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Incidence of afterglow plateaus in GRBs associated with binary NS mergers

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Bologna 7-9 Maggio 2025



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How to identify and accurately localise newly born magnetars



 X-ray plateaus: main properties and possible interpretations

• Our work: computing the incidence of X-ray plateaus in binary driven GRBs

Results and Conclusions

Outline

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Incidence of afterglow plateaus in gamma-ray bursts associated with binary neutron star mergers

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ABSTRACT

One of the most surprising gamma-ray burst (GRB) features discovered with the Swift X-ray telescope (XRT) is a plateau phase in the early X-ray afterglow light curves. These plateaus are observed in the majority of long GRBs, while their incidence in short GRBs (SGRBs) is still uncertain due to their fainter X-ray afterglow luminosity with respect to long GRBs. An accurate estimate of the fraction of SGRBs with plateaus is of utmost relevance given the implications that the plateau may have for our understanding of the jet structure and possibly of the nature of the binary neutron star (BNS) merger remnant. This work presents the results of an extensive data analysis of the largest and most up-to-date sample of SGRBs observed with the XRT, and for which the redshift has been measured. We find a plateau incidence of 18-37% in SGRBs, which is a significantly lower fraction than that measured in long GRBs (>50%). Although still debated, the plateau phase could be explained as energy injection from the spin-down power of a newly born magnetized neutron star (NS; magnetar). We show that this scenario can nicely reproduce the observed short GRB (SGRBs) plateaus, while at the same time providing a natural explanation for the different plateau fractions between short and long GRBs. In particular, our findings may imply that only a minority of BNS mergers generating SGRBs leave behind a sufficiently stable or long-lived NS to form a plateau. From the probability distribution of the BNS remnant mass, a fraction 18-37% of short GRB plateaus implies a maximum NS mass in the range ~2.3-2.35 Mo.

Key words. equation of state - gamma-ray burst: general - stars: magnetars

1. Introduction

Gamma-ray bursts (GRBs) have been a great astrophysical mystery since their discovery in the late 1960s. By the end of the 1990s, their cosmological origin was assessed with the discovery of the afterglow component (e.g. Costa et al. 1997) and the identification of their host galaxies (Metzger et al. 1997). A new breakthrough came with the launch of the Neil Gehrels Swift Observatory (Swift hereafter, Gehrels et al. 2004) in November 2004, which allowed the first observations of the early phases (a few minutes after the burst) of the afterglow emission, leading to the discovery of unexpected features that are thought to encode crucial information on the jet structure and possibly on the nature of the remnant compact object.

telescope (XRT) revealed, in most cases, an initial steep flux decay, likely marking the switching off of the prompt emission, followed by a shallow phase (the so-called plateau), which then (NS) remnant injecting energy into the forward shock (Usov transitions to a characteristic power-law flux decay (Zhang et al. 2006). While the latter is in agreement with the afterglow theory of synchrotron emission by electrons energised in a relativistic shock (Sari et al. 1998), the plateau could not be explained in the same framework, requiring additional physics.

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After almost two decades of Swift/XRT GRB observations, we now know that plateaus occur in the majority of long GRBs, which are those associated with the collapse of massive stars. Short GRBs (SGRBs) associated with binary neutron star (BNS) mergers (and possibly neutron star-black hole mergers), proved harder to study due to their fainter afterglow luminosity with respect to long GRBs: to date, the frequency of plateaus in SGRB afterglows is uncertain (Rowlinson et al. 2013). A precise estimate of plateau incidence in SGRBs is of utmost relevance due to its potential impact on our understanding of their jet morphology and conceivably of the nature of the BNS merger remnant. Indeed, it has been suggested that plateaus could originate from geometrical effects in structured jets1 (e.g. Oganesyan et al. 2020; Beniamini et al. 2022). In this case, the plateau incidence More specifically, the early observations of the Swift/X-ray in short and long GRBs is expected to be comparable, since geometrical effects are of a similar nature in both types. An alternative interpretation invokes the formation of a neutron star 1992; Dai & Lu 1998a,b; Gao & Fan 2006; Metzger et al. 2011; Dall'Osso et al. 2011; Ronchini et al. 2023). A fascinating consequence of this scenario is that the incidence of SGRB plateaus

> See also Dereli-Bégué et al. (2022) for an alternative interpretation (for a small sample of long GRBs) based on a low bulk Lorentz factor and a low-density wind medium

> > A73, page 1 of 13

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Guglielmi, Stratta, Dall'Osso et al. 2024 A&A 692, 73





GRB remnant: BH or NS?

Long GRBs or **Type I GRBs** from collapsars



Credit: ESO

GRB afterglow properties can shed light on the remnant nature and production efficiency Short GRBs or **Type II GRBs** from compact binary mergers



Ascenzi+2011



GRB afterglow



B. Zhang 2019

GRB afterglow



2019 Zhang (

- \bullet following a power-law
- ullet
- At late epochs (~day), $\alpha \geq$ 1.2-1.5 and $\Delta \beta > 0$ \bullet

After min-hrs from the burst -> slow cooling regime of the bulk of the electrons -> X-ray flux decays

Assuming the jet is plunging into a constant density ISM, at early epochs $\alpha \geq 0.7-0.8$

(e.g. Granot & Sari 2002, Zhang et al. 2006)



X-ray afterglow plateaus



Tang et al. 2019

Early X-ray flux evolution (a few x100s-10ks) in the majority of GRBs is too shallow



X-ray afterglow plateaus





Plateau phase strongly challenges the standard sync. radiation scenario

X-ray afterglow plateaus



 $L_{p,x} \propto T_{p,x}^{-1}$

energy reservoir?

from with central engine?

X-ray plateau origin from spin-down NS

- $E_{NS} = 0.5I\omega^2 \sim 3 \times 10^{52}$ erg for P~1ms
- $L_{SD}(t) = \frac{E_{NS}}{\tau(1 + t/\tau)^2}$ (e.g. Zhang & Meszaros 2001) accurately reproduce X-ray plateaus (e.g. Dall'Osso+2011, Rowlinson+2013, Bernardini+2013, Stratta+2018)
- $L_{SD}(\tau) \propto \tau^{-1}$ naturally reproduces the L_{x,p} and T_{x,p} anticorrelation

e.g Lyons+2010; Bernardini+2012; Rowlinson+2013; Bernardini+2013; Stratta+2018 Stratta+2022; <u>Dall'Osso+2023</u>



Dall'Osso+2011

X-ray plateau origin from structured jet effects

- •X-ray plateaus are <u>afterglow emission</u> <u>viewed slightly off-axis</u> with $\theta_{obs} \ge \theta_j$ (e.g. Eichler & Granot 2006, Beniamini+2020)
- •X-ray plateaus are <u>prompt high-latitude</u> <u>emission</u> dominating the afterglow emission (Oganesyan+2020)
- So far, no extensive test of these models on the large available data set has been published yet





X-ray plateaus in GRBs from NS-NS mergers



Adapted from Ascenzi+2011

• X-ray plateaus in the afterglow of GRBs from NS-NS mergers mark the formation of a NS remnant

 Incidence of X-rays plateaus in these GRBs —> proxy of the fraction of NS-NS mergers that generate a NS remnant

X-ray plateaus in GRBs from NS-NS mergers

- Initial sample: all TypeII GRBs with Swift/XRT observed from May 2005 to Dec 2021, with known redshift (based on Fong+2022, O'Connor+2022) -> 85 GRBs
- X-ray light curves from publicly available UK SDC Swift XRT Repository
- To robustly identify a plateau in the Xray afterglow, we requested XRT tot cts >100 (S/N > 10)



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* GRB 150101B: very late XRT observations (>1 day)



"LC fit": Incidence of plateaus

- We find **15/40 (37.5%)** X-ray aflterglow light curves are compatible with a broken power law, with:
 - initial decay index $\alpha < 0.75$ •
 - no evidence of significant spectral evolution before and after the temporal break









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"LC fit": Testing the magnetar scenario

Following Dall'Osso et al. 2011, **13/15** events could be modelled assuming energy injection into the forward shock from a newly-born spinning-down magnetar

GRB name	Inp	ut	Output					
	z	θ_j (deg)	<i>B</i> (10 ¹⁴ G)	P (ms)	$\tau_{\rm sd} \approx 680 \left(\frac{P_{\rm ms}}{B_{15}}\right)^2$ (ks)	χ^2	ν	
051221A	0.5464	6.0	(29 ± 2)	(12.8 ± 0.3)	13.3 ± 2.1	125	77	
060614	0.125	12.6	(37 ± 3)	(24 ± 1)	34 ± 9	1749	465	
070714B	0.923	8.6	(132 ± 40)	(11 ± 2)	0.5 ± 0.3	284	79	
090510	0.903	2.3	(82 ± 7)	(4.5 ± 0.2)	0.20 ± 0.04	80	63	
110402A	0.854	15.0	(96 ± 37)	(14 ± 1)	1.4 ± 1.1	42.0	16	
130603B	0.3568	6.3	(110 ± 2)	(13.2 ± 0.2)	0.98 ± 0.04	137	70	
140903A	0.3529	4.0	(32 ± 4)	(8 ± 0.3)	4.5 ± 1.2	56	36	
150424A	0.3	4.3	(36 ± 4)	(16 ± 1)	13 ± 3	243	115	
151229A	0.63	5.0	(67 ± 9)	(4.3 ± 0.9)	0.3 ± 0.1	142	56	
161001A	0.67	5.0	(47 ± 6)	(4.2 ± 0.2)	0.5 ± 0.1	88	54	
170728B	1.272	3.5	(20 ± 1)	(1.50 ± 0.03)	0.39 ± 0.04	232	193	
210323A	0.733	2.9	(51 ± 11)	(8.7 ± 0.6)	2.0 ± 0.9	62	18	
211211A	0.0763	6.9	(286 ± 26)	(27 ± 1)	0.6 ± 0.1	779	265	





Two discarded events:

- GRB180618A has too short plateau duration to allow model parameter estimates
- GRB061201 provided not plausible B and P values (B>5.6x10¹⁶G, P>38ms)

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"LC fit": Testing the magnetar scenario

Following Dall'Osso et al. 2011, 13/15 events could be modelled assuming energy injection into the forward shock from a newly-born spinning-down magne

GRB name	-	With respect to the whole analysed sample (85): 15%					Construction of the second sec				
051221A 060614 070714B 090510	•	By one	excluc es (25	ling fror low S/N	n the sa + 15010	mpl)1B)	e t and	he "inconclusive" I the EE (40): 32.5%	4 5 6 e from trigger [s]		
110402A 130603B	c										
140903A	0.3529	4.0	(32 ± 4)	(8 ± 0.3)	4.5 ± 1.2	56	36	True disconded success			
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210323A	0.733	2.9	(51 ± 11)	(8.7 ± 0.6)	2.0 ± 0.9	62	18				
211211A	0.0763	6.9	(286 ± 26)	(27 ± 1)	0.6 ± 0.1	779	265	values (B>5.6x10 ¹⁶ G, P>38ms)			



GRB051221A [0.1-30.0 keV]

S/N-reiected (25) GRB 150101B (1 E-rejected (19) LC fit (40)





"EE-rejected": excluding the presence of a magnetar

- If a magnetar remnant is formed, we can define a plausible minimum spin-down luminosity Lmin ~ 3x10⁴⁵ erg/s by assuming:
 - Maximum spin-down timescales of ~10⁵ s
 - Maximum spin period of P ~ 30 ms, corresponding to a rotational energy of ~3x10⁴⁹ erg
- Following Dall'Osso et al. 2023, we also compare the minimum luminosity detected with XRT with the prompt emission luminosity, excluding ratios larger >30
- We found that 9/19 "EE" light curves, clearly show fainter fluxes, incompatible with the presence of a magnetar
- The remainder 10 EE events were declared "inconclusive"



S/N-rejecte
GRB 15010
EE-rejected
LC fit (40)

GRB 080905A



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- 15 GRBs show an X-ray plateau
- For 13 events, plateau is compatible with the magnetar scenario
 - With respect to the whole analysed sample (85): **15%**
 - By excluding from the sample the <u>"inconclusive" ones</u> (25 low S/N + 150101B + 10 "EE") (49): **26.5%**



Compatibility with predicted NS maximum mass



- We computed the **probability** distribution of remnant mass by adopting a double-peaked Gaussian for the mass distribution of NSs in binaries
- The percentage of GRB compatible with a magnetar remnant (15-26%) is reached for a remnant mass of Mrem=(2.31–2.41) M_{\odot}
- Depending on the NS stability Mrem>~M_{TOV} -> compatible with most M_{TOV} estimates so far (e.g. Margalit et al. 2022)



Conclusions

- that a fraction of ~ 1/6 1/4 is compatible with originating a magnetar remnant
- ullet2025)

GRB X-ray afterglow plateaus likely indicate the presence of a magnetar remnant -> potential

FRB sources where to look at given the accurate sky localisation from afterglow MW campaigns

• By analyzing all Swift Type II GRBs detected from 2005 to 2021 with known redshift, we found

Larger samples of identified magnetars can be achieved with future BNS merger and postmerger detections with next generation GW interferometers (as ET), for which accurate sky localisation will be provided through not collimated Kilonova counterparts (e.g. Loffredo et al.







Extra slides



The events with / without plateau, including the faintest ones ("inconclusive"), have similar redshift and energy distribution, suggesting that the lack of plateau evidence is not due to biases against distant or faint events

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$$\kappa = L_{\rm min}/L_{\rm sd} \approx 1.2 \times 10^5 \epsilon P^{5/3} \left(R_6 M_{1.4}^{2/3} \right)^{-1} (\xi/0.5)^{7/2}$$

harbouring a magnetar

Bernardini+2015

EE events with k>30 are incompatible with



