Gravity and nonequilibrium thermodynamics: the origin of evolution equations Máté Pszota in collaboration with

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Overview

- A nonequilibrium thermodynamic theory incorporating gravity
- Evolution equation and flow-frames
- Modified gravity and its applications
- Connection to quantum mechanics \Leftrightarrow

Gravitation and thermodynamics

- Gravitational field coupled to thermodynamic equations of state
	- The internal energy separately contains the field and interaction energy:

$$
u = Ts - pv + \mu = e - \varphi - \frac{\nabla \varphi \cdot \nabla \varphi}{8\pi G \rho}
$$

- Hydrodynamic framework, extensivity is conserved
- Simple heuristic derivation or Liu's procedure for the evolution equations
- The appearance of functional derivatives without variational principles, the classical holographic property and the emergence of quantum phenomena are unexpected general consequences of the II. law of thermodynamics

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VÁN and ABE, Physica A: Statistical Mechanics and its Applications, 588:126505 VÁN, Physics of Fluids 35, 057105 (2023)

Gravitation and thermodynamics

Balances of mass, momentum and internal energy are constraints:

 $\dot{\rho}+\rho\mathbf{\nabla}\cdot\vec{v}=0$ $\rho\dot{\vec{v}} + \nabla\cdot\textbf{P} = 0$ $\rho\dot{e}+\nabla\cdot\vec{q}=-\mathbf{P}:\nabla\vec{\nu}$

 \textcirc II. law of thermodynamics:

 $\rho \dot{s} + \nabla \cdot \vec{j} \geq 0$

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Thermodynamic derivation

$$
\rho \dot{s} + \nabla \cdot \left[\frac{1}{T} \left(\mathbf{q} + \frac{1}{4\pi G} \dot{\varphi} \nabla \varphi \right) \right] =
$$

$$
\left(\mathbf{q} + \frac{\dot{\varphi}}{4\pi G} \nabla \varphi \right) \cdot \nabla \left(\frac{1}{T} \right) + \frac{\dot{\varphi}}{4\pi GT} \left(\Delta \varphi - 4\pi G \rho \right) -
$$

$$
- \left[\mathbf{P} - p \mathbf{I} - \frac{1}{4\pi G} \left(\nabla \varphi \nabla \varphi - \frac{1}{2} \nabla \varphi \cdot \nabla \varphi \mathbf{I} \right) \right] : \frac{\nabla \mathbf{v}}{T} \ge 0
$$

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Thermodynamic forces and fluxes

Constitutive relations

$$
q + \frac{\dot{\varphi}}{4\pi G} \nabla \varphi = \lambda \nabla \left(\frac{1}{T}\right) = -\lambda_F \nabla T
$$

$$
\dot{\varphi} = l_1 \left(\frac{\Delta \varphi}{4\pi G} - \rho\right) + l_{12} \nabla \cdot \nu
$$

$$
\frac{1}{3} \operatorname{Sp}(P) - p + \frac{\nabla \varphi \cdot \nabla \varphi}{24\pi G} = l_{21} \left(\frac{\Delta \varphi}{4\pi G} - \rho\right) + l_2 \nabla \cdot \nu
$$

$$
P - \operatorname{Sp}(P) \frac{I}{3} - \frac{1}{4\pi G} \left(\nabla \varphi \nabla \varphi - \frac{1}{3} (\nabla \varphi)^2 I\right) = -\eta \left(\nabla \nu + (\nabla \nu)^T - \frac{2}{3} \nabla \cdot \nu I\right)
$$

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Flow-frames and evolution equations

 ³ℝ ⁼ 0 = ,

$$
\hat{p} = c \int_{\mathbb{R}^3} p^{\alpha} p^{\beta} f \frac{dp}{p^0} \qquad p^{\mu} = \left(\sqrt{(mc)^2 + |p|^2}, p \right)
$$

Eckart frame – particles: ◈

$$
N^{\alpha} = n_E u_E^{\alpha}, \qquad T^{\alpha\beta} = \frac{e_E}{c^2} u_E^{\alpha} u_E^{\beta} + P_E^{\alpha\beta} + W_E^{\alpha} u_E^{\beta} + W_E^{\beta} u_E^{\alpha}
$$

- The choice of a frame becomes a key problem in relativistic generalisations of theories involving hydrodynamics
- Stability analyses are a topic of current research avenues \diamondsuit
- Idea: the flow is derived by assuming stability (Kovtun)

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L. Gavassino, M. Antonelli, and B. Haskell, Phys. Rev. D 102, 043018 (2020) P. Kovtun, J. High Energ. Phys. 2019, 34 (2019)

Flow-frames comparison

Landau-Lifshitz frame – energy:

$$
N^{\alpha} = n_L u_L^{\alpha} - v^{\alpha}, \qquad T^{\alpha\beta} = \frac{e_L}{c^2} u_L^{\alpha} u_L^{\beta} + P_L^{\alpha\beta}
$$

Eckart frame – particles:

$$
N^{\alpha} = n_E u_E^{\alpha}, \qquad T^{\alpha\beta} = \frac{e_E}{c^2} u_E^{\alpha} u_E^{\beta} + P_E^{\alpha\beta} + W_E^{\alpha} u_E^{\beta} + W_E^{\beta} u_E^{\alpha}
$$

 $\hat{\varphi}$ β -, or thermometer flow-frame – temperature four-vector:

$$
\beta^{\alpha} = \beta(n_T, e_T)u^{\alpha}, \qquad T^{\alpha\beta} = \frac{e_T}{c^2}u_T^{\alpha}u_T^{\beta} + P_T^{\alpha\beta} + W_T^{\alpha}u_T^{\beta} + W_T^{\beta}u_T^{\alpha}
$$

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Akihiko Monnai, Phys. Rev. C 100, 014901, 2019 Byung-Hoon Hwang, Fluid Dyn. Res. 54 015501, 2022

Gravitation and thermodynamics

Cross-effect can arise between 2 of the resulting constitutive equations: \diamondsuit

$$
\dot{\varphi} = l_1 \left(\frac{\Delta \varphi}{4\pi G} - \rho \right) + l_{12} \nabla \cdot \vec{v}
$$

$$
\frac{1}{3} \operatorname{Sp}(\mathbf{P}) - p + \frac{(\nabla \varphi)^2}{24\pi G} = l_{21} \left(\frac{\Delta \varphi}{4\pi G} - \rho \right) + l_2 \nabla \cdot \vec{v}
$$

Dissipative field equation for gravity: \diamondsuit

$$
\frac{\partial \varphi}{\partial t} = \frac{l^2}{\tau} \left(\Delta \varphi - 4\pi G \rho + \frac{l_{12}}{6det(L)} (\nabla \varphi)^2 \right), \text{ where } K := \frac{l_{12}}{6det(L)}
$$

Stationary solution is a modified Poisson's equation: \diamondsuit

 $\Delta \varphi = 4 \pi G \rho + K (\nabla \varphi)^2$

Vacuum solution

$$
\begin{aligned}\n\text{Lef} \cdot g(r) &= -\frac{1}{Kr + Cr^2} \\
\text{Lef} \cdot g(r) &= \frac{1}{K} \ln \left(\frac{r}{K + Cr} \right) + \varphi_0 \\
\text{Lef} \cdot g(r) &= \frac{1}{GM} \\
\text{Lef} \cdot g(r) &
$$

ABE and VÁN, Symmetry 2022, 14(5), 1048

Why do we need Dark Matter?

Astronomical observations suggest an abundance of gravitating matter, more than what can be accounted for:

The dynamics of galaxies and galaxy clusters

Large-scale structure of the universe and CMB-anisotropies

Gravitational lensing

Own work, based on G. Bertone and T. M. P. Tait. Nature, 562(7725):51–56

Is modified gravity better than dark matter?

Regularities in local scaling relations between baryons and dynamics in galaxies pose a \Leftrightarrow problem for dark matter descriptions, but explain a broad range of observations

Right: Baryonic Tully-Fisher relation Stacy McGaugh 14

Left: Bullet Cluster (1E 0657-56), X-ray (pink): NASA/CXC/CfA/M.Markevitch et al.; Optical: NASA/STScI; Magellan/U.Arizona/D.Clowe et al.; Lensing Map (blue): NASA/STScI; ESO WFI; Magellan/U.Arizona/D.Clowe et al.

Galactic rotation curves

- Thermodynamic gravity applied to observed density distributions
- Comparable results to Dark Matter or MOND ◈ approaches (DOI 10.1016/j.dark.2024.101660):

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NGC 3198 rotation curves

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Connection to quantum mechanics

Korteweg fluids – constitutive state variables: (e, ∇e , ρ , $\nabla \rho$, $\nabla^2 \rho$, \mathbf{v} , $\nabla \mathbf{v}$) \diamondsuit

$$
\text{Entropy production: } \rho \dot{s} + \nabla \cdot \mathbf{J} = \mathbf{q} \cdot \nabla \left(\frac{1}{T}\right) - \left[\mathbf{P} - p\mathbf{I} - \frac{\rho^2}{2} \left(\nabla \cdot \frac{\partial s}{\partial \nabla \rho} \mathbf{I} + \nabla \frac{\partial s}{\partial \nabla \rho}\right)\right] : \frac{\nabla \mathbf{v}}{T} \ge 0
$$

- Rigorous methods are needed: Liu's procedure, as before ◈
- The resulting pressure has classical holographic property: ◈

$$
\nabla \cdot \mathbf{P}_K = \rho (\nabla \phi + T \nabla s), \quad \text{where} \quad \phi = \frac{\partial \rho u}{\partial \rho} - \nabla \cdot \frac{\partial (\rho u)}{\partial \nabla \rho} = \delta_\rho(\rho u) \Big|_{\rho s}
$$

Additivity leads to Bohm potential: $\phi_Q(\rho, \nabla \rho, \nabla^2 \rho) = -\frac{\hbar^2}{2m}$ $\Delta\sqrt{\rho}$, recover the Schrödinger equation \diamondsuit $2m²$ $\overline{\rho}$

These fluids thus allow the modelling of quantum-classical transition ◈

Connection to quantum mechanics

- The derivation of evolution equations can be inverted: QM equations can be obtained using this hydrodynamic connection from thermodynamic principles
- Using Liu's procedure, the II. law alone produces holography and dynamics ◈
- QFT, GR can have fluid interpretation: Jackiw et al. (JP A, 2004), Biró and Ván \Leftrightarrow (Foundation of Phys, 2015)
- The connection of thermodynamics and gravity has also been explored by Jacobson (PRL 1995) and Verlinde (JHEP 2011)

Summary

- Particle approach to DM has been a long and unfruitful search, while alternatives keep being developed
- Models using hydrodynamic descriptions have potential for describing galactic dynamics
- Thermodynamic gravity within the hydrodynamic framework offers a potentially powerful method for the description of dynamics on cosmological scales
- Hydrodynamic description greatly benefits from a rigorous foundation and can be a thermodynamic link between QM and gravity

Thank you for your attention!

- Main references: \Leftrightarrow
	- Peter Ván and Sumiyoshi Abe. Emergence of extended Newtonian gravity from thermodynamics. Physica A: Statistical Mechanics and its Applications, 588:126505, 2022
	- P. Ván, Holographic fluids: A thermodynamic road to quantum physics, Physics of Fluids 35, 057105, 2023
	- L. Gavassino, M. Antonelli, and B. Haskell, When the entropy has no maximum: A new perspective on the instability of the first-order theories of dissipation, Phys. Rev. D 102, 043018, 2020
	- ◈ S. Abe and P. Ván. Crossover in Extended Newtonian Gravity Emerging from Thermodynamics. Symmetry, 14(5), 2022
	- M. Pszota and P. Ván. Field equation of thermodynamic gravity and galactic rotational curves. DOI \Diamond 10.1016/j.dark.2024.101660, accepted in Physics of the Dark Universe, (arXiv:2306.01825v3)

Cosmological and galactic simulations

- Evolution of ΛCDM halos in the fluid approximation (left)
- N-body simulation (right)
- Periodic boundary conditions
- Scalar Field DM, Bose-Einstein condensate (superfluid) description

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FOIDL, RINDLER-DALLER, and ZEILINGER, PHYS. REV. D 108, 043012 (2023)