

Xrays from Massive-Stars

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Collisional shock-heated X-ray emission

Rough rule of thumb:

$$\frac{T_x}{10MK} \approx \frac{E_x}{1kev} \approx \left(\frac{V}{1000\text{km/s}} \right)^2$$

$E_x \sim 2\text{-}10 \text{ keV} \Rightarrow V \sim 3000 \text{ km/s} \Rightarrow$ **Colliding wind binaries**

$E_x \sim 0.5\text{-}1 \text{ keV} \Rightarrow V \sim 700 \text{ km/s} \Rightarrow$ **Intrinsic instabilities**

$E_x \sim 1\text{-}3 \text{ keV} \Rightarrow V \sim 1000 \text{ km/s} \Rightarrow$ **Magnetic confinement**

Radiative cooling

$$\frac{t_{cool}}{t_{exp}} \approx \frac{V_8 d_{12}}{\dot{M}_{-7}}$$

Stevens+ 1992

$$V_8 \equiv \frac{V}{10^8 \text{cm/s}} = \frac{V}{1000 \text{km/s}}$$

$$d_{12} \equiv \frac{d}{10^{12} \text{cm}} \approx \frac{d}{0.1 \text{au}}$$

$$\dot{M}_{-7} \equiv \frac{\dot{M}}{10^{-7} M_{\odot} / \text{yr}}$$

Colliding Wind Shock X-ray emission

$$\chi \equiv \frac{t_{cool}}{t_{exp}} \approx \frac{l}{r} \gg 1$$

Adiabatic shock for binary separation D

$$L_x \sim \int \rho^2 dV \sim \left(\frac{\dot{M}}{V_\infty} \right)^2 \int_R^D \frac{r^2 dr}{r^4} \sim \frac{1}{D}$$

$$\chi \equiv \frac{t_{cool}}{t_{exp}} \approx \frac{l}{r} \ll 1$$

Radiative shocks: insensitive to D

$$L_x \lesssim L_{wind} \equiv \dot{M} v_\infty^2 / 2 \text{ for 1-D spherical, laminar shock}$$

COLLIDING WINDS FROM EARLY-TYPE STARS IN BINARY SYSTEMS

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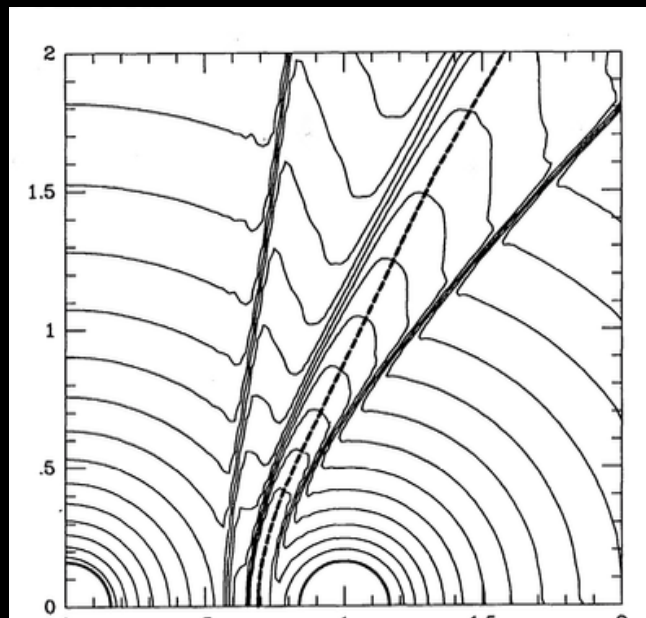
A. M. T. POLLOCK

Computer & Scientific Co., Ltd., 34 Westwood Road, Sheffield S11 7EY, England

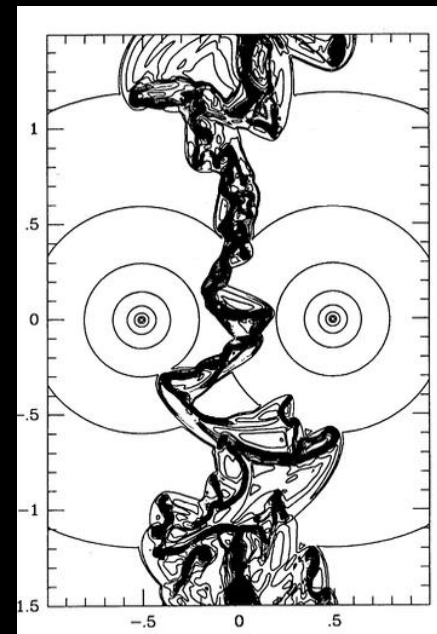
Received 1991 May 23; accepted 1991 August 15

$$\chi = \frac{t_{\text{cool}}}{t_{\text{esc}}} \approx \frac{v_8^4 d_{12}}{\dot{M}_{-7}},$$

$\chi \gg 1$; adiabatic



$\chi \ll 1$; radiative



Suppression of X-rays from radiative shocks by their thin-shell instability

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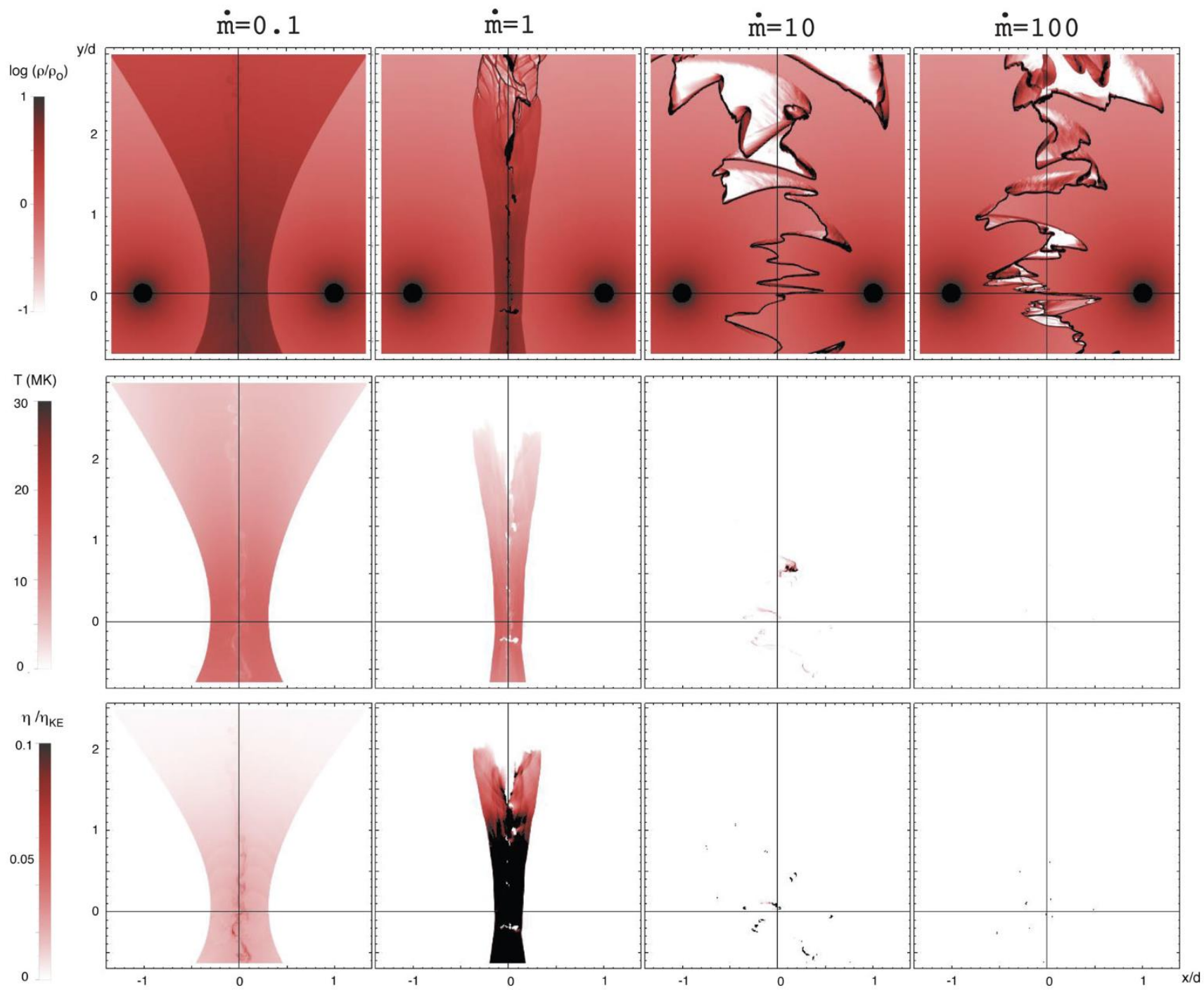
²Penn State Worthington Scranton, Dunmore, PA 18512, USA

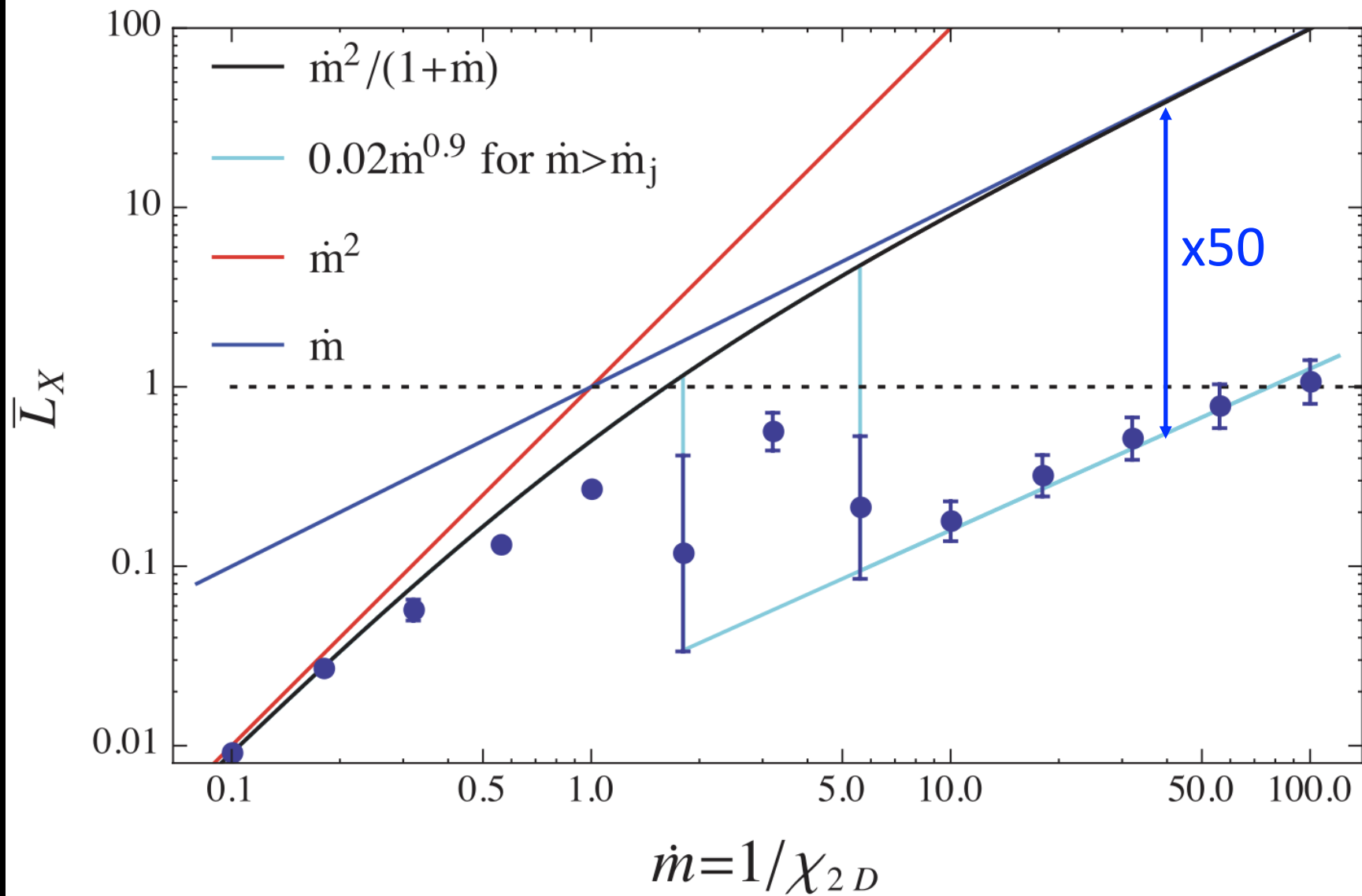
for planar
shock, expect $\frac{L_x}{L_{KE}} \approx \frac{\dot{m}^2}{1 + \dot{m}}$

$$\dot{m} \equiv 1/\chi$$

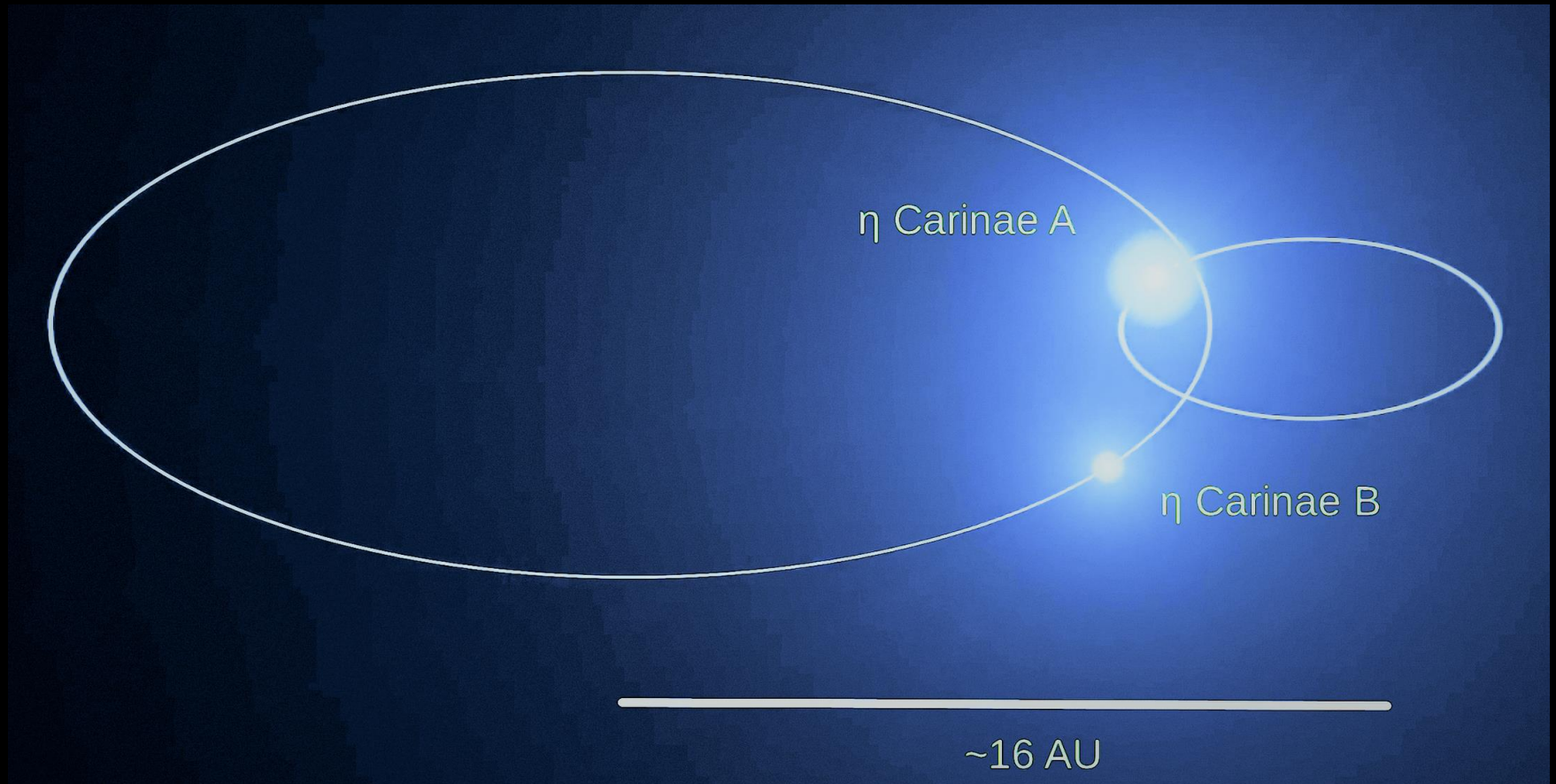
$$\chi = \frac{t_{\text{cool}}}{t_{\text{esc}}} \approx \frac{v_8^4 d_{12}}{\dot{M}_{-7}},$$

but for $\dot{m} \gg 1$,
thin-shell gives $\frac{L_x}{L_{KE}} \approx \frac{\dot{m}}{50}$





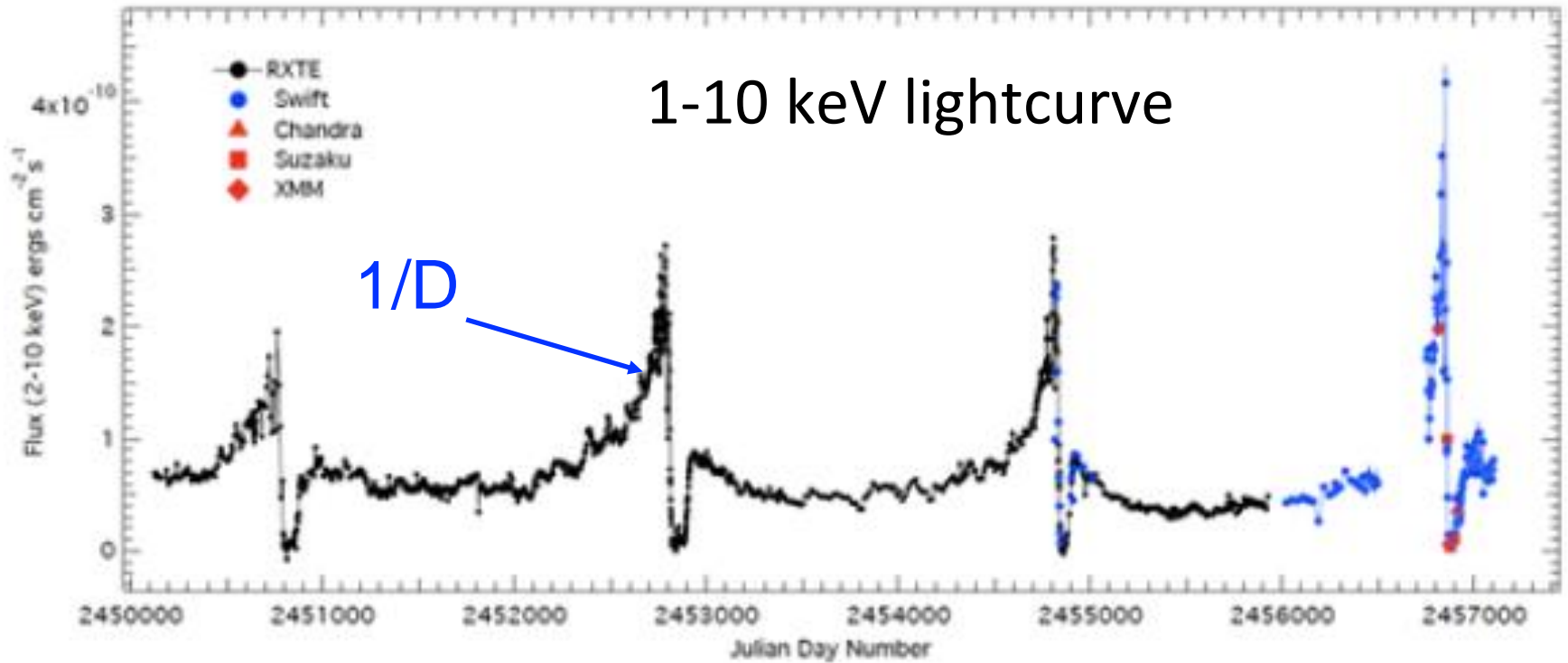
Colliding winds in eta Carinae



5.54 yr (2023 d) binary, $a \sim 16 \text{ au}$, $e = 0.9$

$M_A + M_B \sim 130 M_\odot$

5.54 yr (2023 d) binary



1998

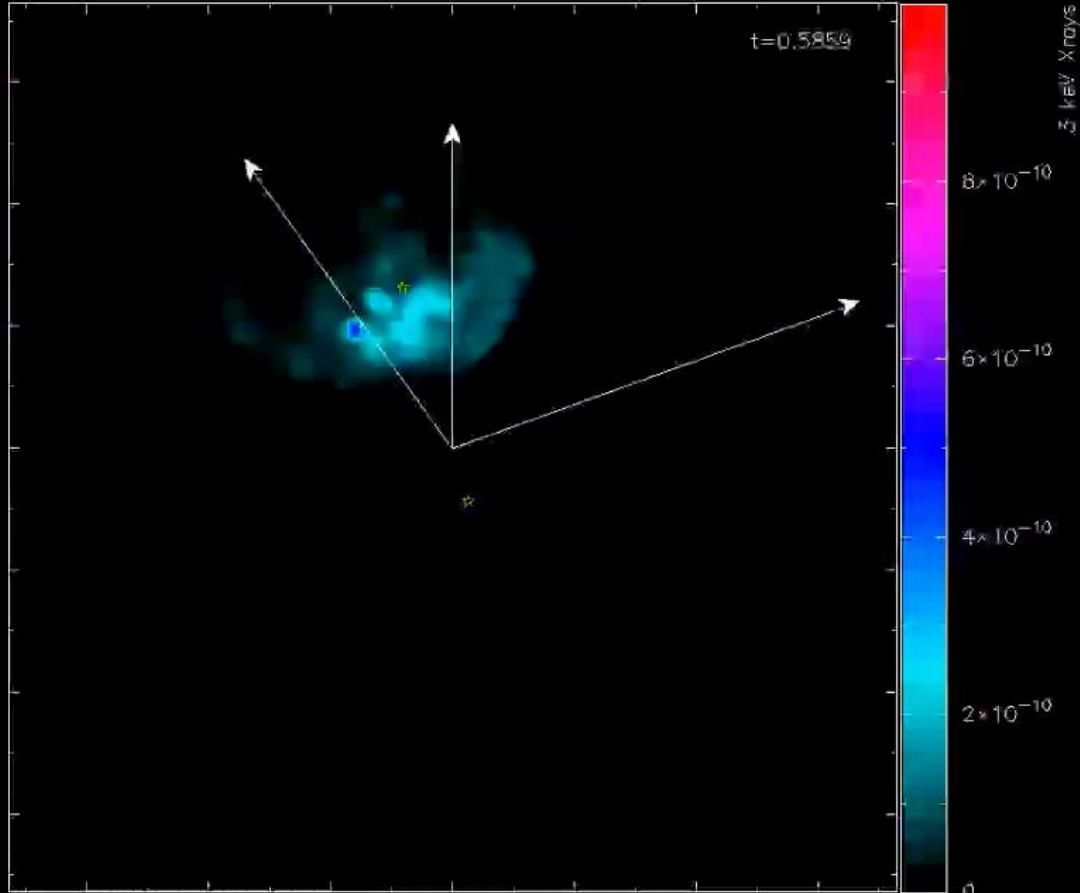
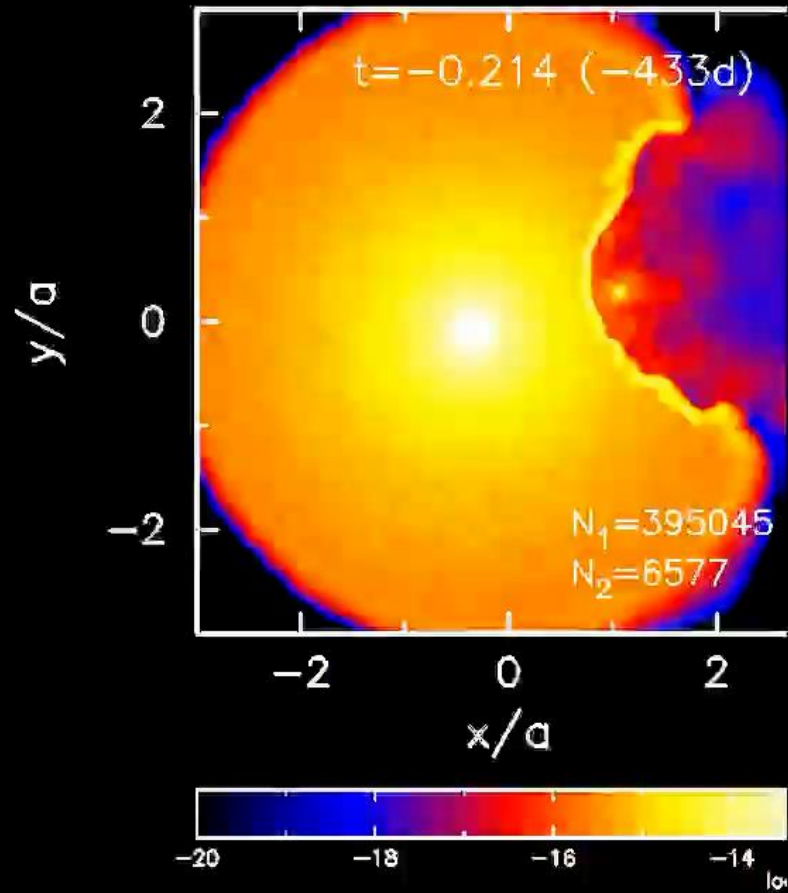
2003

2009

2014

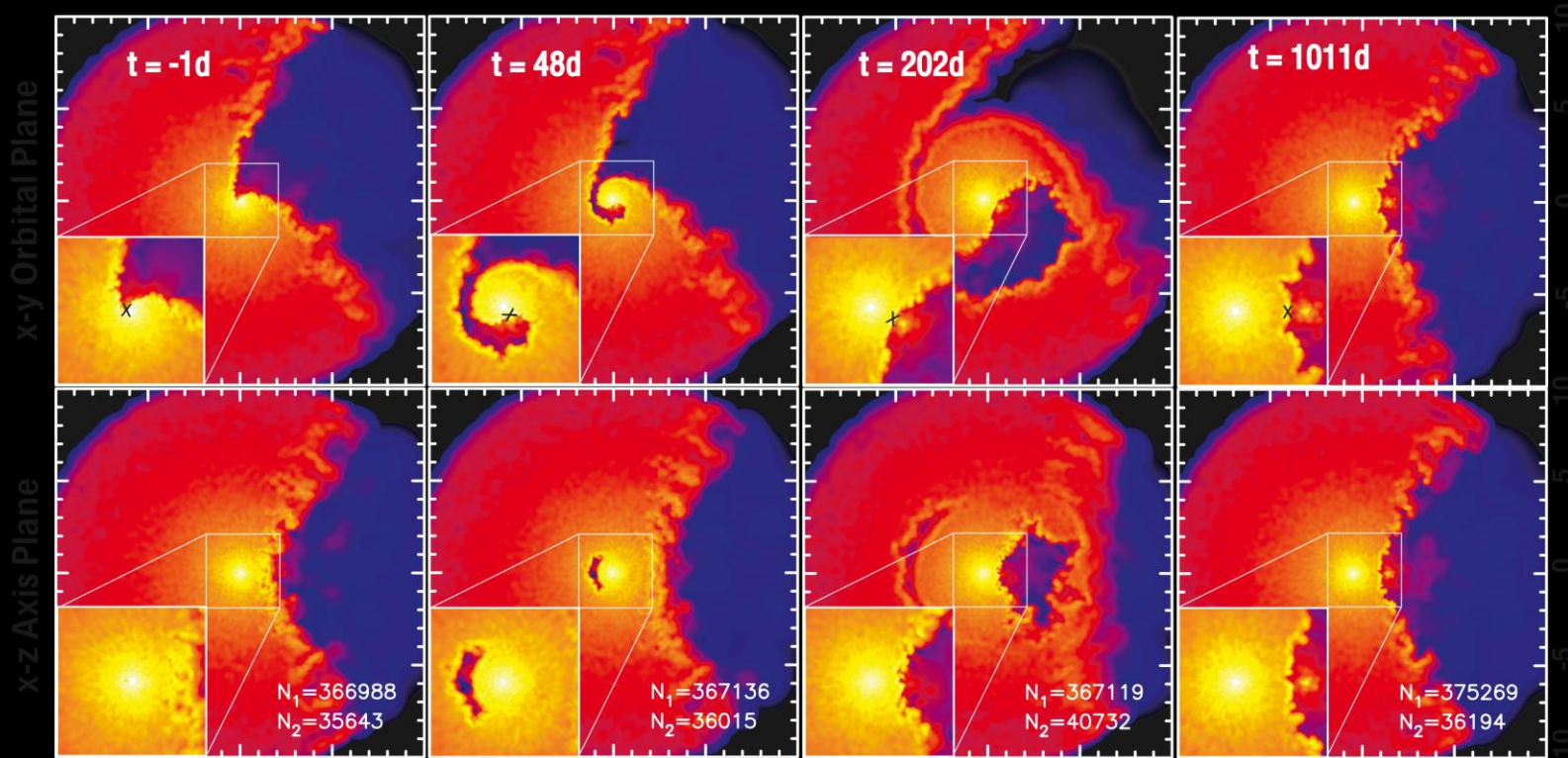
density

X-rays

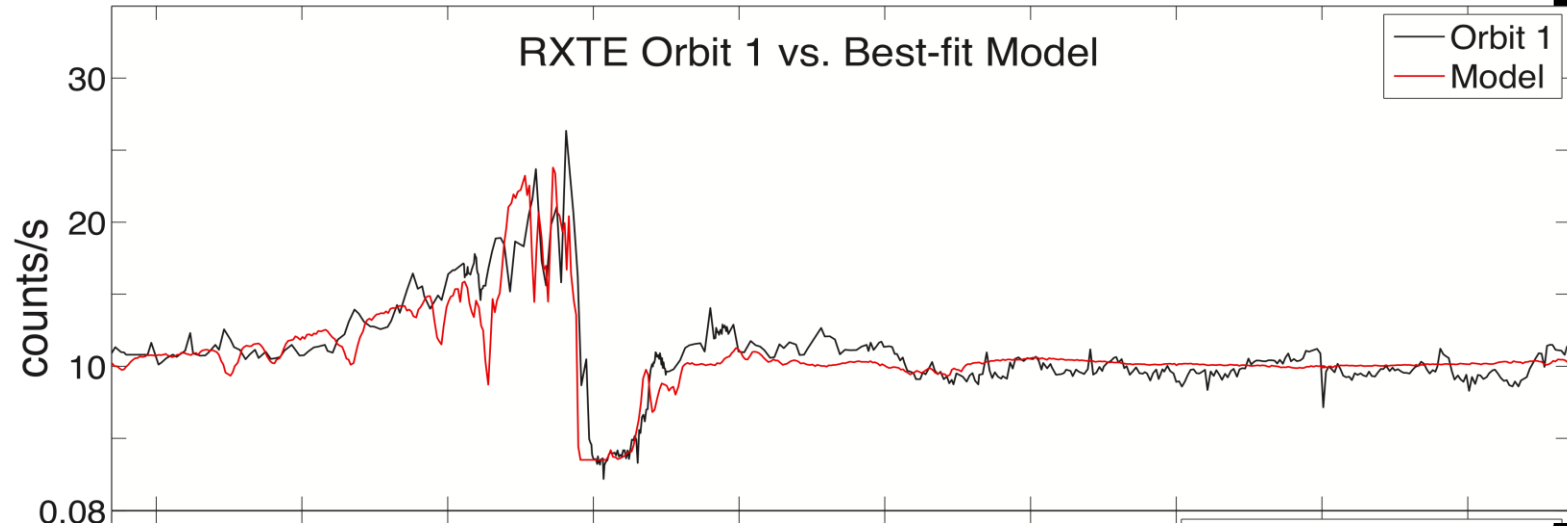


courtesy C. Russel

3D SPH sim of η Car CWB

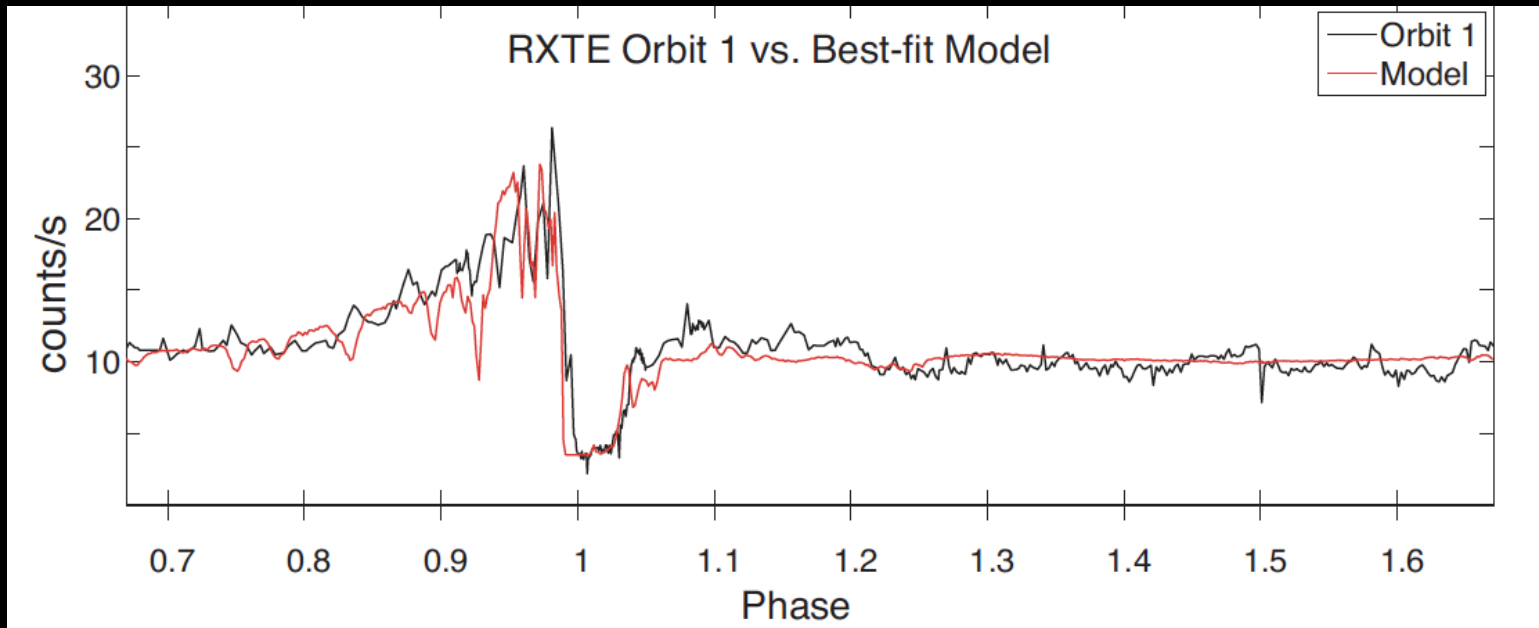


RXTE Orbit 1 vs. Best-fit Model



1-10 keV X-ray lightcurve from 5.54 yr (2023 d) binary

C. Russel, UDel Ph.D. thesis 2014



$$\dot{M}_A = 10^{-3} \frac{M_\odot}{\text{yr}}$$

$$\dot{M}_B = 10^{-5} \frac{M_\odot}{\text{yr}}$$

Star B likely
a stripped
WR star

$$V_A = 600 \frac{\text{km}}{\text{s}}$$

$$V_B = 3000 \frac{\text{km}}{\text{s}}$$

$$L_A = 5 \times 10^6 L_\odot$$

$$L_B < 10^6 L_\odot$$

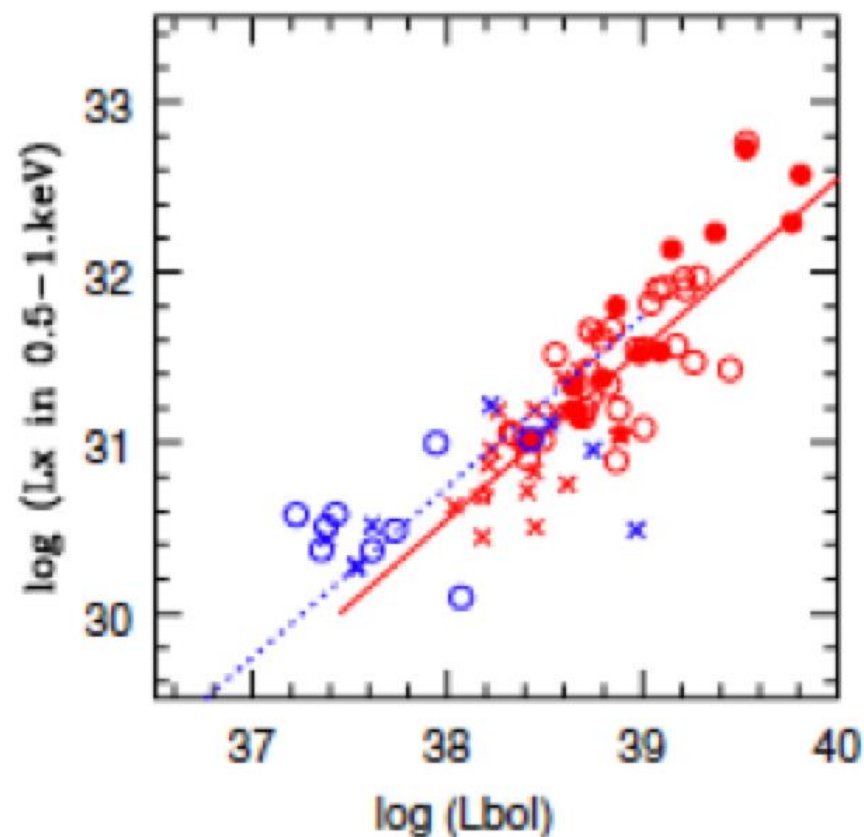
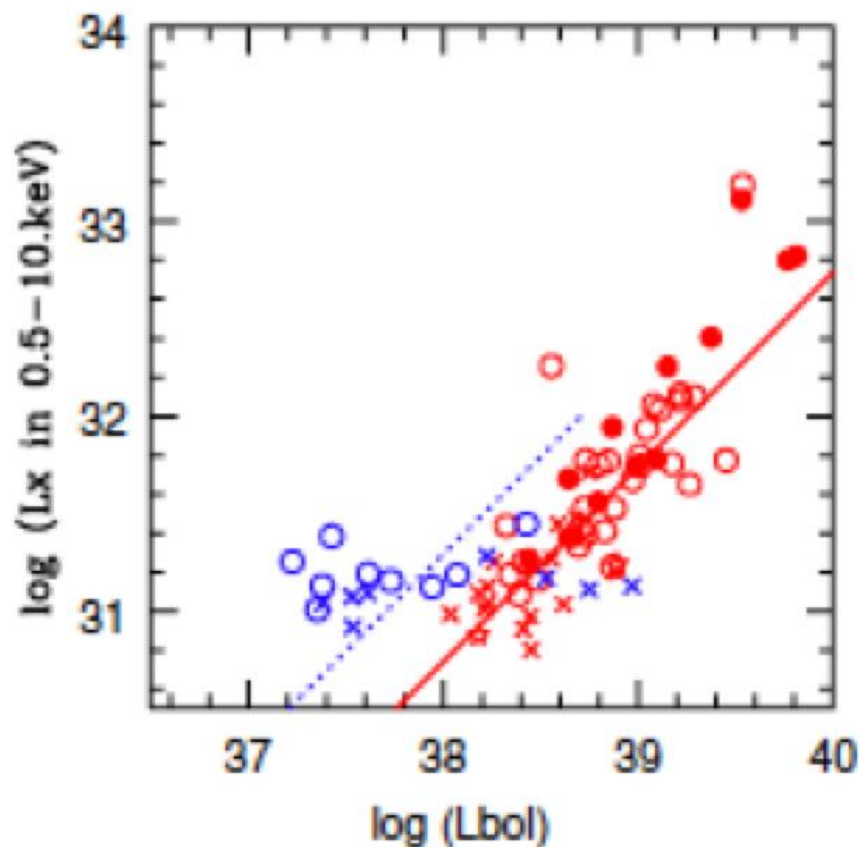
X-rays from single OB Stars

- Single O-stars emit soft (~ 0.5 keV) X-rays
- Thought to arise from **Embedded Wind Shocks**
- **EWS** arise from intrinsic **Line-Driving Instability**
- Observed scaling is $L_x \sim 10^{-7} L_{\text{bol}}$
- EWS theory predicts $L_x \sim \dot{M} / V_\infty \sim (L_{\text{bol}})^{1.7}$
- Reconcile here with “**thin shell mixing**” of shocks

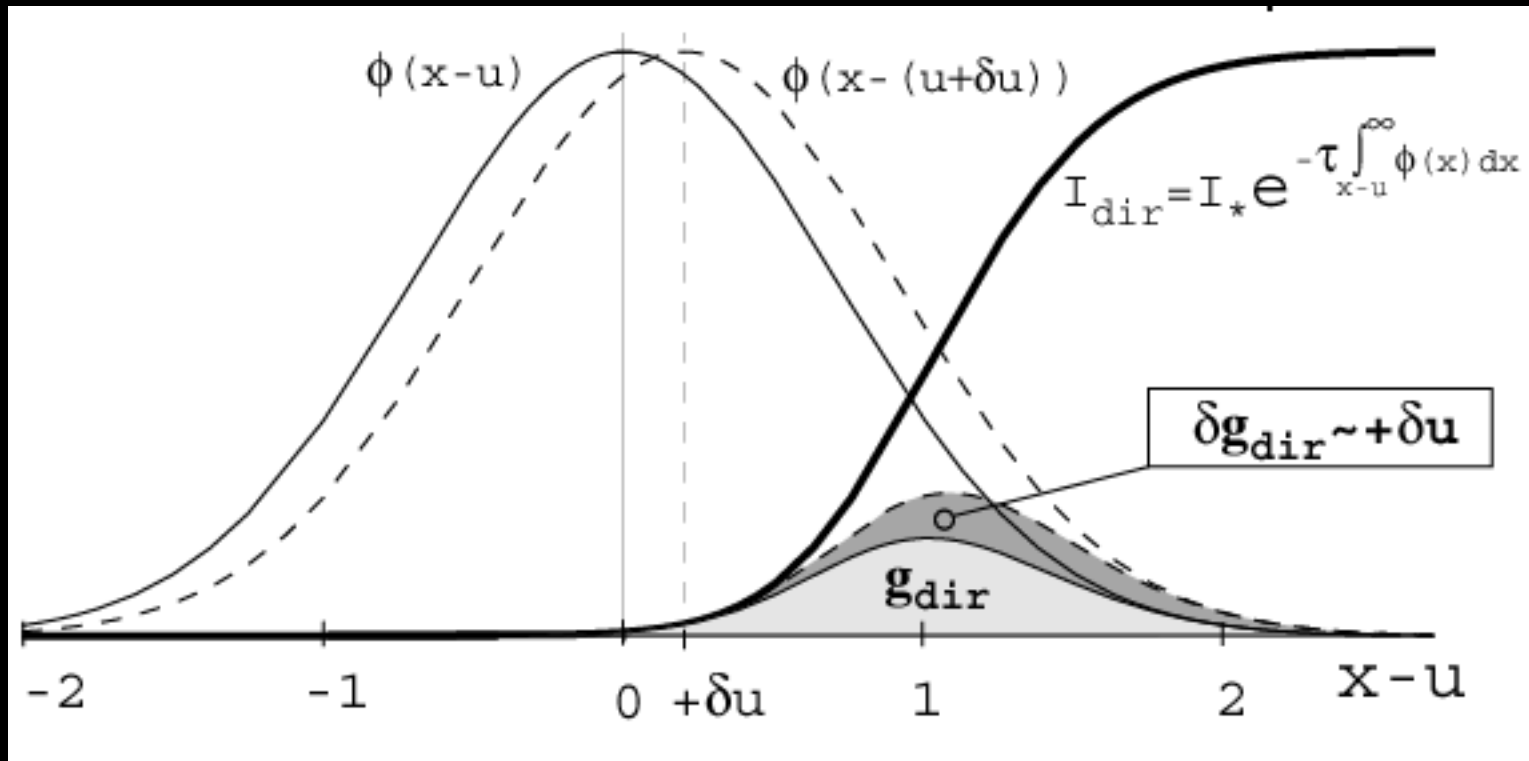
$L_x \sim 10^{-7} L_{bol}$ for Chandra observations of Carina OB stars

Naze et al. 2011

O (red dots) and B (blue dots) stars



Line-Deshadowing Instability



for $\lambda < L_{\text{sob}}$:

$$i\omega = \delta g / \delta v$$

$$= +g_0 / v_{\text{th}} = \Omega$$

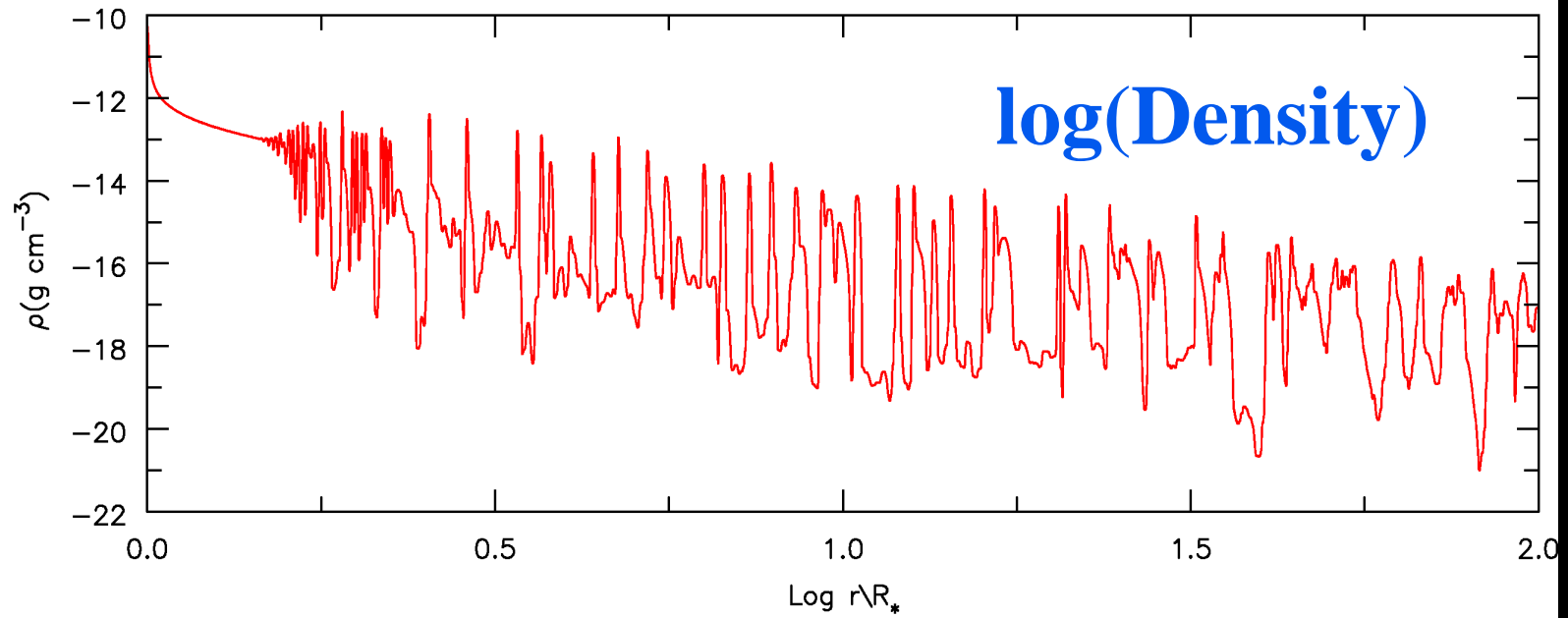
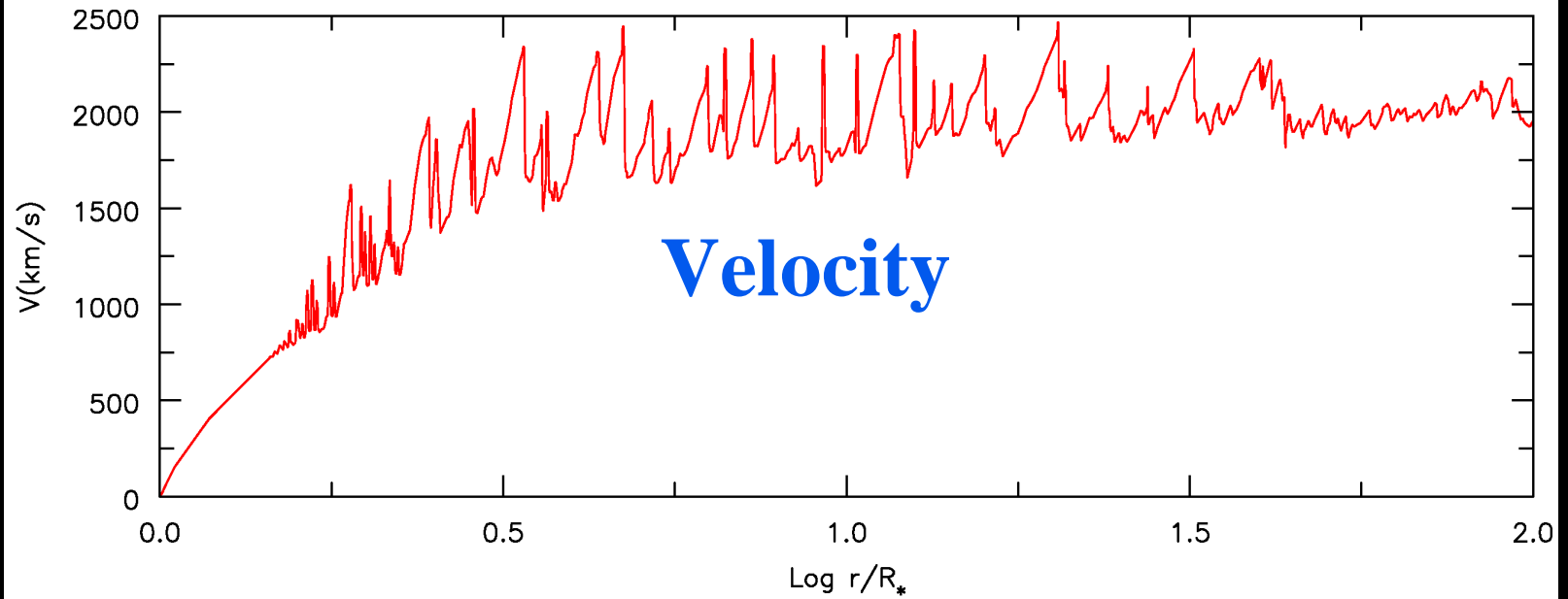


Instability with growth rate

$$\Omega \sim g_0 / v_{\text{th}} \sim v v' / v_{\text{th}} \sim v / L_{\text{sob}} \sim 100 v / R$$

e^{100} growth!

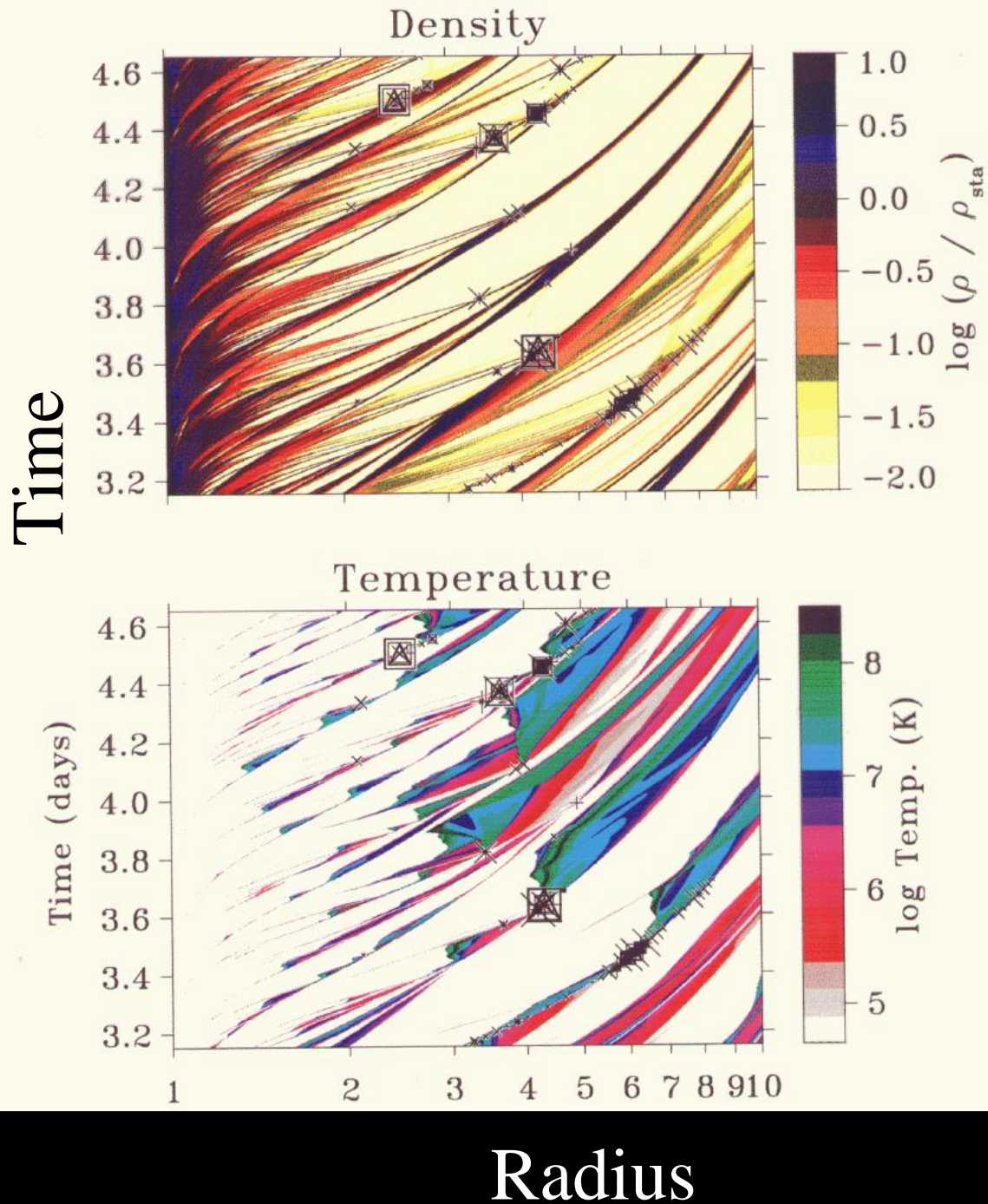
Time snapshot of wind structure vs. radius



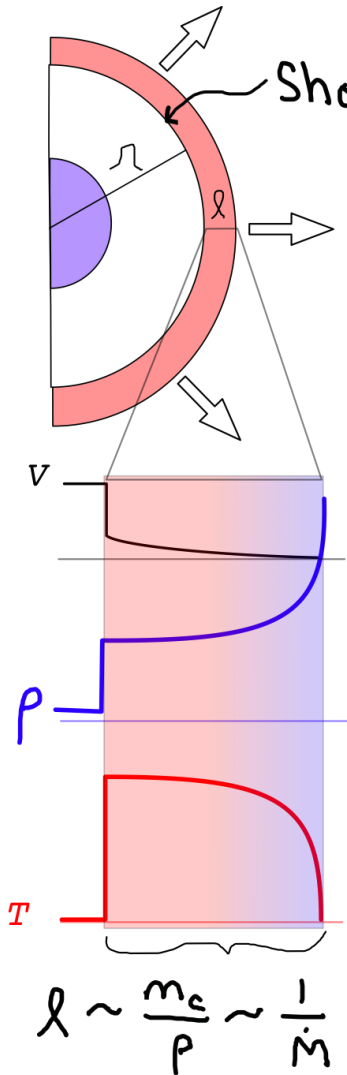
Turbulence-seeded clump collisions

Enhances V_{disp}
and thus X-ray
emission

Feldmeier et al.
1997



Shock cooling & thin-shell mixing



$$\dot{M} \sim L_{\text{Bol}}^{1/2} \sim L_{\text{Bol}}^{1.7}$$

$$L_x \sim f_x \dot{M} \Delta V^2$$

adiabatic $l \gg r$

$$f_x \sim \frac{r}{l} \Rightarrow$$

$$L_x \sim \dot{M}^2 \sim L_{\text{Bol}}^{3.4}$$

radiative $l \ll r$

$$f_x \sim \text{const.}$$

$$L_x \sim \dot{M} \sim L_{\text{Bol}}^{1.7}$$

+
thin-shell mixing

$$f_x \sim \left(\frac{l}{r}\right)^m$$

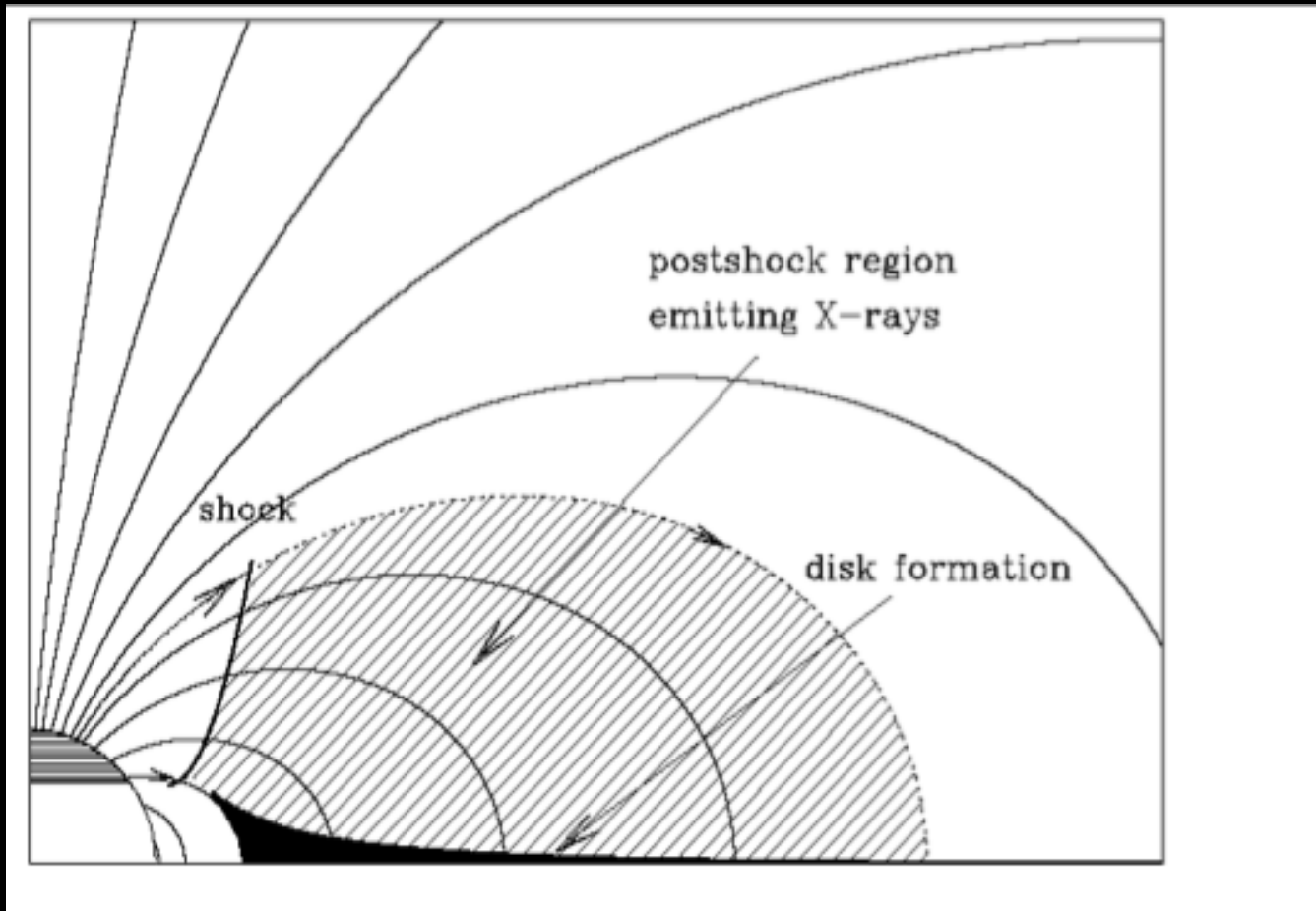
$$L_x \sim \dot{M}^{1-m} \sim L_{\text{Bol}}^{\frac{1-m}{1.7}}$$

if $m \approx 1 - \alpha \approx 0.4$

$$L_x \sim L_{\text{Bol}}^{\frac{1-m}{1.7}} \sim L_{\text{Bol}}$$

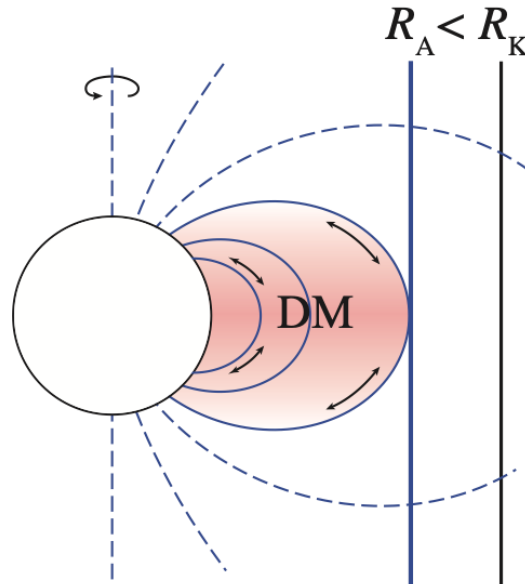
X-rays from Magnetically Confined Wind-Shocks

Babel & Montmerle 1997



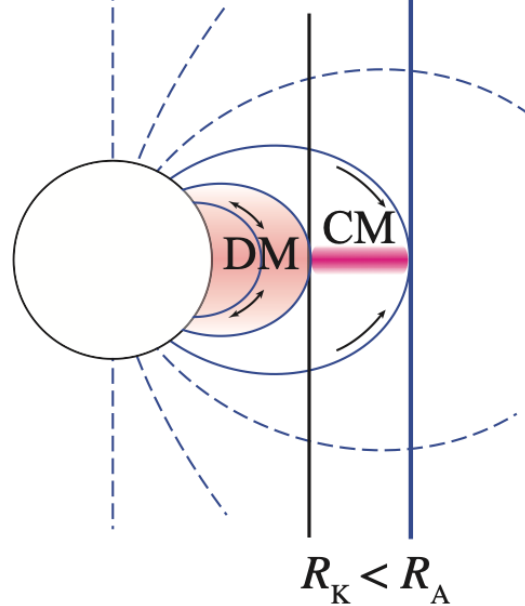
Dynamical Magnetospheres

$$R_A < R_K$$



Centrifugal Magnetospheres

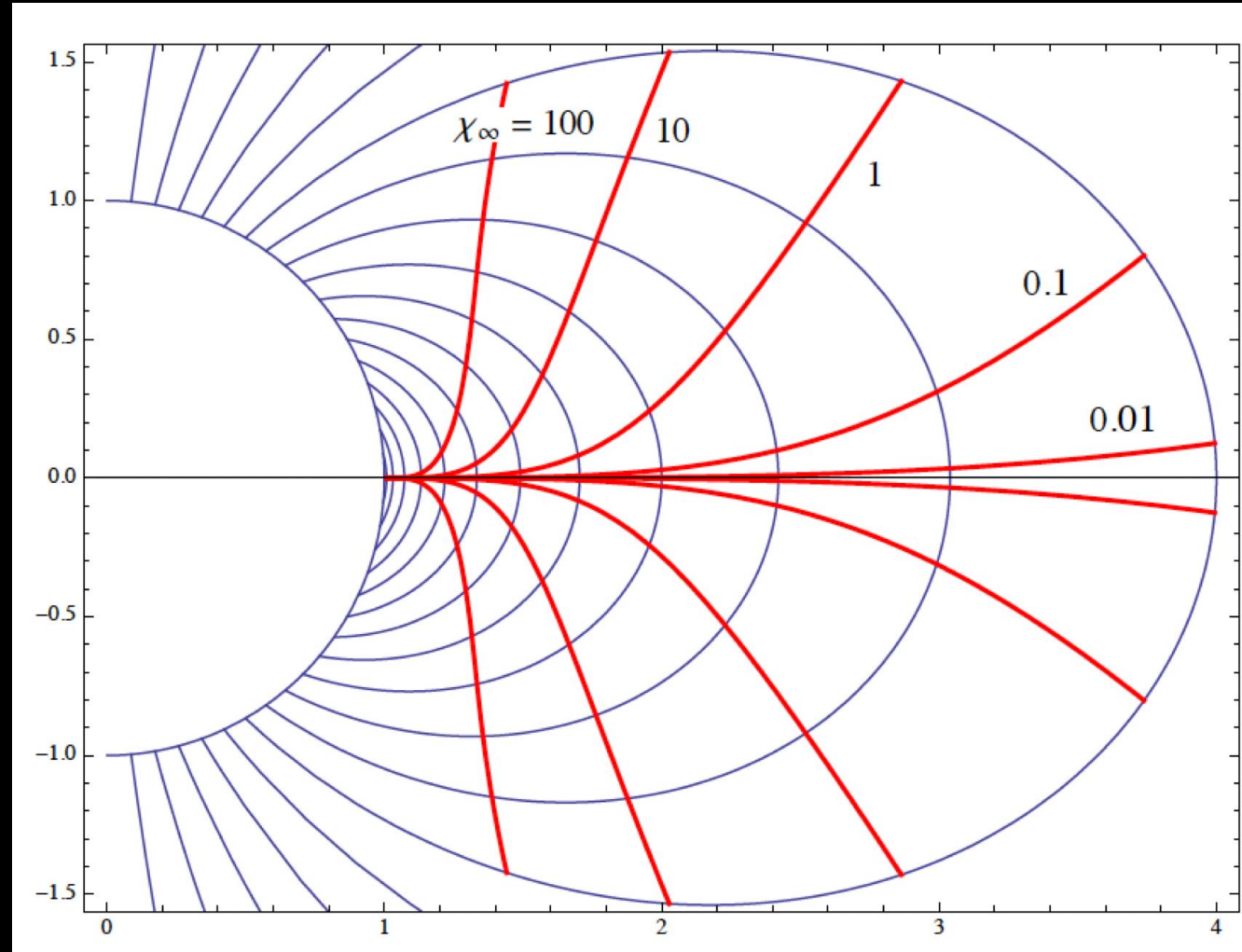
$$R_K < R_A$$



Shock retreat in a DM

χ_∞ measures in cooling **inefficiency**

$$\chi_\infty \equiv \frac{15\pi}{128} \frac{V_\infty^4 R_*}{\dot{M} \Lambda_m}$$
$$\approx 0.034 \frac{V_8^4 R_{12}}{\dot{M}_{-6}}$$



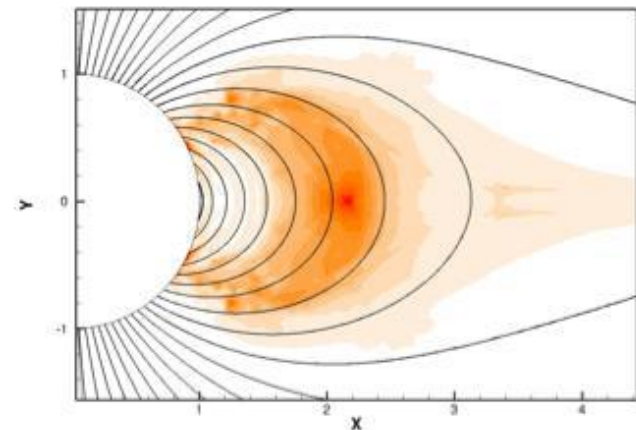
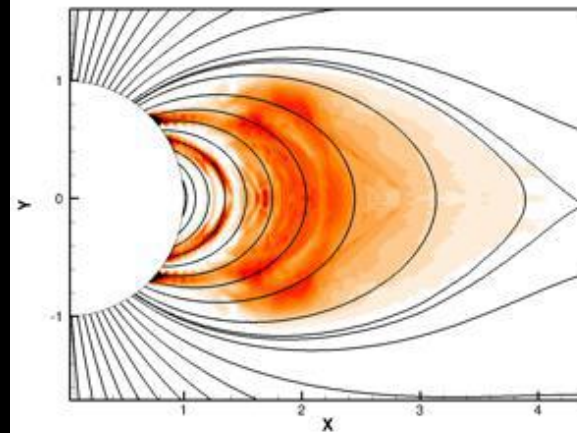
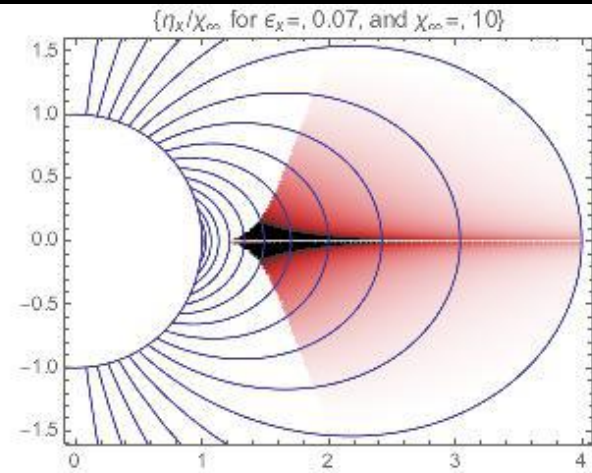
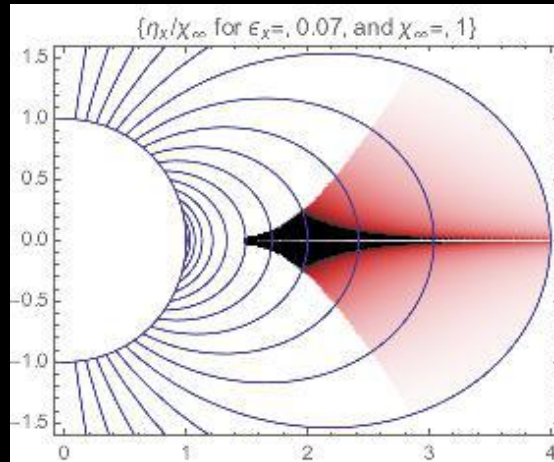
ADM vs. MHD: X-ray emissivity

moderate cooling

weak cooling

$$\chi_{\infty} = 1$$

$$\chi_{\infty} = 10$$



ADM

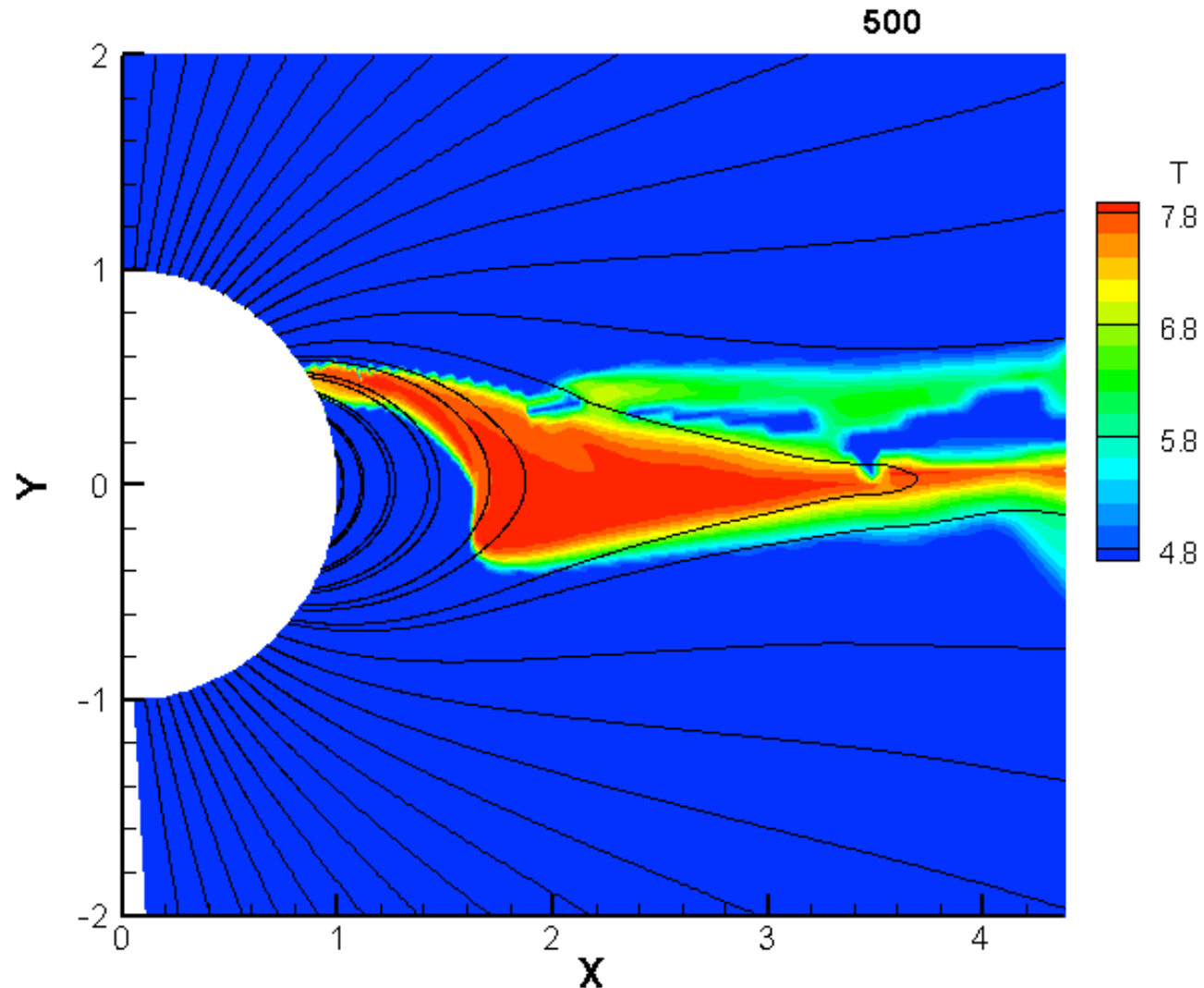
MHD

Magnetically Confined Wind Shocks

$\log T$

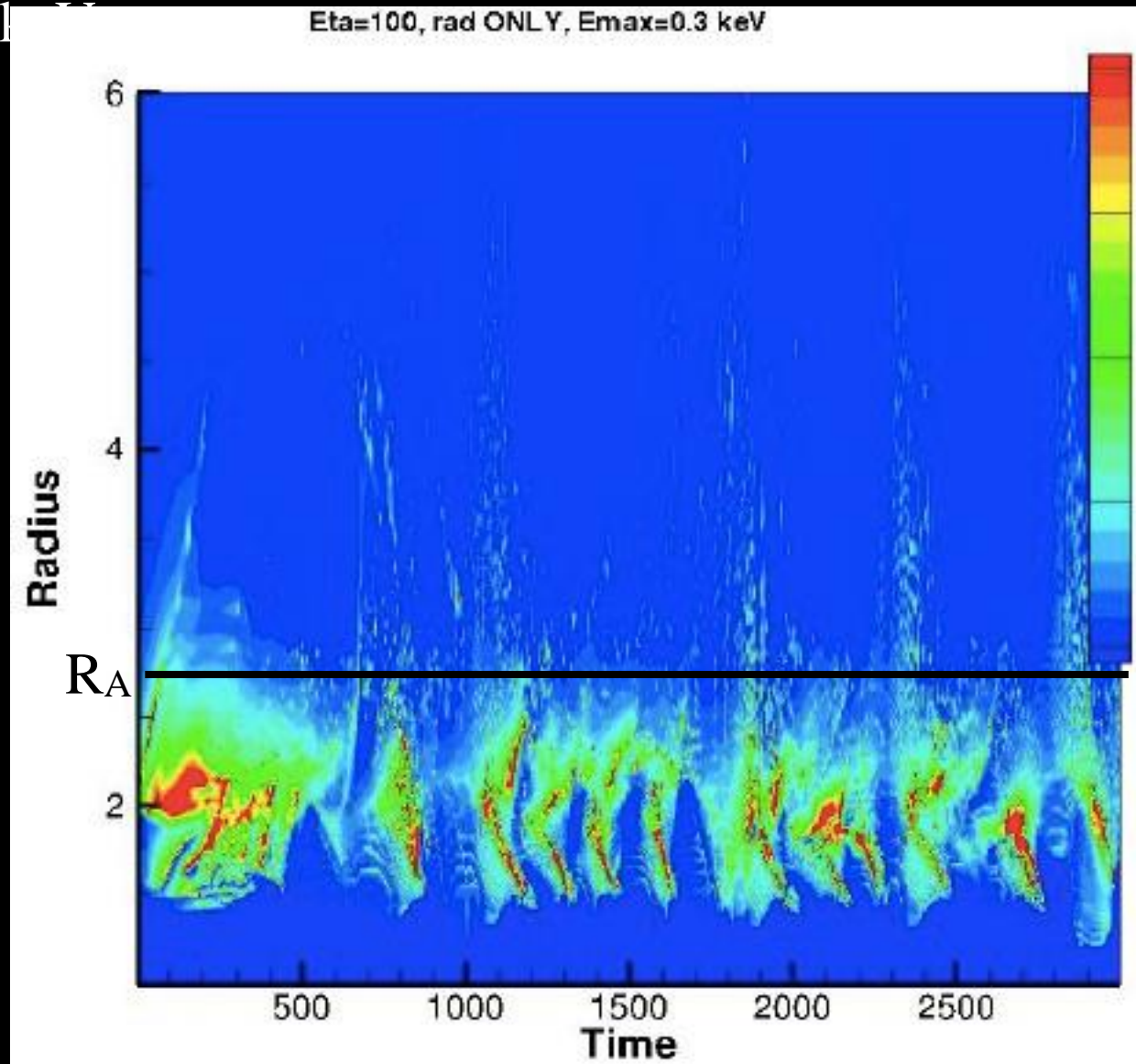
no
rotation
but now with
**Radiative
Cooling**

~ 2 keV
X-rays fit
Chandra
spectrum
for Θ^1 Ori C



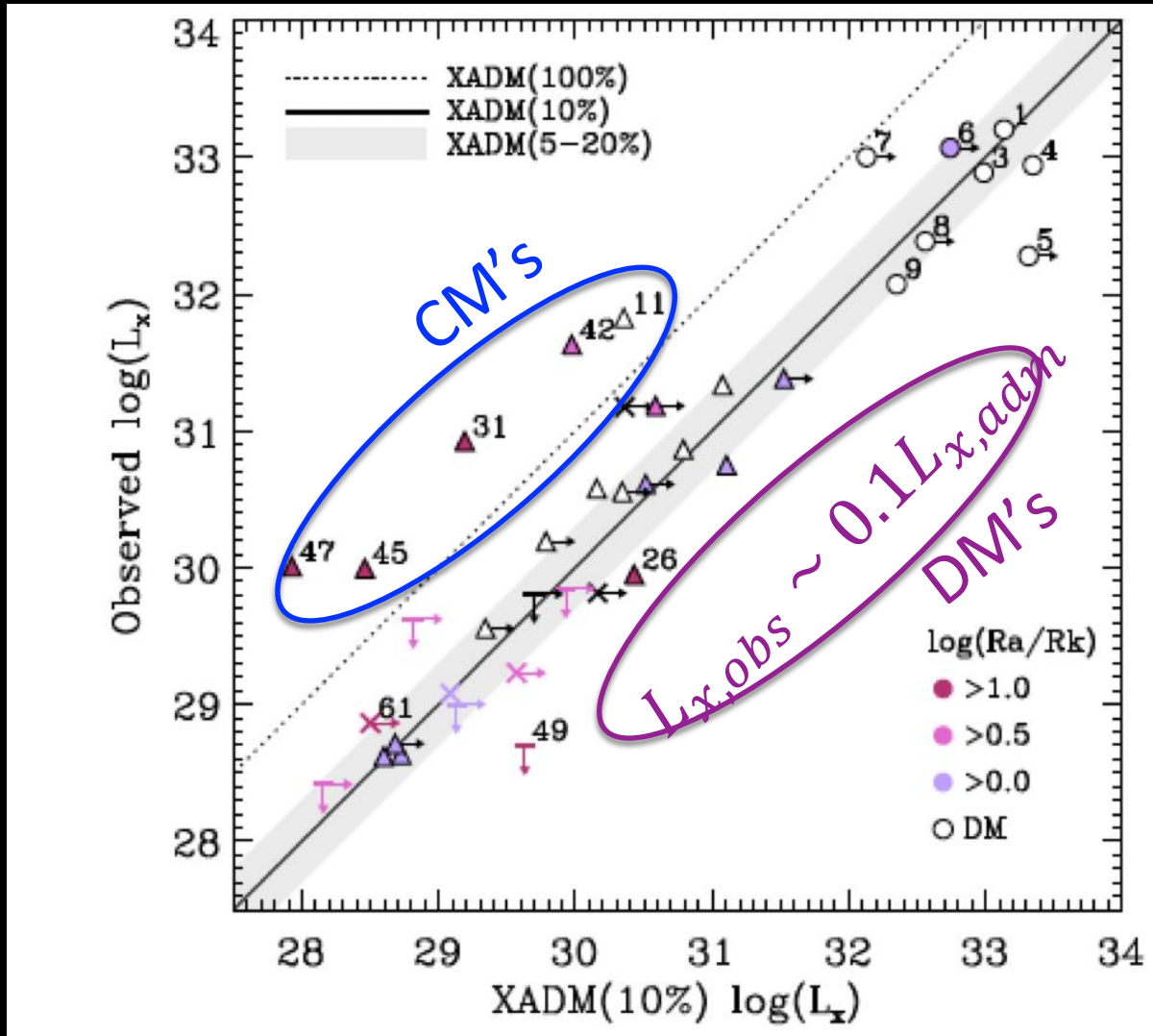
2D MHD: energy w/ rad cooling

r, t variation of (theta-averaged) X-rays > 0.3



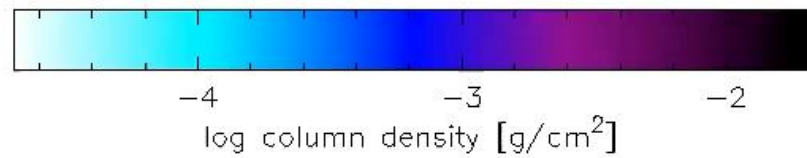
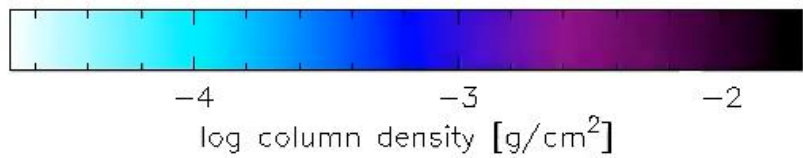
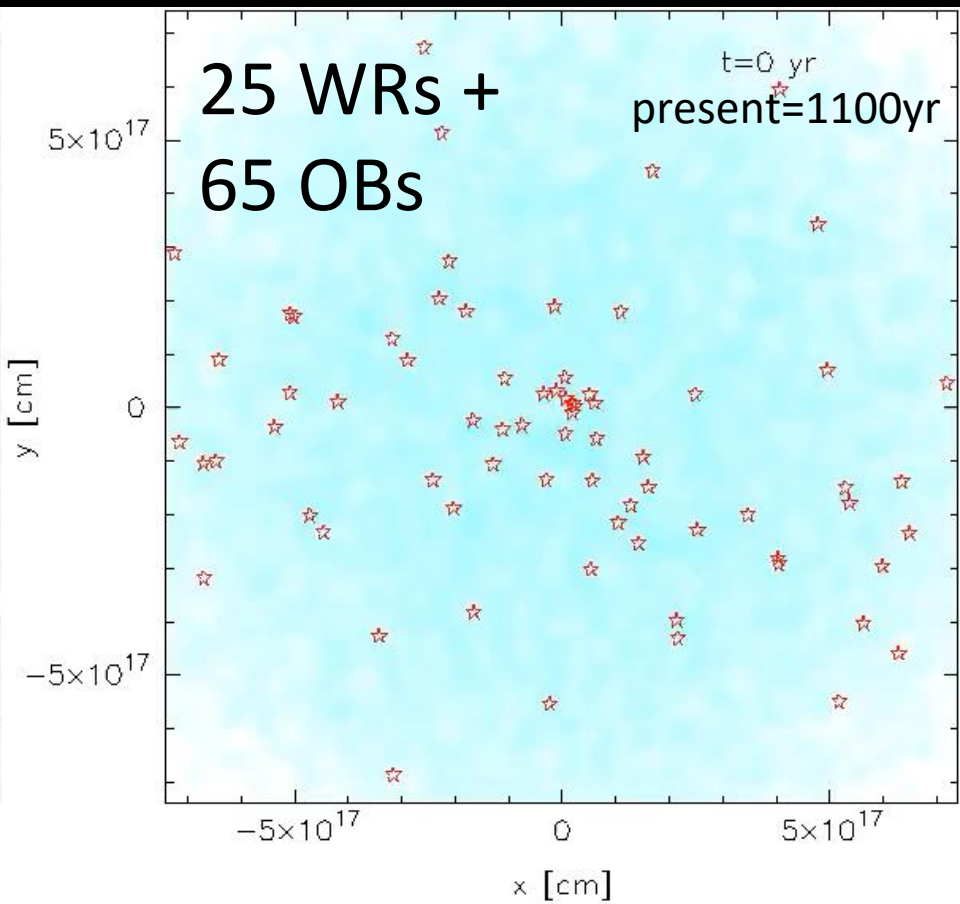
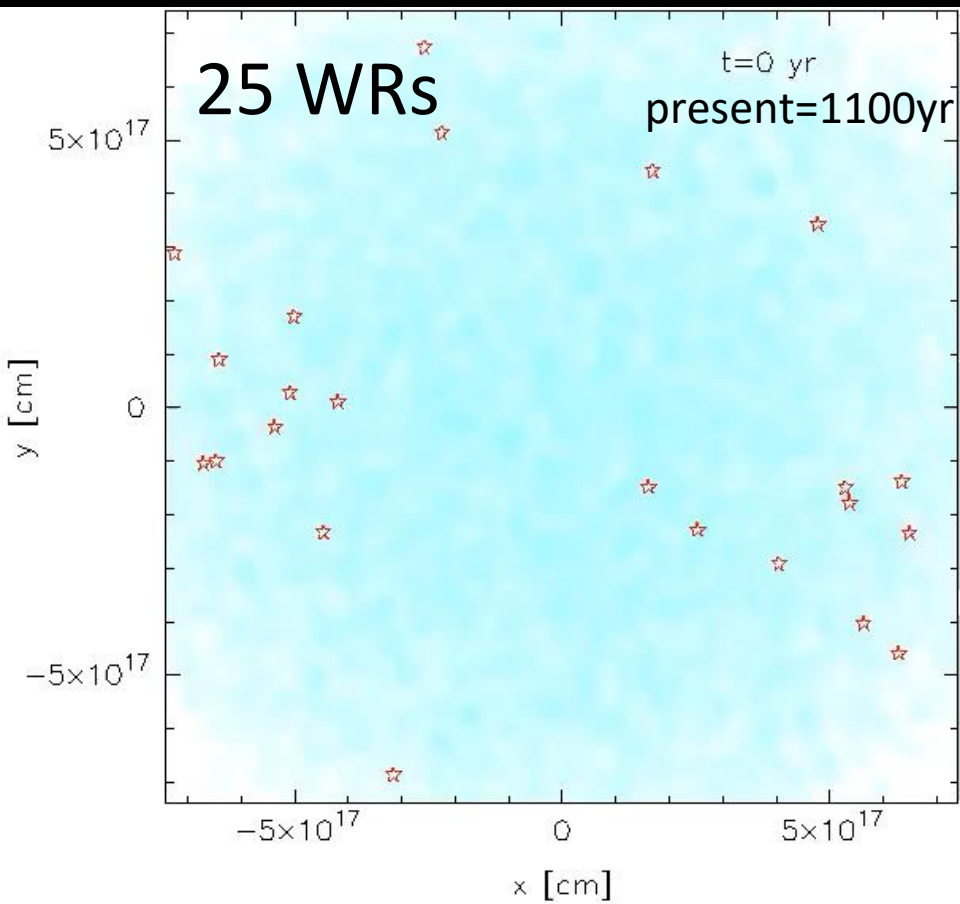
Lx vs. Mdot: obs. vs. XADM theory

Naze' et al. 2014



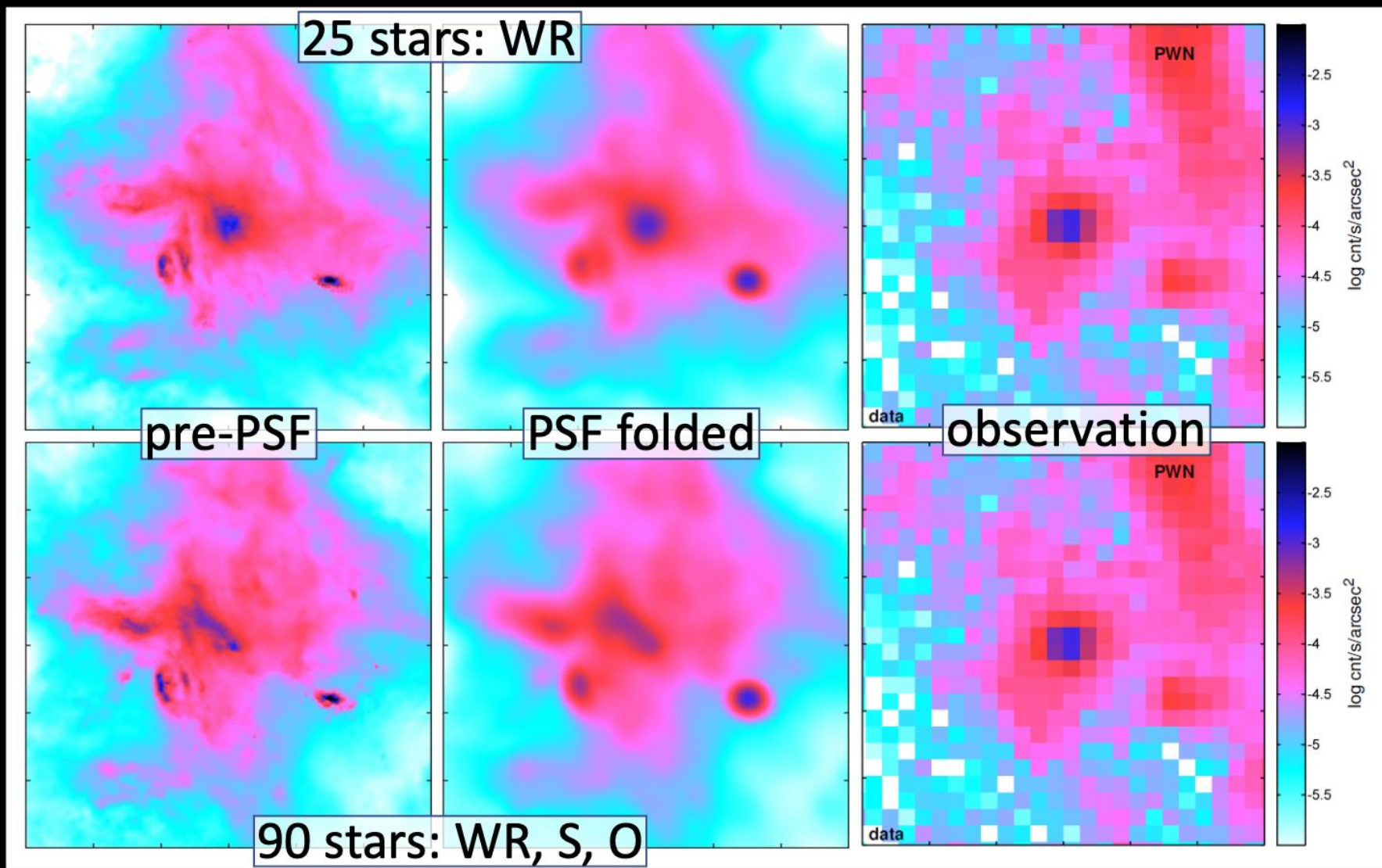
Hydro models of wind collisions in **gc cluster**

- ~30 Wolf-Rayet stars
 - 60+ earlier massive stars
- } smorgasbord of colliding winds
- Input constraints: orbital and wind properties from infrared
Paumard+06, Martin+07, Gillessen+17
Outbursts from Sgr A* Pont+10, Clavel+13, Chuard+18
 - Hydro simulation → Result is **ρ and T in central parsec**
 - Output constraints: thermal X-ray emission from Chandra
Wang+13



X-rays – 25 vs 90 stars

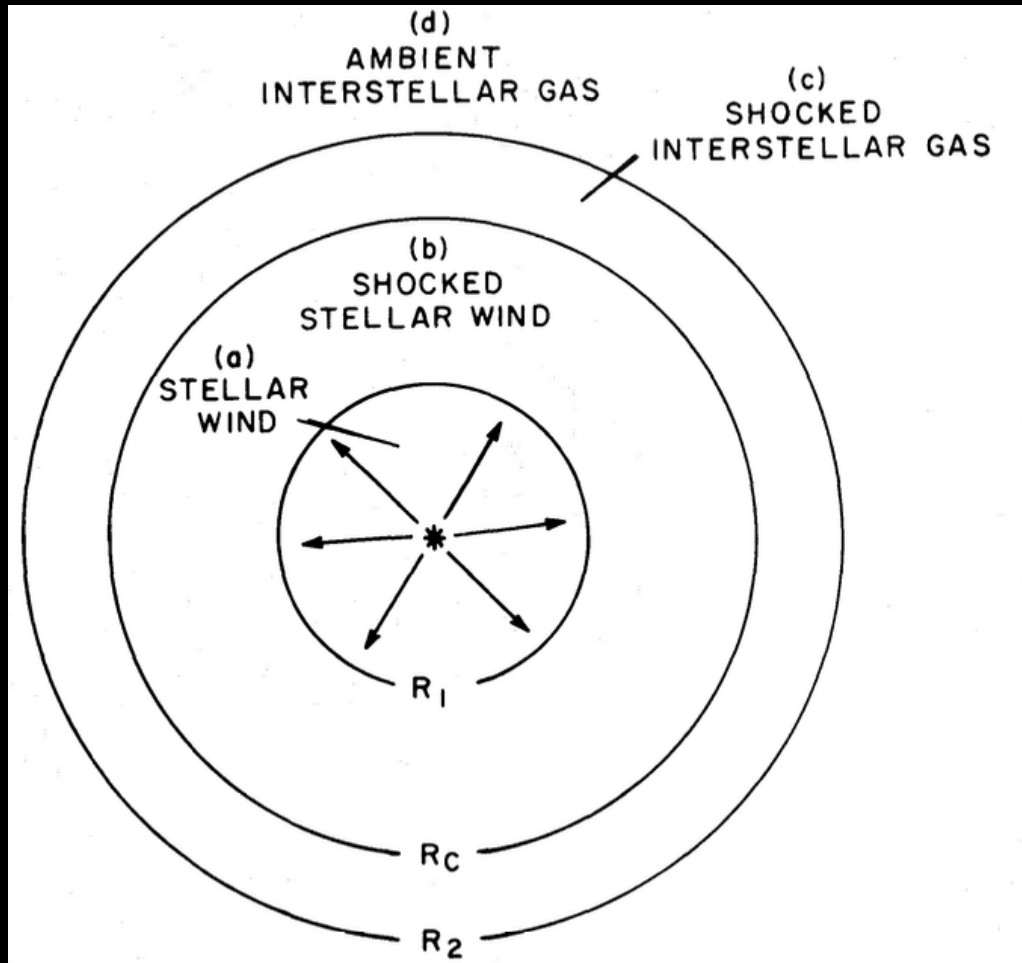
4-9 keV HETG 0th-order
1pc x 1pc



Summary

- Colliding wind X-rays
 - $E_x \approx \text{keV}(V/1000)^2$; $L_x \lesssim L_{wind} \equiv \dot{M}V_\infty^2/2$
- Embedded wind shock X-rays
 - $E_x \approx 0.5\text{keV}$; $L_x \sim 10^{-7}L_{bol} \sim L_{bol}^{(1-m)/\alpha}$
- Magnetically Confined Wind Shocks
 - $E_x \approx (1 - 2)\text{keV}$; $L_x \sim 0.1L_{XADM}$
- Galactic center wind cluster X-rays

Weaver+ 1977



Questions for cluster wind models

- How do multiple wind collisions turn into “superwind”
 - $T_x \approx 10\text{MK}(V/1000)^2$; $V_w \approx 3000\text{km/s}$
 - driven by radiation or gas pressure?
 - latter should be adiabatically cooled
 - origin of B-field; O vs WR
 - slow rotation weakens Parker spiral
- Role of Radiative Cooling
 - Adiabatic shocks: “Energy driven”
 - Radiative shocks: “Momentum driven”
 - But how can $\chi \equiv t_{cool}/t_{exp} \approx v_8^4 d_{12}/\dot{M}_{-7}$ be <1 ?
 - e.g. $d \sim \text{pc} \sim 3 \times 10^{18} \text{ cm} \Rightarrow \chi \approx 3000 v_8^4 / \dot{M}_{-4}$
 - mass loading from clumpy ISM??