Simulations of the polarized Galactic sky for the Epoch of Reionization observations

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## EoR with the 21 cm line: the challenge of foregrounds

- Extragalactic Point Sources (PS) radio galaxies, AGNs, ..
- Galactic and Extragalactic free-free
bremsstrahlung radiation from electron-ion collisions
- Galactic synchrotron (dominant foreground) cosmic ray electrons interacting with the galactic magnetic field.
e.g. Santos et al 2005, Jelic et al 2008, Geil et al 2011


EPOCH OF REIONIZATION

EXTRAGALACTIC FOREGROUNDS

GALACTIC
FOREGROUNDS
credit: LOFAR

Synchrotron is linearly polarised.

## Foreground cleaning

Heritage form CMB analysis (more complex with the 3D 21 cm signal)

$$
T(\nu, \hat{\mathbf{n}})=\sum_{k=1}^{N_{\mathrm{fg}}} f_{k}(\nu) S_{k}(\hat{\mathbf{n}})+T_{\text {cosmo }}(\nu, \hat{\mathbf{n}})+T_{\text {noise }}(\nu, \hat{\mathbf{n}})
$$

- Blind subtraction: spectral polynomial fitting (LOS), principal component analysis (PCA), ..
- model the foregrounds with a physically motivated functions (log-polynomial in Global signal analysis)
foreground sources are expected to be spectrally smooth, while the cosmological EoR signal is expected to fluctuate rapidly with frequency


## Foreground Avoidance

Alternative strategy: use data only from the EoR window (in the cylindrical $k_{\perp}-k_{\|}$Fourier plane) that are not contaminated by foregrounds


- Smooth foregrounds are expected at small $k_{\|}$
- going to higher $k_{\perp}$, due to the instrument response, foregrounds leak out to higher $k_{\|}$(foreground wedge)

Liu et al. (2014)

## Polarised Synchrotron emission



## Why is it important?

- non trivial frequency structure
- can leak into the unpolarized part
- can contaminate the 21 cm analysis:

1. preventing component separation to work properly
2. mimic the high $k_{\|}$EoR emission, scattering power in otherwise clean EoR window Moore et al. (2013)

## Synchrotron generalities

- Depends on $B_{\perp}$ to the LOS modulated by the density of cosmic electrons
- Diffuse polarised emission:

$$
\begin{aligned}
& P=\Pi_{0} I e^{2 i \phi} \text { with } \phi=\phi_{0}+\psi \lambda^{2} \text { faraday rotation } \\
& \text { given by } B_{\|} \text {and the presence of thermal electrons } \\
& \qquad \psi \propto \int_{\mathrm{LOS}} n_{e} B_{\|} d r
\end{aligned}
$$

At Eor frequencies P simulations are difficult:

- lack of correlation with total intensity
- not a lot of polarised data at low frequencies
- depolarisation effects prevent extrapolation from higher frequencies


## Simulation strategy

Spinelli, Bernardi, Santos (2018) Use RM-synthesis framework ( $\psi$ and $\lambda^{2}$ as a Fourier pair) Bretjens\& Bruyn (2005) Heald, Brown \& Edmonds (2009):

- full-sky gaussian $\mathrm{Q}, \mathrm{U}$ maps in $\psi$ space with specific power spectrum:

$$
\begin{gathered}
\tilde{Q}(\psi, \hat{\mathbf{n}})=\sum_{\ell m} \tilde{q}_{\ell m}(\psi) Y_{\ell m}(\hat{\mathbf{n}}) \\
\left\langle\tilde{q}_{\ell m}(\psi) \tilde{q}_{\ell^{\prime} m^{\prime}}^{*}(\psi)\right\rangle=(2 \pi)^{2} A(\psi) \ell^{-\alpha(\psi)} \\
\left\langle\tilde{q}_{\ell m}(\psi) \tilde{q}_{\ell m}^{*}\left(\psi^{\prime}\right)\right\rangle \propto \rho(\Delta \psi, \ell)
\end{gathered}
$$

We use MWA data to constrain free parameters (from Bernardi et al 2013 but we can use other data)

- transform back to frequency space using the Fourier relation between $\psi$ and $\lambda^{2}$

$$
Q\left(\lambda^{2}, \hat{\mathbf{n}}\right)=\int \tilde{Q}(\psi, \hat{\mathbf{n}}) e^{2 \pi \lambda^{2} \psi} d \psi
$$

## Full-sky polarized maps

- complex frequency structure
- a worst case scenario for depolarization (MWA @189 MHz)
- no ionosphere (yet)
- simulations in the range $50-200 \mathrm{MHz}$ publicly available at UWC-CRC drive folder
- option to select frequency and spatial resolution




## Some applications

- foreground avoidance: an example from PAPER
$21-\mathrm{cm}$ power spectrum polarization contamination EoR
$120-180 \mathrm{MHz}$
- foreground cleaning: (simulated) EDGES 21-cm Global signal analysis in presence of polarized foregrounds
Cosmic Dawn $50-100 \mathrm{MHz}$

PAPER: Karoo desert (South Africa)


EDGES: Murchison
Radio-astronomy Observatory
(Western Australia)


## "Cleaning" in Global-signal studies for Cosmic Dawn

$$
T(t, \nu) \propto \int_{\Omega} T_{s k y}\left(t, \nu, \hat{n}^{\prime}\right) A\left(\nu, \hat{n}^{\prime}\right) d \hat{n}^{\prime}+T_{N}
$$



Bernardi et al. (2016)

$$
T(t, \nu)=T_{f}(t, \nu)+T_{\mathrm{HI}}(\nu)
$$

Foregrounds:
$\log _{10} T_{f}\left(\nu_{j}\right)=$
$\sum_{n=0}^{N} p_{n}\left[\log _{10}\left(\frac{\nu_{j}}{\nu_{0}}\right)\right]^{n}$
Signal:
$T_{H I}\left(\nu_{j}\right)=$
$A_{H I} e^{-\frac{\left(\nu_{j}-\nu_{H I}\right)^{2}}{2 \sigma_{H I}^{2}}}$
or the flattened Gaussian Bowman et al. (2018)

## What if there is a polarized sky contamination?

A polarized signal can leak into the measurements but the analysis assumes it is weak.
Can EDGES absorption profile have local astrophysical origin?

- observations every h from $l s t=0$ to $l s t=8$
- frequency range: $50-100 \mathrm{MHz}$
- beam model Dowell (2011)
- Bayesian pipeline for foreground removal: HIBayes Zwart et al. (2016) based on MultiNest Feroz et al. (2009)


Spinelli, Bernardi, Santos in prep

## Contamination

- Two types of simulations:

1. all $\phi$
2. $\phi<5 \mathrm{rad} / \mathrm{m}^{2}$

- 1000 sims for each case
- high contamination of the order of $\sim 200 \mathrm{mK}$
- single dipole will measure $I+Q(I-Q)$
- I: 5th degree log-polynomial
- Q: typical cases from our simulations




## Possible scenarios

- current simulation typical Gaussian:
the measured profile is unlikely to be due to (our) astrophysical foreground
- confirmed EDGES signal:
$+Q$ (solid) and $-Q$ (dashed) contamination could bias the result and shift the deep up and down




## Conclusions

- polarized synchrotron can in principle be one of the main challenges to unveil the 21 cm signal both for foreground avoidance and foreground cleaning
- we developed full-sky simulations based on MWA data in the range $50-200 \mathrm{MHz}$.
- we have investigated the impact of a (unaccounted for) polarized sky component in global-signal analysis for Cosmic Dawn:

1. EDGES results are unlikely to have a local astrophysical origin (assuming our simulated sky)
2. there is still a possible bias due to polarization in the extraction of the cosmic signal

Backup

## Foreground avoidance with PAPER

- Sky brightness I,Q,U,V through the telescope: $\left.\mathbf{s}^{\prime}(\hat{\mathbf{r}}, \nu)\right)=\mathbf{A}(\hat{\mathbf{r}}, \nu) s(\hat{\mathbf{r}}, \nu)$ (A Mueller matrix)
- $\tau=\frac{\mathbf{b} \cdot \hat{\mathbf{r}}}{c}$ geometric delay between antenna pairs
- delay-transform:

Parsons \& Backer (2009) localizes foreground emission in "delay" space

- compute $P\left(k_{\perp}, k_{\|}\right)$


$$
k_{\perp} \propto|\mathbf{b}|, k_{\|} \propto \tau
$$

Formalism (and code) adapted from Nunhokee et al 2017

## Looking for a clean window



Q, U and leaked I at $z=8.5$ for a 30 m baseline

## Polarization leakage in Global Signal Analysis






Spinelli, Bernardi, Santos in prep

## EDGES results I

Bowman et al. (2018)


## EDGES results II

Bowman et al. (2018)


## (Some of the) EDGES criticism

## Hills et al. (2018)



## Rotation Measure (RM) synthesis

## Bretjens\& Bruyn (2005) Heald, Brown \& Edmonds (2009)

Use Fourier relation between polarised surface brightness ( P ) and surface brightness per unit of Faraday depth (F)

$$
P\left(\lambda^{2}\right)=\int_{-\infty}^{+\infty} F(\psi) e^{i 2 \psi \lambda^{2}} d \psi
$$

Inverting this formula:

- only positive $\lambda$ have physical meaning
- incomplete sampling in $\lambda^{2}$

Need to define a RM transfer function (RMTF) that gives the resolution in Faraday depth:

$$
\text { FWHM } \sim\left(\Delta \lambda^{2}\right)^{-1} \text { total bandwidth }
$$

lack of sensitivity to structures extended in Faraday depth

## MWA data

G. Bernardi et al. 2013


- MWA 32 element 2400 degrees
- RM synthesis
cube of polarised images at selected faraday depth
$-50<\mathrm{RM}<+50 \mathrm{rad} \mathrm{m}{ }^{\wedge}-2$
in step of $1 \mathrm{rad} \mathrm{m}^{\wedge}-2$
RMTF $4.3 \mathrm{rad} \mathrm{m}^{\wedge}-2$
describe MWA statistical behaviour and extend it to full-sky
- CONs: fine and local structures impossible to catch
- PROs: using genuine polarisation data instead of intensity


## Characterization of MWA data

- At fixed $\psi$, the data can be approximated with a Rayleigh distribution $R(\sigma(\psi))$
- retain only maps with $S / N>2$ : the interval
$-18<\psi<+23$
- Power Spectrum reconstruction with HEALPIX (Gorski et al. 2005) and MASTER (Hivon et al. 2002)
- Fit a power law considering cosmic variance on large scale and noise on small scales (Tegmark 1997)




## Mimicking the correlations



- Correlation decreses with $\ell$
- No dependecy on $\Delta \psi$
- Residual correlation still present at high $\ell$

In the simulations:

$$
\begin{aligned}
& \vec{q}_{\ell m}=\frac{1}{\sqrt{2}} N\left(0, \Sigma_{\ell}\right)+\frac{i}{\sqrt{2}} N\left(0, \Sigma_{\ell}\right) \\
& \quad \text { with } \\
& \Sigma_{\ell}^{i j}=\rho(\ell, \Delta \psi)\left(\ell^{\alpha\left(\psi_{i}\right)+\alpha\left(\psi_{j}\right)}\right)^{1 / 2}
\end{aligned}
$$

- $N_{\psi} \times N_{\psi}$ matrix $\forall \ell$
- the model reproduce the data well for $\ell>200$.
- At lower $\ell$ more complex situation (demasking?)


## A model for the beam

$$
A(\nu, \theta, \phi)=\sqrt{\left[p_{E}(\nu, \theta) \cos \phi\right]^{2}+\left[p_{H}(\nu, \theta) \sin \phi\right]^{2}}
$$

Taylor et al. (2012), Ellingson et al. (2013), Dowell (2011)

$$
\begin{gathered}
p_{i}(\nu, \theta)= \\
{\left[1-\left(\frac{\theta}{\pi / 2}\right)^{\alpha_{i}(\nu)}\right](\cos \theta)^{\beta_{i}(\nu)}+} \\
\gamma_{i}(\nu)\left(\frac{\theta}{\pi / 2}\right)(\cos \theta)^{\gamma_{i}(\nu)}
\end{gathered}
$$




