

Extragalactic (VHE) Gamma-Ray Sources and Hadronic Models

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**National
Research
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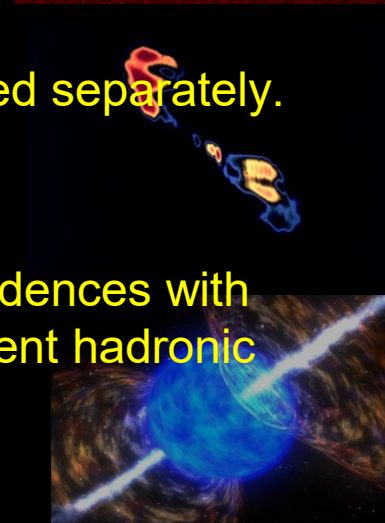
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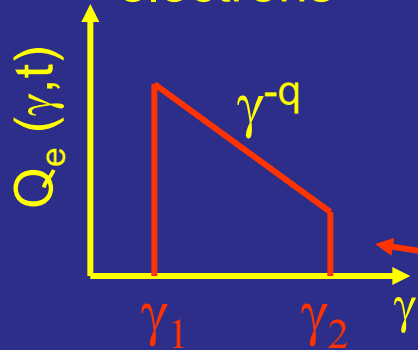
Extragalactic VHE Gamma-Ray Sources

- Blazars
 - Most numerous source class, promising neutrino sources → Focus of this talk
- Radio galaxies
 - 6 detected; probably same physics as blazars → Not covered separately.
- Gamma-ray bursts
 - 5 detected so far (incl. the BOAT, GRB 221009A); no coincidences with neutrino events so far → No compelling evidence of prominent hadronic interactions.
- Starburst galaxies
 - 2 detected (M82, NGC 253). Most likely hadronic cosmic-ray interactions. Not covered here due to lack of time.
- Low-Luminosity AGN (LLAGN)
 - Only NGC 4278 → Giulia Migliori's talk Monday



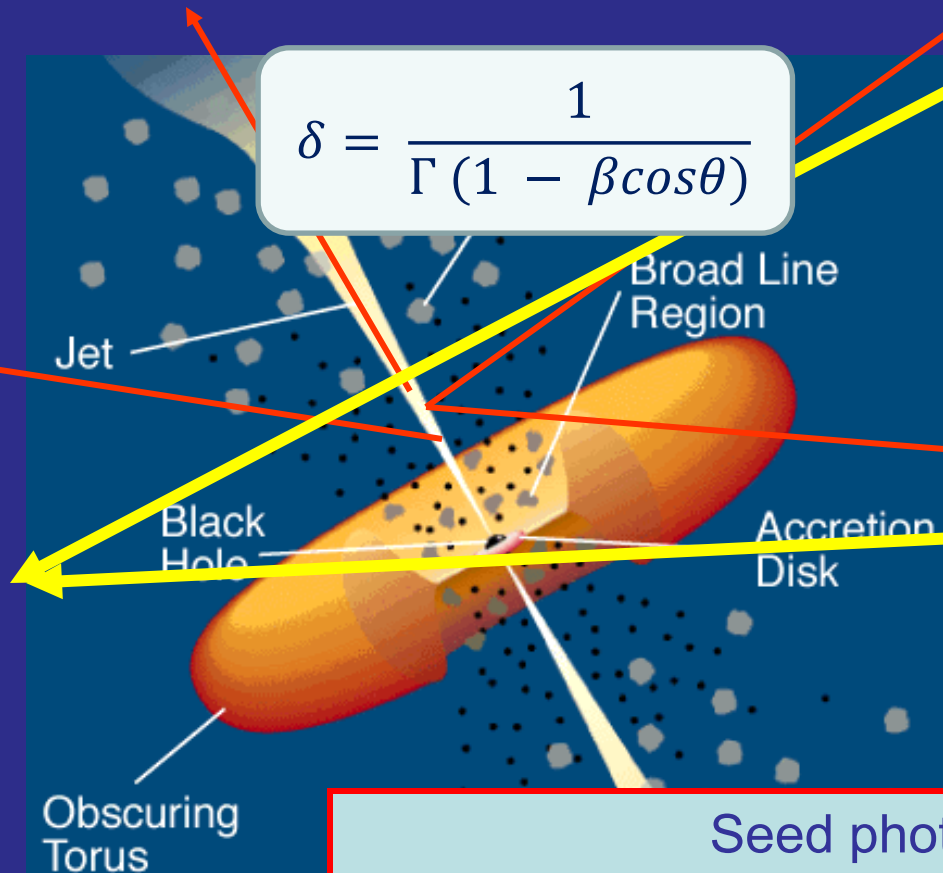
Leptonic Blazar Model

Injection,
acceleration of
ultrarelativistic
electrons

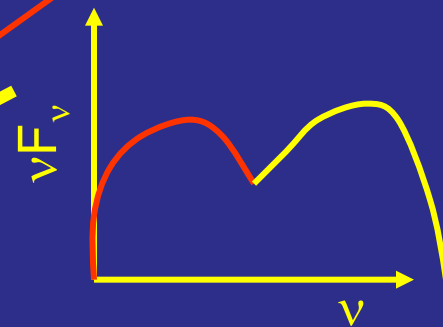


Relativistic jet outflow with $\Gamma \approx 10$

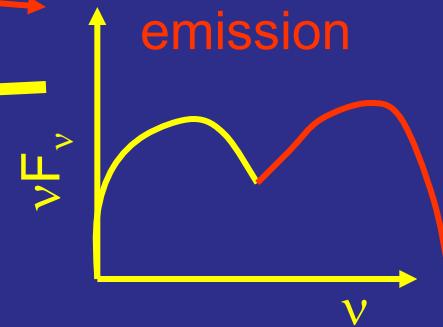
$$\delta = \frac{1}{\Gamma (1 - \beta \cos \theta)}$$



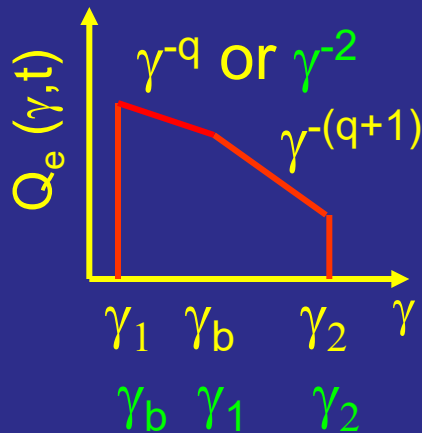
Synchrotron
emission



Compton
emission



Radiative cooling
 $\leftrightarrow$ escape $\Rightarrow$

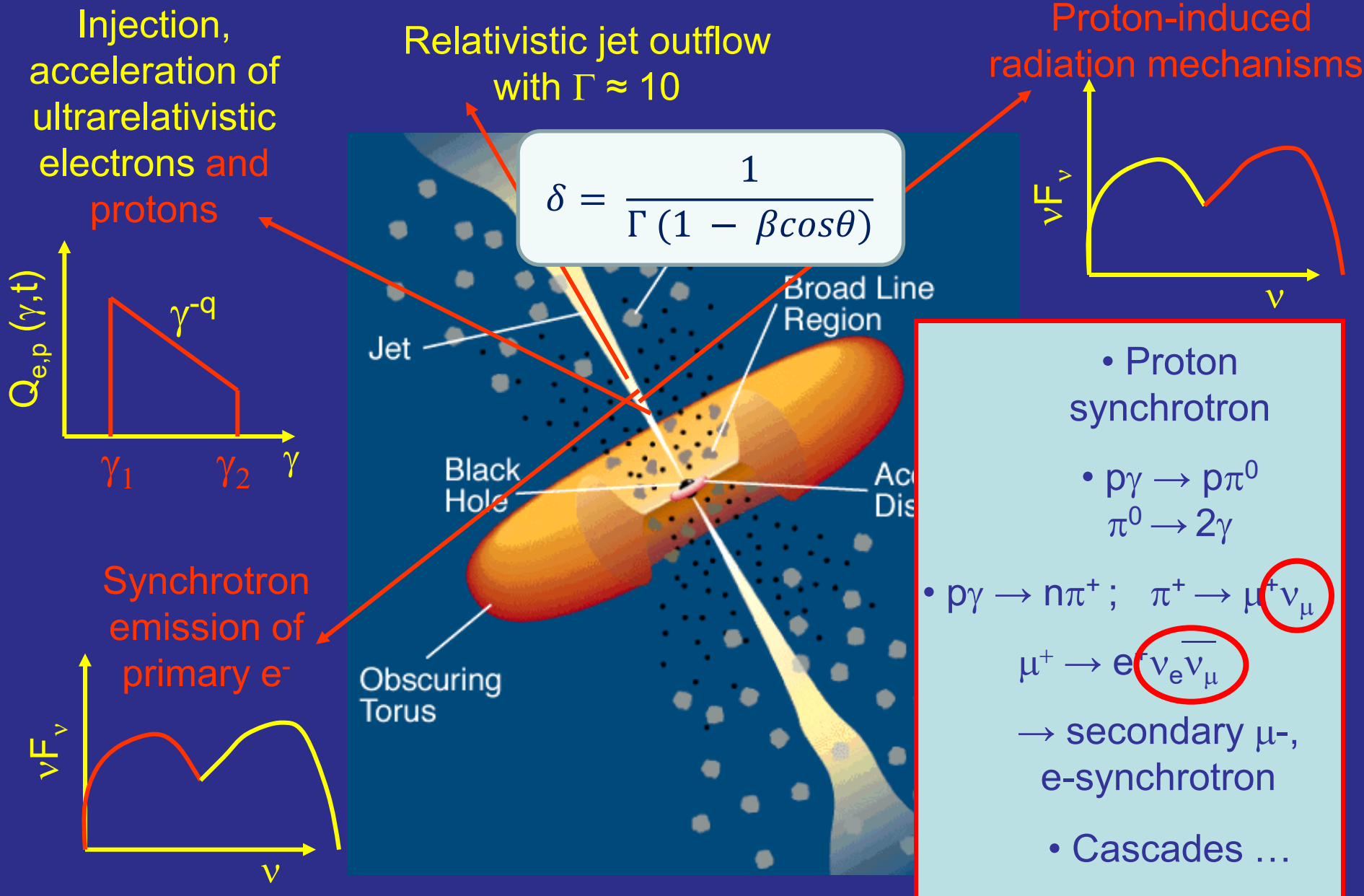


$$\gamma_b: \tau_{\text{cool}}(\gamma_b) = \tau_{\text{esc}}$$

Seed photons:

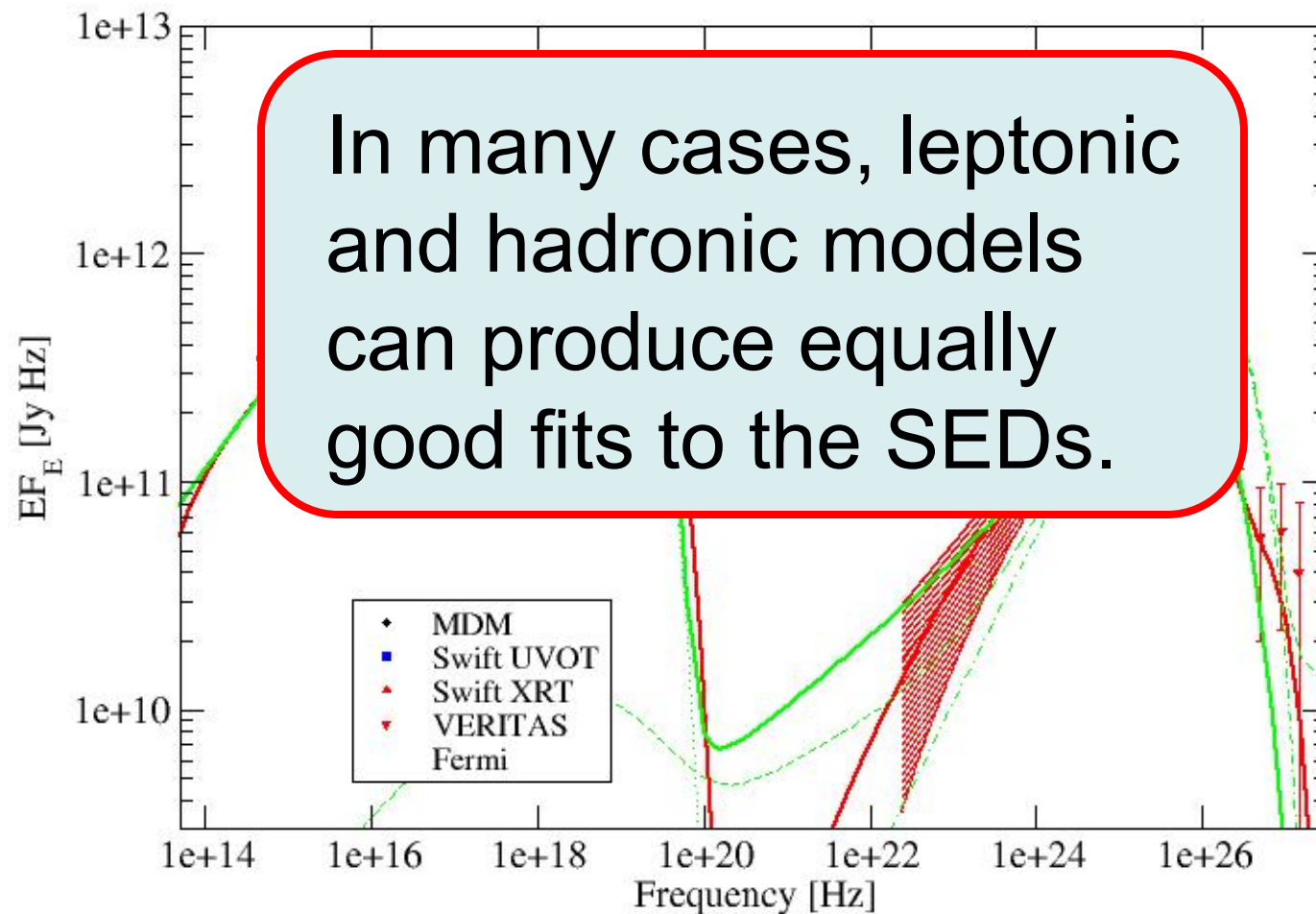
Synchrotron (within same region [SSC] or external sources, e.g., accretion disk, BLR, dust torus (EC = External Compton))

Hadronic Blazar Models



Lepto-Hadronic Model Fits to Blazar SEDs

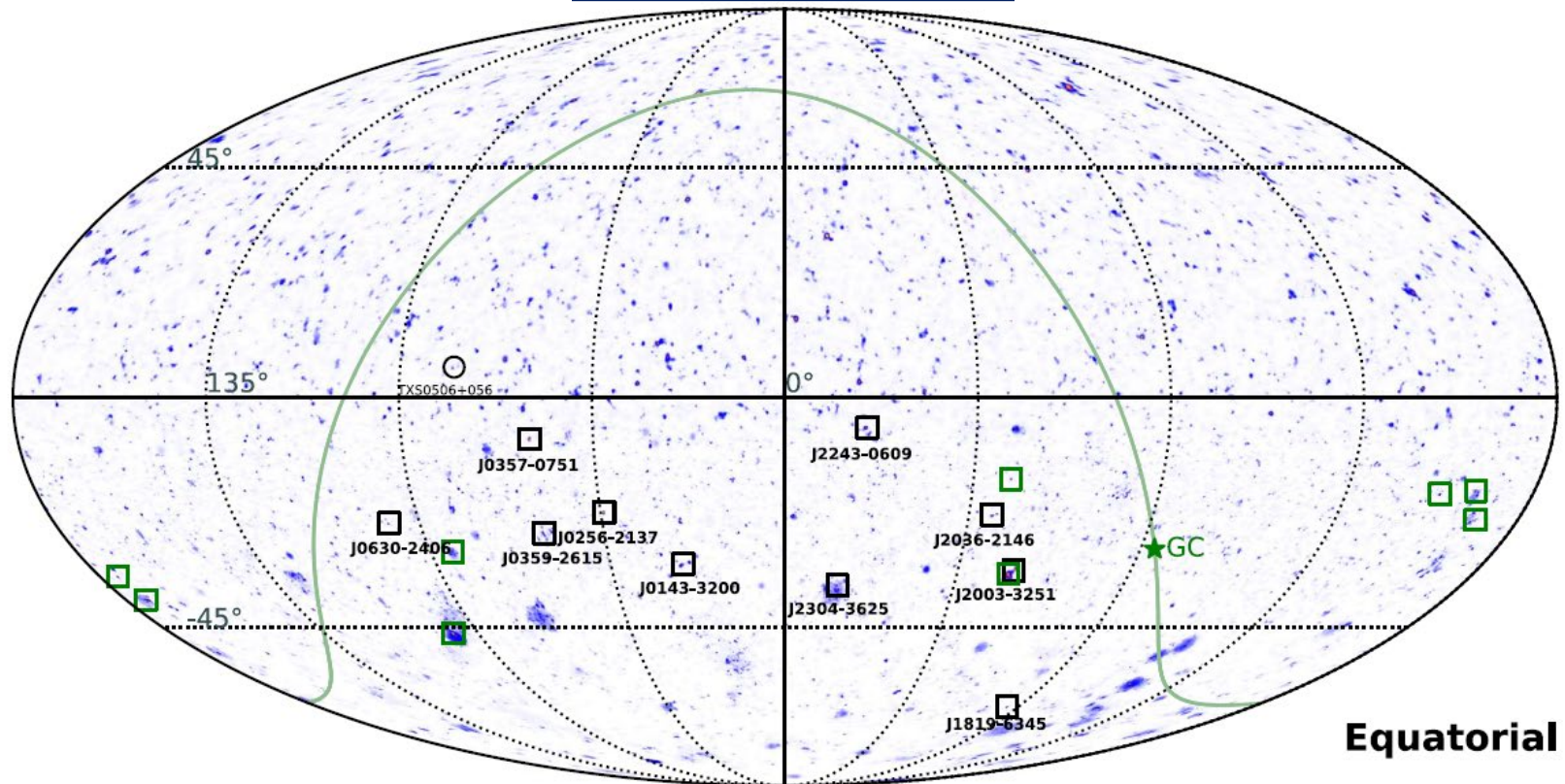
RGB J0710+591 (HBL)



Possible
Diagnostics to
distinguish:

- Neutrinos
- Polarization
- Variability

Origin of IceCube-Detected Neutrinos

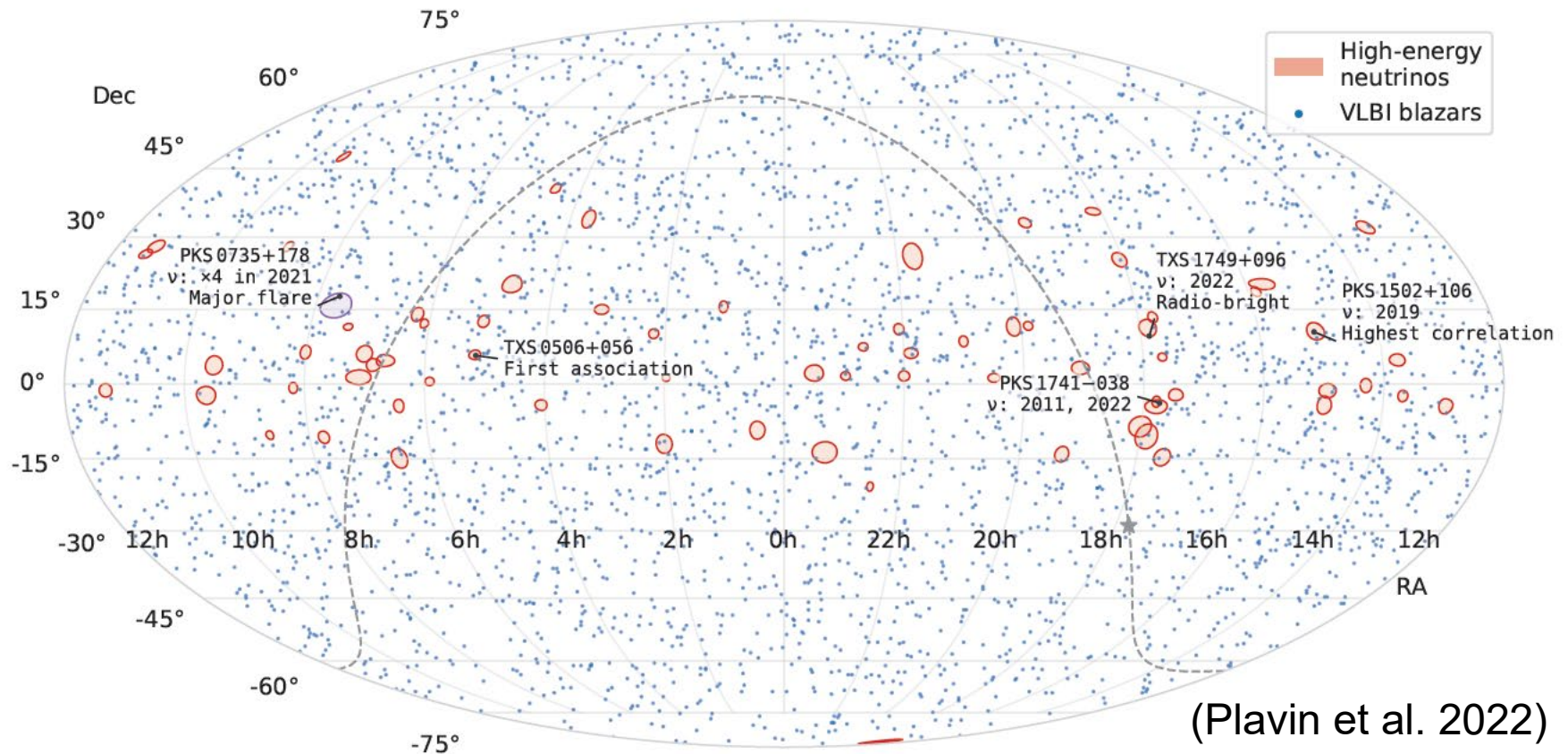


(Buson et al. 2022)

Significant correlation of IceCube neutrinos with γ -ray (Fermi-LAT) **blazars** (chance coincidence probability $p = 6 \cdot 10^{-7}$)

– but can not be responsible for all IceCube neutrinos (e.g., Murase et al. 2018)

Origin of IceCube-Detected Neutrinos



Also: Correlation with radio-loud blazars
(chance-coincidence probability of $p = 3 \cdot 10^{-4}$)
(Plavin et al. 2020, 2021, 2022)

Basic ingredients of hadronic jet multi-messenger emission models

- Particle acceleration/injection (usually modelled as instantaneous injection; 2nd-order [stochastic] acceleration included in some models/codes)
- Particle interaction / radiation modules (synchrotron, SSC, EC standard, but differences in $p\gamma$ interactions and cascades)
- Secondary particle (pions, muons, neutrons, electron-positron pair) production (MC vs. semi-analytical templates)
- Radiative transport (emission / absorption / scattering) (geometry-dependent in multi-zone / extended jet models; polarization?)
- Evolution of particle distributions (re-acceleration / cooling / escape) (usually through solutions to Fokker-Planck equations)

Basics of Neutrino Production

- $p + p (N) \rightarrow p + p/n + n_0\pi^0 + n_+\pi^+ + n_-\pi^-$ ($\sigma_{pp} \sim 0.1 \text{ mb}$)

- $p + \gamma \rightarrow p + \pi^0 \quad (2/3)$
or $n + \pi^+ \quad (1/3)$ ($\sigma_{p\gamma} \sim 0.6 \text{ mb}$)

$$\pi^0 \rightarrow 2\gamma$$

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$$\pi^- \rightarrow \mu^- + \nu_\mu$$

$$\mu^+ \rightarrow e^+ + \nu_\mu + \nu_e$$

$$\mu^- \rightarrow e^- + \nu_\mu + \nu_e$$

$$\tau = 2.55 \times 10^{-8} \text{ s}$$

$$\tau = 2.2 \times 10^{-6} \text{ s}$$

Neutrino Production in AGN Jets

$$n_{ph} \sim \frac{L_{sy}}{\delta^4 \langle \epsilon \rangle m_e c^2 4\pi R^2 c} \sim 3 \times 10^{18} \epsilon_{-6}^{-1} R_{16}^{-2} L_{sy,44} \delta_1^{-4} \text{cm}^{-3}$$

$$n_p \leq \frac{L_j}{\Gamma^2 m_p c^2 \pi R^2 c} \sim 10^4 R_{16}^{-2} \Gamma_1^{-2} L_{j,46} \text{cm}^{-3} \quad \epsilon = \frac{E_{ph}}{m_e c^2}$$

$$\frac{t_{pp}}{t_{p\gamma}} \sim \frac{n_{ph}}{n_p} \sim 3 \times 10^{14} \frac{\Gamma_1^2 L_{sy,44}}{\epsilon_{-6} \delta_1^4 L_{j,46}}$$



Photo-Pion Production Cross Section

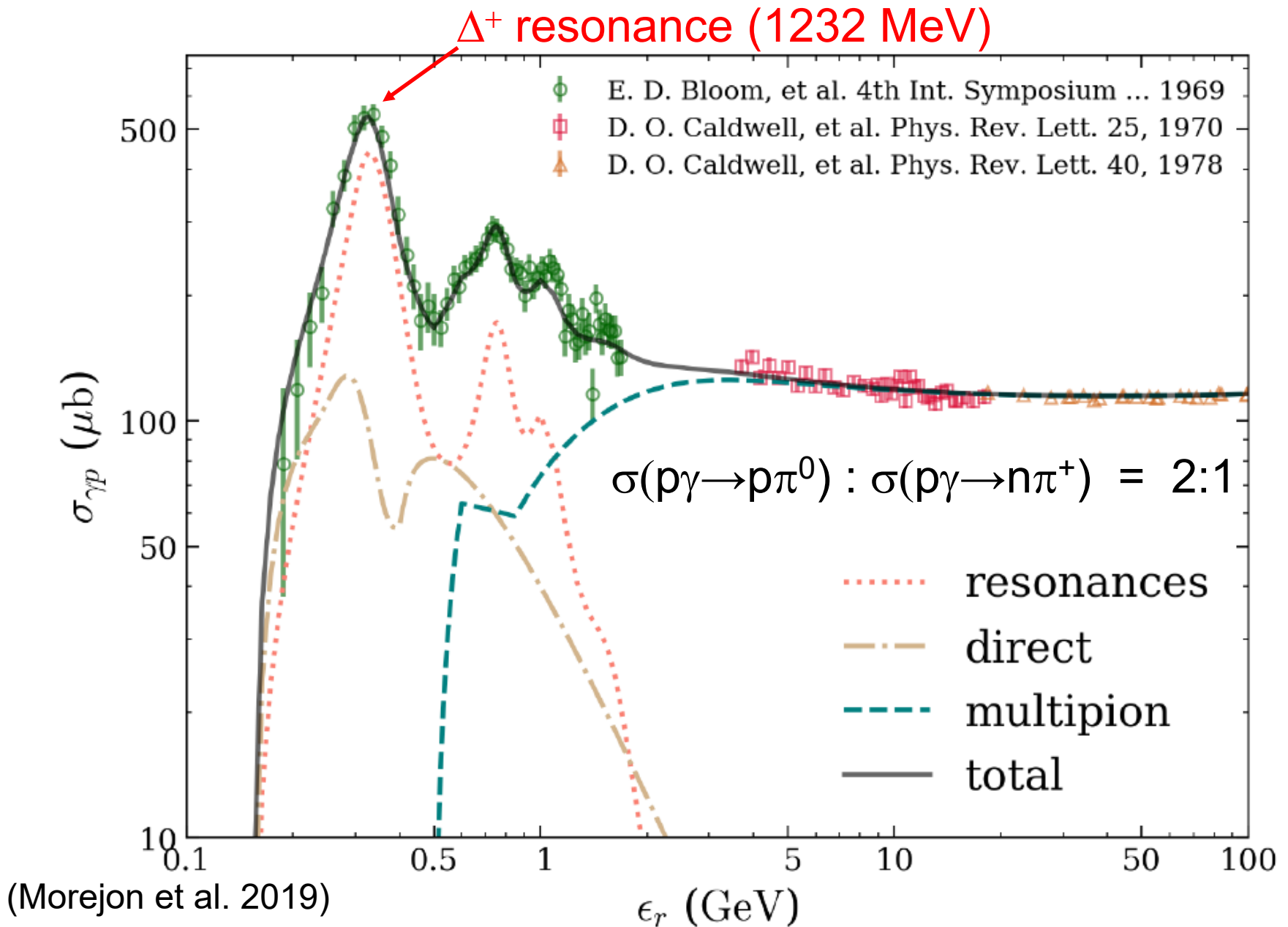
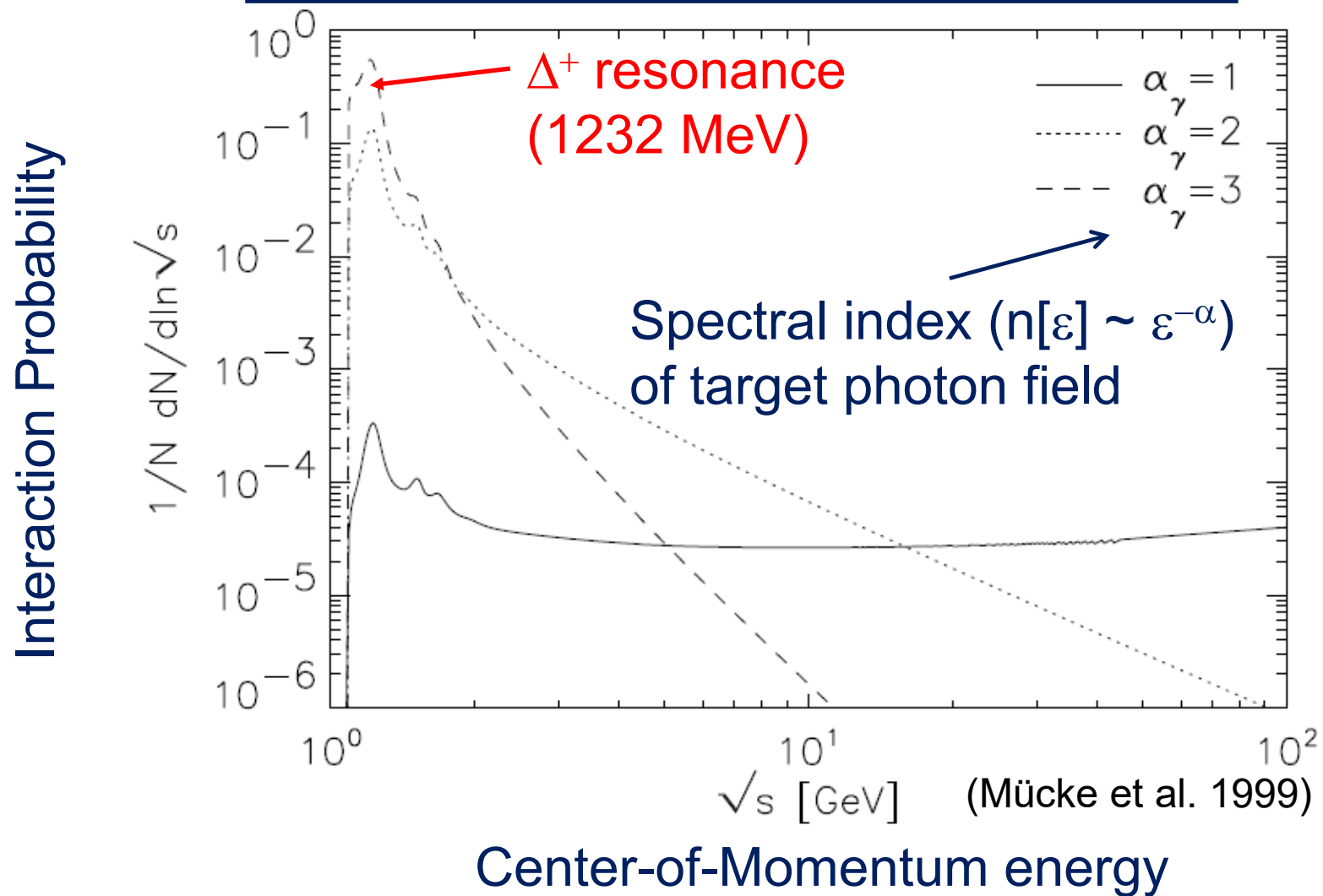
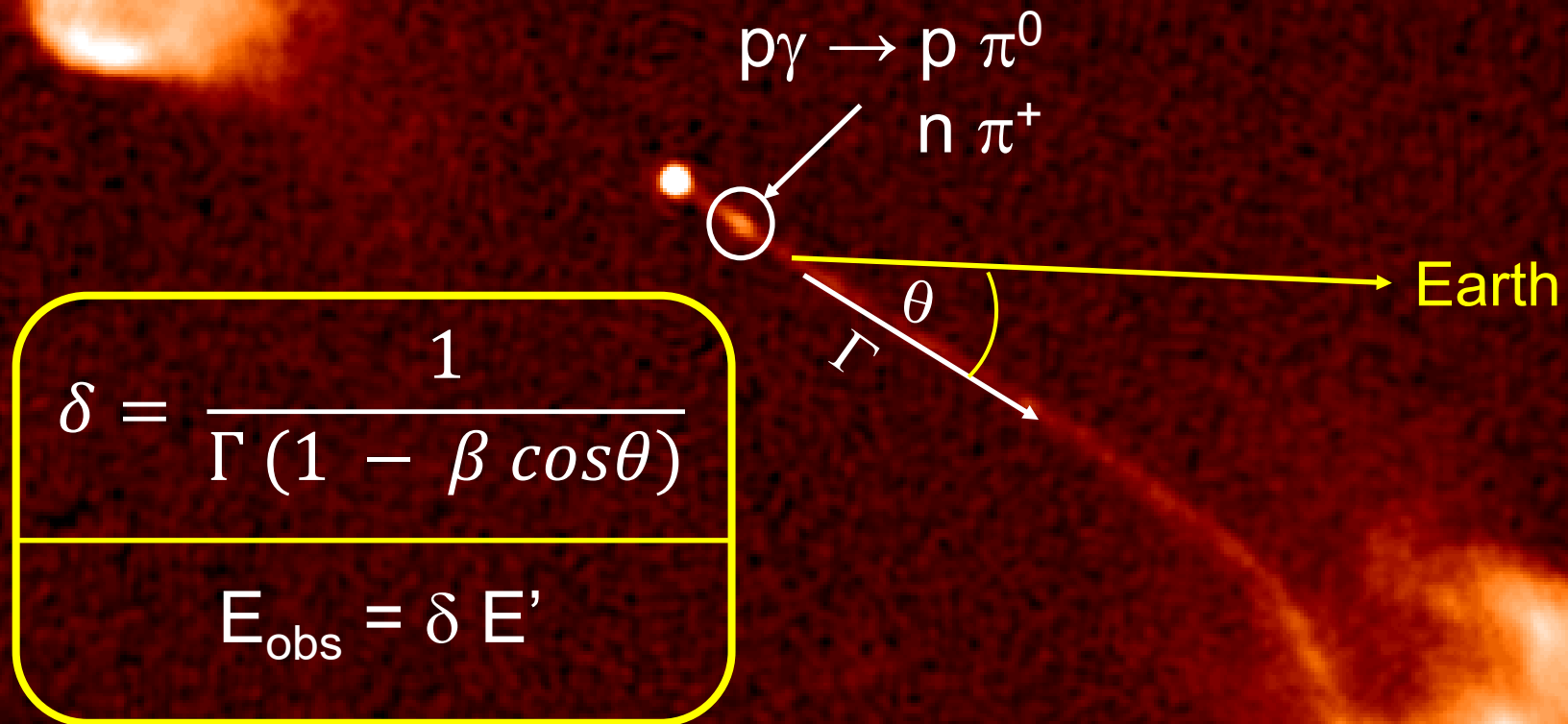


Photo-Pion Production



For realistic target photon fields, most interactions occur near threshold (at Δ^+ resonance).

Photo-pion induced neutrino production in relativistic jets



Quasar 3C175

VLA 6cm image (c) NRAO 1996

Photo-pion production - Energetics

p- γ threshold:
$$E_p^{\text{thr}} = \frac{m_p m_\pi c^4}{2 E_{\text{ph}}} \left(1 + \frac{m_\pi}{2 m_p} \right) \sim 10^{17} \text{ eV } E_{\text{t,eV}}^{-1}$$

At Δ^+ resonance:

$$s = E_p' E_t' (1 - \beta_p' \mu) = E_{\Delta^+}^2 = (1232 \text{ MeV})^2$$

Each neutrino takes about $\sim 5 \%$ of the proton's energy

$\Rightarrow$ To produce IceCube neutrinos ($\sim 100 \text{ TeV} \rightarrow E_\nu = 10^{14} E_{14} \text{ eV}$):

Need protons with

$$E_p' \sim 200 E_{14} \delta_1^{-1} \text{ TeV}$$

and target photons with

$$E_t' \sim 1.6 E_{14}^{-1} \delta_1 \text{ keV}$$

Photo-pion production – Origin of Target Photons

To produce IceCube neutrinos ($\sim 100 \text{ TeV} \rightarrow E_\nu = 10^{14} E_{14} \text{ eV}$):

Need protons with

$$E'_p \sim 200 E_{14} \delta_1^{-1} \text{ TeV}$$

and target photons with

$$E'_t \sim 1.6 E_{14}^{-1} \delta_1 \text{ keV} \Rightarrow \text{X-rays!}$$

(At least) two possible scenarios:

a) Target photons co-moving with the emission region

$$\Rightarrow E_t^{\text{obs}} \sim 16 E_{14}^{-1} \delta_1^2 / (1+z) \text{ keV}$$

$\Rightarrow$ Observed as Doppler-boosted hard X-rays

Tightly constrained by
observed hard X-ray flux $\rightarrow$
Energetics constraints.

b) Target photons stationary in the AGN frame

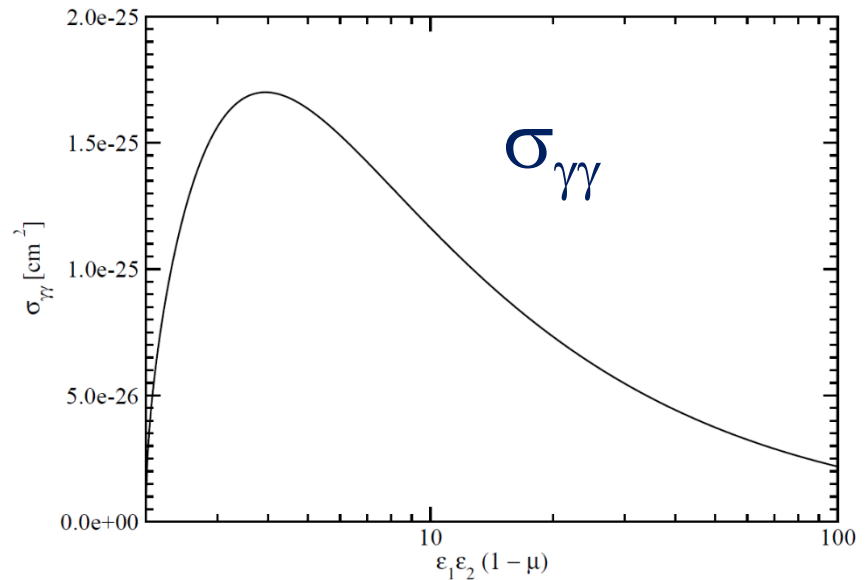
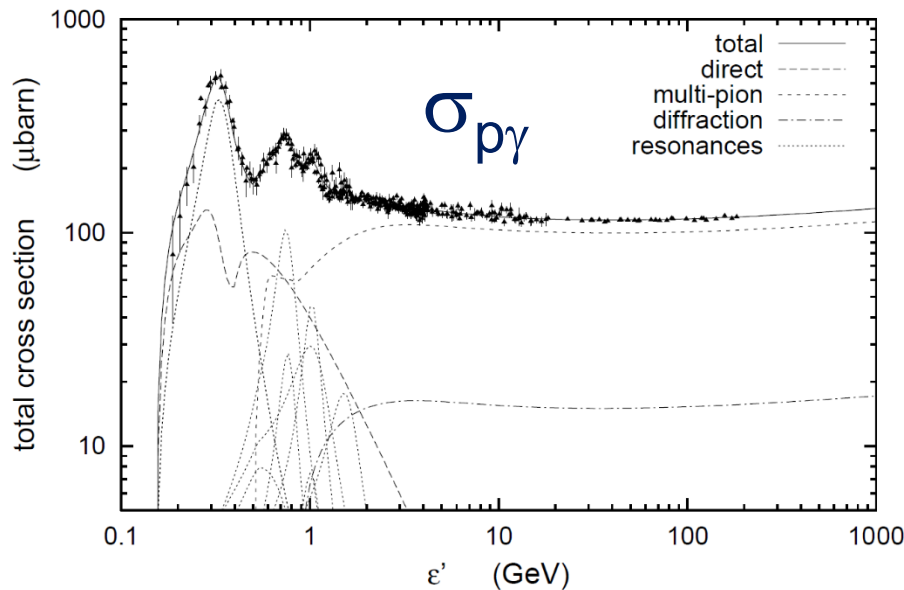
$$\Rightarrow E_t^{\text{obs}} \sim 160 E_{14}^{-1} / (1+z) \text{ eV}$$

$\Rightarrow$ Observed as UV / soft X-rays, Doppler boosted into the emission-region frame

Much more relaxed
energetics constraints.

The $p\gamma$ Efficiency Problem

- Efficiency for protons to undergo $p\gamma$ interaction $\sim \tau_{p\gamma} = \ell_{\text{esc}} \sigma_{p\gamma} n_{\text{ph}}$
- Likelihood of γ -ray photons to be absorbed $\sim \tau_{\gamma\gamma} = R \sigma_{\gamma\gamma} n_{\text{ph}}$



$$\frac{\tau_{p\gamma}}{\tau_{\gamma\gamma}} = \frac{\sigma_{p\gamma} \ell_{\text{esc}}}{\sigma_{\gamma\gamma} R} \approx \frac{1}{300} \frac{\ell_{\text{esc}}}{R}$$

ℓ_{esc} = average length
travelled by protons
until escape

at $E_{\gamma} \sim \frac{m_e^2 c^4}{E_t} \sim 3.3 \times 10^{-5} E_{\nu} \longleftarrow \sim \text{few to tens of GeV } \gamma\text{-rays}$

The $p\gamma$ Efficiency Problem

- Efficiency for protons to undergo $p\gamma$ interaction $\sim \tau_{p\gamma} = \ell_{\text{esc}} \sigma_{p\gamma} n_{\text{ph}}$

- Likely

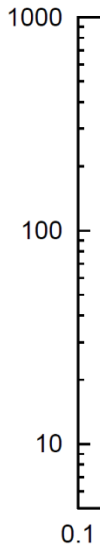
$\Rightarrow$ Efficient neutrino production sites are likely to be optically thick to gamma-rays

$\Rightarrow$ Expect no correlation between GeV – TeV gamma-ray and neutrino activity!

$\Rightarrow$ Cascades! $\rightarrow$ Expect correlation with X-ray / soft (MeV) γ -ray and possibly multi-TeV gamma-ray activity!

$\Rightarrow$ Promising targets for ASTRI / SSTs

total cross section (μbarn)



at $E_\gamma \sim \frac{m_e^2 c^4}{E_t} \sim 3.3 \times 10^{-5} E_\nu \longleftarrow \sim \text{few to tens of GeV } \gamma\text{-rays}$

until escape

length
ons

Modeling Multi-messenger Emissions from photo-pion production

- Most accurate method: Monte-Carlo simulations of individual reactions (e.g., SOPHIA: Mücke et al. 2000: Comp. Phys. Comm., 124, 290)
- Most often used method: Semi-analytical approximations of spectra of final decay products, e.g.,

$$\frac{dN_l(E_l)}{dE_l} = \int \frac{dE_p}{E_p} f_p(E_p) \int d\varepsilon f_{ph}(\varepsilon) \Phi_l(\eta, x)$$

With $\eta = \frac{4 \varepsilon E_p}{(m_p c^2)^2}$, $x = E_l/E_p$, and $\Phi_l(\eta, x)$ tabulated in Kelner & Aharonian (2008: Phys. Rev. D., 78, 034013)

- Integrates out over intermediate (π, μ) particles, i.e., neglects π, μ synchrotron emission (and losses).

=> Only applicable if $\tau_{\text{sy}}(\pi, \mu) \gg \gamma_p \tau'_{\text{decay}}$

=> Requires $\gamma_{p, \text{max}} B \ll 5.6 \times 10^{10} \text{ G}$

Modeling Multi-messenger Emissions from photo-pion production

- More generally applicable method: Production (injection) rates of pions $Q_\pi(E_\pi)$ based on semi-analytical templates from Hümmer et al. (2010: ApJ, 721, 630)
- Solve coupled Fokker-Planck equations for the evolution of all particle species, e.g.,

$$\frac{\partial n_\pi(\gamma, t)}{\partial t} = \frac{\partial}{\partial \gamma} \left[\frac{\gamma^2}{(a+2)t_{\text{acc}}} \frac{\partial n_\pi(\gamma, t)}{\partial \gamma} \right] - \frac{\partial}{\partial \gamma} (\dot{\gamma}_{\text{rad}} n_\pi(\gamma, t)) + Q_\pi(\gamma, t) - \frac{n_\pi(\gamma, t)}{t_{\text{esc}}} - \frac{n_\pi(\gamma, t)}{\gamma t'_{\text{decay}}}$$

Stochastic acceleration
Radiative (synchrotron) losses
Pion production
Escape
Pion decay: $\pi^+ \rightarrow Q_\mu(E_\mu)$
 $\pi^0 \rightarrow Q_\gamma(E_\gamma)$

(e.g., Diltz et al. 2016, Zacharias et al. 2019, 2021)

(Lepto-)Hadronic codes “on the market”

- AM³ (Gao et al. 2017; Klinger et al. 2024)
 - Atheva (Mastichiadis & Kirk 1995; Mastichiadis et al. 2005; Dimitrakoudis et al. 2012; Petropoulou et al. 2014)
- B13 (Böttcher et al. 2013)
 - LeHa-Paris (Cerruti et al. 2015)
 - LeHaMoC (Stathopoulos et al. 2024)
 - OneHaLe (Zacharias et al. 2019, 2021)

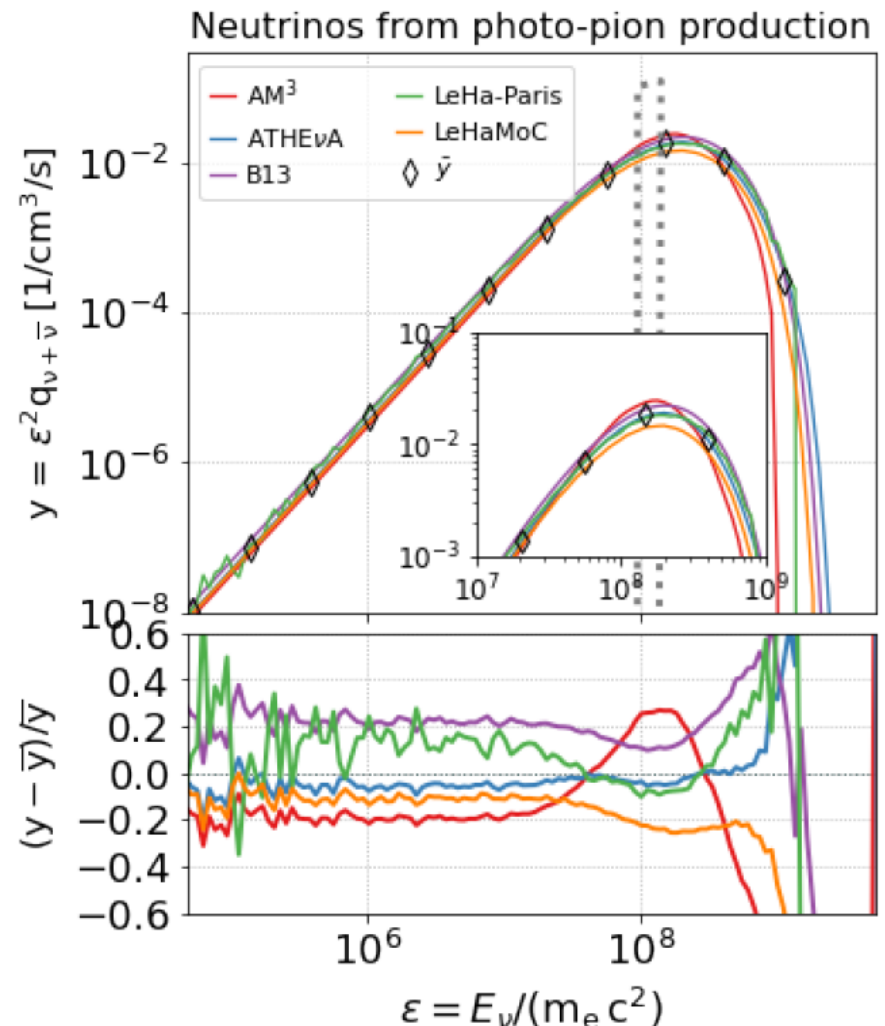
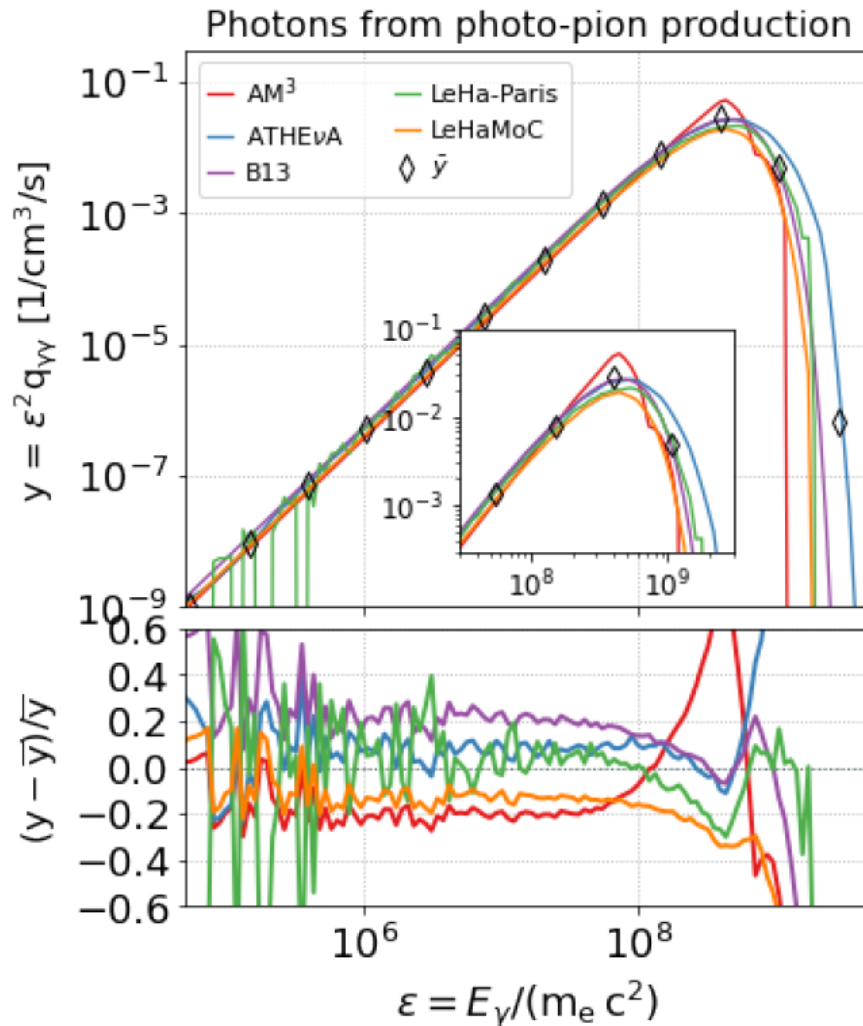
Table 1
Physical Processes Included in the Numerical Codes

Physical Processes	Codes				
	AM ³	ATHE ν A	B13	LeHa-Paris	LeHaMoC
Electron synchrotron radiation	✓	✓	✓	✓	✓
Synchrotron self-absorption	✓	✓	✓	✓	✓
Electron inverse Compton scattering	✓	✓	✓	✓	✓
Electron–positron annihilation	✗	✓	✓	✗	✗
Photon–photon pair production	✓	✓	✓	✓	✓
Triplet pair production	✗	✓	✗	✗	✗
Proton synchrotron radiation	✓	✓	✓	✓	✓
Proton inverse Compton scattering	✓	✗	✗	✗	✗
Proton–photon pair production	✓	✓	✗	✓	✓
Proton–photon pion production	✓	✓	✓	✓	✓
Proton–proton pion production	✗	✗	✗	✓	✓
Neutron–photon pion production	✓	✓	✗	✗	✗
Neutron decay	✗	✓	✗	✗	✗
Kaon synchrotron radiation	✗	✓	✗	✗	✗
Pion synchrotron radiation	✓	✓	✗	✗	✗
Muon synchrotron radiation	✓	✓	✗	✓	✗

(Cerruti et al. 2026: Hadronic code comparison project)

(Lepto-)Hadronic codes “on the market”

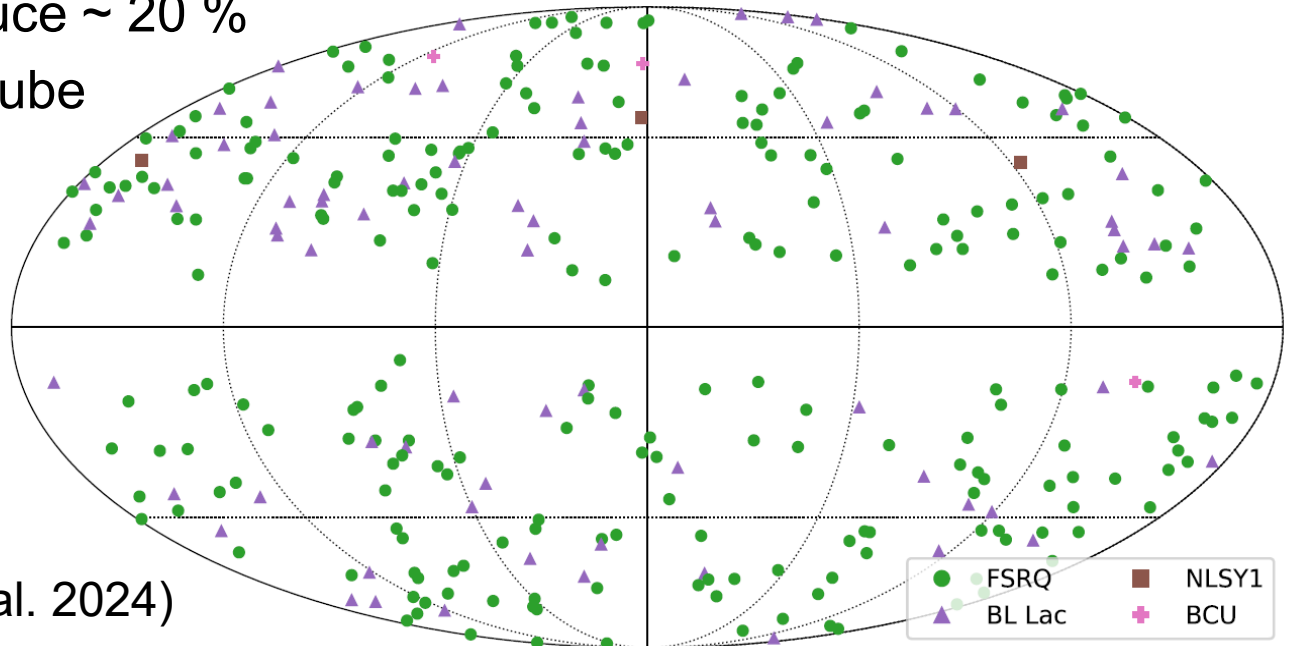
Various codes agree with each other to within a few %.



(Cerruti et al. 2026: Hadronic code comparison project)

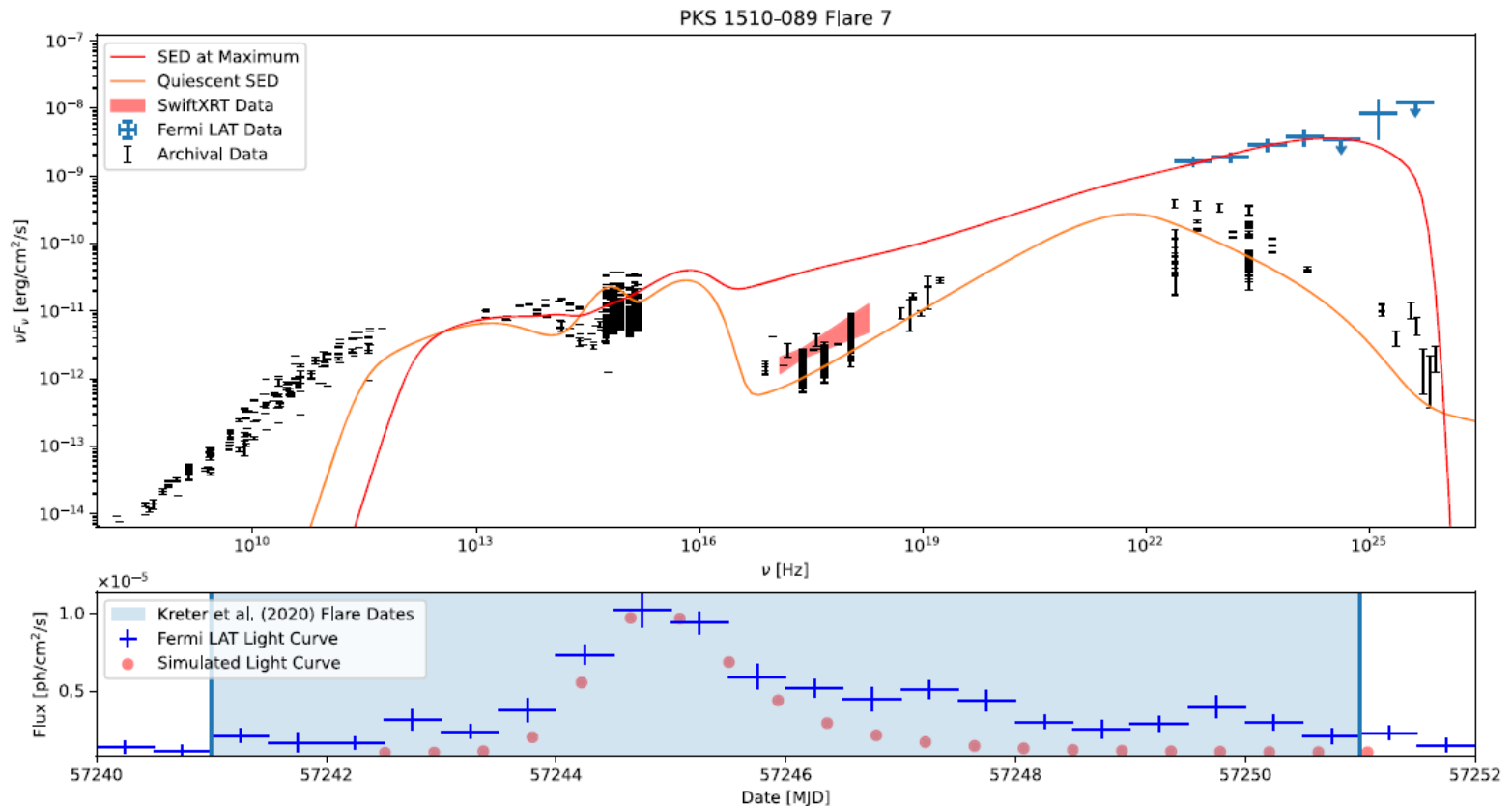
Recent candidate neutrino blazar modelling results

- Large body of publications on candidate neutrino blazar modelling using various codes.
- Arguably most comprehensive: Rodrigues et al. (2024): Systematic modelling (AM³) of the SEDs of 324 bright blazars :
 - In most cases, leptonic-dominated models suffice;
 - only ~ 33 % benefit from an additional hadronic component.
- Blazars may produce ~ 20 % of the diffuse IceCube neutrino flux, in agreement with stacking limits.



Neutrino detection potential from bright blazar flares

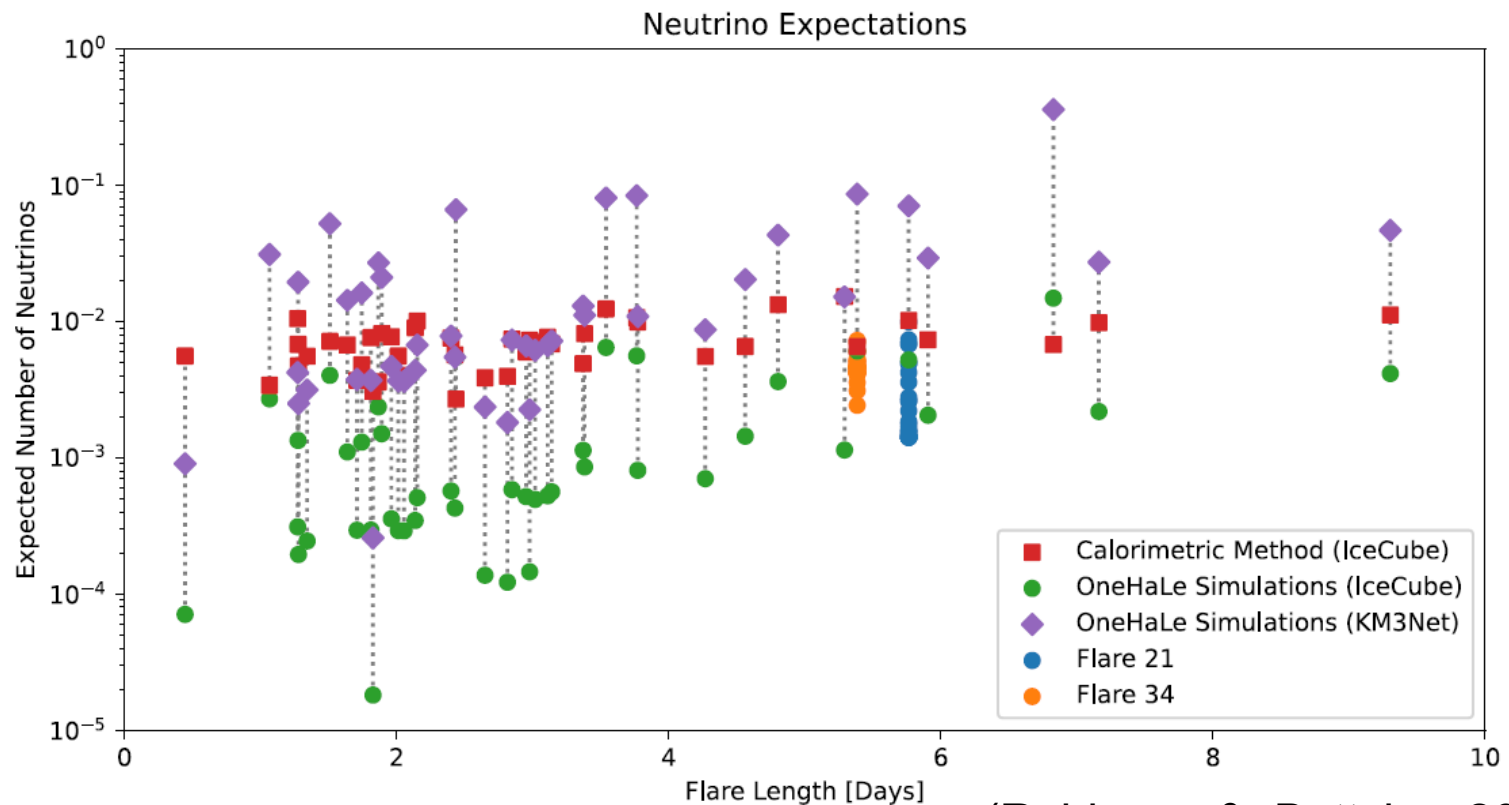
- Using OneHaLe for hadronically-dominated modelling of SEDs and light curves of the 20 brightest Fermi-LAT blazar flares studied by Kreter et al. (2020) with a simple calorimetric method (neutrino output proportional to bolometric high-energy flux).



(Robinson & Böttcher 2024)

Neutrino detection potential from bright blazar flares

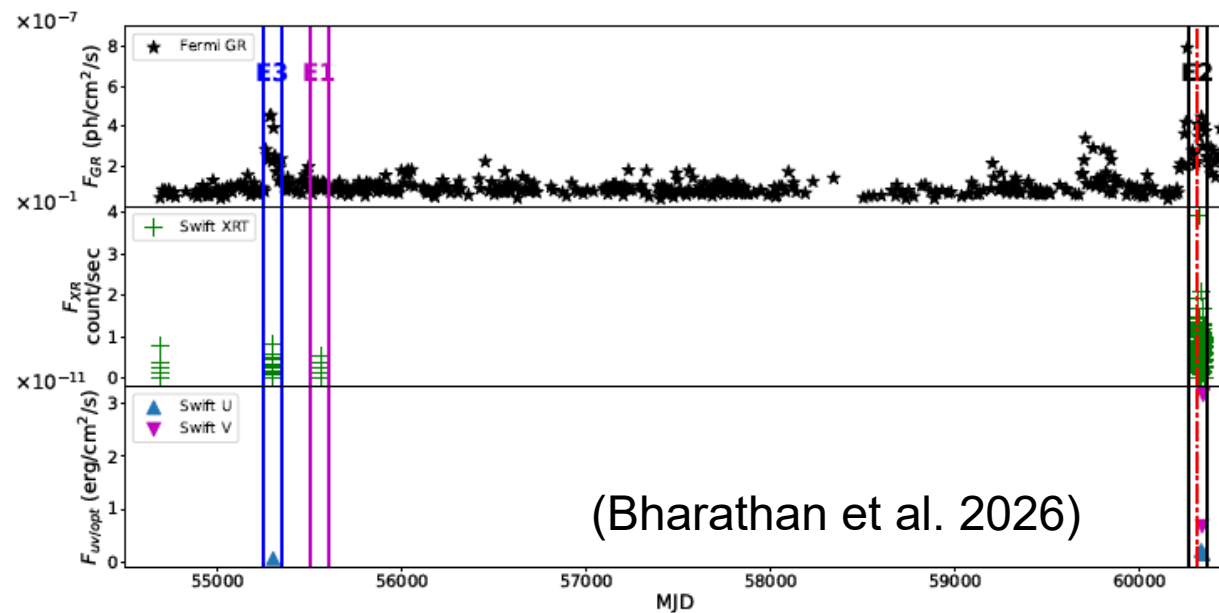
- Even the brightest blazar flares expected to yield < 0.1 neutrinos even for KM3NeT design configuration.
- Calorimetric approximation tends to over-estimate neutrino expectations by about an order of magnitude.



(Robinson & Böttcher 2024)

Dedicated SED modelling of 4 neutrino blazars

- Bharathan et al. (2026): Selected 4 recent blazar-neutrino coincidences with good broadband coverage:
 - PKS 0446+112 ($z = 2.153$; IC 240105)
 - TXS 0506+056 ($z = 0.3365$; IC170922A)
 - PKS 1424-418 ($z = 1.522$; IC 121204)
 - PKS 1502+106 ($z = 1.8378$; IC 190730A)
- Selected 3 different states for each source:
 - E1: quiescent
 - E2: around neutrino alert
 - E3: bright γ -ray flare
- Lepto-hadronic modelling with Böttcher et al. (2013) code.



Dedicated SED modelling of 4 neutrino blazars

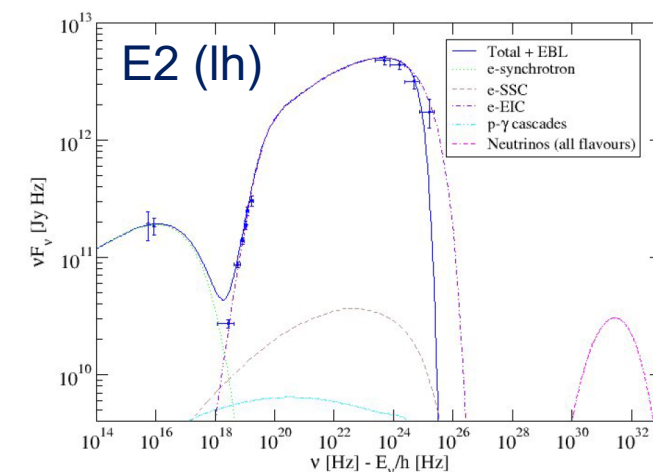
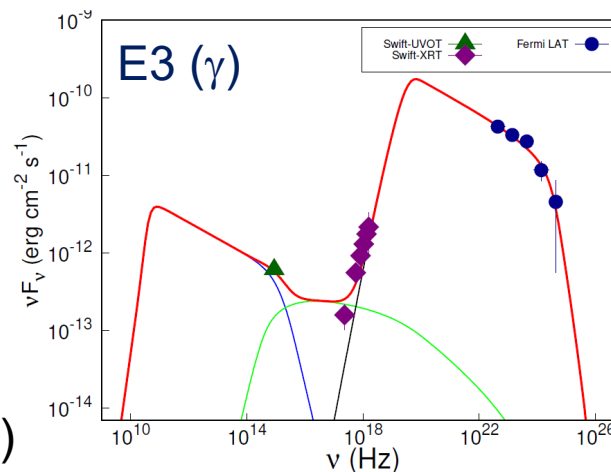
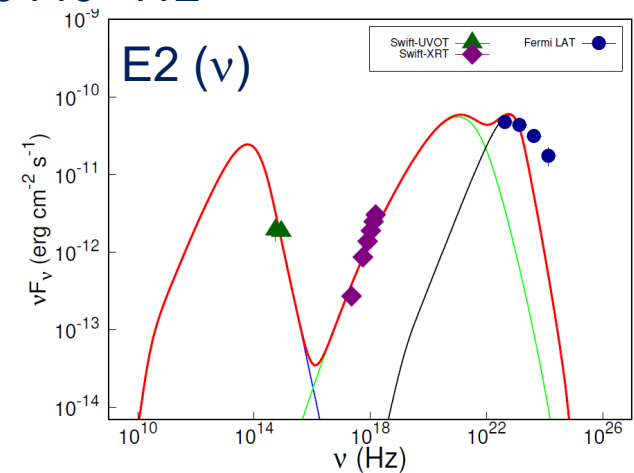
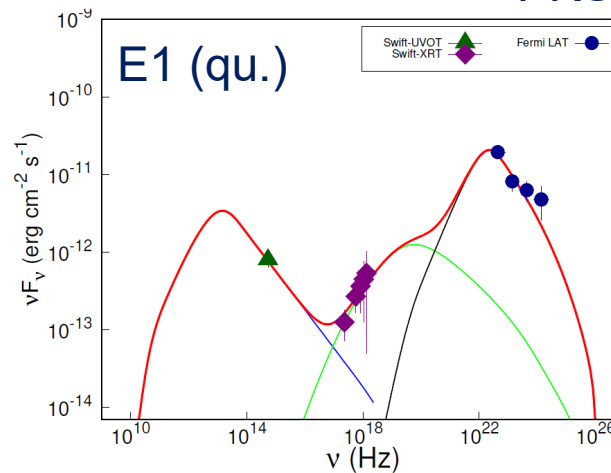
- Most SEDs well fitted with purely leptonic models; hadronically ($p\gamma$) dominated models excluded by hard X-ray spectra.

- 2 exceptions where leptonic models fail:

- PKS 0446+112 during neutrino detection, but $p\gamma$ cascade component still sub-dominant (see also Khatoun et al. 2026)
- PKS 1502+106 during neutrino detection.

(Bharathan et al. 2026)

PKS 0446+112



Summary

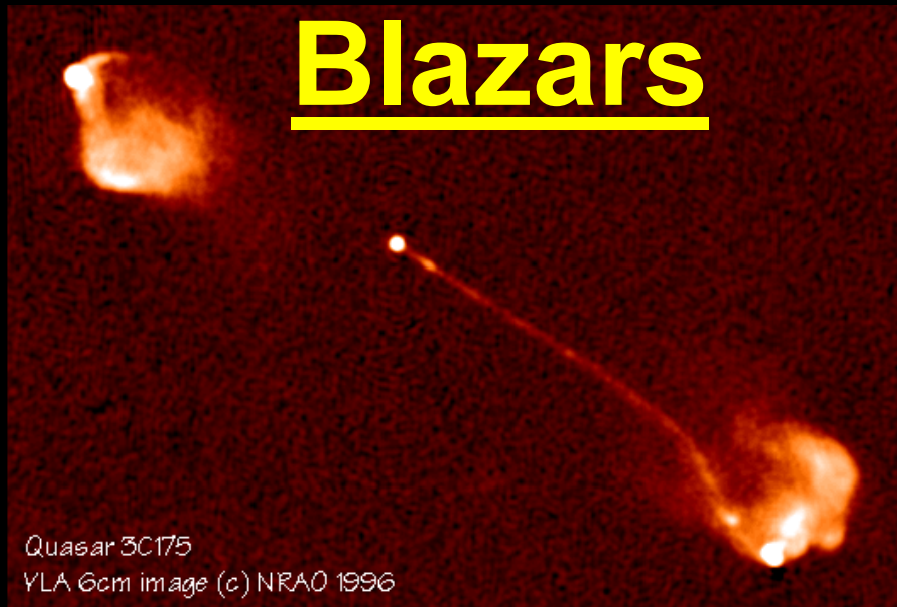
1. Association of IceCube neutrinos with γ -ray and radio bright blazars becomes more and more significant.
2. Energetics constraints prefer external target radiation fields for $p\gamma$ interactions.
3. SED modeling of blazars generally prefers leptonically dominated broadband emission with at most a sub-dominant contribution of hadronic processes to X-ray – γ -ray emission.
4. Even the brightest γ -ray flares of blazars are unlikely to yield a neutrino detection probability of $> 10\%$ even for KM3NeT design sensitivity.

A scenic view of a church in a mountain valley. The church is a light-colored building with a tall, slender bell tower topped with a dark dome. The tower has two clock faces. The church is situated on a grassy slope, with a cemetery in front of it. In the background, there are steep, rocky mountains under a blue sky with scattered white clouds. The text "Thank you!" is overlaid in the center of the image.

Thank you!

Backup

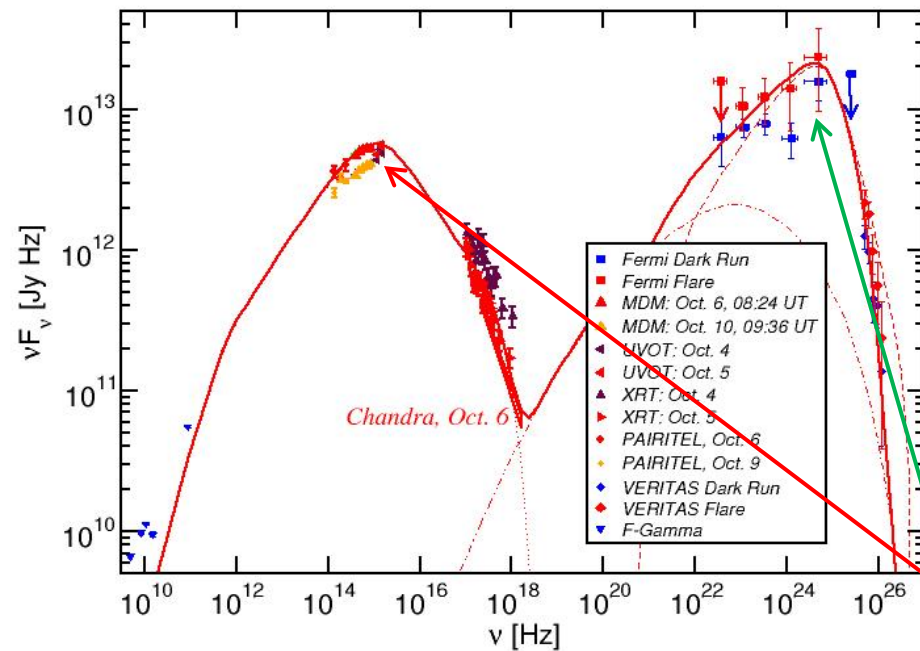
Blazars



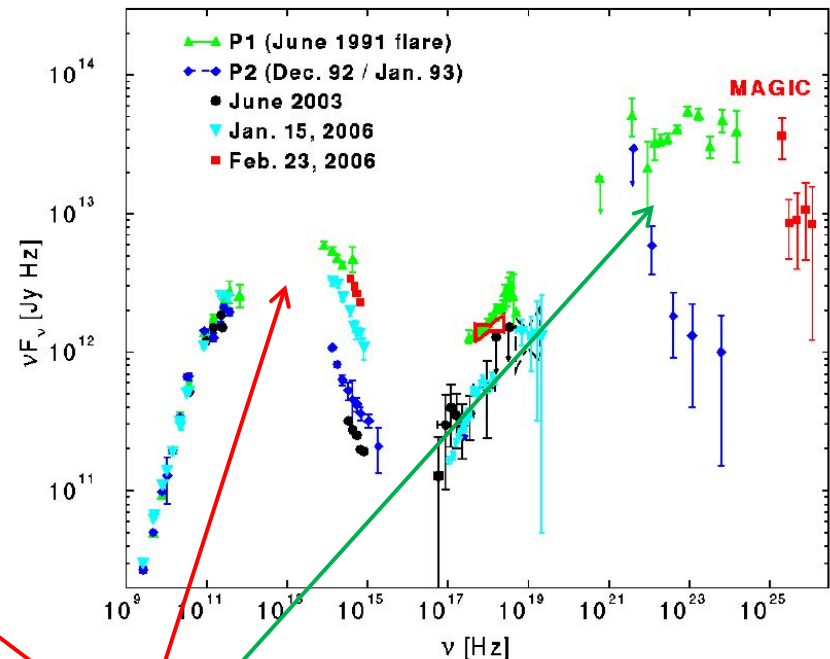
- Class of AGN consisting of BL Lac objects and gamma-ray bright quasars with relativistic jets pointing close to our line of sight
- Rapidly (often intra-day) variable
- Strong gamma-ray sources
- Radio knots often with superluminal motion
- Radio and optical (and X-ray) polarization

Blazar Spectral Energy Distributions (SEDs)

3C66A



3C279

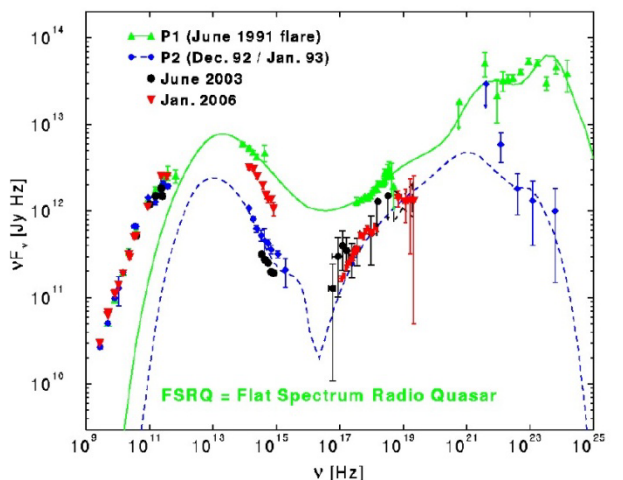


Non-thermal spectra with
two broad bumps:

- Low-energy (probably synchrotron):
radio-IR-optical(-UV-X-rays)
- High-energy (X-ray – γ -rays)

Blazar Classification

3C279



(Hartman et al. 2000)

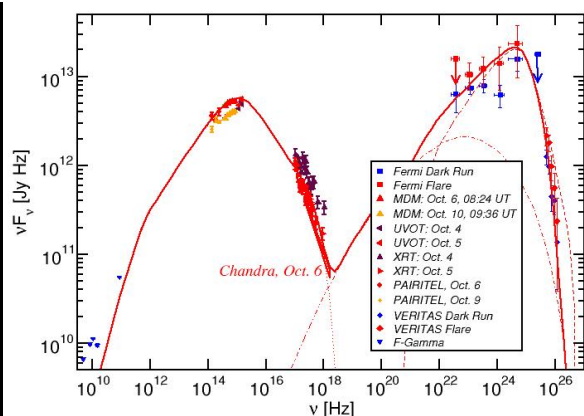
Low-Synchrotron Peaked (LSP): Quasars (FSRQs)/ Low-frequency peaked BL Lac Objects (LBLs)

Low-frequency component from radio to optical/UV,

$$\nu_{\text{sy}} \leq 10^{14} \text{ Hz}$$

High-frequency component from X-rays to γ -rays, often dominating total power

3C66A



(Abdo et al. 2011)

Intermediate-Synchrotron Peaked (ISP): Intermediate BL Lacs (IBLs):

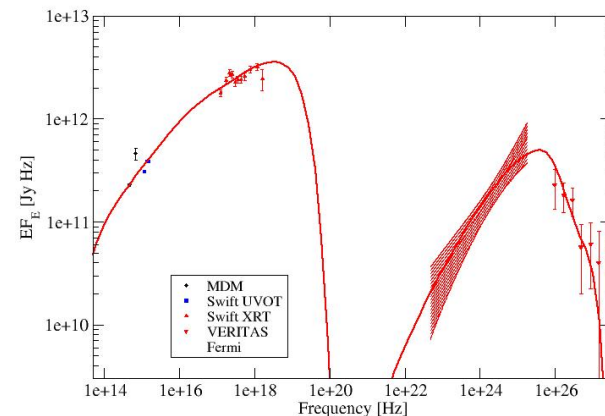
Peak frequencies at IR/Optical and GeV gamma-rays,

$$10^{14} \text{ Hz} < \nu_{\text{sy}} \leq 10^{15} \text{ Hz}$$

Intermediate overall luminosity

Sometimes γ -ray dominated

RGB J0710+591



(Acciari et al. 2009)

High-Synchrotron Peaked (HSP): High-frequency peaked BL Lacs (HBLs):

Low-frequency component from radio to UV/X-rays,

$$\nu_{\text{sy}} > 10^{15} \text{ Hz}$$

often dominating the total power

High-frequency component from hard X-rays to high-energy gamma-rays

Flux and Polarization Variability

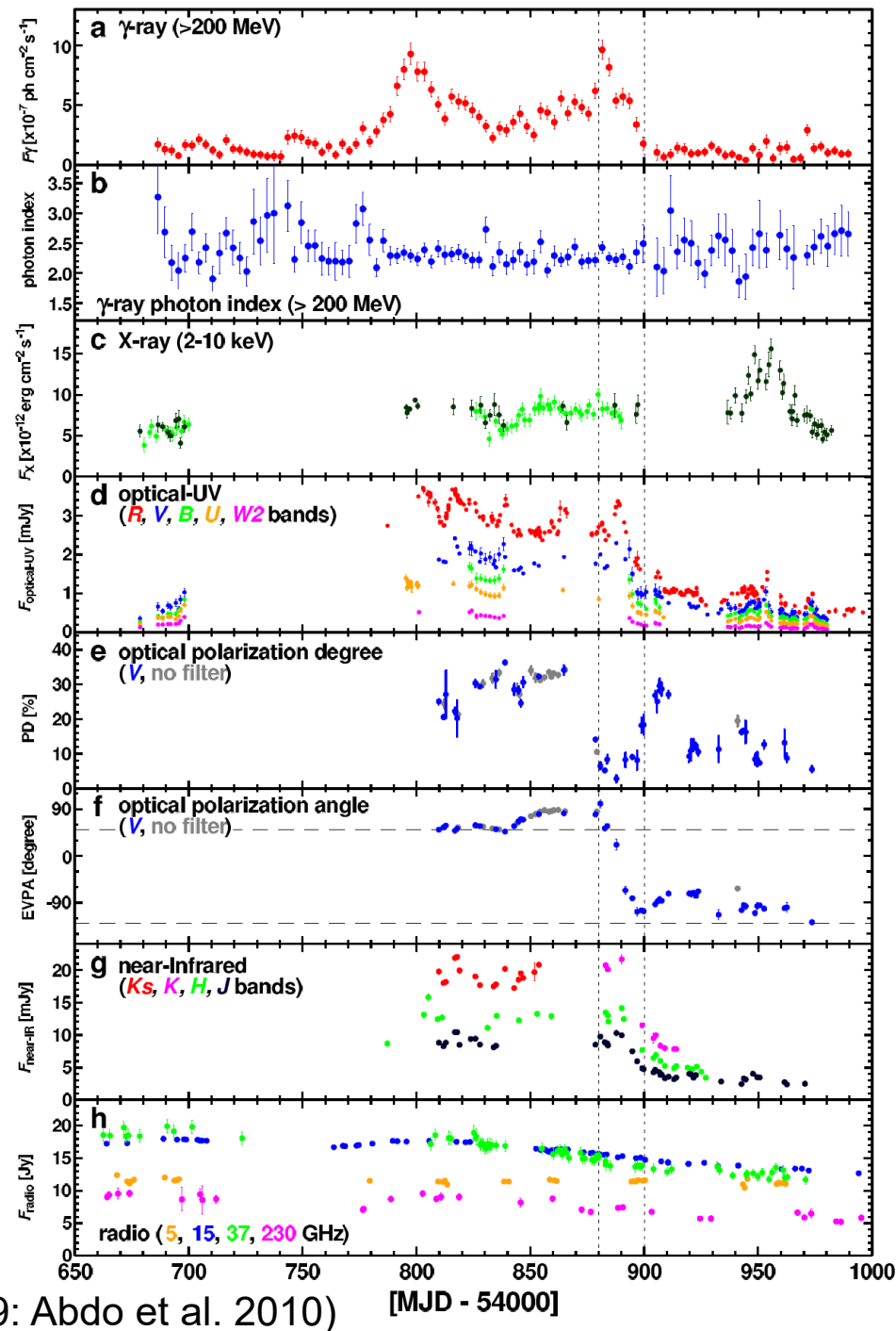
Multi-wavelength variability on various time scales (months – minutes)

Sometimes correlated, sometimes not

Observed optical polarization degrees

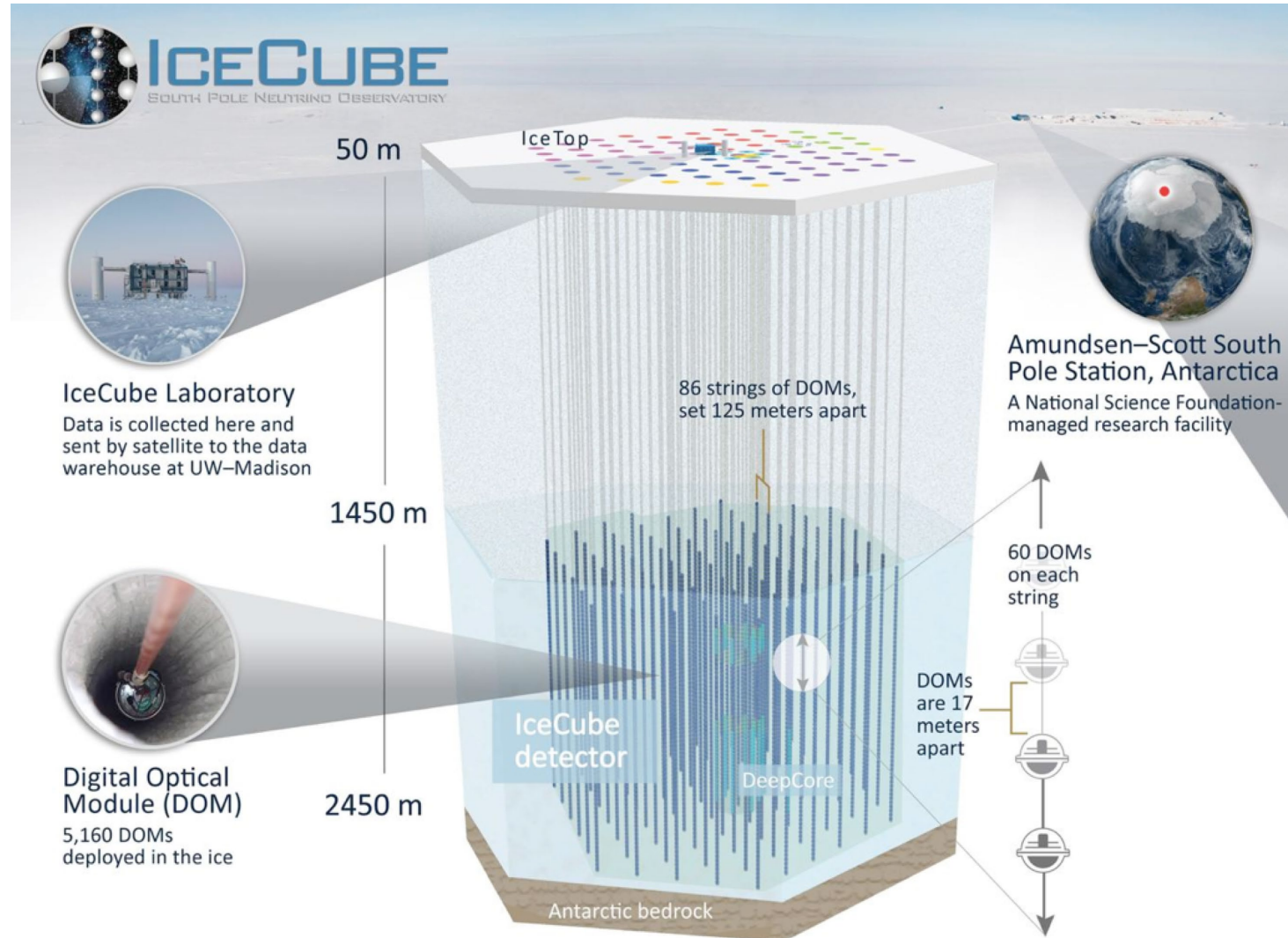
$$\Pi_{\text{opt}} < \sim 30 \%$$

Both degree of polarization and polarization angles vary.



(3C279: Abdo et al. 2010)

The IceCube Neutrino Detector at the South Pole



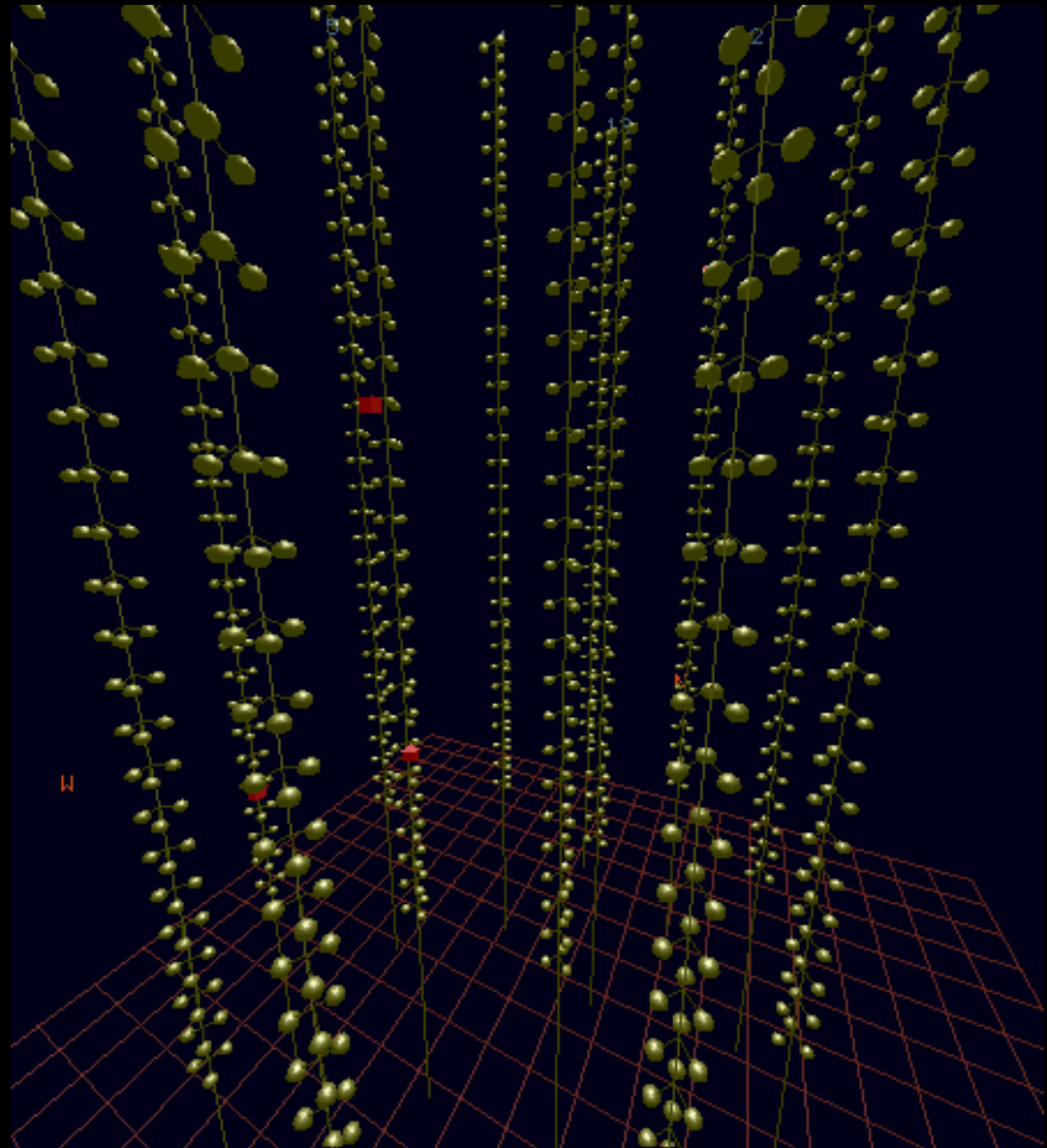
Fully operational since 2010.

High-Energy Neutrino Detectors

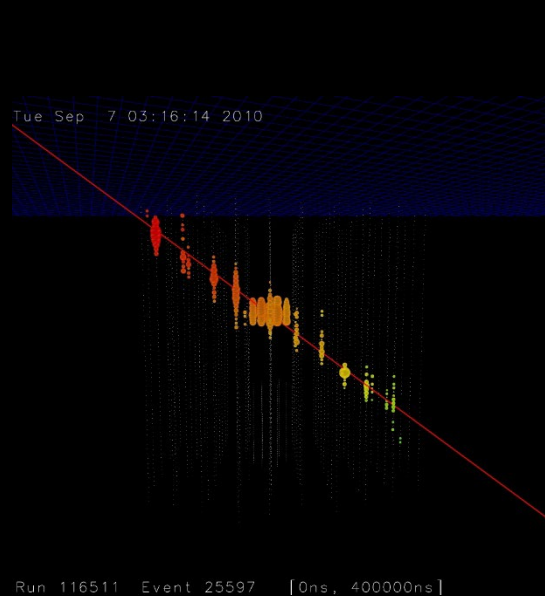
IceCube:

$E_\nu \sim 100 \text{ TeV} - \text{few PeV}$

ν -matter scattering,
followed by particle
cascades in ice/water
→ Cherenkov light
detection.

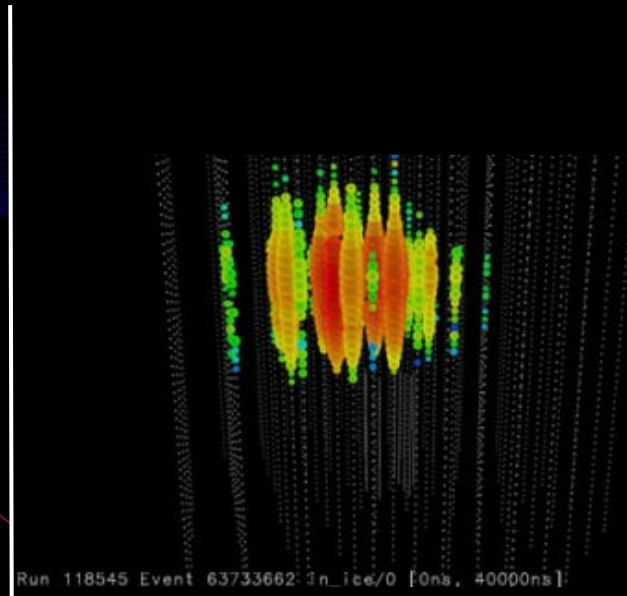


Neutrino Event Types



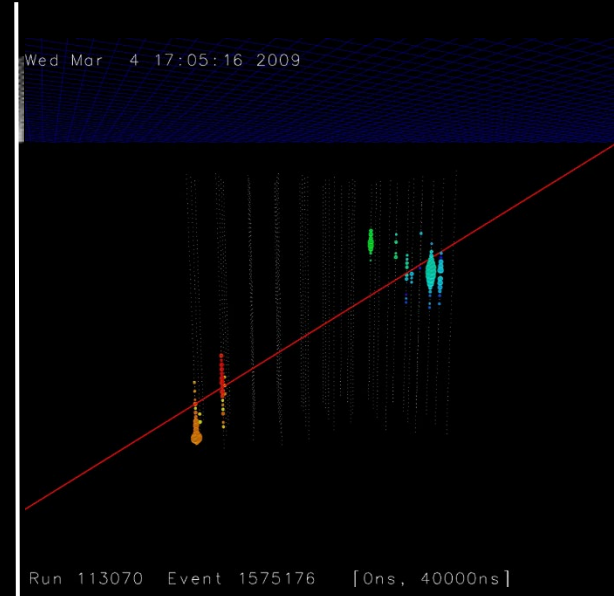
Tracks $\leftarrow \nu_{\mu}$

ν_{μ} interaction
produces a muon
→ Long track
→ Localization down
to $\sim 0.4^{\circ}$.



Cascades $\leftarrow \nu_e$

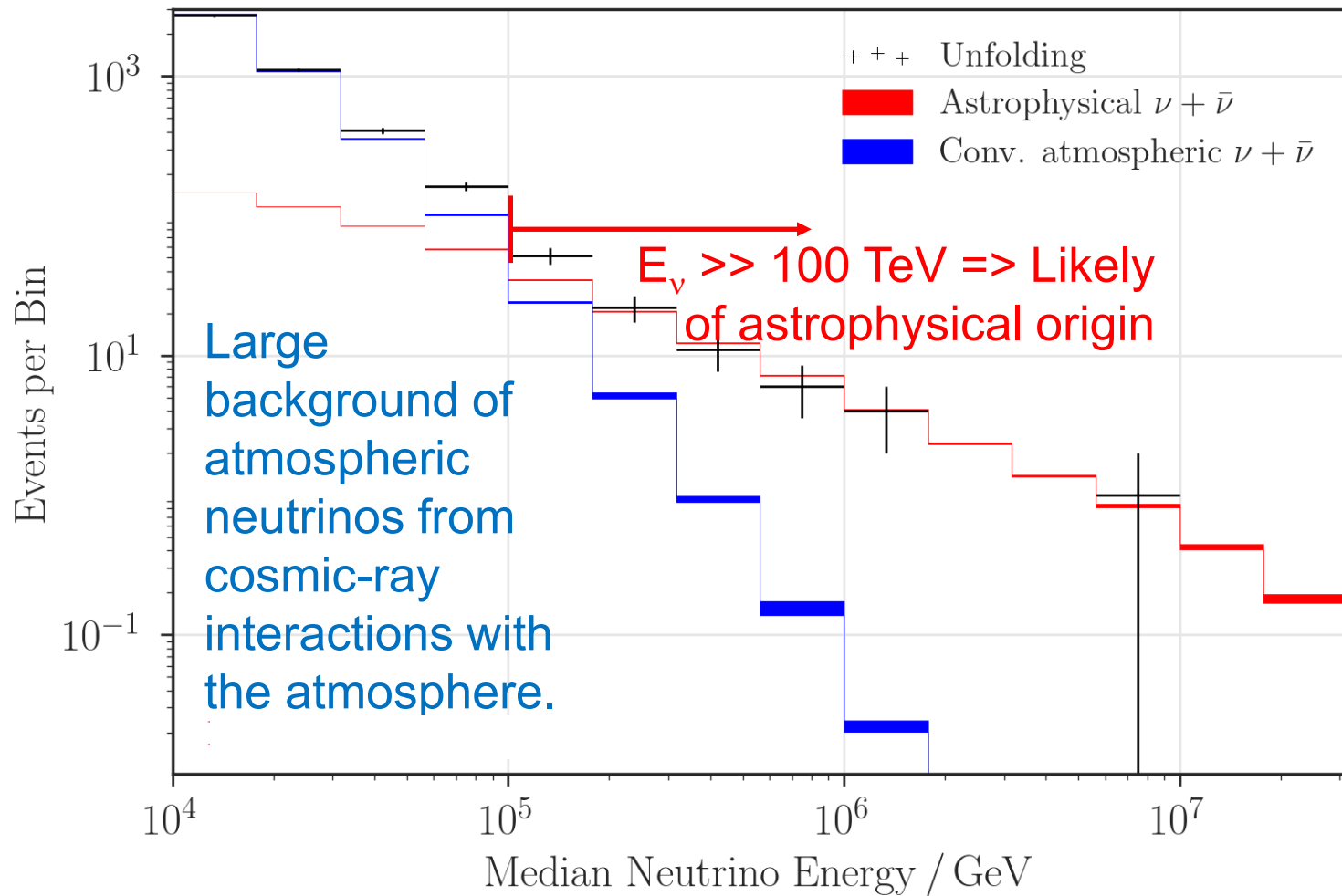
ν_e interaction produces
an electron
→ loses energy quickly
in a blob-like cascade.
Localization to $\sim 8^{\circ}$.



Double Bang $\leftarrow \nu_{\tau}$

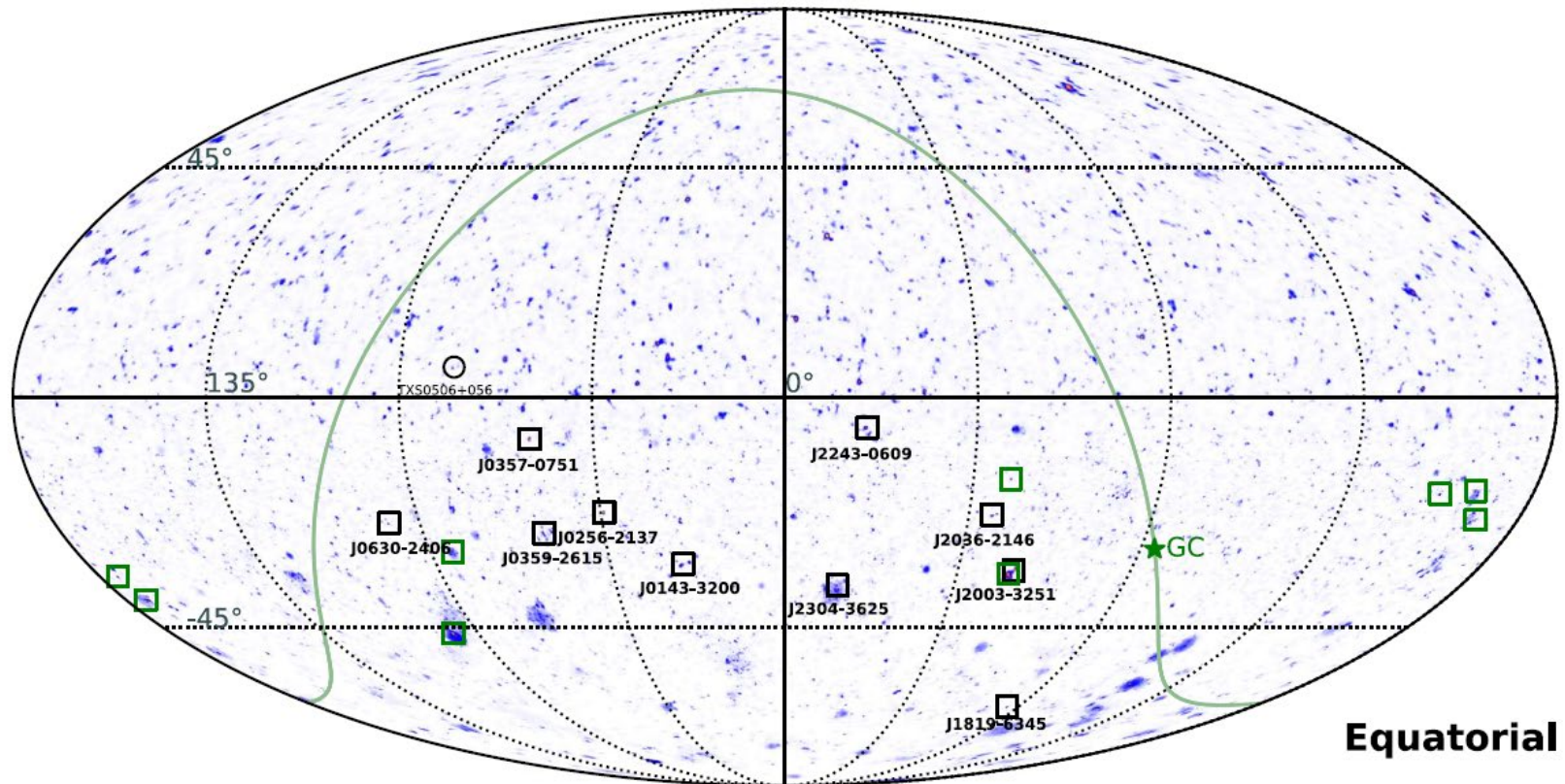
ν_{τ} interaction
produces a tau lepton
→ 2 cascades: One at
tau production, one at
tau decay.

The IceCube Neutrino Spectrum



First evidence for astrophysical neutrinos published in 2013.

Origin of IceCube-Detected Neutrinos

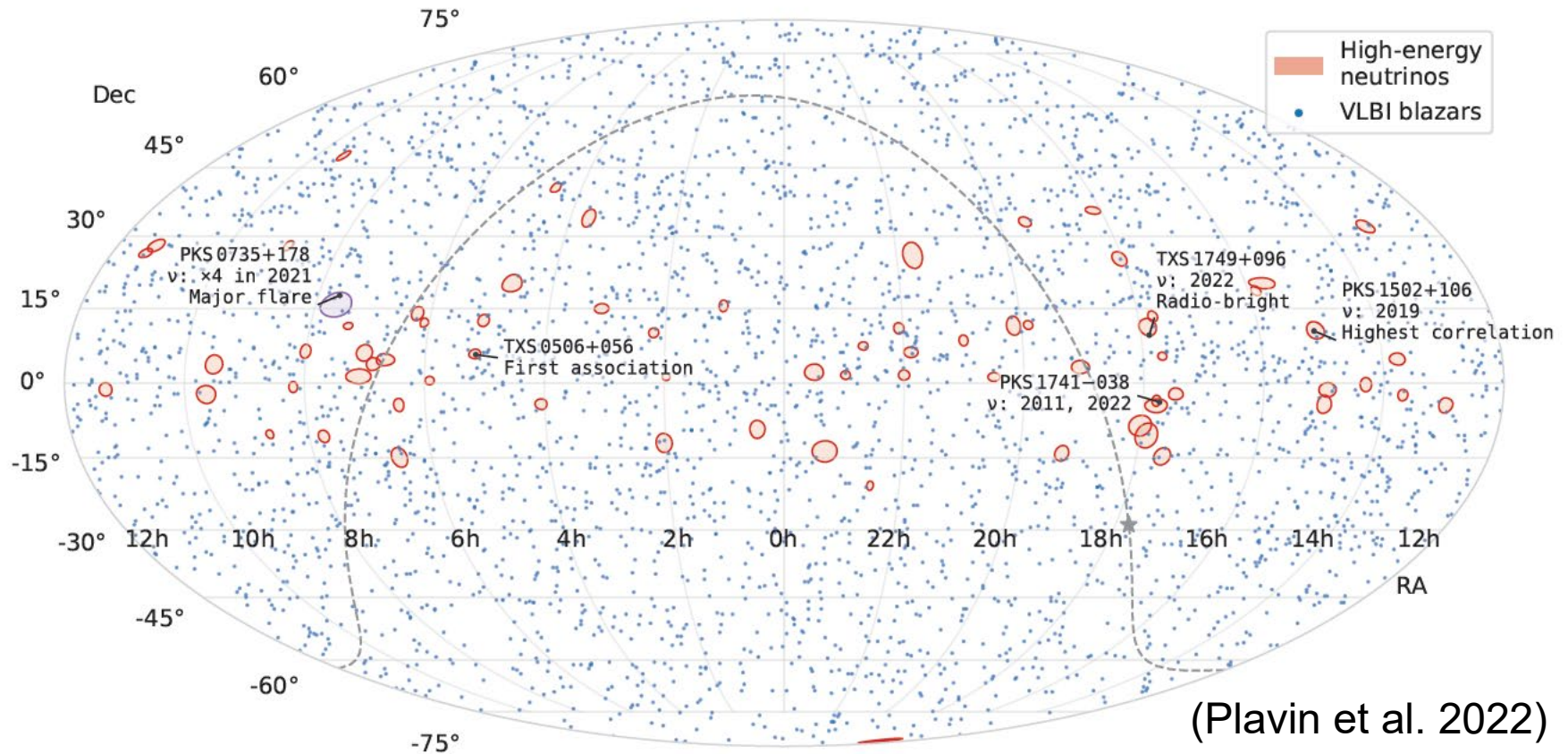


(Buson et al. 2022)

Significant correlation of IceCube neutrinos with γ -ray (Fermi-LAT) **blazars**
(chance coincidence probability $p = 6 \cdot 10^{-7}$)

– but can not be responsible for all IceCube neutrinos (e.g., Murase et al. 2018)

Origin of IceCube-Detected Neutrinos



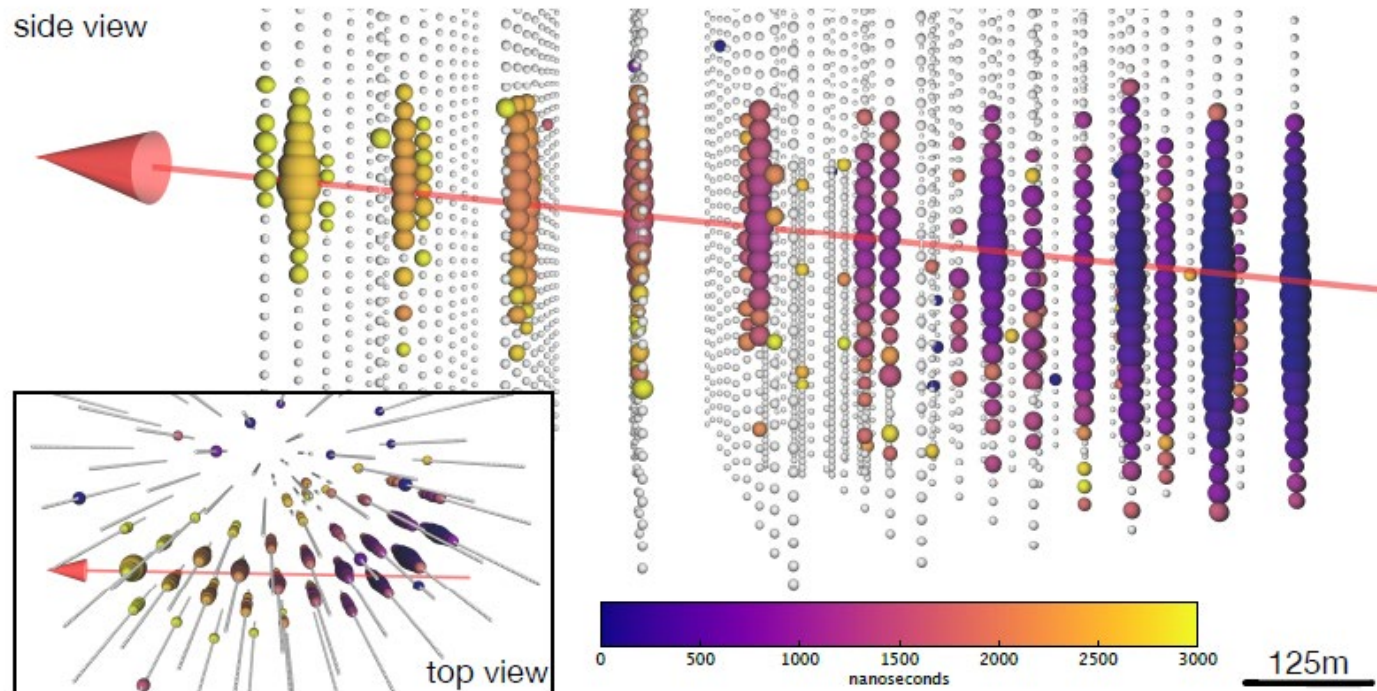
Also: Correlation with radio-loud blazars
(chance-coincidence probability of $p = 3 \cdot 10^{-4}$)
(Plavin et al. 2020, 2021, 2022)

The Eddington Bias

- Most modelled blazar – neutrino coincidences based on 1 single neutrino.
- ~ 100 other blazar sources with similar properties (same neutrino flux?)
- Models only need to produce ~ 0.01 neutrinos from any individual source! (Strotjohann et al. 2019)

Tentative Associations of IceCube Neutrinos with Blazars

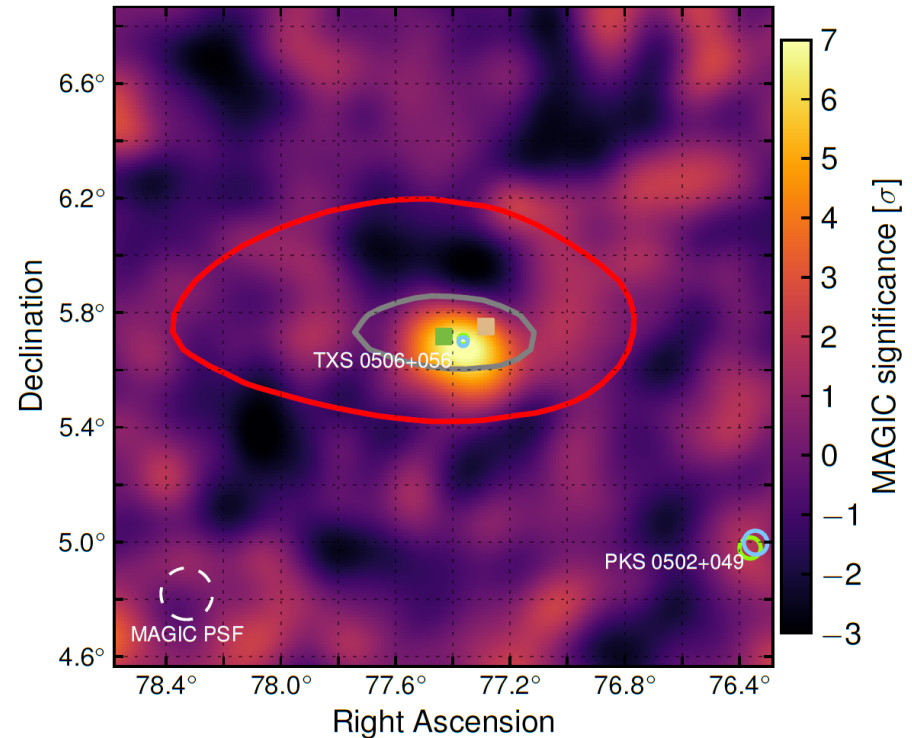
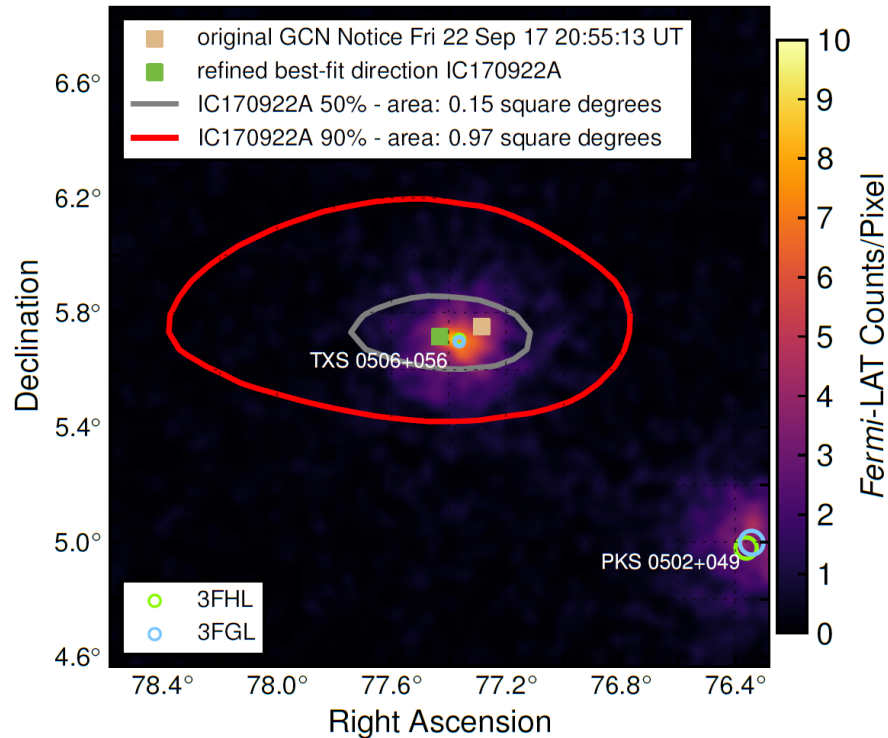
IceCube-170922A



(IceCube et al. 2018a)

Sept. 22, 2017: neutrino-induced muon track with $E_\nu \sim 290$ TeV
~ 50 % probability of being astrophysical

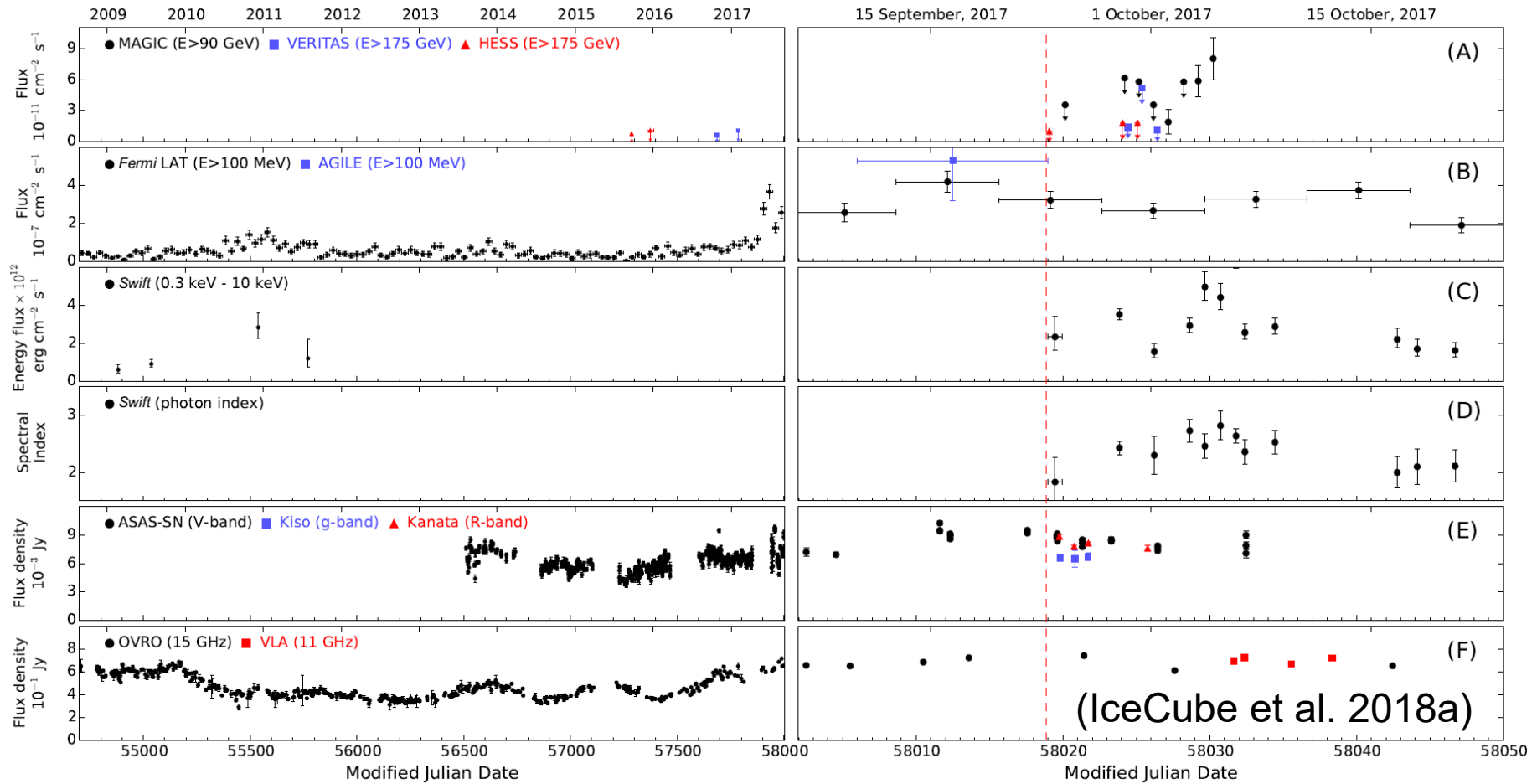
IceCube-170922A and TXS 0506+056



(IceCube et al. 2018a)

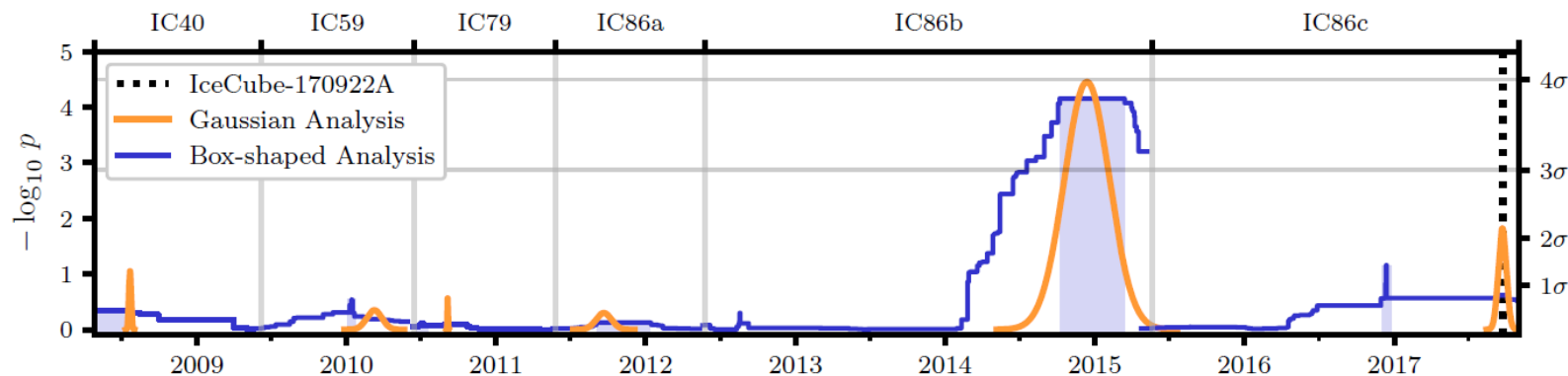
Position consistent with BL Lac object TXS 0506+056
($z = 0.3365$) during a ~ 4 weeks long γ -ray high state
($\sim 3 \sigma$ probability of association)

IceCube-170922A and TXS 0506+056



Position consistent with BL Lac object TXS 0506+056 ($z = 0.3365$) during a ~ 4 weeks long γ -ray high state.
Flare detection by MAGIC several days after neutrino

The Neutrino Flare from TXS 0506+056



(IceCube et al. 2018b)

Search in archival data => Evidence
for 13 ± 5 excess lower-energy
(< 100 TeV) neutrinos from the
direction of TXS 0506+056
in 2014 – 2015

(~ 4 months around December 2014)

=> Possible to estimate flux and
spectrum!

