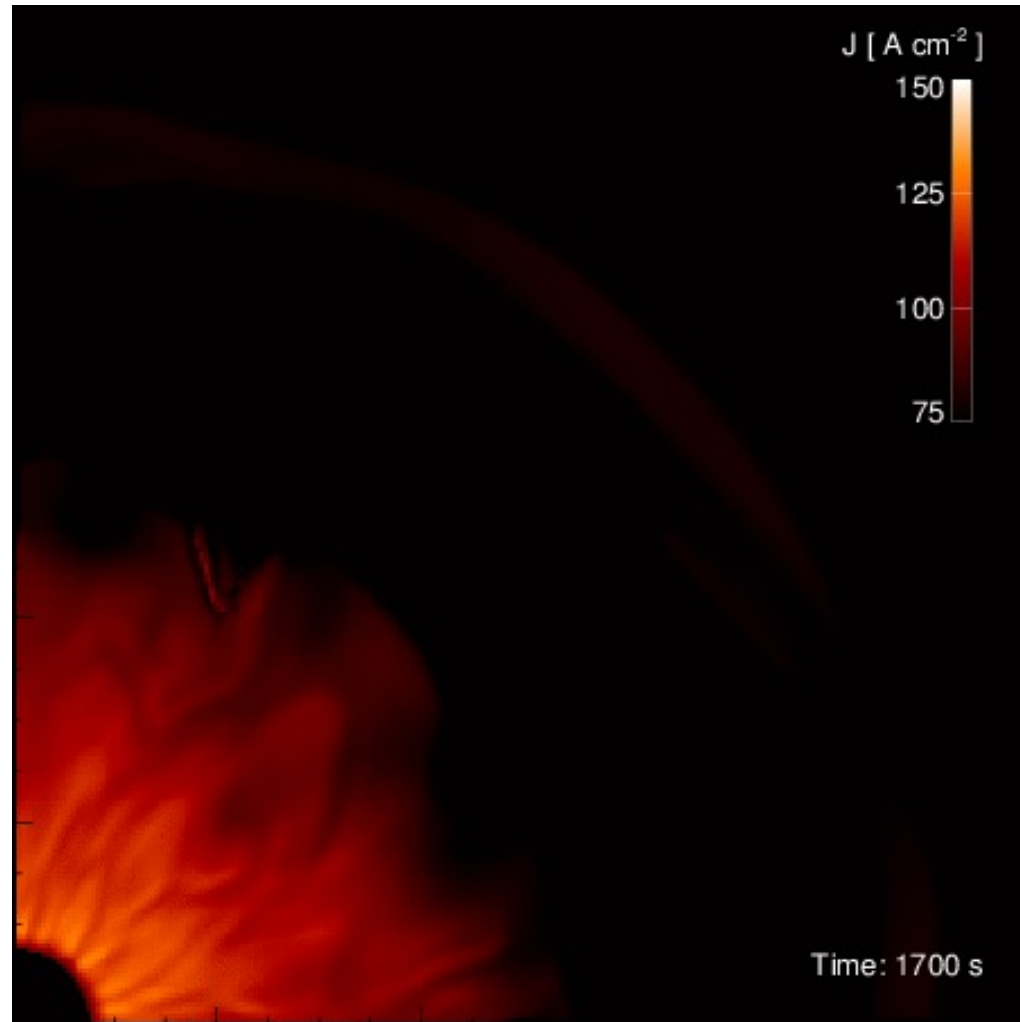
- Larger currents and heating released mostly at the footpoints (*where the magnetic field is more intense because of TAPERING*)



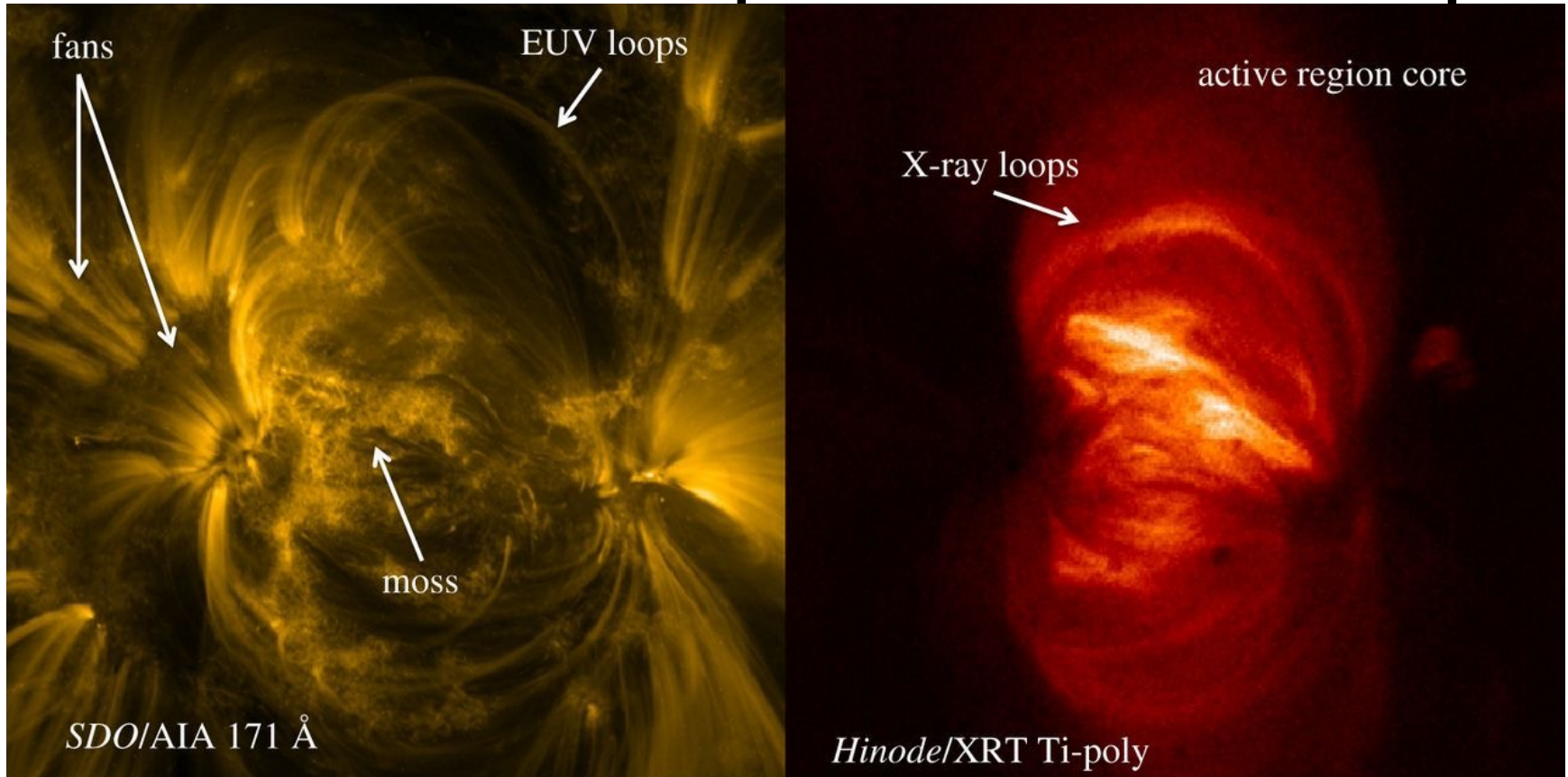
The footpoints are the key



Dissipation of currents



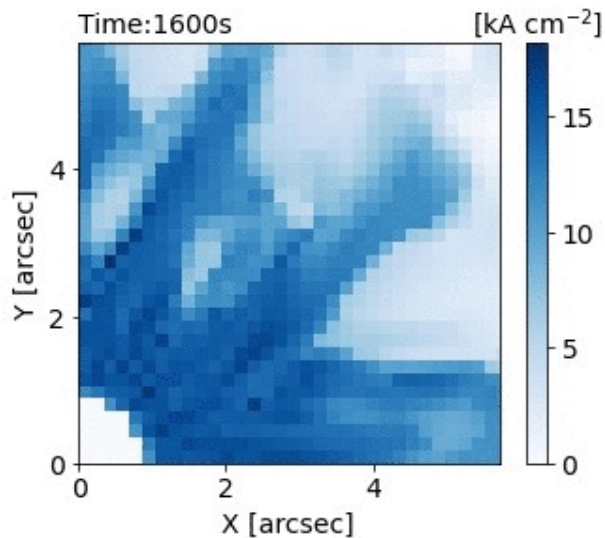
In some EUV lines (e.g., *FeIX 171A*) we observe footpoints of hot loops



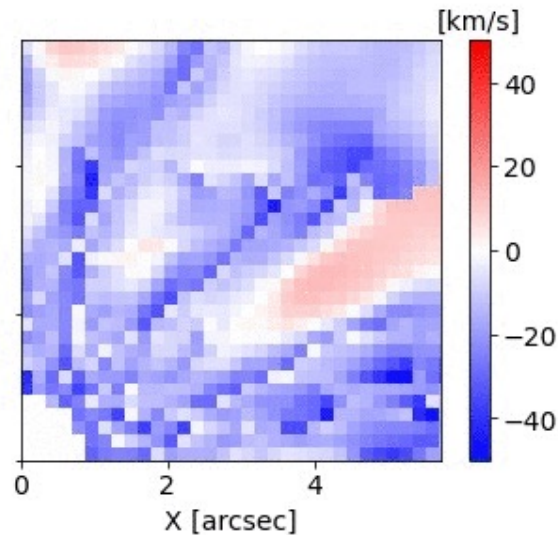
Schmelz & Winebarger 2015

High resolution spectroscopy can monitor the fine structure of currents!

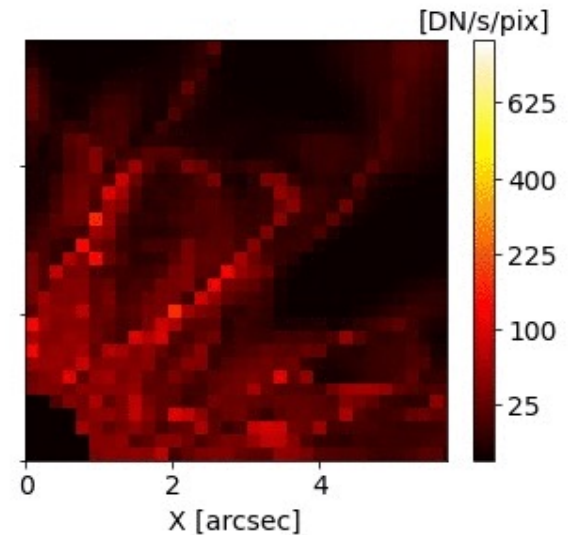
Current density



FeIX 171a: Doppler shift



FeIX 171a: Intensity



Pixel size: 0.17''

Conclusions

- Field stressing by braiding/twisting can provide loop heating
- Higher current dissipation at the footpoints because of loop tapering
- We show that EUV high resolution spectroscopy (FeIX 171A) can track the fine structure of currents and dissipation