



## Prospects for gamma-ray detection of Galactic Stellar Clusters

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#### **PeVatrons**

#### Accelerators of Galactic Cosmic Rays to ${\sim}10^{15}\,eV$

Suitable PeVatron candidates are those that meet the Hillas criterion:

Supernova remnants are prime candidates as a Galactic source

 $\sim 10^{51}$  erg per explosion,

- ~10% into CR acceleration,
- ~3 per century in the Milky Way
- -> would be sufficient to power Cosmic Rays

BUT -> no active PeVatron supernovae seen to date

Only act as accelerators for a short period of time?









### Gamma-ray emission from Stellar Clusters



- Stellar clusters can drive a collective wind into the surrounding medium
- Particle acceleration may occur at the wind termination shock up to PeV?
- Analytical model for CR acceleration from G. Morlino et al. 2021
  - apply this model to stellar clusters listed in Gaia DR2
  - estimate the expected CR flux and resulting gamma-ray emission

See AM, Morlino, Celli, Menchiari, Specovius https://arxiv.org/abs/2403.16650 (submitted)





Relevant sizes can be approximated based on cluster properties:

Bubble radius  $R_b$ :

$$R_b(t) = 112 \left(\frac{n}{\text{cm}^{-3}}\right)^{-0.2} \left(\frac{L_{\text{w,c}}}{10^{37} \text{ erg s}^{-1}}\right)^{0.2} \left(\frac{t}{10 \text{ Myr}}\right)^{0.6} \text{ pc}$$

 $R_{cd} = 0.95 R_b$  (contact discontinuity)

Wind termination shock radius  $R_s$ :

$$R_{s}(t) = 26 \left(\frac{n}{\text{cm}^{-3}}\right)^{-0.3} \left(\frac{\dot{M}_{c}}{10^{-4}\text{M}_{\odot}\text{yr}^{-1}}\right)^{0.3} \left(\frac{v_{w}}{2000 \,\text{km}\,\text{s}^{-1}}\right)^{0.1} \left(\frac{t}{10 \,\text{Myr}}\right)^{0.4} \,\text{pc}$$

Cluster radius  $R_c$ :

use the angular radius containing 50% of the member stars as reported in the Gaia DR2 stellar cluster catalogue, converted into a physical size using the most recent distance estimate (i.e. from Gaia DR3 if available).



#### Gamma-ray emission from Stellar Clusters



- Model from Morlino et al. 2021
- Estimate the resulting gamma-ray emission from the wind-blown bubble around stellar clusters as listed in the Gaia DR2
- Particle injection spectrum depends on location within the system:
- In the bubble:

$$f_2(r,p) = f_s(p)e^{\alpha(r)} \frac{1 + \beta(e^{\alpha(R_b)}e^{-\alpha(r)} - 1)}{1 + \beta(e^{\alpha(R_b)} - 1)} + f_{\text{Gal}}(p) \frac{\beta(e^{\alpha(r)} - 1)}{1 + \beta(e^{\alpha(R_b)} - 1)}$$

- In the surrounding ISM:

$$f_3(r,p) = f_b(p)\frac{R_b}{r} + f_{\text{Gal}}(p)\left(1 - \frac{R_b}{r}\right)$$





Particle injection spectrum from the termination shock:

$$f_s(p) = \frac{\chi_{\rm CR}}{2} \frac{\sigma n_1 u_1^2}{4\pi \Lambda_p (m_p c)^3 c^2} \left(\frac{p}{m_p c}\right)^{-s} e^{-\Gamma(p)}$$

and where 
$$e^{-\Gamma(p)} \simeq \left[1 + A_0 \left(\frac{p}{p_{\max,1}}\right)^{A_2}\right] \exp\left[-A_1 \left(\frac{p}{p_{\max,1}}\right)^{A_3}\right]$$
,

is an approximation that depends on the diffusion scenario being considered.

Note that 
$$\alpha(r,p) = \frac{u_2 R_s}{D_2(p)} \left(1 - \frac{R_s}{r}\right)$$
 and  $\beta(p) = \frac{D_3(p) R_b}{u_2 R_s^2}$ 

also depend on the diffusion properties of the medium.





Again, region 1 within the wind and region 2 within the wind-blown bubble.

The wind velocity takes the following form: 
$$u(r) = \begin{cases} u_1 & r < R_s \\ u_2 \left(\frac{R_s}{r}\right)^2 & R_s < r < R_b \\ 0 & r > R_b \end{cases}$$

Where  $u_2 = u_1/4$  due to a shock compression ratio of  $\sigma = 4$  being assumed.

#### Density profile in the wind and bubble



The density of the wind depends on the mass-loss rate from the cluster

$$n_{\rm w}(r) = \frac{M_{\rm c}}{4\pi r^2 v_{\rm w}} \qquad R_{\rm c} \le r < R_{\rm s}$$

The density of the bubble is: 
$$n_{\rm b} = \frac{M_{\rm b}}{m_{\rm p}\frac{4}{3}\pi(R_{\rm cd}^3 - R_{\rm s}^3)}$$

where the mass of the bubble:

$$M_{\rm b} = (\dot{M}_{\rm sh} + \dot{M}_{\rm w})T_{\rm age} - \frac{\dot{M}_{\rm w}}{v_{\rm w}(R_{\rm s} - R_{\rm c})}$$

And the evaporated mass

$$\dot{M}_{\rm sh} = 2 \times 10^{-4} \left( \frac{L_{\rm w,c}}{10^{37} \,{\rm erg/s}} \right)^{\frac{27}{35}} \left( \frac{n}{10 \,{\rm cm}^{-3}} \right)^{-\frac{2}{35}} \left( \frac{t}{1 \,{\rm Myr}} \right)^{\frac{6}{35}} M_{\odot} \,{\rm yr}^{-1}$$



... in the case of a clumpy medium, this mass can be redistributed.

~10% of the shell material may fragment to form uniformly distributed clumps of material.

The bubble density is then modified as:  $n_b^* = n_b^0 + 0.117 n_{ism}$  and the shell density:  $n_{sh}^* = 0.9 n_{sh}^*$ 



Now we have most ingredients of the physical model (skipping some details...)

The particle distribution is converted into a gamma-ray flux using Kelner et al. 2006:

$$\Phi_{\gamma,\nu}(E_{\gamma,\nu}, R', t') = cn \int_{E_{\gamma,\nu}}^{\infty} \sigma_{\rm pp}(E) f(E, R', t') K_{\gamma,\nu}\left(\frac{E_{\gamma,\nu}}{E}, E\right) \frac{dE}{E}$$

The flux detected at Earth is therefore:  $F(E_{\gamma,\nu}, t) = \Phi_{\gamma,\nu}(E_{\gamma,\nu}, t)/(4\pi D_c^2)$ 

Where  $D_c$  is the distance to the cluster.

Consider a representative (hypothetical) stellar cluster to explore how free parameters affect the model.

Parameter	Value	Parameter	Value
YCR	10%	t NG	10 Myr
$\chi_{\rm B}$	10%	M Vw c	$10^{-1} M_{\odot} / yr$ 3000. km /s
$L_{\rm coh}$	2 pc	$\delta$	0.5 (Kraichnan)
$\sigma$	4	$D_0 (10 {\rm GeV})$	$3.0 \times 10^{28} \text{ cm}^2/\text{s}$
		n <sub>ism</sub>	$10 \mathrm{cm}^{-3}$





Massive Wolf-Rayet stars contribute with their strong stellar winds

-> contribution added explicitly to each cluster based on Gaia DR3 catalogue of Wolf-Rayet stars.

Below an age ~ 3Myr, no supernovae will have occurred.

Above this age, a number of supernovae are expected to have occurred, depending on the stellar cluster mass and age.

For each cluster, we estimate the number of supernovae already occurred. -> All SNe will contribute material to increase the density of the bubble region -> Only SNe within 1 advection time will contribute to the cluster wind luminosity -> We assume the contribution from SNe does not affect the bubble size

The contribution from SNe does, however, enhance the flux.





We find that the SNe contribution enhances the flux as

$$\langle f_{\rm SNR}(E) \rangle = 1.74 N_{\rm SN}(t_{\rm esc}) \left(\frac{E_{\rm SN}}{10^{51} {\rm erg}}\right) \left(\frac{L_w}{10^{37} {\rm erg/s}}\right)^{-1} \left(\frac{t_{sc}}{3 {\rm Myr}}\right)^{-1} \times f_s(10 \,{\rm GeV}) \left(\frac{E}{10 \,{\rm GeV}}\right)^{-2.3} \exp(-E/E_{\rm max})$$

Which contribution dominates the gamma-ray flux depends on the age of the cluster

After ~3Myr, supernovae dominate







The expressions so far have assumed that the cluster mass  $M_c$  (mass-loss rate  $\dot{M}_c$ ) and wind luminosity  $L_{w,c}$  are available as free parameters of the model and required as input.

How do we determine them for known stellar clusters?

Turned out to not be straight forward.

Developed a multi-step procedure, outlined in:

"On the mass and wind luminosity of young Galactic open clusters in Gaia DR2" Celli, Specovius, Morlino, AM, Menchiari, A&A 686 (2024) A118 https://doi.org/10.1051/0004-6361/202348541



1) Extract the number and magnitude of stars in each cluster in the catalogue

We adopt the initial mass function:

$$\xi(M) = k \times \begin{cases} \left(\frac{M}{M_{\min}}\right)^{-\beta_1} & M_{\min} \le M < M_0 \\ \left(\frac{M_0}{M_H}\right)^{-\beta_1} \left(\frac{M}{M_0}\right)^{-\beta_2} & M_0 \le M < M_1 & (1) \\ \left(\frac{M_0}{M_H}\right)^{-\beta_1} \left(\frac{M_1}{M_0}\right)^{-\beta_2} \left(\frac{M}{M_1}\right)^{-\beta_3} & M_1 \le M < M_{\max} \end{cases}$$
  
with  $\beta_1 = 1.30, \beta_2 = 2.30, \beta_3 = 2.35, M_{\min} = 0.08 M_{\odot}, M_0 = 0.50 M_{\odot}, and M_1 = 1.00 M_{\odot}, M_{\odot}$  being the Sun mass.  
$$N^* \text{ per cluster}$$

30

2) Perform a kernel density fit to the G-band magnitude distribution of each cluster to obtain Gmin & Gmax In place of the maximum magnitude, we use the **mode** magnitude —> i.e. completeness limit.

Convert to bolometric magnitude:  $\mathcal{M}_{bol} = \mathcal{M}_G + BC_G$  with Bolometric correction  $BC_G(T_{eff}) = \sum_{i=0}^{4} \alpha_i \left(T_{eff} - T_{eff,\odot}\right)^i$ 



3) Convert into intrinsic luminosity accounting for the extinction correction via:

$$L_s = 10^{0.4(\mathcal{M}_{\odot,\text{bol}} - \mathcal{M}_{\text{bol}} + A_G)} L_{\odot}$$

4) Convert from luminosity to stellar mass via the mass-luminosity relation:

$$L_{\rm s}(M) = \begin{cases} L_{b1} \left(\frac{M}{M_{b1}}\right)^{\gamma_1} \left[\frac{1}{2} + \frac{1}{2} \left(\frac{M}{M_{b1}}\right)^{1/\Delta_1}\right]^{\Delta_1(-\gamma_1 + \gamma_2)} & M/M_{\odot} < 12 \\ \\ \epsilon L_{b2} \left(\frac{M}{M_{b2}}\right)^{\gamma_2} \left[\frac{1}{2} + \frac{1}{2} \left(\frac{M}{M_{b2}}\right)^{1/\Delta_2}\right]^{\Delta_2(-\gamma_2 + \gamma_3)} & M/M_{\odot} \ge 12 \end{cases}$$

-> Yields the **minimum**  $M^*_{\min}$  and **maximum**  $M^*_{\max}$  mass of *observed* member stars

5) Derive the normalisation k of the stellar mass function 
$$\xi(M)$$
 by inverting the integral:  $N^* = \int_{M_{\min}^*}^{M_{\max}^*} \xi(M) \, dM$ 



6) Then the total cluster mass can be obtained via  $M_{\rm c} = \int_{M_{\rm min}}^{M_{\rm max}} \xi(M) M \, dM$ 

- this time the bounds of the integral are given by theoretical considerations:  $M_{\rm min} = 0.08 M_{\odot}$  and  $M_{\rm max}$  depends on the age of the stellar cluster.



We expect a systematic underestimate of a factor ~2-3, based on clusters where the mass can be determined independently.

The cluster mass-loss rate can be estimated from the mass-loss rates of individual stars  $M_s$ 

$$\dot{M}_{\rm c} = \int_{M_{\rm min}}^{M_{\rm max}} \xi(M) \dot{M}_{\rm s}(M) \, dM$$

Similarly, via momentum conservation

$$\dot{M}_{c}v_{w,c} = \int_{M_{min}}^{M_{max}} \xi(M)\dot{M}_{s}(M)v_{w,s}(M) dM$$
, where  $v_{w,s}$  is the stellar wind velocity.

Finally, the cluster wind luminosity can be obtained via:

$$L_{\rm c,w} = \frac{1}{2} \dot{M}_{\rm c} v_{\rm w,c}^2 = \frac{1}{2} \frac{\left(\int_{M_{\rm min}}^{M_{\rm max}} \xi(M) \dot{M}_{\rm s}(M) v_{\rm w,s}(M) dM\right)^2}{\int_{M_{\rm min}}^{M_{\rm max}} \xi(M) \dot{M}_{\rm s}(M) dM}$$



### Free parameters of the model

...and their default values

Properties of the cosmic ray efficiency and particle diffusion -> CR Properties of the stellar cluster (age, mass, wind etc.) -> Cluster Number of supernovae occurring in the stellar cluster -> SNR





	Parameter	Value
CR	ξcr ξ	10%
	$L_{\rm coh}$	2 pc
	$\sigma$	4
	<i>s</i> <sub>wts</sub>	4.0
	$\delta$	0.5 (Kraichnan)
	$D_0 (10 {\rm GeV})$	$3.0 \times 10^{28} \mathrm{cm}^2 \mathrm{/s}$
Cluster	tage	3 Myr
	$M_{\rm sc}$	$10^4 M_{\odot}$
	М	$10^{-5} { m M}_{\odot}/{ m yr}$
	$v_w$	2500 km s <sup>-1</sup>
	n <sub>ISM</sub>	$10  {\rm cm}^{-3}$
	d	1 kpc
	$\eta_{ m fr}$	0.1
SNR	$N_{ m SN}$	5
	$M_{ m ej}$	$5 M_{\odot}$
	$E_{ m sn}$	$10^{51}  {\rm erg}$
	s <sub>sn</sub>	4.3

### Comparison to data

Gamma-ray detected stellar clusters in Gaia DR2

We apply the model to known stellar clusters with gamma-ray emission

Shaded band - allowed variation via reasonable changes to the model









Spectral energy distributions adjusted to data by varying (e.g.) the average ambient density in the surrounding ISM and the CR efficiency

Shape of Bohm diffusion curve is more consistent with data than Kraichnan

(No fitting —> visual comparison only)



Westerlund 1 —> Smaller observed bubble size ( $\sim$ 60pc) & revised distance estimate are consistent with red line Provided a higher ISM ambient density is assumed —> also consistent with observations.





NGC 3603 —> based on Fermi-LAT data, the expected emission in the TeV range is well above HESS sensitivity

Detection of Danks 1 & 2 also recently reported based on Fermi-LAT data.





Select young stellar clusters (< 30 Myr) for which necessary information is available

Determine  $v_w$  based on the aforementioned estimates for  $\dot{M}_c$  and  $L_w \rightarrow coupled$  via:  $L_w = \frac{1}{2}\dot{M}_c v_w^2$ 

Angular size given by  $R_b$  and distance.

Compare predicted gamma-ray / neutrino flux to instrument sensitivity







Many wind-blown bubbles around stellar clusters are simply large, which decreases the surface brightness.

Instrument sensitivity is also reduced at larger angular sizes.

Solid and dashed lines indicate the plausible range for each stellar cluster







Using the most optimistic scenario for the model (upper bound / dashed line curves)

Red circles = LHAASO ultra-high-energy sources & Cyan points = 4FGL unidentified sources

Many bubbles have predicted sizes that are ~multi-degree scale and often multiple overlapping along the line of sight





Due to either supernovae or the stellar cluster wind —> we wish to identify not only those clusters likely to be detectable, but also those for which this indicates PeVatron activity

Usually, supernovae dictate the maximum energy

$$E_{\text{max}} = 48 \left(\frac{B}{10\,\mu\text{G}}\right) \left(\frac{L_{\text{coh}}}{2\,\text{pc}}\right)^{-1} \left(\frac{u_{st}}{5000\,\text{km/s}}\right)^{\frac{1}{2}} \left(\frac{R_{st}}{10\,\text{pc}}\right)^{\frac{1}{2}} \text{TeV}$$

Whilst the maximum momentum determined via confinement in the wind is:

$$cp_{\max,1}^{Bohm} = 3u/c eBR_s$$
  
 $cp_{\max,1}^{Kra} = (3u/cR_s)^2 eB/L_{coh}$ 





In the vicinity of stellar clusters, molecular clouds may locally enhance the expected gamma-ray / neutrino flux.

Clouds — from catalogues based on CO data — are paired with clusters out to a distance  $2R_b$  from the cluster centre.

An important consideration here is also the fraction of the flux due to the stellar cluster wind vs illumination by the sea of galactic cosmic rays.

Define  $\zeta = F_{\gamma}^{\rm sc}/F_{\gamma}^{\rm sc+gal}$ 









- Detailed modelling of specific clusters, especially those with detected emission
- Addressing some assumptions of the model (e.g. losses, transport, composition, density & B-field profiles...)
- Follow-up analyses and/or observation proposals for clusters with the brightest predicted fluxes
- Predicting the anticipated enhancement factor for flux from nearby clouds
- Selection of which molecular cloud catalogue / gas distribution maps are most suitable





# Thank you for your attention

Any questions?



Total flux increases with age, with wind velocity and with the mass loss rate.

The bubble contributes most to the total flux.

Note: below 100 GeV, we use the expression  $\frac{dN_{\gamma}}{dE_{\gamma}} = 2 \int_{E_{\min}}^{\infty} \frac{q_{\pi}(E_{\pi})}{\sqrt{E_{\pi}^2 - m_{\pi}^2}} dE_{\pi} \quad \text{with } E_{\min} = E_{\gamma} + m_{\pi}^2/4E_{\gamma}$ 

and 
$$q_{\pi}(E_{\pi}) = \tilde{n} \frac{c n_H}{K_{\pi}} \sigma_{\text{inel}} \left( m_p + \frac{E_{\pi}}{K_{\pi}} \right) J_p \left( m_p + \frac{E_{\pi}}{K_{\pi}} \right)$$
.





For neutrinos, we assume that the neutrino flux at source is:

$$\Phi^*_{
u} = \phi^*_{
u_e} + \phi^*_{
u_{\mu}}$$

Such that the flux arrive at Earth is modified due to neutrino oscillations and distributed over three neutrino flavours:

$$\begin{split} \phi_{\nu_{e}} &= & P_{ee}\phi_{\nu_{e}}^{*} + P_{\mu e}\phi_{\nu_{\mu}}^{*} \\ \phi_{\nu_{\mu}} &= & P_{e\mu}\phi_{\nu_{e}}^{*} + P_{\mu\mu}\phi_{\nu_{\mu}}^{*} \\ \phi_{\nu_{\tau}} &= & P_{e\tau}\phi_{\nu_{e}}^{*} + P_{\mu\tau}\phi_{\nu_{\mu}}^{*} \end{split}$$

For comparing to neutrino detector sensitivity, we use the muon neutrino flux only as the most promising channel for detection



Stellar mass-loss rate:

$$\dot{M}_{\rm s}(M) \simeq 9.55 \times 10^{-15} \left(\frac{L_{\rm s}(M)}{L_{\odot}}\right)^{1.24} \left(\frac{M}{M_{\odot}}\right)^{0.16} \left(\frac{R_{\rm s}(M)}{R_{\odot}}\right)^{0.81} \frac{M_{\odot}}{\rm yr}$$

The stellar radius depends on mass as:

 $R_{\rm s}(M) = 0.85 R_{\odot} (M/M_{\odot})^{0.67}$ .

The stellar wind velocity is given by:

$$v_{\rm w,s}(M) = C(T_{\rm eff}) \sqrt{\frac{2G_{\rm N}M}{R_{\rm s}(M)} \left(1 - \frac{L_{\rm s}(M)}{L_{\rm Edd}(M)}\right)}$$