

Direct neutron capture experiments on stable and unstable isotopes (for the s- and i-processes)

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Instituto de Física Corpuscular (Valencia, Spain)
n_TOF Collaboration (CERN)

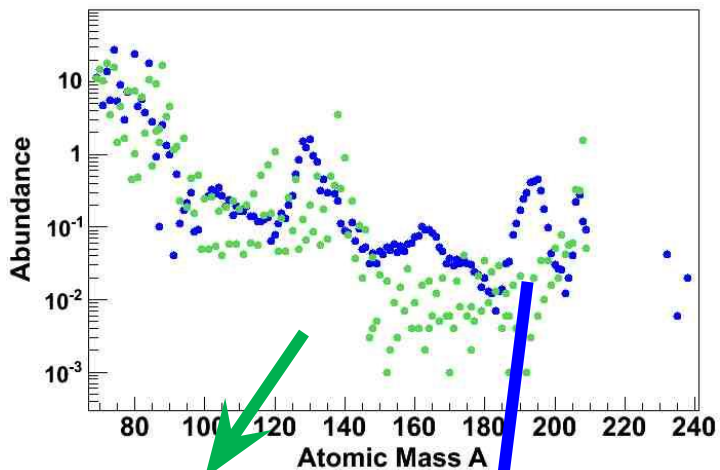
s, i & r Element Nucleosynthesis (sirEN)
CONFERENCE

Giulianova (Italy), 8-13 June 2025

Dedicated to the memory of Prof. Roberto Gallino

- **The synthesis of (heavy) elements in the universe**
- **Nucleosynthesis & neutron capture: data needs for unstable and stable nuclei**
- **Direct neutron capture reactions: from stars to the lab**
- **State-of-the art & future of direct neutron capture reactions**

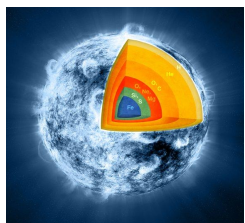
Neutrons capture (n,g) reactions play a key role in the synthesis of heavy elements



s-process

$T = 10^8 - 10^9$ K

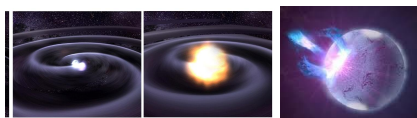
$N_n = 10^6 - 10^{12} \text{ cm}^{-3}$



r-process

$T = 10^8 - 10^{10}$ K

$N_n = 10^{20} - 10^{27} \text{ cm}^{-3}$

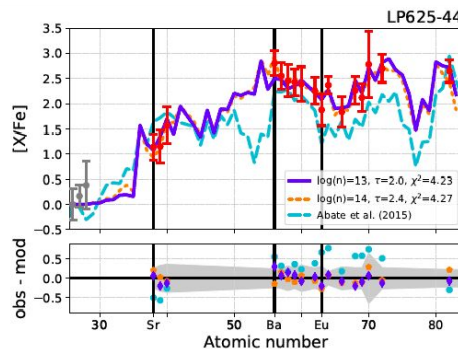
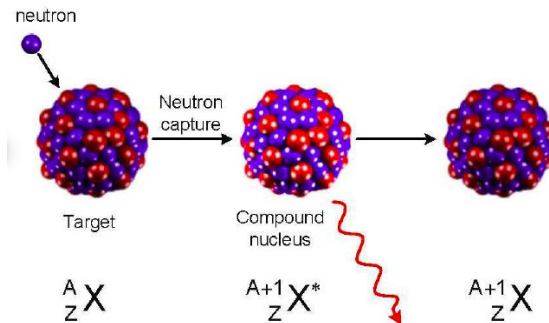
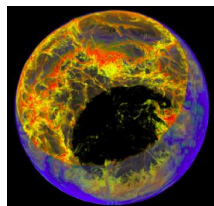


s + i + r EN

i-process

$T > 10^8$ K

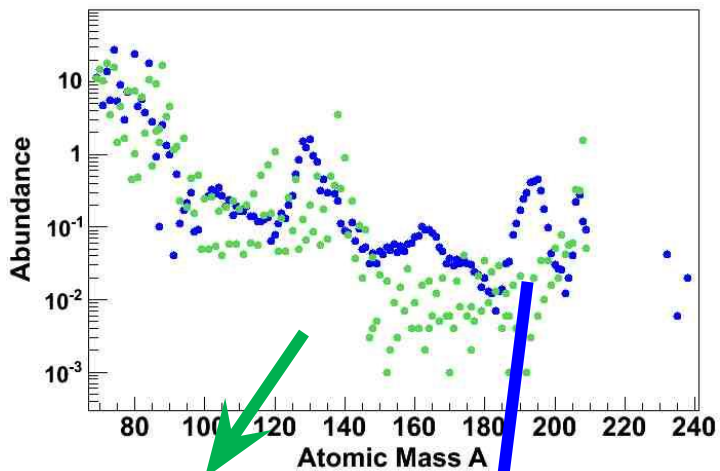
$N_n = 10^{13} - 10^{16} \text{ cm}^{-3}$



i-process:
abundance
pattern in
CEMP stars

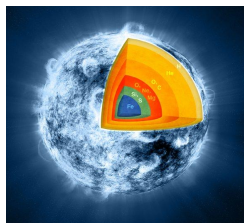
Falk Herwig et al 2011 ApJ 727 89
Melanie Hampel et al 2019 ApJ 887 11

Neutrons capture (n,g) reactions play a key role in the synthesis of heavy elements



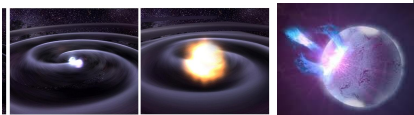
s-process

$T = 10^8 - 10^9 \text{ K}$
 $N_n = 10^6 - 10^{12} \text{ cm}^{-3}$



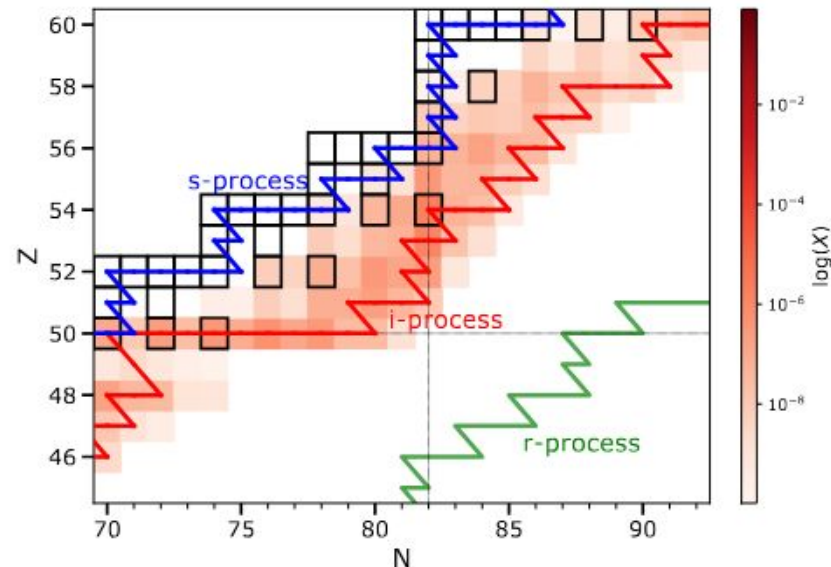
r-process

$T = 10^8 - 10^{10} \text{ K}$
 $N_n = 10^{20} - 10^{27} \text{ cm}^{-3}$



i-process

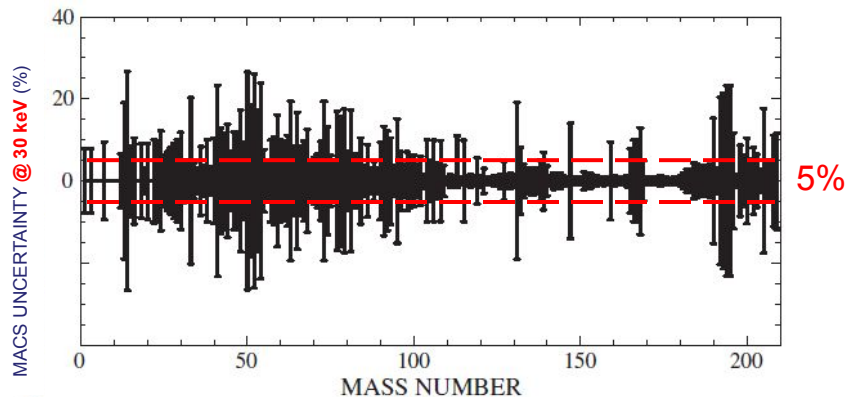
$T > 10^8 \text{ K}$
 $N_n = 10^{13} - 10^{16} \text{ cm}^{-3}$



A. Choplin, EPJ Web of Conferences 279, 07001 (2023)

s-process: stable or long-lived unstable → **Direct (n,g)** 
i-process: few units from stability → shorter-lived → **Direct (n,g) in future** 
r-process: n-rich (veery short lived) → **Direct (n,g)** 

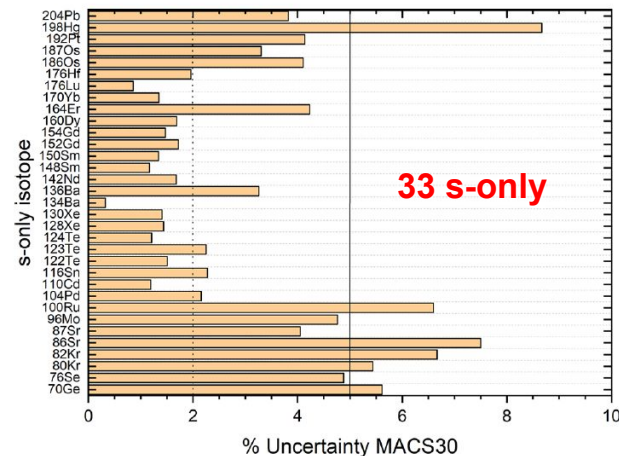
- The synthesis of (heavy) elements in the universe
- **Nucleosynthesis & neutron capture: data needs for unstable and stable nuclei**
- Direct neutron capture reactions: from stars to the lab
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F. Käppeler et al. [Rev. Mod. Phys. \(2011\)](#)

s-only isotopes:
Benchmark for s-process
and galactic-chemical
evolution
(GCE) models,

**Aim: Accuracy MACS
down to 2% unc.**

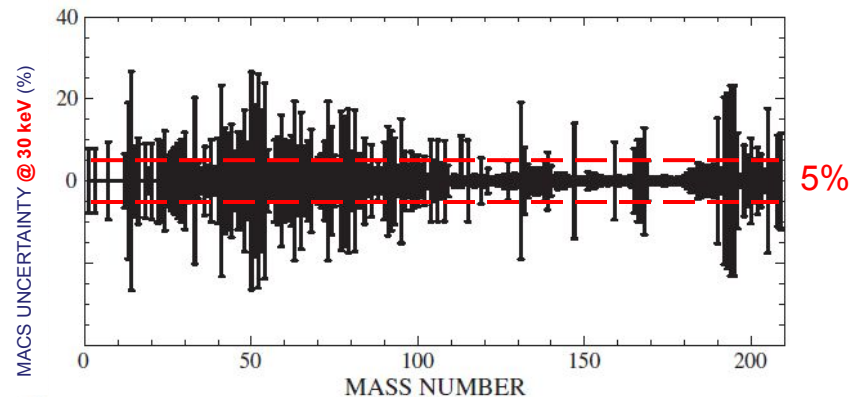


I. Dillmann et al., [Eur. Phys. J. A 59, 105 \(2023\)](#)

MC Sensitivity studies:

AGB: G. Cescutti et al., *MNRAS*, 478, 4101(2018).

Massive: Nishimura et al., *MNRAS*, 469, 1752 (2017)

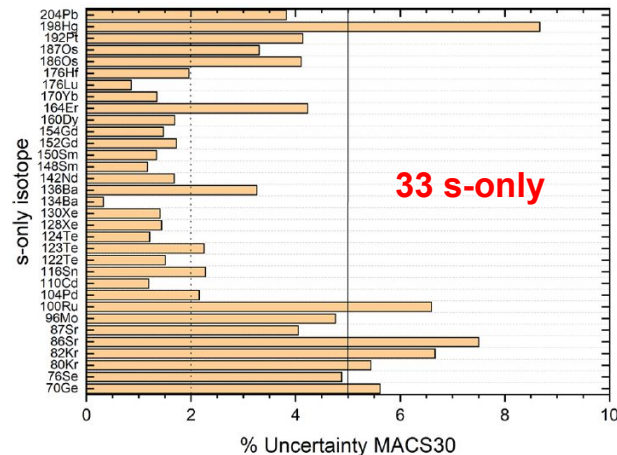


F. Käppeler et al. [Rev. Mod. Phys. \(2011\)](#)

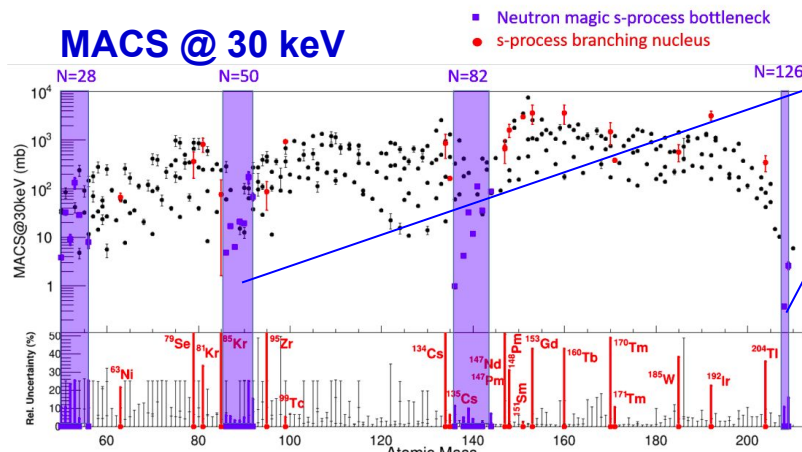
s-only isotopes:
Benchmark for s-process
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**Aim: Accuracy MACS
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**n-magic nuclei:
s-process bottlenecks**



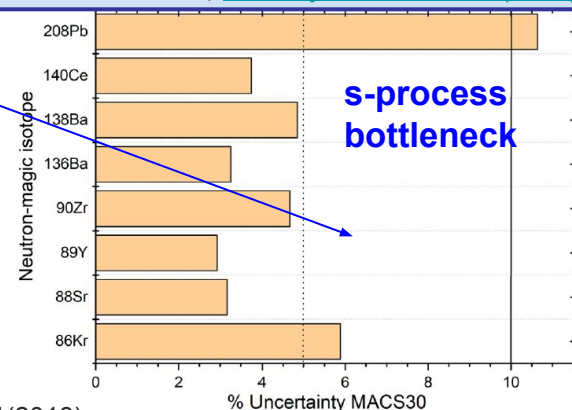
MACS @ 30 keV



Very small (n,g) XS

**High accuracy
(n,g) data on
stable nuclei
still needed!**

I. Dillmann et al., [Eur. Phys. J. A 59, 105 \(2023\)](#)



MC Sensitivity studies:

AGB: G. Cescutti et al., *MNRAS*, 478, 4101(2018).

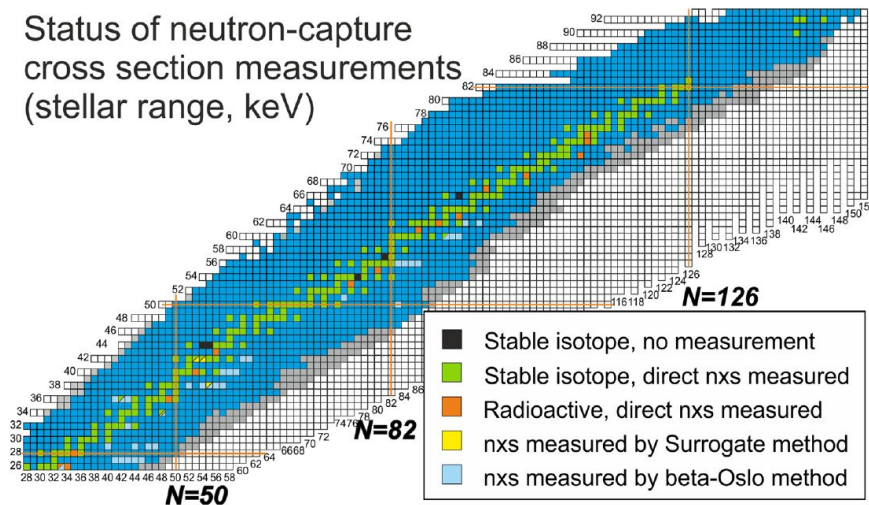
Massive: Nishimura et al., *MNRAS*, 469, 1752 (2017)

Key unstable cases: branching points

$^A_Z(n,\gamma)$ competes with β decay: (n,γ) cross sections $\rightarrow$ conditions stellar environment

I. Dillmann et al., *Eur. Phys. J. A* **59**, 105 (2023).

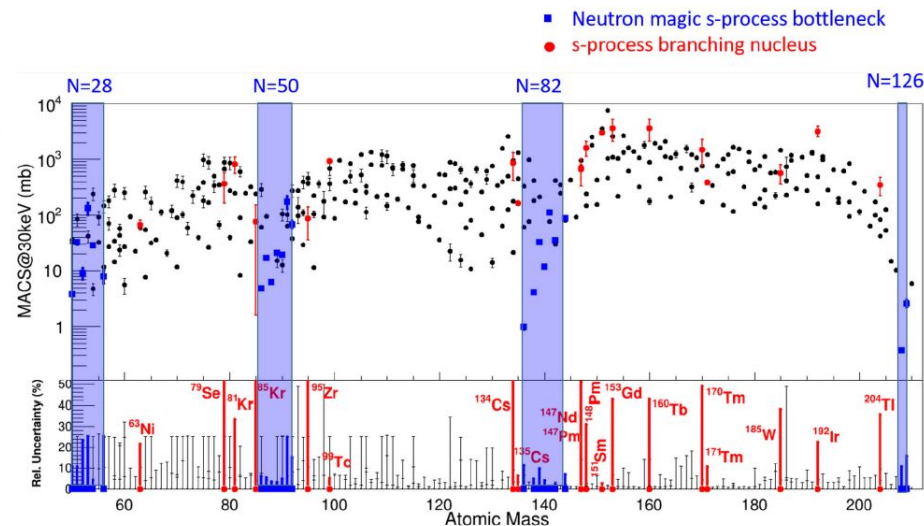
Status of neutron-capture cross section measurements (stellar range, keV)



Status Unstable isotopes:

Very little direct (n,γ) data,
Very large uncertainties
Incomplete neutron energy range

C. Domingo-Pardo et al. *Eur. Phys. Jour. A* **59**, 8 (2023)

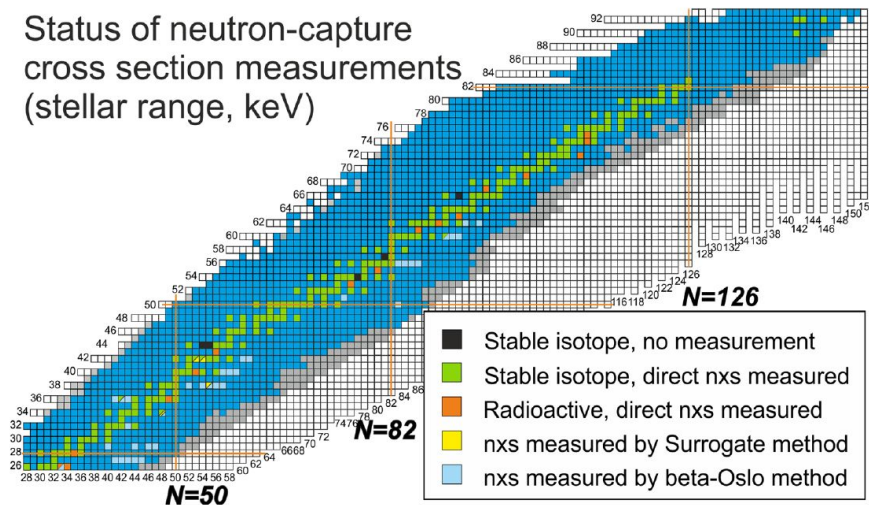


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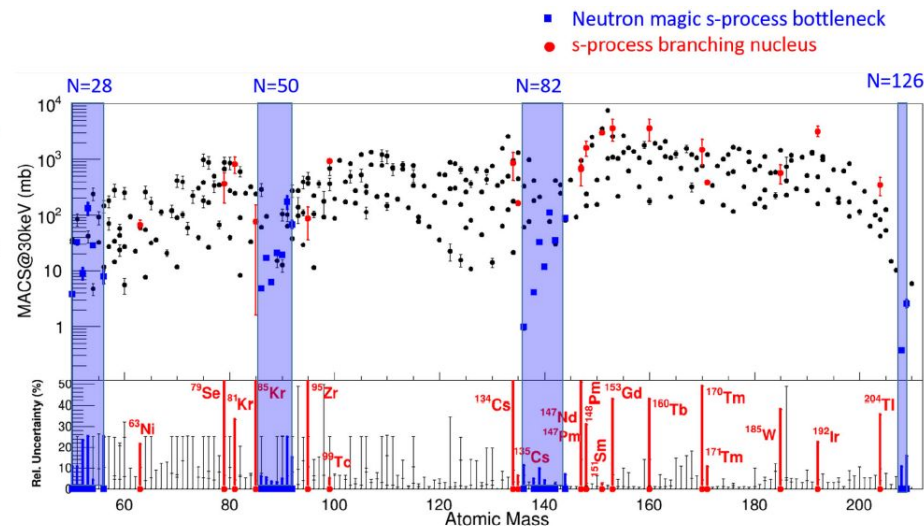
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C. Domingo-Pardo et al. *Eur. Phys. Jour. A* **59**, 8 (2023)



WHY? Challenging measurements:

Radioactive isotopes, low masses

The impact of (n,γ) reaction rate uncertainties on the predicted abundances of i-process elements with $32 \leq Z \leq 48$ in the metal-poor star HD94028

John E. McKay,^{1,2†} Pavel A. Denissenkov,^{1,3*†} Falk Herwig,^{1,3*†}
Georgios Perdikakis^{3,4,5} and Hendrik Schatz^{3,5,6†}

The impact of (n,γ) reaction rate uncertainties of unstable isotopes on the i-process nucleosynthesis of the elements from Ba to W

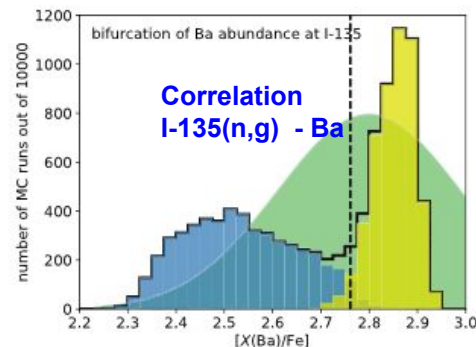
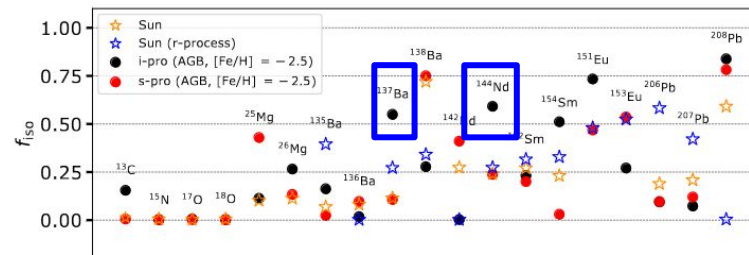
Pavel A. Denissenkov,^{1,2,3*†} Falk Herwig,^{1,2*†} Georgios Perdikakis,^{2,4,5} and Hendrik Schatz^{2,5,6†}

Signatures of i-process: Ba-137, Nd-144:

^{137}Cs (**30 y**)(n,γ) , ^{144}Ce (**285 d**)(n,γ) compete with production

Very high impact (n,g) in i-process abundances:

^{66}Ni (**55h**)(n,γ), ^{72}Zn (**46.5h**)(n,γ), ^{75}Ga (**2.10m**)(n,γ),
 ^{135}I (**6.6h**)(n,γ), ^{141}Ba (**18.3 m**)(n,γ), ^{141}La (**3.92h**)(n,γ)



Reaction	Element
^{66}Ni 54.6 h	Zn
	Ge
	As
^{72}Zn 46.5 h	Ge
^{75}Ga 2.10 m	As
Element	Reaction
Ba	^{135}I 6.6 h

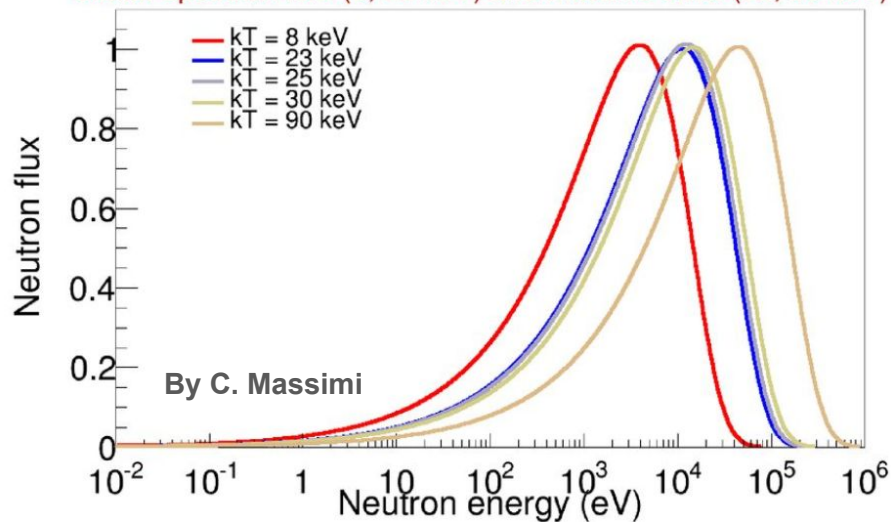
(n,g) for i-process: involves n-rich nuclei → even more challenging than s-process

1. Half-life is **in general too short (minutes-hours)** to allow a (days/weeks-long) measurement.
2. To be produced in **exotic beam facilities**
3. **Requires: new neutron facilities & experimental techniques**

**Currently: indirect
(n,g) (A. Spyrou)**

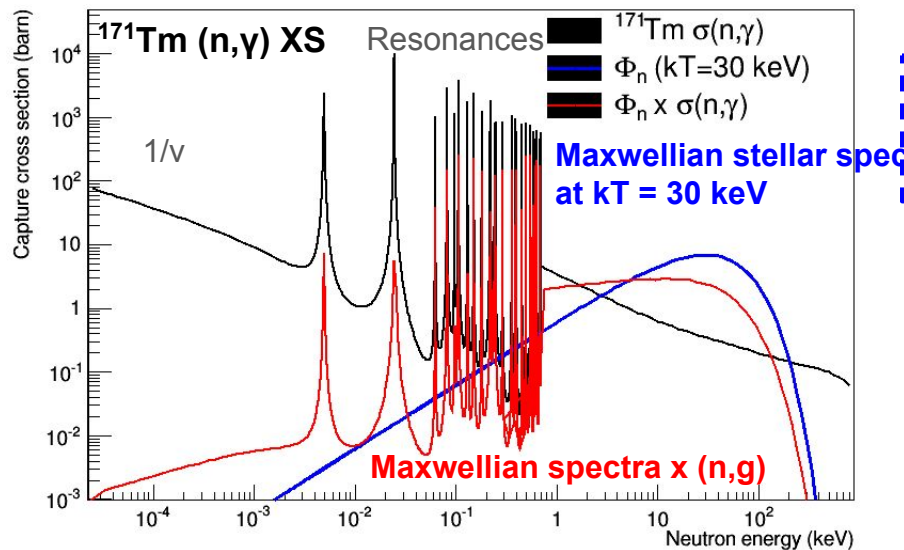
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Stellar spectra: AGB (8, 23 keV) and Massive stars (25, 30, 90 keV)



Maxwellian stellar spectra

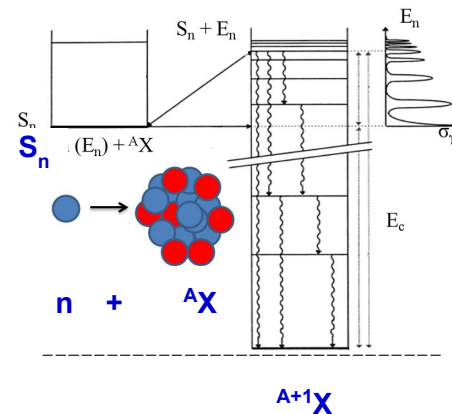
$$\phi_{MB}(v)dv = \phi_{MB}(E)dE = \frac{2}{\sqrt{\pi}} \frac{1}{(kT)^{3/2}} \sqrt{E} e^{-\frac{E}{kT}} dE$$

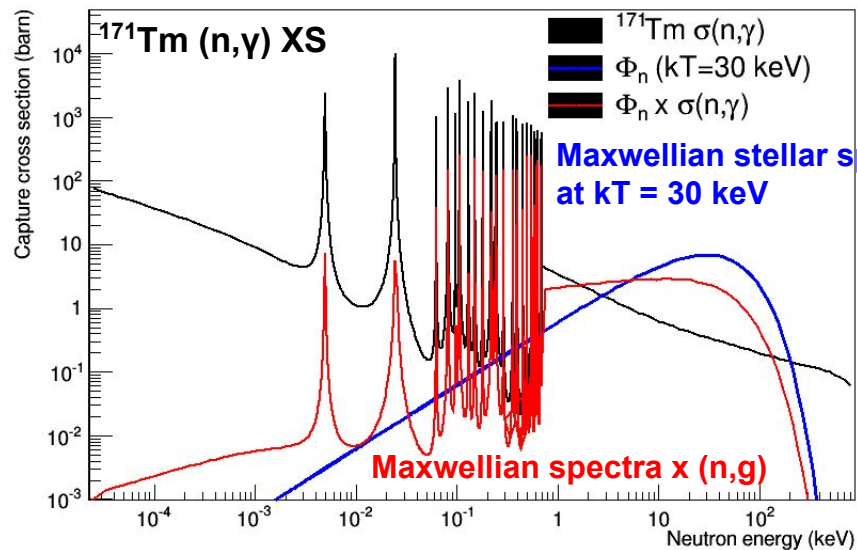


Maxwellian stellar spectra

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Resonances <-> Compound nucleus





Maxwellian stellar spectra

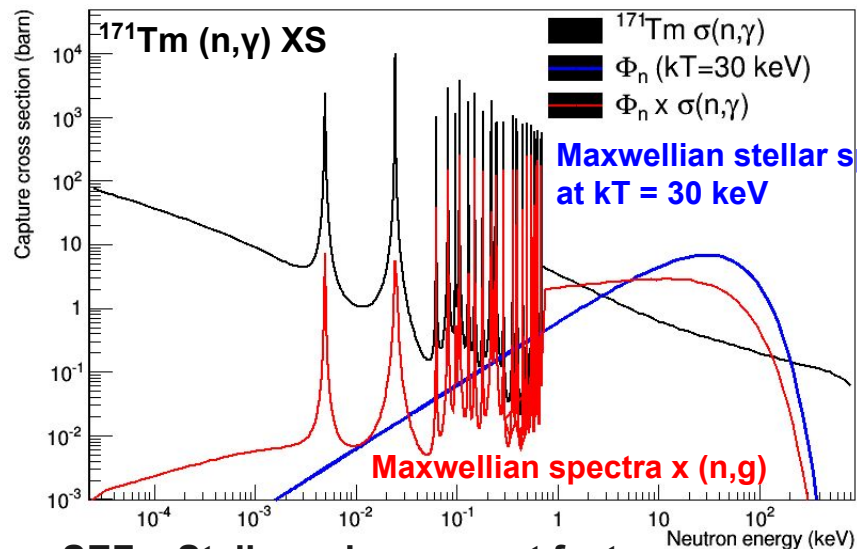
$$\phi_{MB}(v)dv = \phi_{MB}(E)dE = \frac{2}{\sqrt{\pi}} \frac{1}{(kT)^{3/2}} \sqrt{E} e^{-\frac{E}{kT}} dE$$

Reaction rate per particle pair

$$\langle \sigma v \rangle = \left(\frac{8}{\pi \mu} \right)^{1/2} \frac{1}{(kT)^{3/2}} \int_0^\infty E \sigma(E) e^{-\frac{E}{kT}} dE.$$

MACS = Maxwellian Averaged Cross Section

$$\text{MACS} = \frac{\langle \sigma v \rangle}{v_T} = \frac{2}{\sqrt{\pi} (kT)^2} \int_0^\infty E \sigma(E) e^{-\frac{E}{kT}} dE.$$



SEF = Stellar enhancement factor

For some isotopes:

$$SEF(T) = \frac{\langle \sigma \rangle^{\text{star}}}{\langle \sigma \rangle_{kT}},$$

Maxwellian stellar spectra

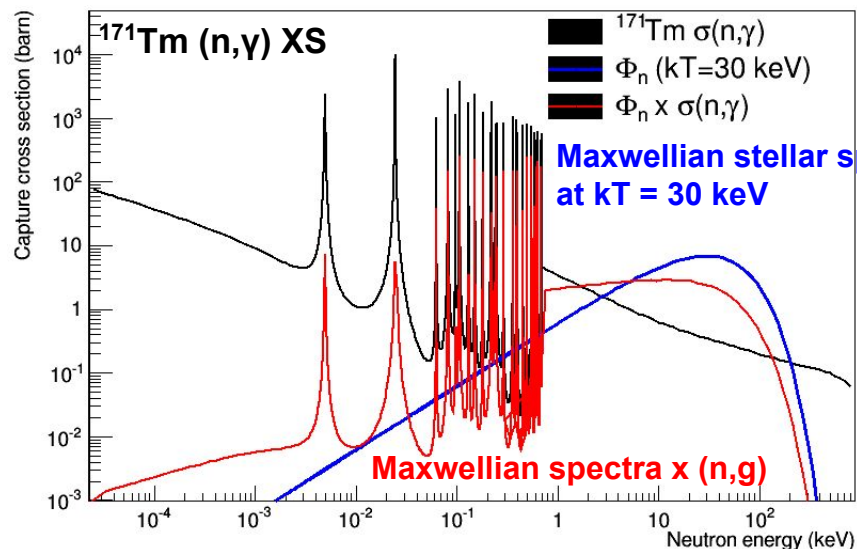
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Maxwellian stellar spectra

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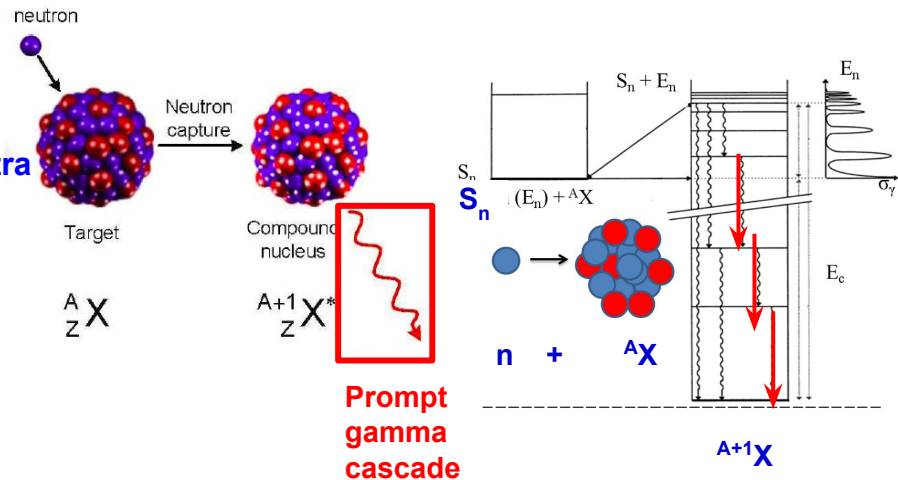
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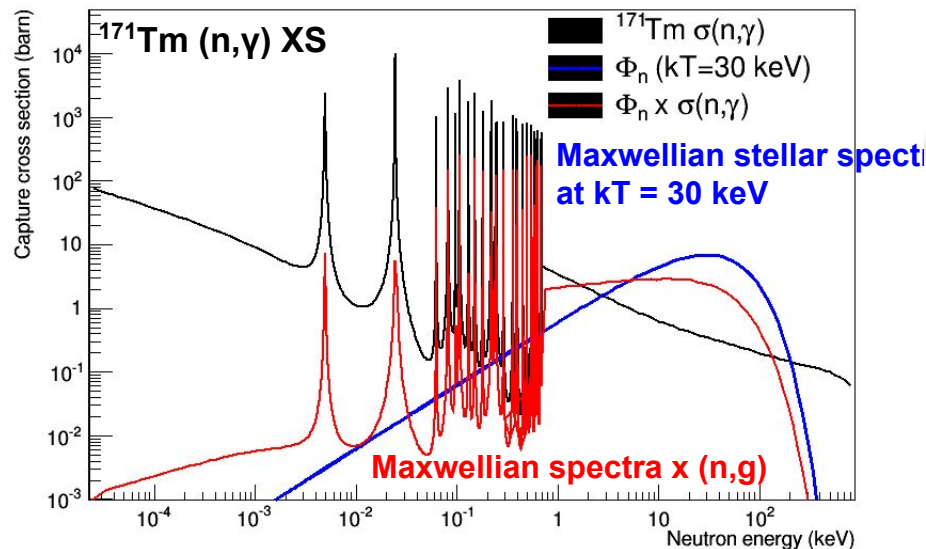
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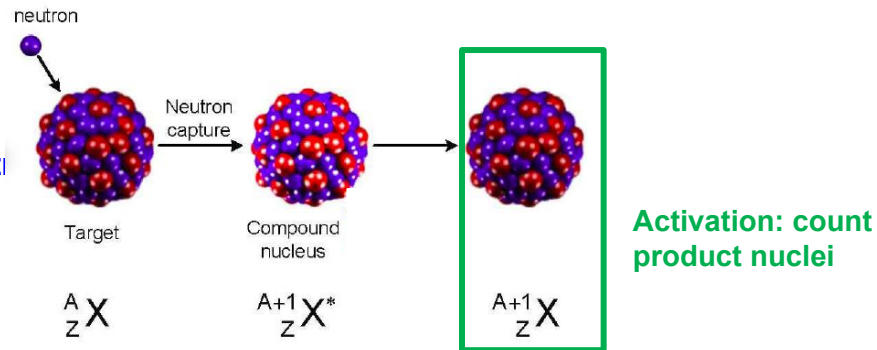
Is it possible to replicate a **Maxwellian neutron spectra** in the lab?
How do we measure **energy-dependent cross sections**?



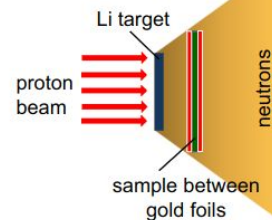
Beam: Pulsed white neutron beam + Time-of-flight technique
Technique: Measure (n,g) prompt g-ray yield vs neutron energy
Outcome: Cross section vs neutron energy → **MACS different kT**



Low energy accelerator + $^7\text{Li}(p,n)^7\text{Be}$

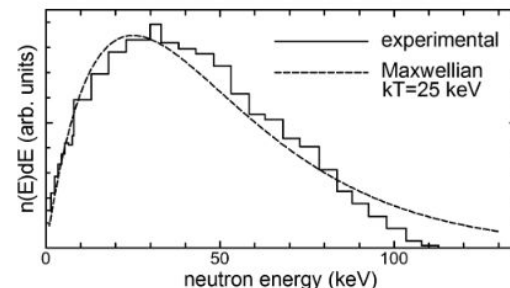



$^7\text{Li}(p,n)^7\text{Be}$



$E_p = 1912 \text{ keV}$

Quasi-maxwellian @ 25 keV



A. Wallner, [Phys. Rev. C 93, 045803 \(2016\)](#)

OPTION 2:

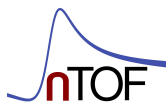
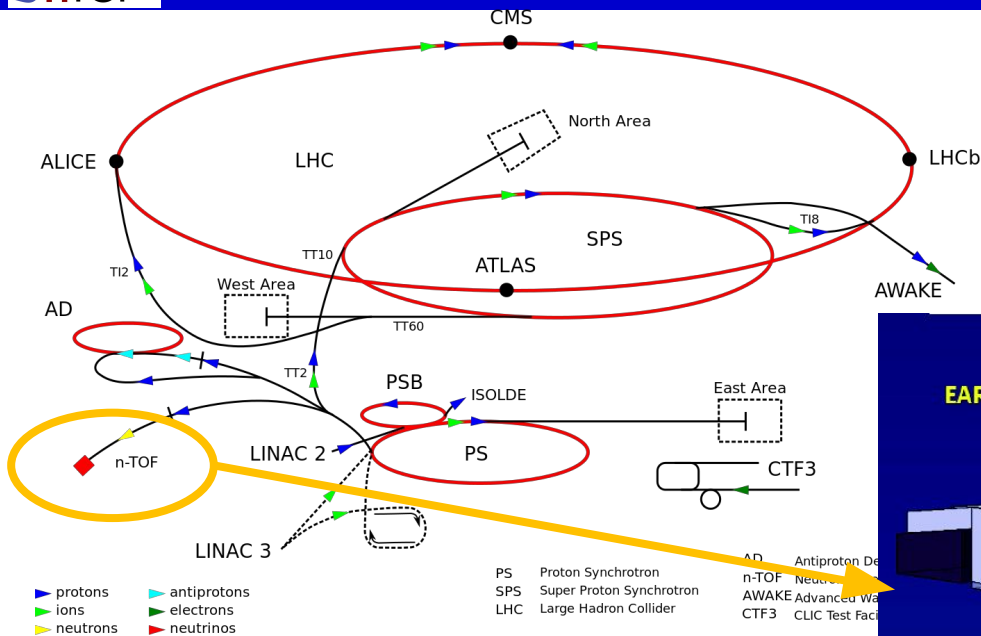
Beam: Quasi-maxwellian spectra

Technique: Count (n,g) product nuclei (activation)

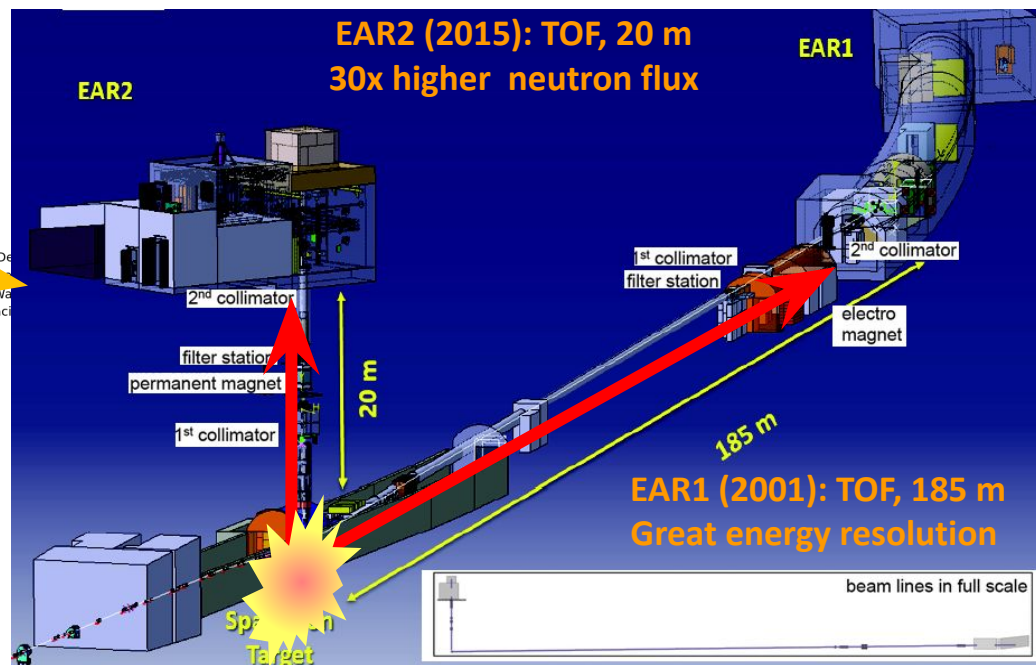
Outcome: Integral XS (SACS) → **MACS at kT = 25 keV**

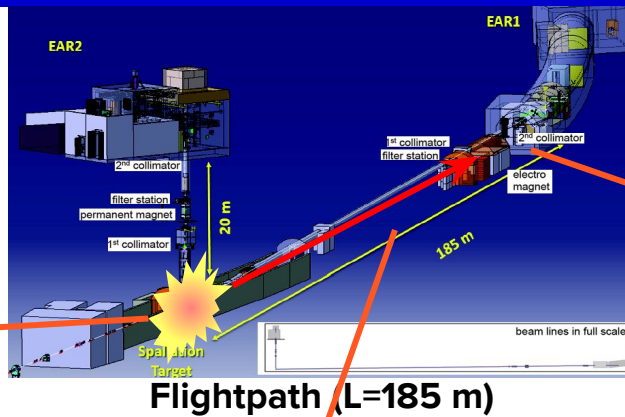


White neutrons sources
 "Stellar" neutron beams
 Fast neutrons



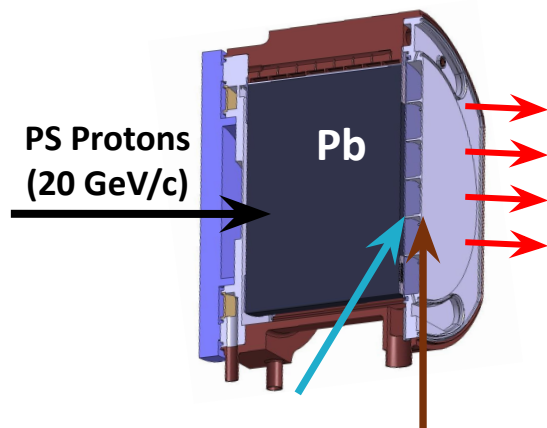
中国散裂中子源
China Spallation Neutron Source



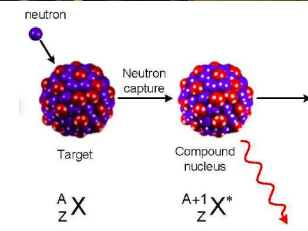
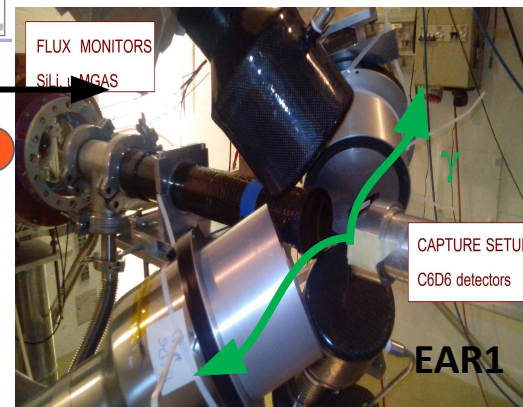


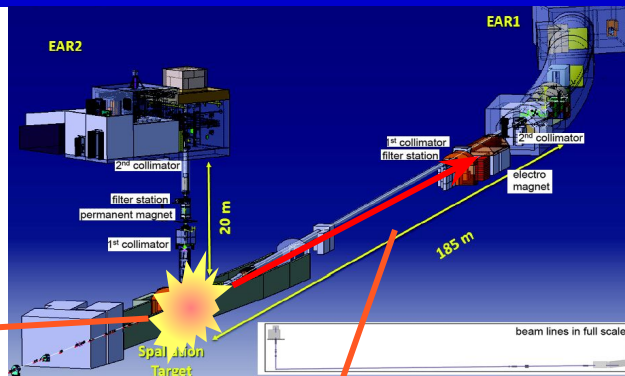
EXPERIMENTAL AREA

SPALLATION TARGET

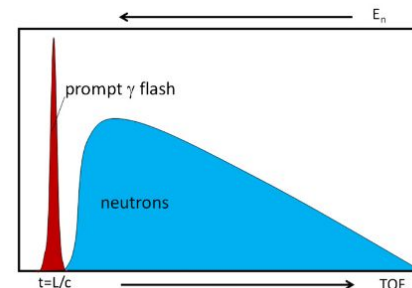


Neutron energies (meV to GeV)

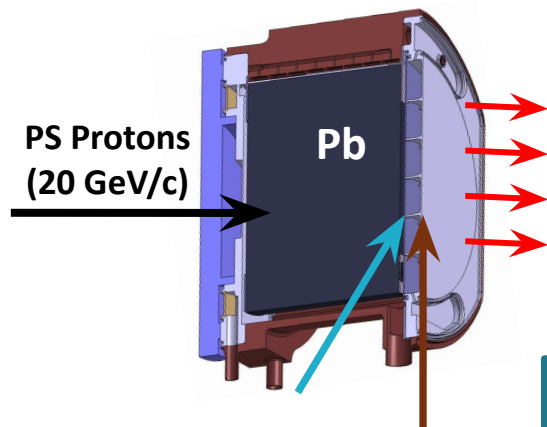




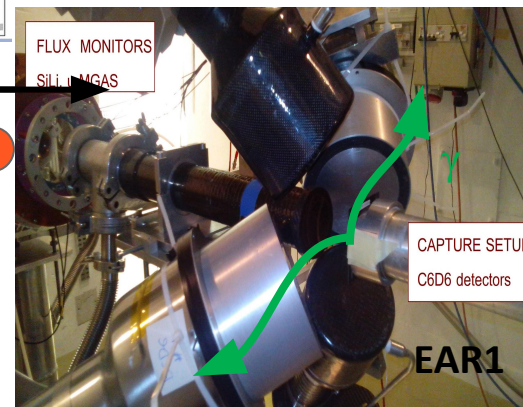
Flightpath (L=185 m)



**SPALLATION
TARGET**



**Neutron energies
(meV to GeV)**



**Time to energy
conversion**

$$E_n = \frac{1}{2}mv^2 = K^2 \frac{L^2}{t^2}$$

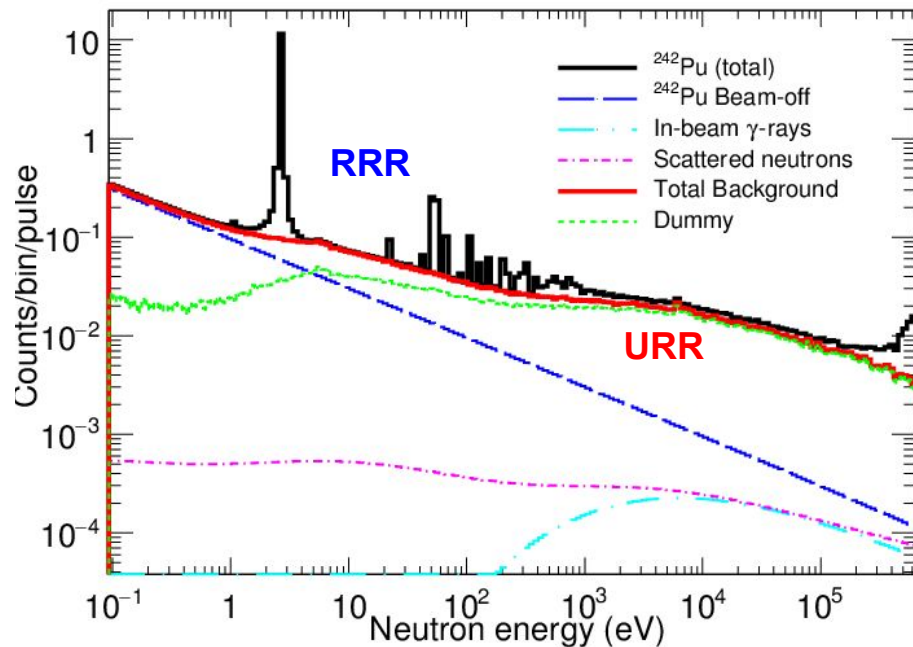
**Neutron
time-of-flight**

$$t = t_m - (t_\gamma - L/c)$$

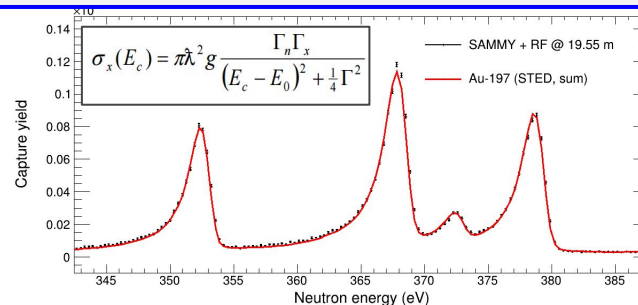
**Time
reference: t_γ**
Arrival of
prompt γ -rays

$$Y(E_n) = F_{norm}^{thr} \cdot \frac{C(E_n) - B(E_n)}{\Phi(E_n) \cdot \varepsilon_c}$$

Capture yield:
Many points + experimental effects

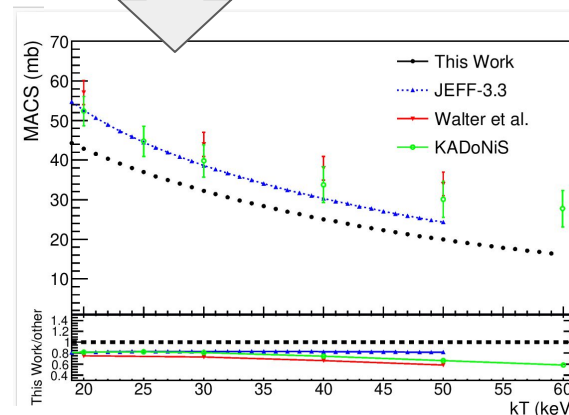


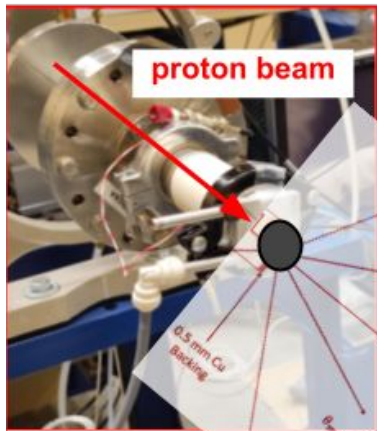
Cross-section: R-Matrix Analysis
Fit yield & extract cross section with Few parameters per resonance (E, Jp, Γ)



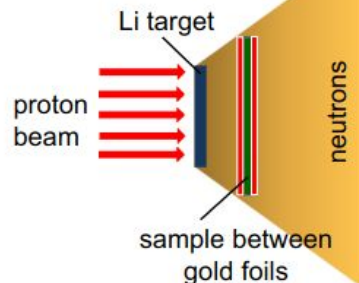
MACS from:

- Resonance (RRR)
- XS vs En (URR)



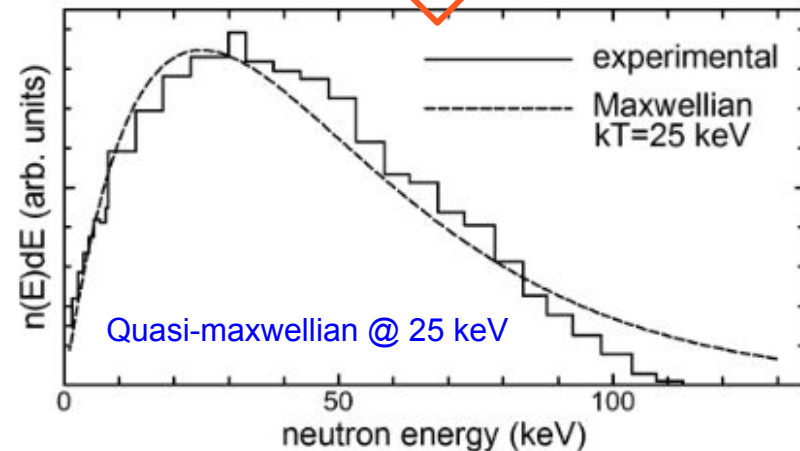


Low E (MeV) proton accelerators

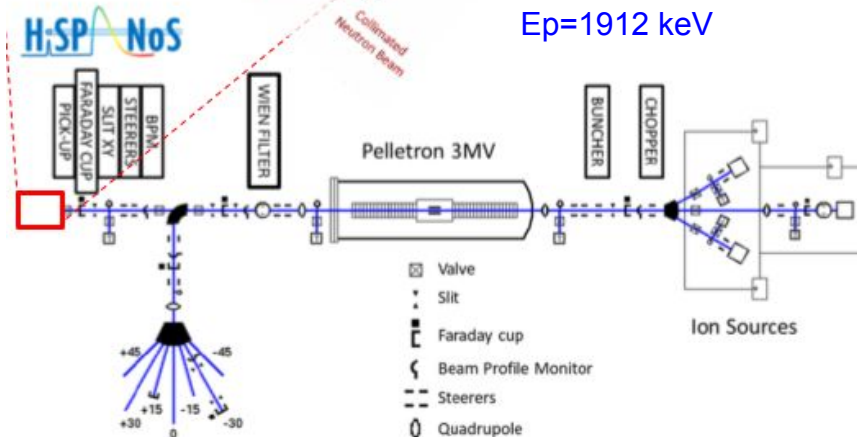


Stellar spectra: AGB (8, 23 keV) and
Massive stars (25, 90 keV)

STANDARD METHOD



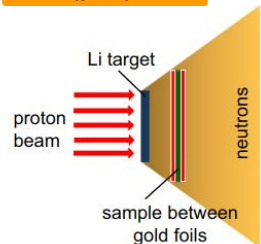
A. Wallner, [Phys. Rev. C 93, 045803 \(2016\)](#)



Adapted from P. Pérez-Matoto, [PhD Thesis \(U. Sevilla, 2024\)](#)

Low E (MeV) proton accelerators

${}^7\text{Li} (p,n) {}^7\text{Be}$

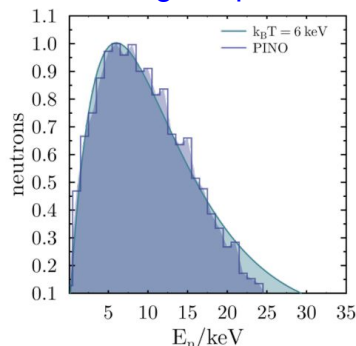


+ RECENT:
STELLAR
ENERGIES

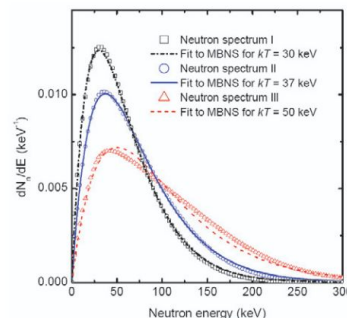
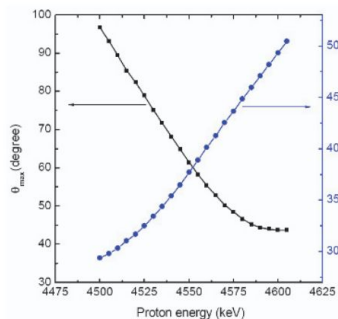
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LeNoS
Legnaro Neutron Source

Lower s-process temperatures
with ring-shaped samples



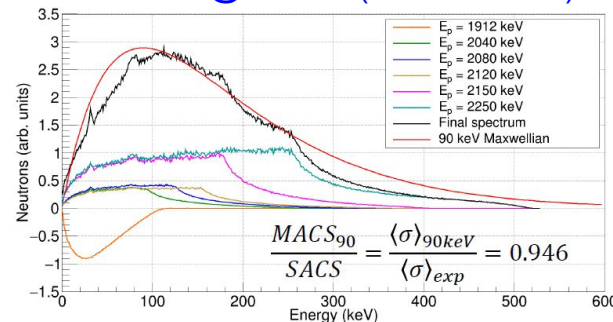
Tuning the kT with the angle at higher proton energies



P. Mastinu, Physics Procedia 26, p.261-273 (2012)

H₂SP NoS

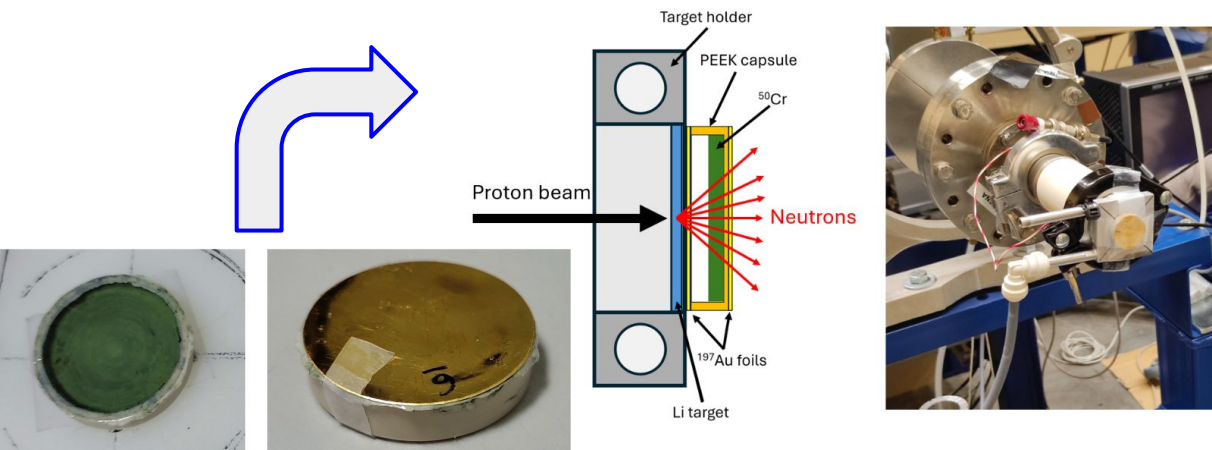
Different E_p combined $\rightarrow$
MACS @ 90 keV (massive stars)



P. Pérez-Maroto, EPJ Web of Conferences 294, 01004 (2024)

T. Hefrich, EPJ Web of Conferences 279, 06011 (2023)

2) Irradiation in the neutron beam

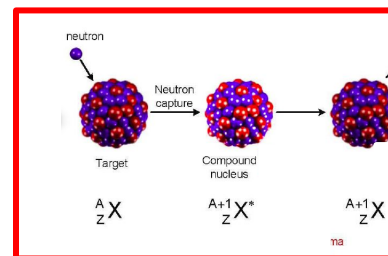


1) Main sample + reference/normalization samples

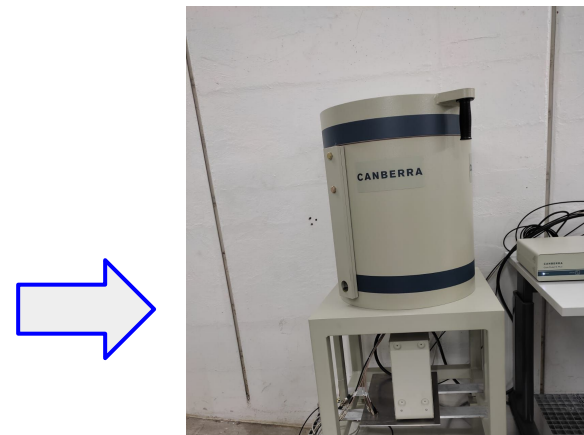
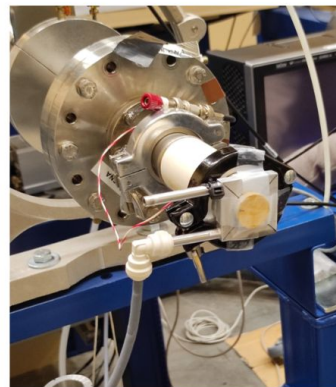
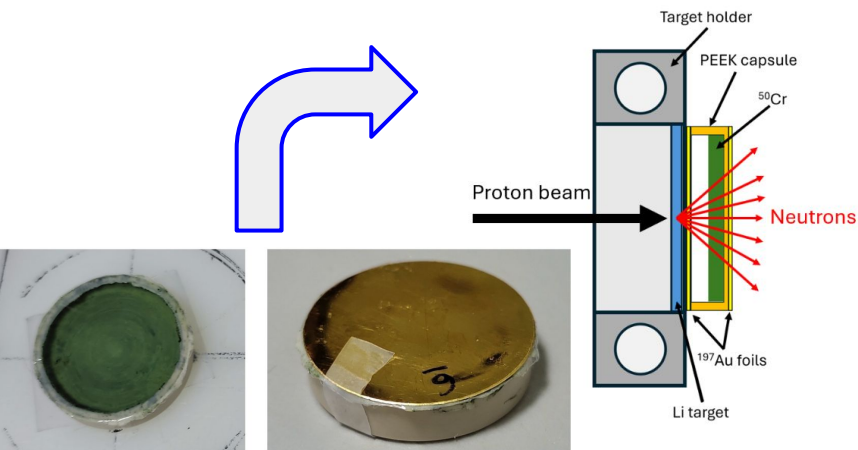
P. Pérez-Maroto, [PhD Thesis \(U. Sevilla, 2024\)](#)

Activation:

1) Irradiation in neutron beam $\rightarrow {}^A_Z\text{X}(n,g){}^{A+1}_Z\text{X}$



2) Irradiation in the neutron beam



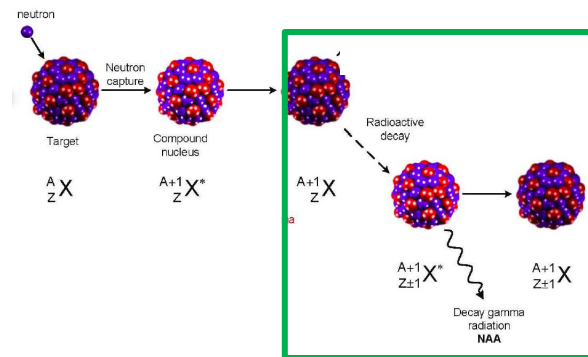
1) Main sample and reference/normalization samples sample

P. Pérez-Maroto, [PhD Thesis \(U. Sevilla, 2024\)](#)

Activation:

1) Irradiation in neutron beam $\rightarrow {}^A_Z\text{X}(n,g){}^{A+1}_Z\text{X}$

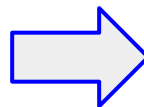
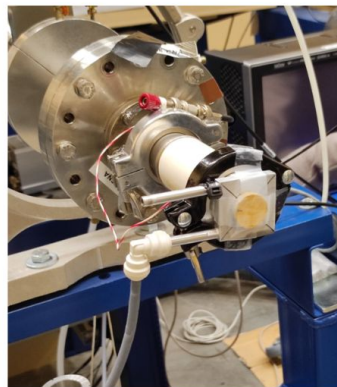
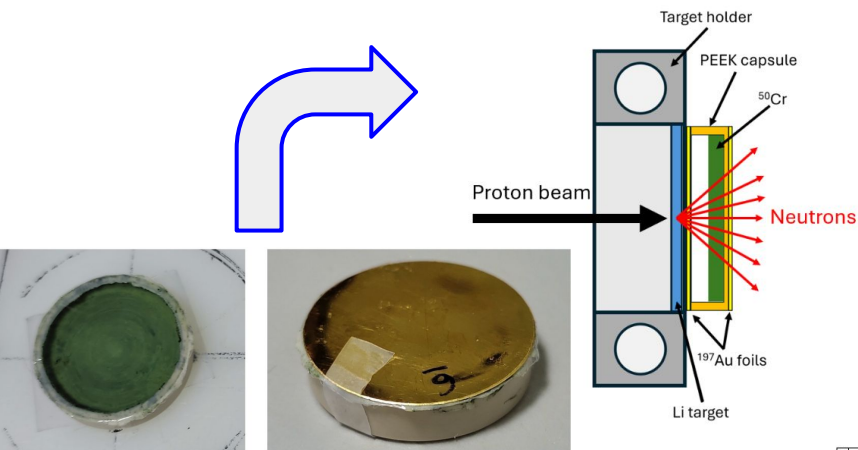
2) Count number of (n,g) reactions from the decay of ${}^{A+1}_Z\text{X}$



Only applies for unstable ${}^{A+1}_Z\text{X}$!

3) Decay measurement at: High resolution station (HPGe, LaBr3)

2) Irradiation in the neutron beam



3) Decay measurement at: resolution station (HPGe, LaBr3)



4) MACS derived from number of decays (relative to reference MACS, e.g. Au-197)

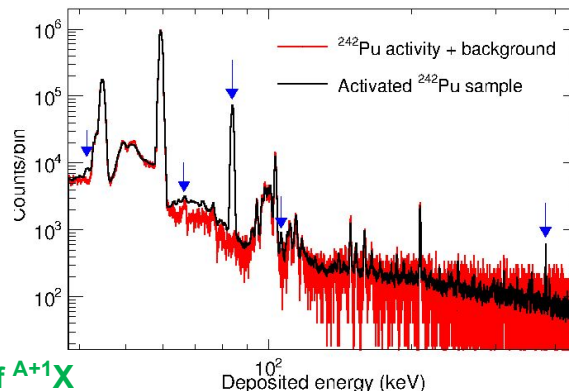
1) Main sample and reference samples (Au-197)

P. Pérez-Maroto, [PhD Thesis \(U. Sevilla, 2024\)](#)

Activation:

1) Irradiation in neutron beam $\rightarrow {}^AX(n,g){}^{A+1}X$

2) Count number of (n,g) reactions from the decay of ${}^{A+1}X$



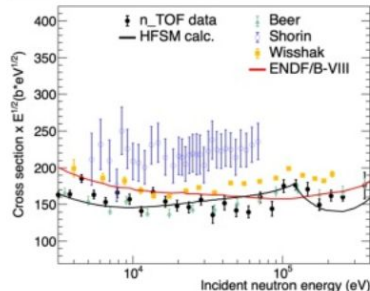
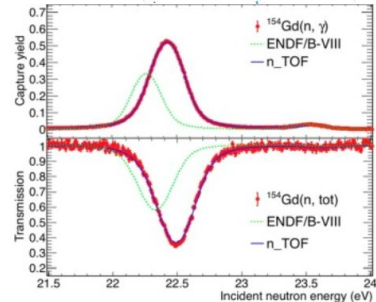
- The synthesis of (heavy) elements in the universe
- Nucleosynthesis & neutron capture: data needs for unstable and stable nuclei
- Direct neutron capture reactions: from stars to the lab
- State-of-the art & future of direct neutron capture reactions

s-only isotopes: Benchmark for s-process models

Physics Letters B 804 (2020) 135405

Measurement of the $^{154}\text{Gd}(n,\gamma)$ cross section and its astrophysical implications

A. Mazzone^{a,b}, S. Cristallo^{c,d}, O. Aberle^e, G. Alaerts^f, V. Alcañe^g, S. Amaducci^{h,i},



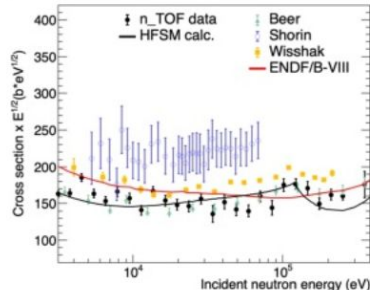
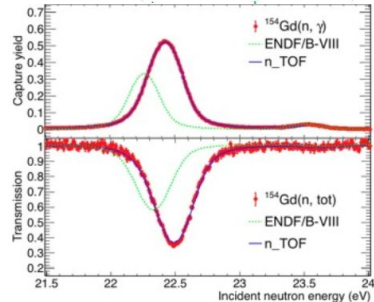
Gd-154(n,g): New cross section reduces discrepancy between observed and calculated solar s-only abundances

s-only isotopes: Benchmark for s-process models

Physics Letters B 804 (2020) 135405

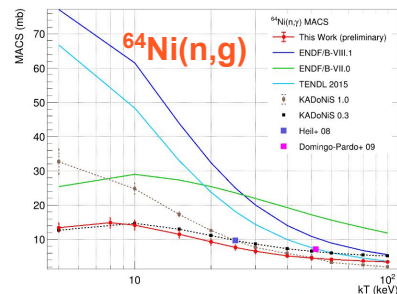
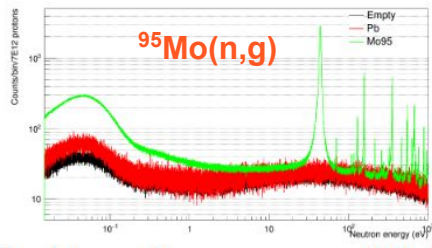
Measurement of the $^{154}\text{Gd}(n,\gamma)$ cross section and its astrophysical implications

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Gd-154(n,g): New cross section reduces discrepancy between observed and calculated solar s-only abundances

Ni, and Mo-isotopic anomalies in presolar SiC grains



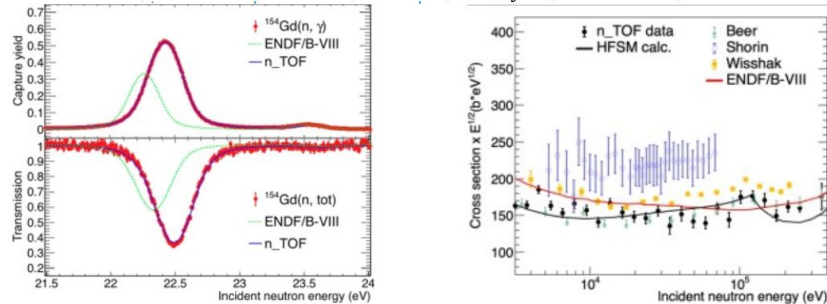
+ Details: Talks by M. Spelta & R. Mucciola

s-only isotopes: Benchmark for s-process models

Physics Letters B 804 (2020) 135405

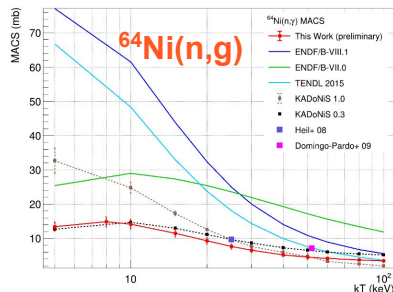
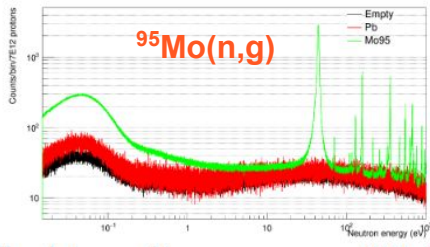
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Gd-154(n,g): New cross section reduces discrepancy between observed and calculated solar s-only abundances

Ni, and Mo-isotopic anomalies in presolar SiC grains



+ Details: Talks by M. Spelta & R. Mucciola

n-magic nuclei: s-process bottlenecks, very small (n,g) XS

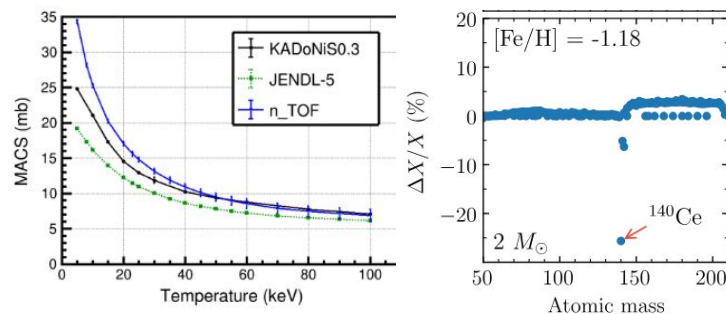
PHYSICAL REVIEW LETTERS 132, 122701 (2024)

Editors' Suggestion

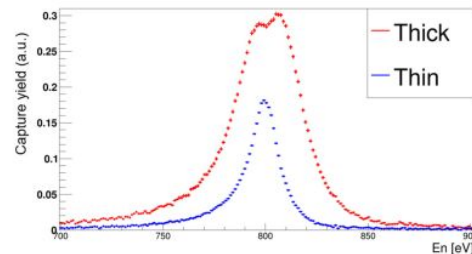
Featured in Physics

Measurement of the $^{140}\text{Ce}(n,\gamma)$ Cross Section at n_TOF and Its Astrophysical Implications for the Chemical Evolution of the Universe

S. Amaducci¹, N. Colonna², L. Cosentino¹, S. Cristallo^{3,4}, P. Finocchiaro¹, M. Krtićka⁵, C. Massimi^{6,7}



Ce-140(n,g): The new XS leads to an even smaller Ce-abundance



One of the lowest
MACS (25 keV)
2.89 mb!

Bi-209(n,g): repeat w/ high precision less in-beam g-rays (EAR2) and thinner sample (MS) and higher sensitivity setup (sTED)

REVIEW OF MODERN PHYSICS, VOLUME 83, JANUARY-MARCH 2011

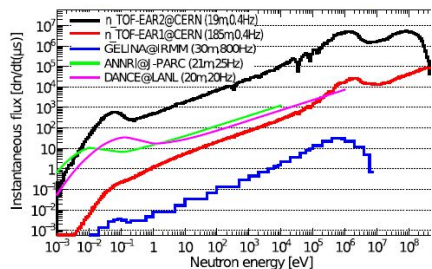
Sample	Half-life (yr)	Q value (MeV)	Comment
⁶³ Ni	100.1	β^- , 0.066	C. Lederer <i>et al.</i> , Phys. Rev. Lett. 110, 022501 (2013)
⁷⁹ Se	2.95×10^5	β^- , 0.159	Important branching, constrains s -process temperature in massive stars
⁸¹ Kr	2.29×10^5	EC, 0.322	Part of ⁷⁹ Se branching
⁸⁵ Kr	10.73	β^- , 0.687	Important branching, constrains neutron density in massive stars
⁹⁵ Zr	64.02 d	β^- , 1.125	Not feasible in near future, but important for neutron density low-mass AGB stars
¹³⁴ Cs	2.0652	β^- , 2.059	Important branching at $A = 134, 135$, sensitive to s -process temperature in low-mass AGB stars, measurement not feasible in near future
¹³⁵ Cs	2.3×10^6	β^- , 0.269	So far only activation measurement at $kT = 25$ keV by Patronis <i>et al.</i> (2004)
¹⁴⁷ Nd	10.981 d	β^- , 0.896	Important branching at $A = 147/148$, constrains neutron density in low-mass AGB stars
¹⁴⁷ Pm	2.6234	β^- , 0.225	Part of branching at $A = 147/148$
¹⁴⁸ Pm	5.368 d	β^- , 2.464	Not feasible in the near future
¹⁵¹ Sm	90	β^- , 0.076	U. Abbondanno <i>et al.</i> , Phys. Rev. Lett. 93, 161103 (2004)
¹⁵⁴ Eu	8.593	β^- , 1.978	Complex branching at $A = 154, 155$, sensitive to temperature and neutron density
¹⁵⁵ Eu	4.753	β^- , 0.246	So far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1995)
¹⁵³ Gd	0.658	EC, 0.244	Part of branching at $A = 154, 155$
¹⁶⁰ Tb	0.198	β^- , 1.833	Weak temperature-sensitive branching, very challenging experiment
¹⁶³ Ho	4570	EC, 0.0026	Branching at $A = 163$ sensitive to mass density during s process, so far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1996b)
¹⁷⁰ Tm	0.352	β^- , 0.968	Important branching, constrains neutron density in low-mass AGB stars
¹⁷¹ Tm	1.921	β^- , 0.098	Part of branching at $A = 170, 171$
¹⁷⁹ Ta	1.82	EC, 0.115	Crucial for s -process contribution to ¹⁸⁰ Ta, nature's rarest stable isotope
¹⁸⁵ W	0.206	β^- , 0.432	Important branching, sensitive to neutron density and s -process temperature in low-mass AGB stars
²⁰⁴ Tl	3.78	β^- , 0.763	Determines ²⁰⁵ Pb/ ²⁰⁵ Tl clock for dating of early Solar System

F. Käppeler *et al.*,
[Rev. Mod. Phys 83, 157](#)
(2011)

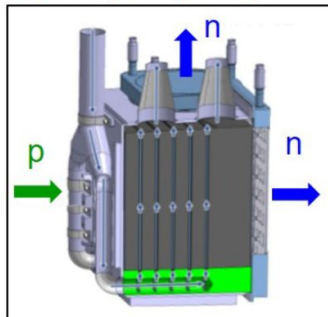
**Challenging
measurements:**
Radioactive
isotopes, low
masses

Before 2015: only **2/21** of the key s -process isotopes measured by TOF

Upgrade I: instantaneous neutron flux $\rightarrow$ n_TOF EAR2 (+ Target upgrade)



Target#3 (2021-present)



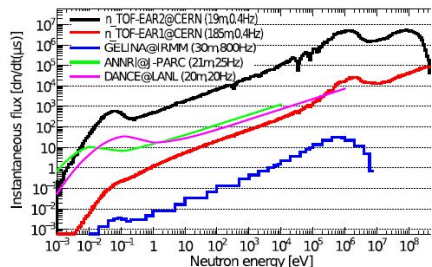
n-TOF EAR2: highest
Instantaneous flux
(FOM for
neutron-to-activity
ratio)

x300 instantaneous
flux of n_TOF-EAR1!

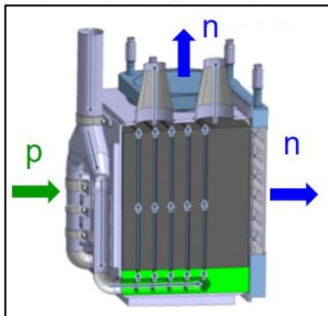
J. Lerendegui-Marco et al., *Eur. Phys. J. A* **52**, 100 (2016).

J.A. Pavón-Rodríguez et al., *Eur. Phys. J. A* (arxiv) (2025)

Upgrade I: instantaneous neutron flux $\rightarrow$ n_TOF EAR2 (+ Target upgrade)



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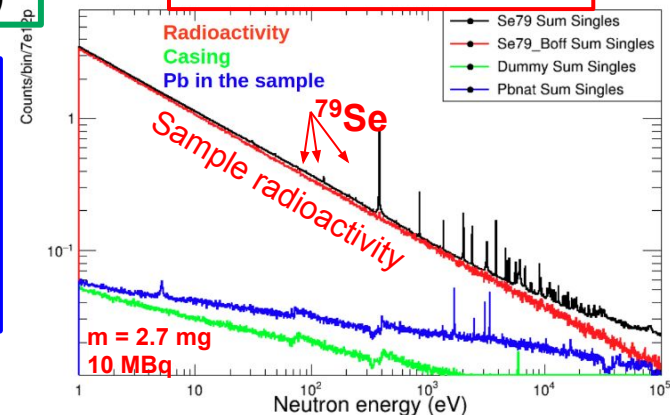
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x300 instantaneous
flux of n_TOF-EAR1!

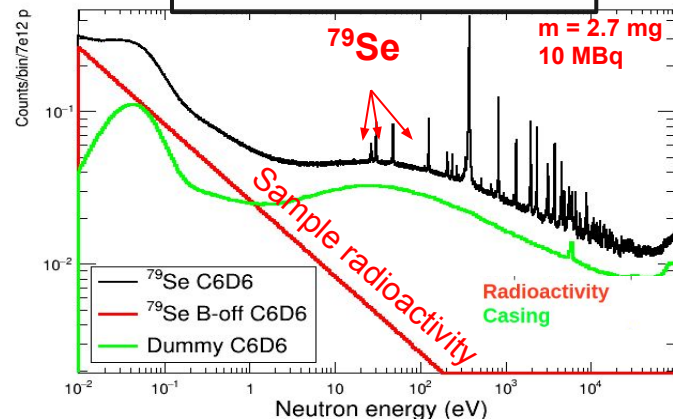
J. Lerendegui-Marco et al., *Eur. Phys. J. A* **52**, 100 (2016).

J.A. Pavón-Rodríguez et al., *Eur. Phys. J. A* (arxiv) (2025)

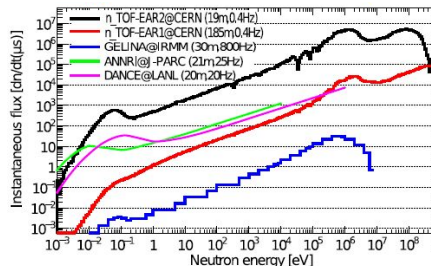
$^{79}\text{Se}(n,\gamma)$ @ n_TOF EAR1



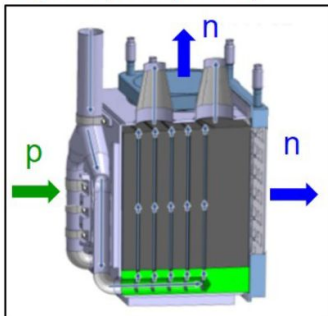
$^{79}\text{Se}(n,\gamma)$ @ n_TOF EAR2



Upgrade I: instantaneous neutron flux $\rightarrow$ n_TOF EAR2 (+ Target upgrade)



Target#3 (2021-present)



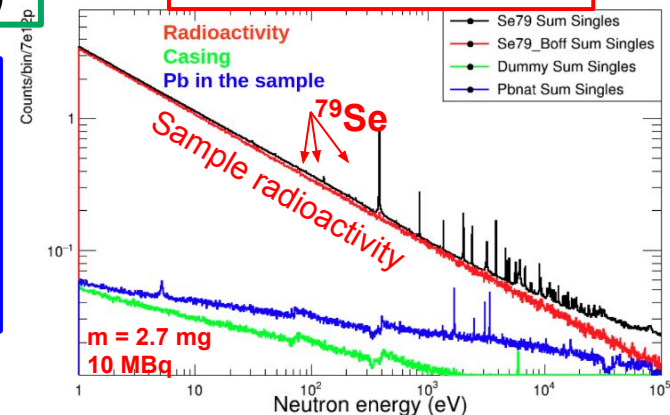
n-TOF EAR2: highest
Instantaneous flux
(FOM for
neutron-to-activity
ratio)

x300 instantaneous
flux of n_TOF-EAR1!

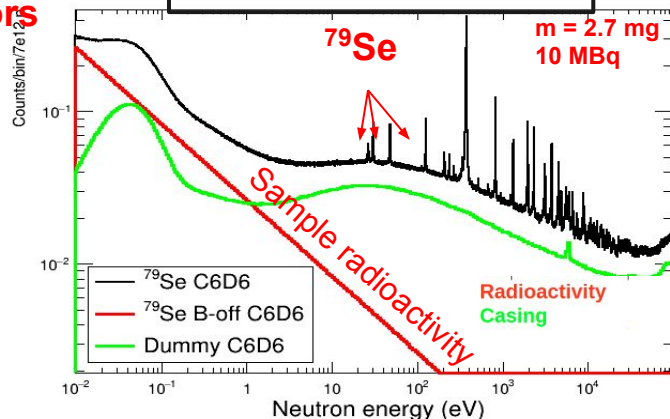


New generation of detectors

$^{79}\text{Se}(n,\gamma)$ @ n_TOF EAR1



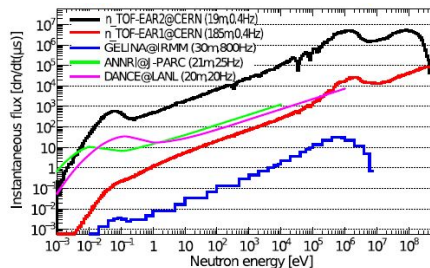
$^{79}\text{Se}(n,\gamma)$ @ n_TOF EAR2



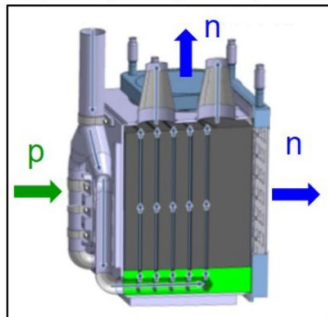
J. Lerendegui-Marco et al., *Eur. Phys. J. A* **52**, 100 (2016).

J.A. Pavón-Rodríguez et al., *Eur. Phys. J. A* (arxiv) (2025)

Upgrade I: instantaneous neutron flux $\rightarrow$ n_TOF EAR2 (+ Target upgrade)



Target#3 (2021-present)



n-TOF EAR2: highest
Instantaneous flux
(FOM for
neutron-to-activity
ratio)

x300 instantaneous
flux of n_TOF-EAR1!

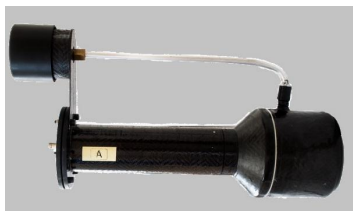


New generation of detectors

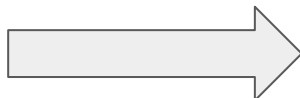
J. Lerendegui-Marco et al., *Eur. Phys. J. A* **52**, 100 (2016).

J.A. Pavón-Rodríguez et al., *Eur. Phys. J. A* (arxiv) (2025)

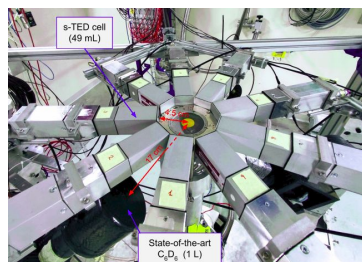
Upgrade II: detectors for high C. rates (+ enhanced signal-to-bckg)



Large TED (~1L)
C.R. max ~ 20 MHz

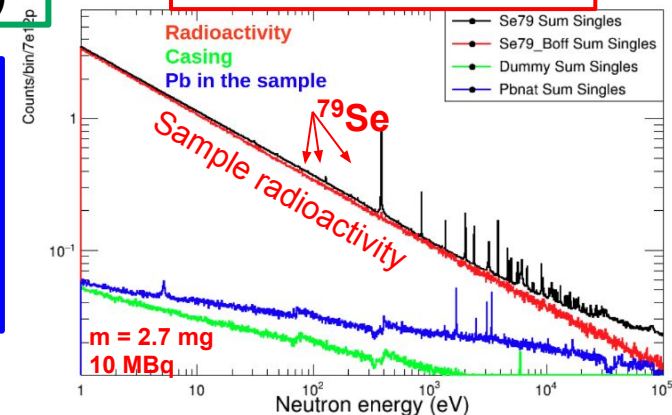


s-TED
(~50mL)
C.R. max ~1 MHz

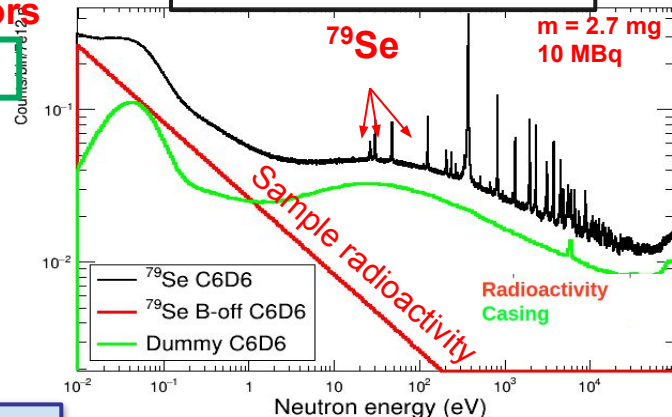


V. Alcayne, *EPJ Web of Conferences* **284**, 01043 (2023)

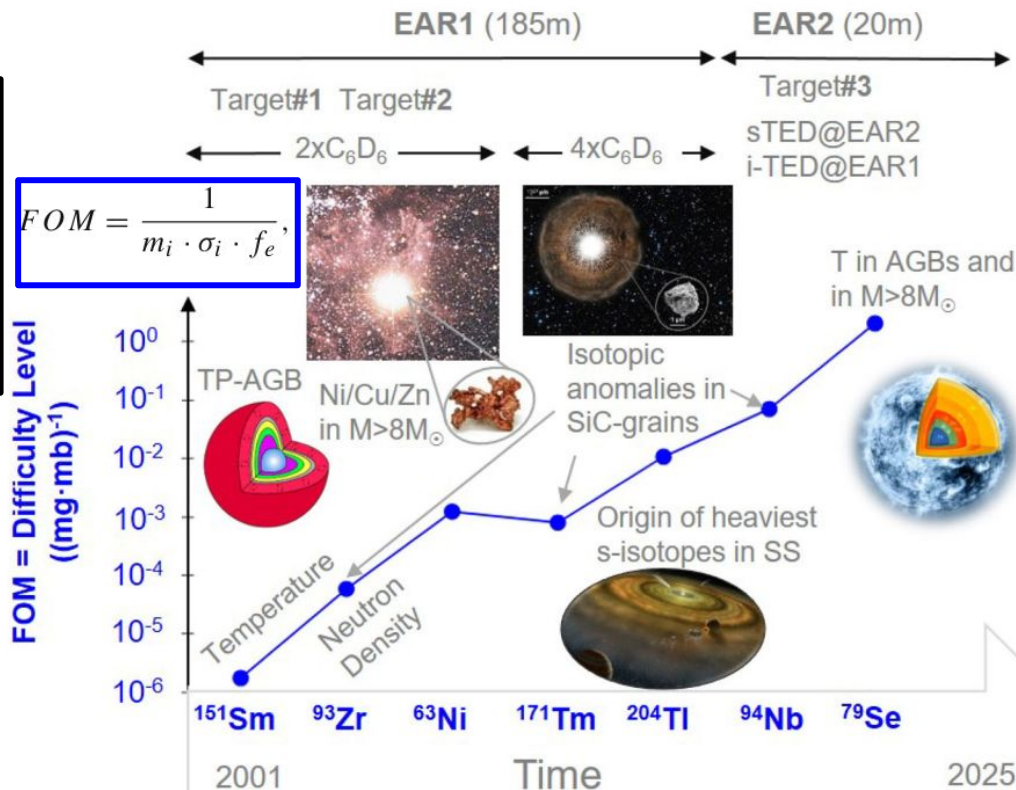
$^{79}\text{Se}(n,\gamma)$ @ n_TOF EAR1



$^{79}\text{Se}(n,\gamma)$ @ n_TOF EAR2

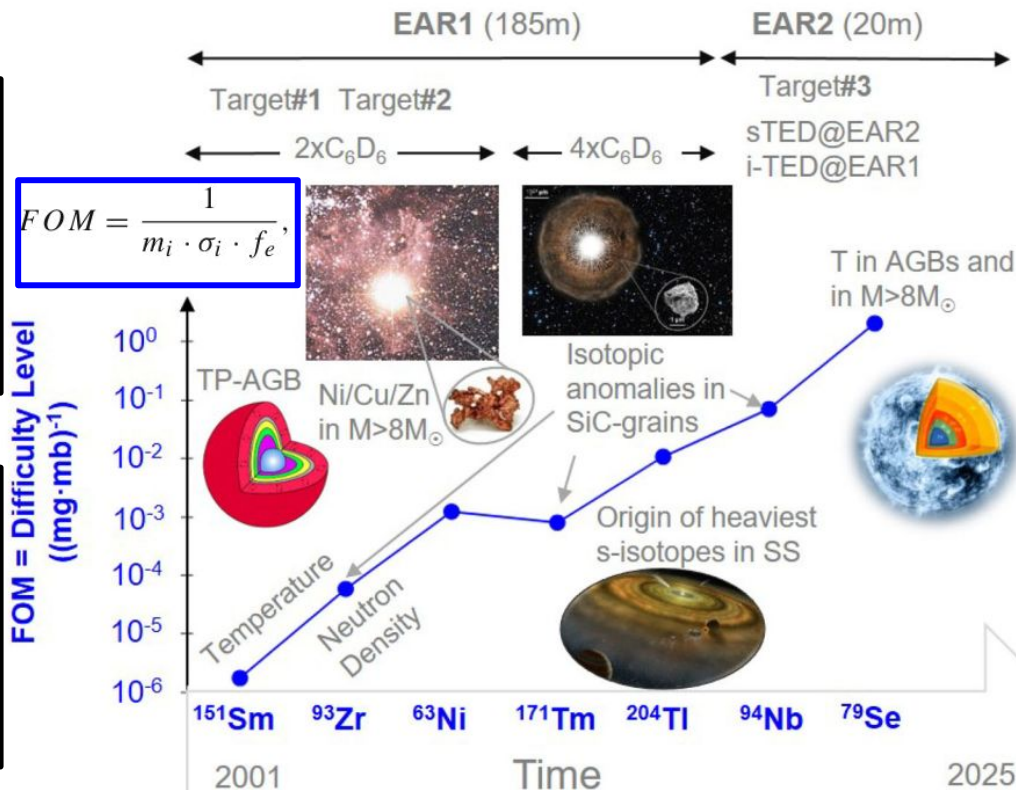


Facility upgrades
 (EAR2, new spallation target, ...)
 +
Advanced detectors
 =
ENHANCED DIFFICULTY LEVEL



Facility upgrades
(EAR2, new spallation target, ...)
+
Advanced detectors
=
ENHANCED DIFFICULTY LEVEL

Before 2015: only **few** key s-process
branchings measured by TOF:
 ^{151}Sm , ^{93}Zr , ^{63}Ni
2015-2018: ^{171}Tm (TOF & Act), ^{204}Tl
(TOF), ^{147}Pm (Act)
Very recent progresses: ^{79}Se , ^{94}Nb

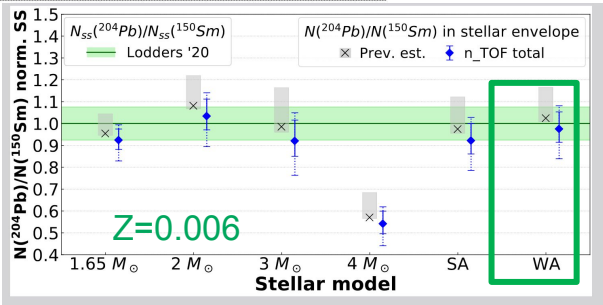
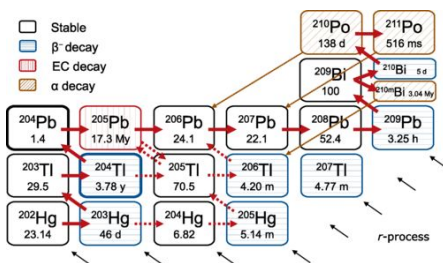
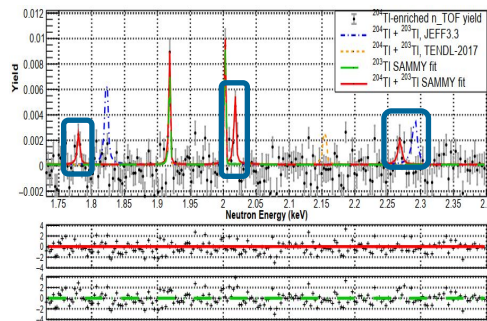


Tl-204(n,γ)

$m = 9 \text{ mg}$
 $(+200 \text{ mg } ^{203}\text{Tl})$
 $\sigma = 260(90) \text{ mb}$

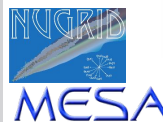
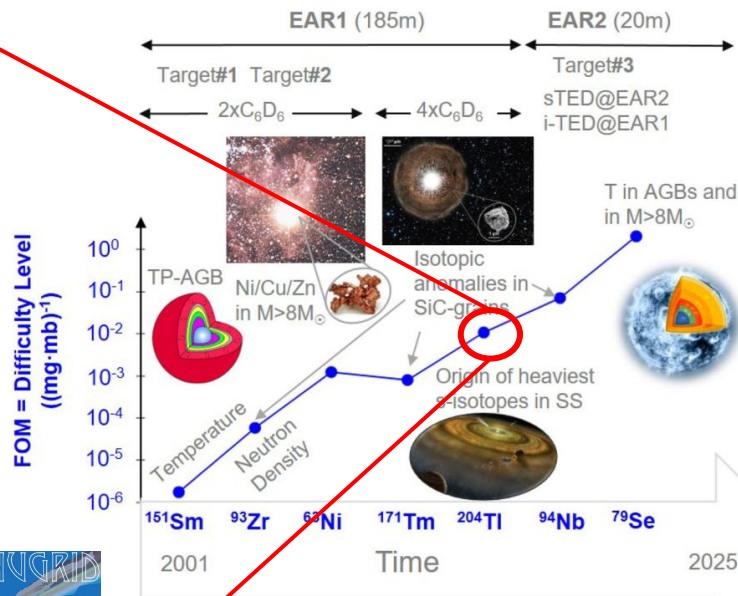


First measurement ever
TOF-unique case



Uncertainty in the abundance of the heaviest s-only 204Pb has been reduced from ~30% down to +8%/-6%

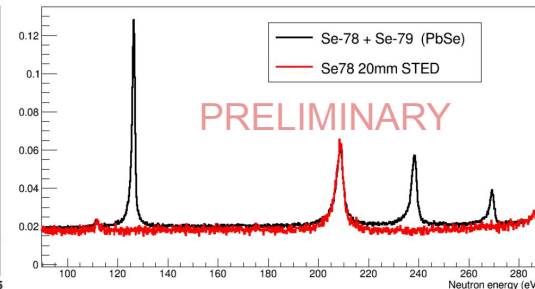
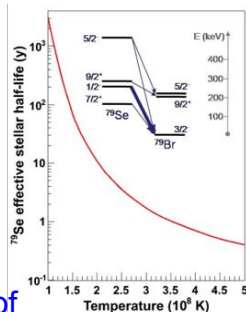
A. Casanovas, [Phys. Rev. Letters 133, 052702 \(2024\).](#)



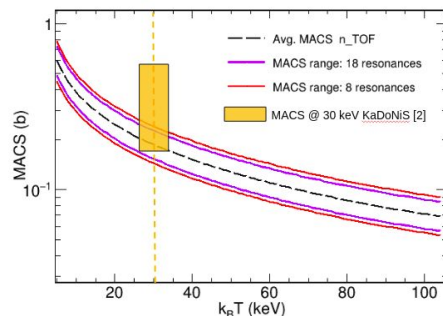
Se-79(n,γ)

First measurement ever
TOF-unique case

$m = 2.6 \text{ mg}$
 $(+3 \text{ g } ^{208}\text{Pb}^{78}\text{Se})$
 $\sigma = 263(43) \text{ mb}$



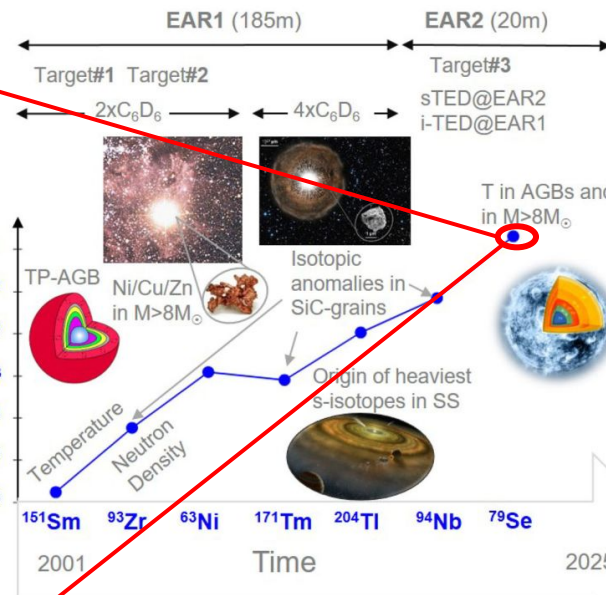
PRELIMINARY



Expected: 10-15 Resonances →

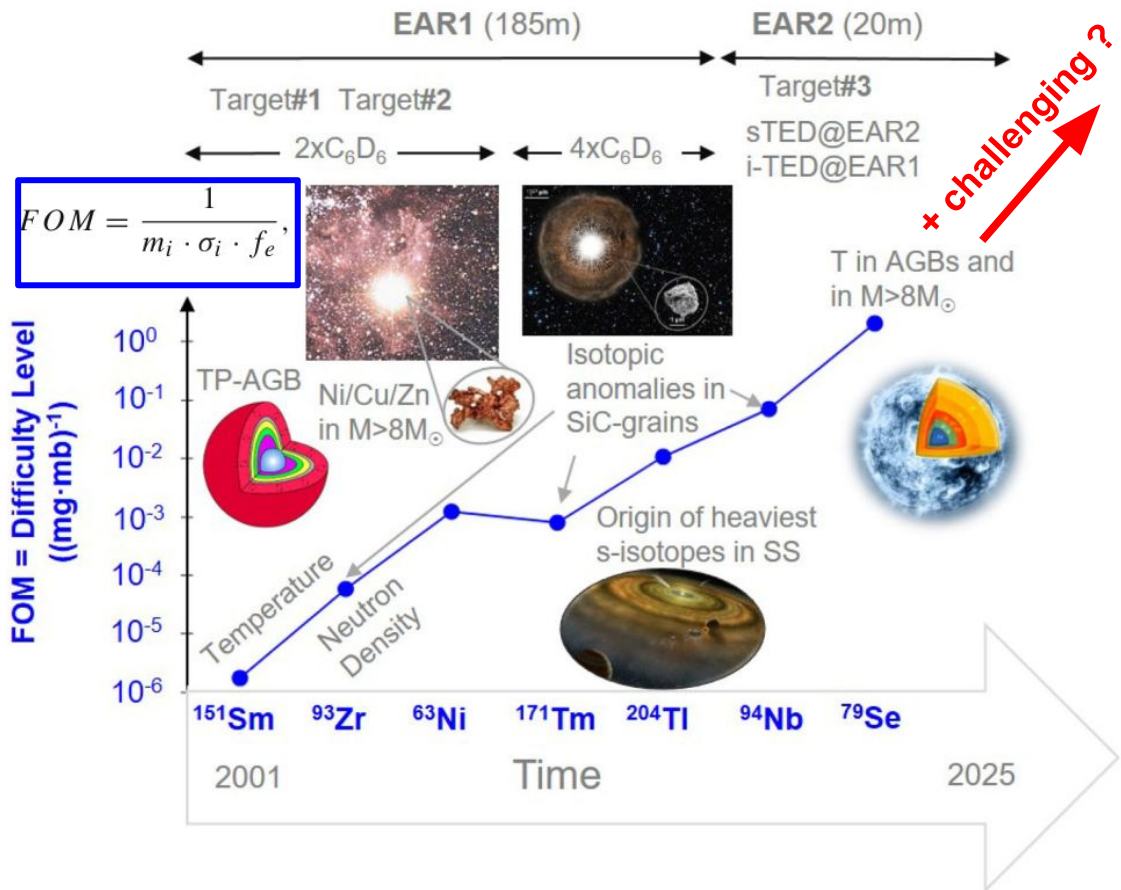
experimental constraint of MACS down to ~30%

Lerendegui-Marco, J. et al., [EPJ Web Conf. 279, 13001 \(2023\)](https://doi.org/10.1051/epjconf/202327913001).



Challenges unstable s-process branching points:

1. TOF requires mass $\geq 5 \times 10^{17} - 10^{18}$ atoms (EAR2)
2. Mass/Activity: **radiation hazard & activity background**
 → Many (n,g) reactions still not accessible via TOF

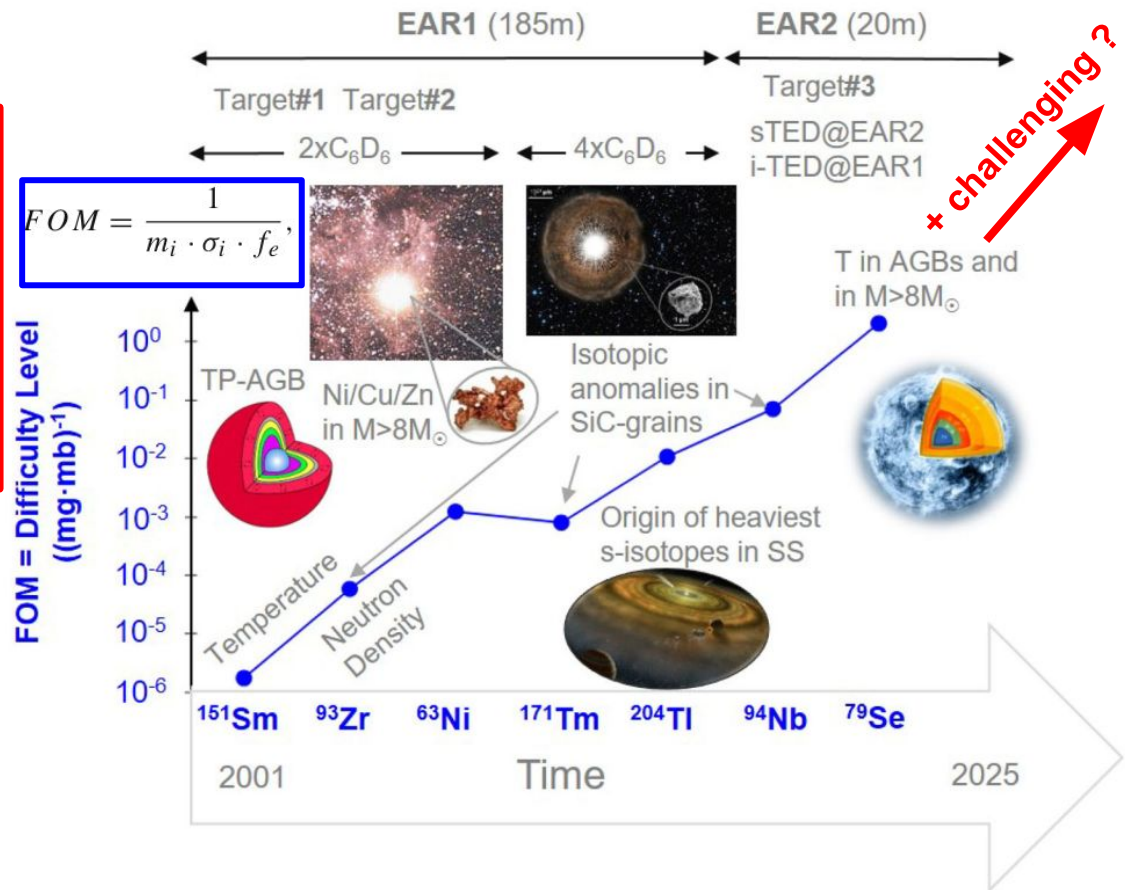


Challenges unstable s-process branching points:

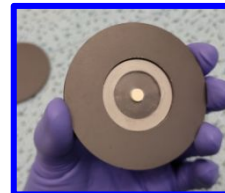
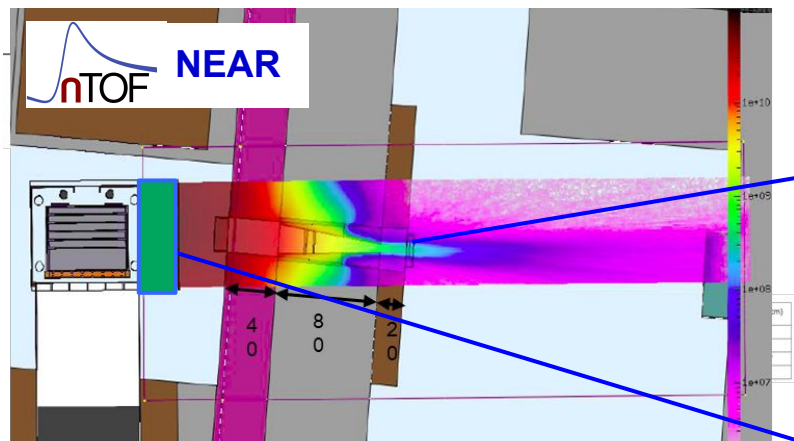
1. TOF requires mass $\geq 5 \times 10^{17} - 10^{18}$ atoms (EAR2)
2. Mass/Activity: **radiation hazard & activity background**
 → Many (n,g) reactions still not accessible via TOF

SOLUTION:

Activation in high flux facilities
(HIGHER SENSITIVITY → LOWER MASS REQUIRED)

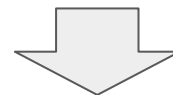


n_TOF NEAR: (n,g) activations with very high flux (~ 100 EAR2): small masses, unstable isotopes

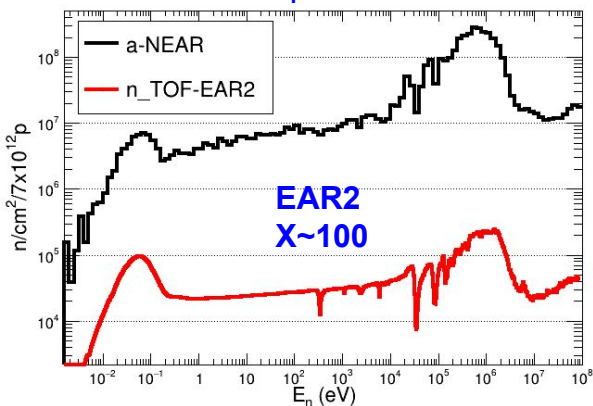


Different Stellar
temperatures
kT: 0.1~500 keV

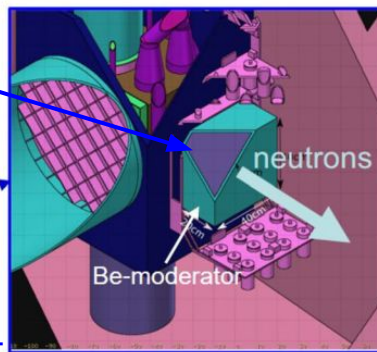
$^{10}\text{B}_4\text{C}$ Filters



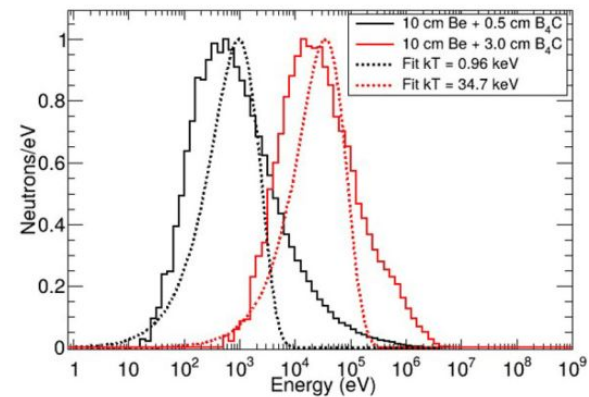
9×10^8 n/cm²/pulse $\rightarrow$ Mass limit $\sim 10^{15}$ atoms



White neutron
beam moderated +
filtered
 $\rightarrow$ "quasi-Stellar"

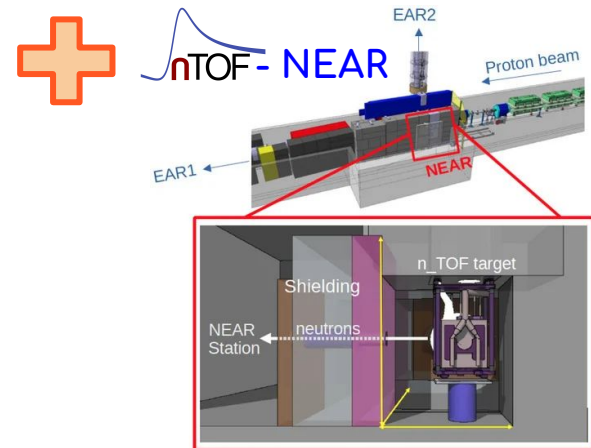
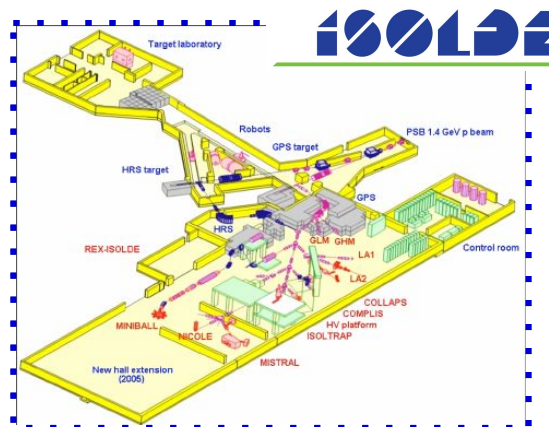


Be / Al₂O₃ moderator

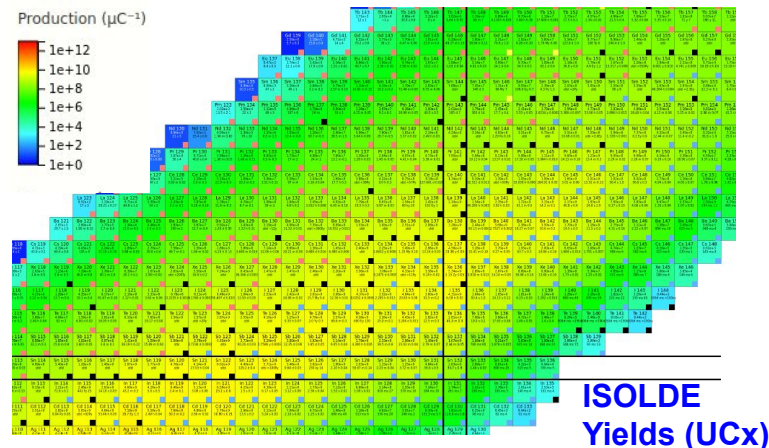


Synergy RIB-neutron facilities

Goal: Produce samples of relevant unstable nuclei at **RIB** & measure MACS at **high-flux facility**

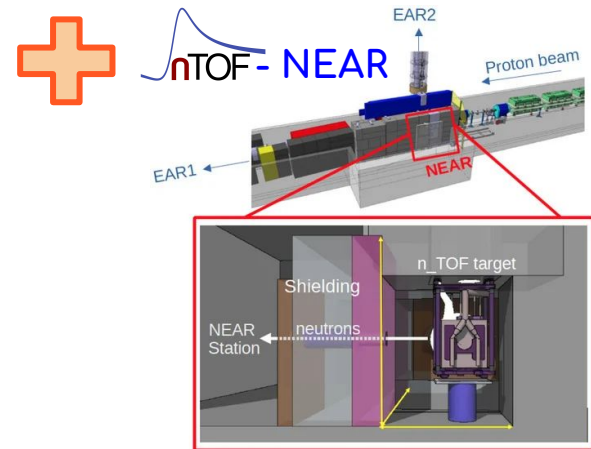
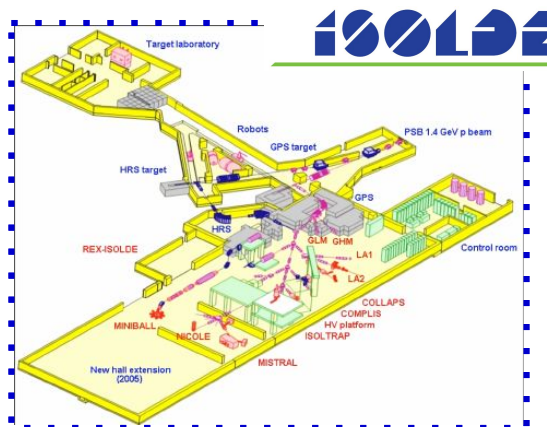


- **Example at CERN: ISOLDE + n_TOF-NEAR**
- **High flux + proximity:** Smaller masses & shorter-lived isotopes would be accessible
- **Examples:**
 - 59Fe, 134Cs, 135Cs, 148Pm, 154Eu, 155Eu, 160Tb, 170Tm, and 181Hf (s-process),
 - Cs-137, Ce-144, 66Ni, 72Zn (i-process)



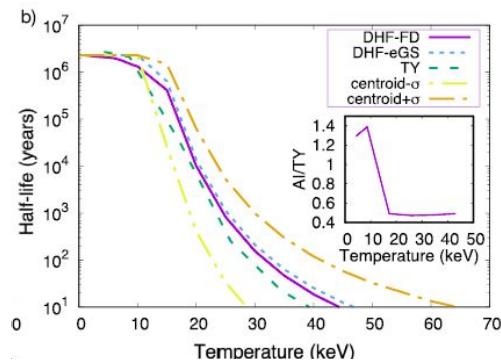
Synergy RIB-neutron facilities

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- **Example at CERN: ISOLDE + n_TOF-NEAR**
- **High flux + proximity:** Smaller masses & shorter-lived isotopes would be accessible
- **Examples:**
 - 59Fe, 134Cs, **135Cs**, 148Pm, 154Eu, 155Eu, 160Tb, 170Tm, and 181Hf (s-process),
 - Cs-137, Ce-144, 66Ni, 72Zn (i-process)

J. Lerendegui, S. Carollo et al.,
[CERN-INTC-2022-040; INTC-P-641](#)



Cs-135: branching s-process, New stellar beta decay rates (Talioli)

TOF facility in astro-range: FRANZ (?)

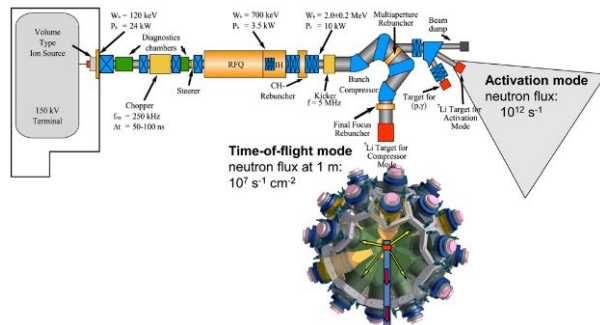
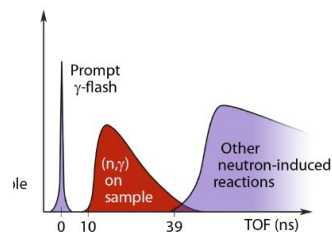


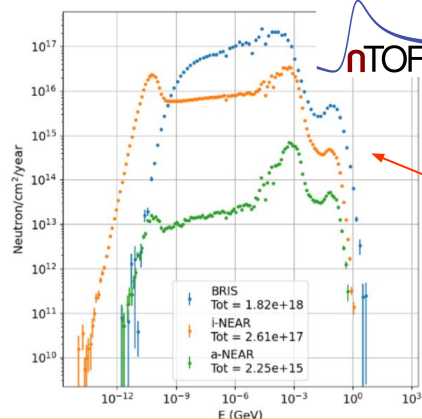
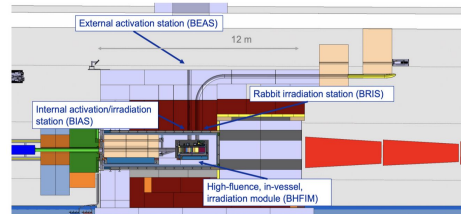
Figure 10. Sketch of the neutron facility FRANZ, currently under construction at Goethe University Frankfurt/Germany.



TOF (L=1m):
 $\sim 1e7 \text{ n/cm}^2/\text{s}$

50 mA p-beam + high repetition rate $\rightarrow$
Integrated flux $\sim x5$ (n_TOF-EAR2)

Activation facility nAct @ BDF/SHIP (SPS,CERN) (>2030)



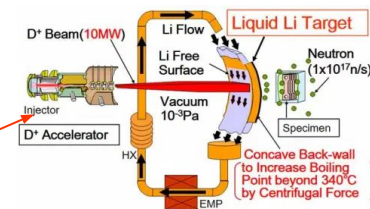
NEAR-like
spectrum
shaping

Flux (a-NEAR x1000) + fast rabbit
(T1/2 > min-h)

(n,g) at IFMIF-DONES? (>2033)



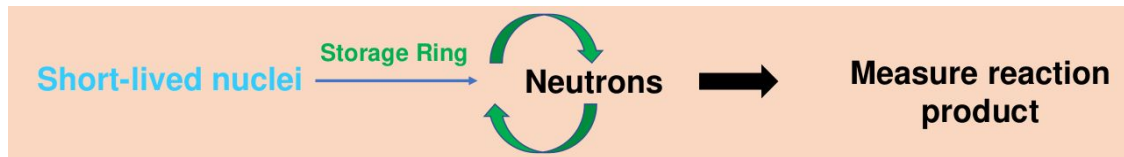
10 MW 40 MeV d-beam:
($10^{13} \text{ n/cm}^2/\text{s}$) @ 14 MeV
:UNIQUE FACILITY



Ancillary science:
quasi-stellar flux
 $\sim 1e9 - 4e9 \text{ n/cm}^2/\text{s}$
($x \sim 100$ a-NEAR,
 $x \sim 200$ HISPANoS)

DIRECT (N,G) IN INVERSE KINEMATICS

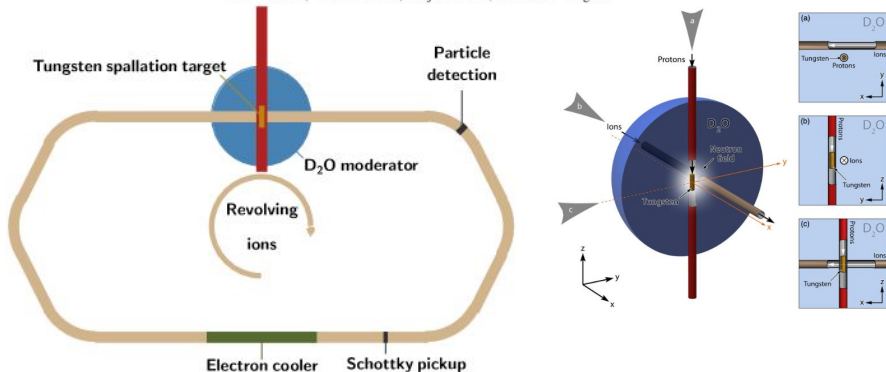
Aim: short-lived (s- and i-process)



PHYSICAL REVIEW ACCELERATORS AND BEAMS **20**, 044701 (2017)

Spallation-based neutron target for direct studies of neutron-induced reactions in inverse kinematics

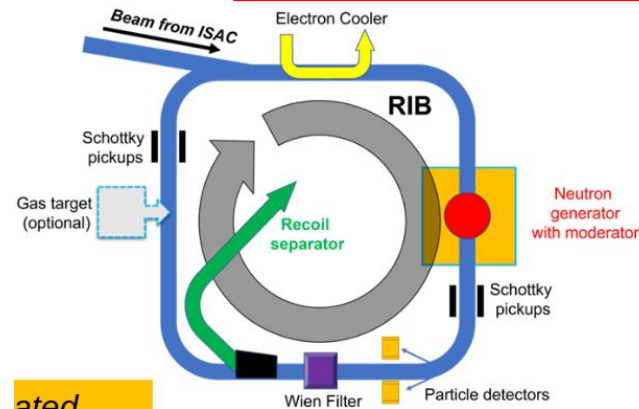
René Reifarth,^{*} Kathrin Göbel, Tanja Heftrich, and Mario Weigand



Neutron target facility @ LANL

800 MeV p, 100 mA + spallation target + moderation
n density = $8 \cdot 10^9$ n/cm²

- s-process: ^{134}Cs ($t_{1/2} = 2.06$ y), ^{147}Nd ($t_{1/2} = 11$ d), ^{148}Pm ($t_{1/2} = 5.37$ d), ^{153}Gd ($t_{1/2} = 240.4$ d), ^{160}Tb ($t_{1/2} = 72.3$ d), ^{170}Tm ($t_{1/2} = 128.6$ d), and ^{204}Tl ($t_{1/2} = 3.78$ y).
- i-process: ^{66}Ni ($t_{1/2} = 54.6$ h), ^{75}Ga ($t_{1/2} = 126$ s), ^{135}I ($t_{1/2} = 6.58$ h), ^{141}Ba ($t_{1/2} = 18.27$ m), and ^{141}La ($t_{1/2} = 3.92$ h).



TRISR @ ARIEL-TRIUMF

Compact high-flux neutron generator:
 10^8 n/s/cm²

- **Accurate neutron capture cross sections** are key **the s- and i-process of stellar nucleosynthesis**, for validating and constraining stellar nucleosynthesis models.
- **Data needs for nucleosynthesis studies:**
 - Stable nuclei: s-only, s-process bottlenecks
 - Unstable nuclei: s-process branching points & i-process
- **Beams & techniques for direct neutron capture measurements**
 - TOF in a white neutron beam: En-dependent cross section in a wide range
 - Activation in quasi-stellar beam: MACS @ 25 keV
- **State-of-the-art**
 - Improving accuracies of key stable isotopes
 - Many first-time (n,g) on unstable s-process branchings in the last decade due to upgrades in facilities and detectors
 - Towards more unstable isotopes (s-process, i-process) → **Current limit for TOF: signal-to-background**
- **Future**
 - Higher sensitivity: High flux facilities for activation (e.g. n_TOF NEAR)
 - Faster/closer sample production: Synergy between neutron beam and RIB facilities (e.g. NEAR & ISOLDE)
 - Future very bright neutron sources: FRANZ, nACT @ BDF, IFMIF-DONES
 - Towards very short-lived (i-process): Direct (n,g) in inverse kinematics: neutron target + exotic ion storage ring

Thank you!

Any questions?

J. Lerendegui-Marco,

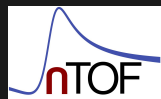
Instituto de Física Corpuscular (Valencia, Spain) / n_TOF Collaboration (CERN)

jorge.lerendegui@ific.uv.es

**s, i & r Element Nucleosynthesis (sirEN)
CONFERENCE**

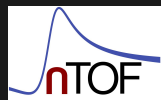
Giulianova (Italy), 8-13 June 2025

Dedicated to the memory of Prof. Roberto Gallino



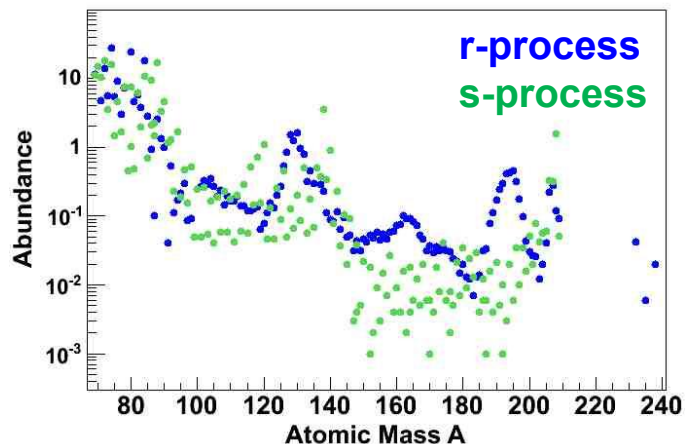
BACK-UP





s- / i-process





high neutron fluxes , short time scales, very n-rich nuclei

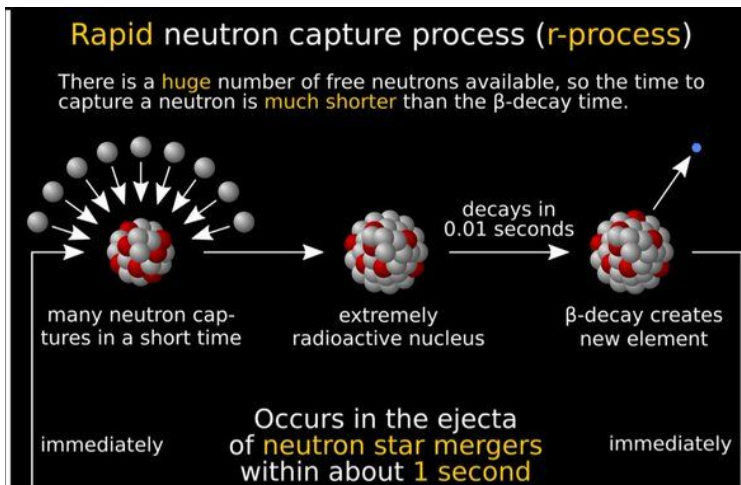


SN, N-Star Mergers, etc ? [r-process path \(gif\)](#)

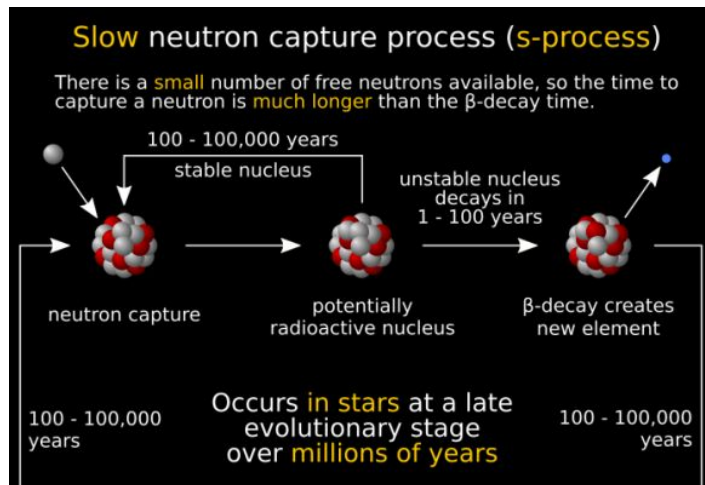
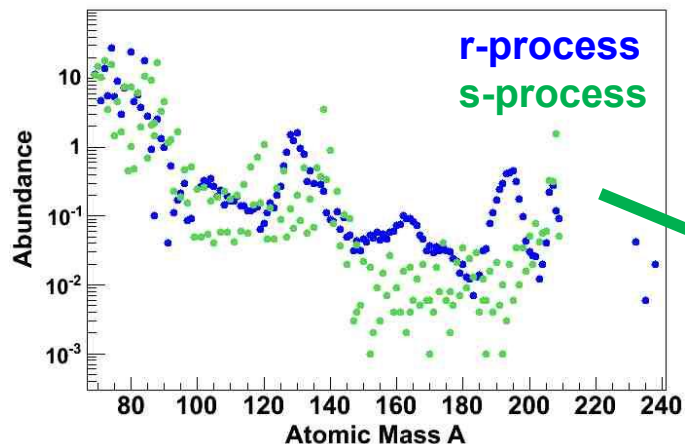
rapid (r)-process

$$T = 10^8 - 10^{10} \text{ K}$$

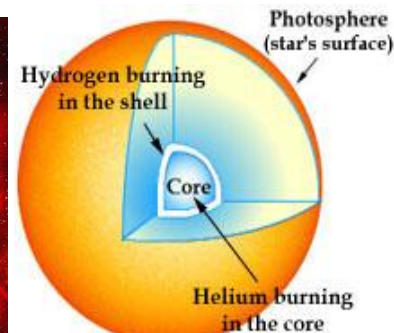
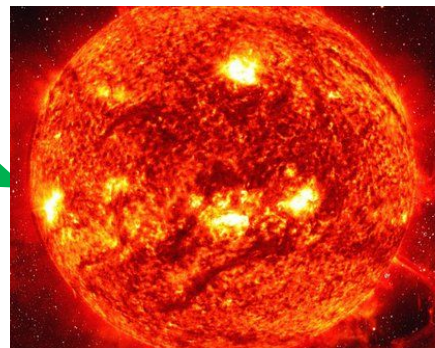
$$N_n = 10^{20} - 10^{27} \text{ cm}^{-3}$$



r-process : extremely short-lived nuclei →
Neutron capture cannot be directly studied.



Low neutron fluxes , very long time scales,
Stable or long-lived unstable nuclei



AGB (red giant) stars, massive stars

slow (s)-process

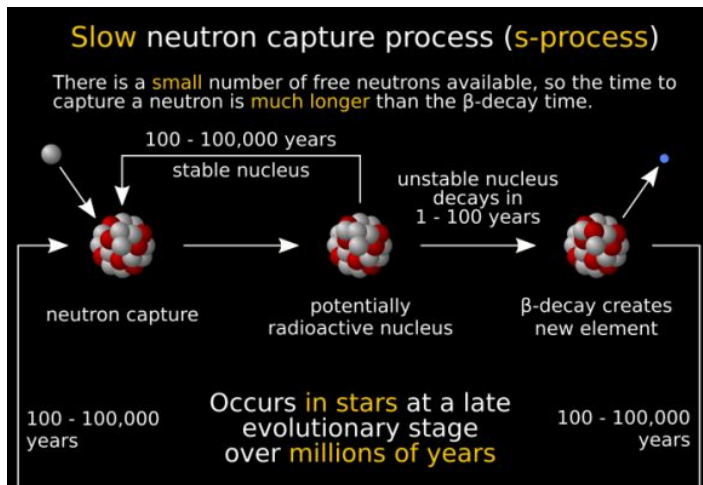
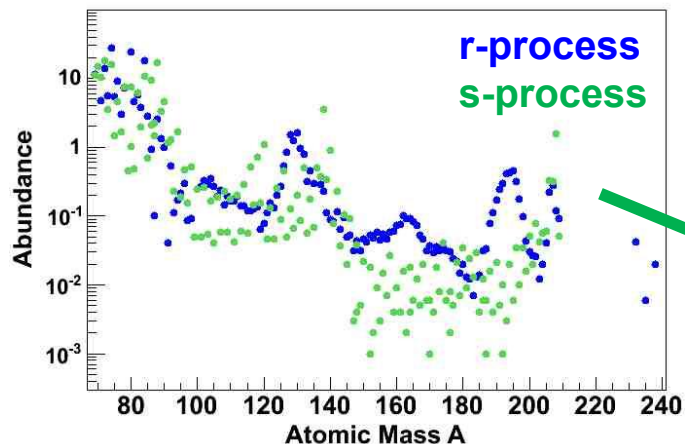
$$T = 10^8 - 10^9 \text{ K}$$

$$N_n = 10^6 - 10^{12} \text{ cm}^{-3}$$

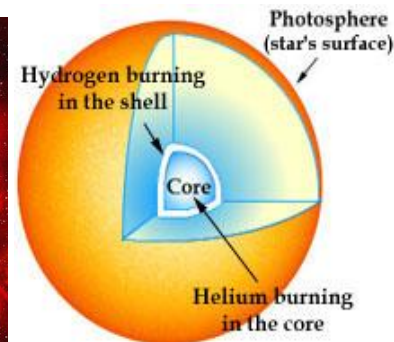
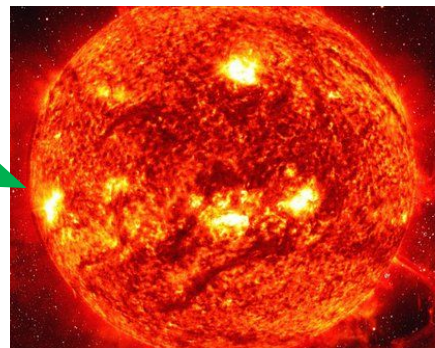
NEUTRON SOURCES (AGB)

$$^{22}\text{Ne}(\alpha, n): 10^{12} \text{ n/cm}^3, kT \sim 30 \text{ keV}$$

$$^{13}\text{C}(\alpha, n): 10^7 \text{ n/cm}^3, kT \sim 5 \text{ keV}$$



Low neutron fluxes , very long time scales,
Stable or long-lived unstable nuclei



AGB (red giant) stars, massive stars

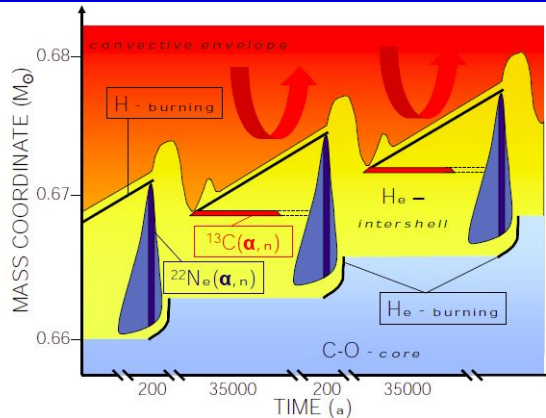
slow (s)-process

$$T = 10^8 - 10^9 \text{ K}$$

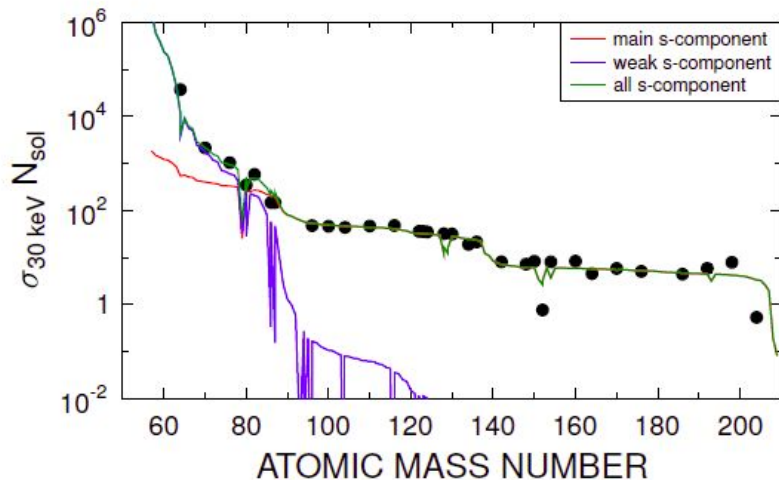
$$N_n = 10^6 - 10^{12} \text{ cm}^{-3}$$

s-process: along stability valley $\rightarrow$

Direct neutron capture reactions in
existing neutron beams



Thermally pulsing
low-mass AGB stars



S-PROCESS COMPONENTS

MAIN S-PROCESS:

He shell burning phase in AGB STARS

$^{22}\text{Ne}(\alpha, n)$: 10^{12} n/cm^3 , $kT \sim 30 \text{ keV}$

$^{13}\text{C}(\alpha, n)$: 10^7 n/cm^3 , $kT \sim 5 \text{ keV}$

$90 < A < 209$

WEAK S-PROCESS:

Massive stars

$^{22}\text{Ne}(\alpha, n)$

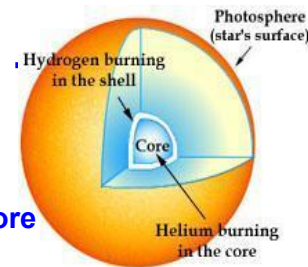
He core burning: 10^6 n/cm^3 , $kT \sim 25 \text{ keV}$

C-shell burning: 10^{12} n/cm^3 , $kT \sim 90 \text{ keV}$

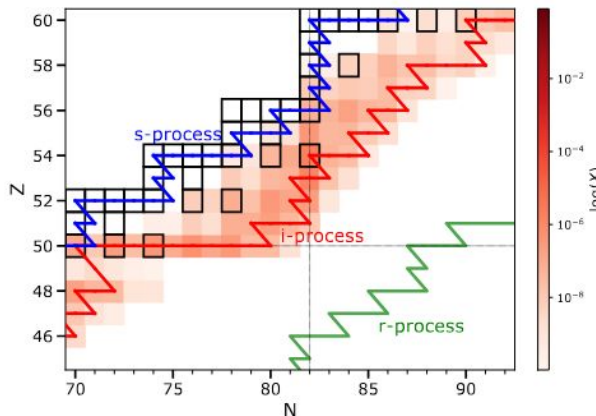
$56 < A < 90$

Energy -
Temperature
 $90 \text{ keV} \sim 1 \text{ GK}$

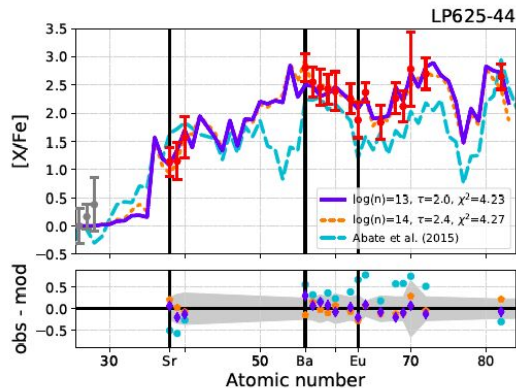
He-burning core
Massive star



i-process



Required to reproduce experimental abundance pattern in CEMP stars



Intermediate neutron fluxes and time scales,
short(er)-lived unstable nuclei (hours-days)

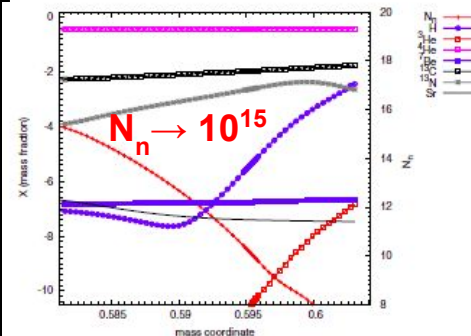
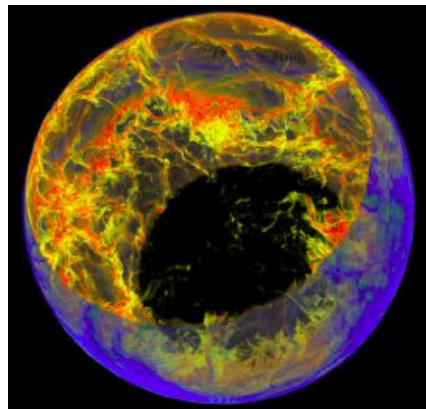


Figure 7. Abundance profiles snapshot (RUN103), just before the mixing split is imposed, demonstrates the simultaneous action of nucleosynthesis and mixing on similar timescales.

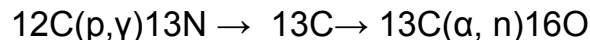
Low metallicity, low mass AGB

Intermediate (i)-process

$T > 10^8 \text{ K}$

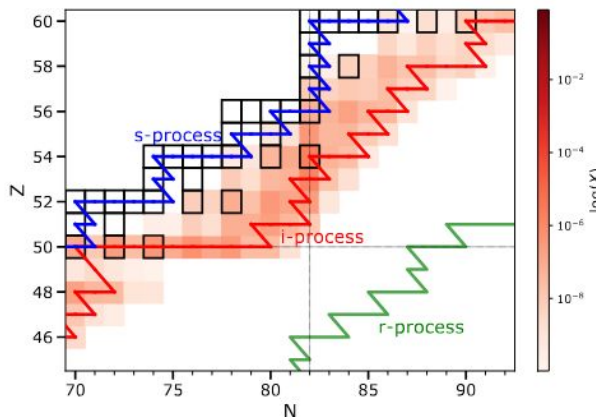
$N_n = 10^{13} - 10^{16} \text{ cm}^{-3}$

NEUTRON SOURCE

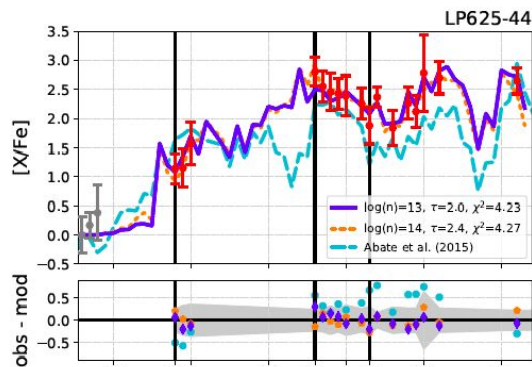


Falk Herwig et al 2011 ApJ 727 89
Choplin, EPJ Web of Conferences 279, 07001 (2023)
Melanie Hampel et al 2019 ApJ 887 11

i-process



Required to reproduce experimental abundance pattern in CEMP stars



Intermediate neutron fluxes and time scales,
short(er)-lived unstable nuclei (hours-days)

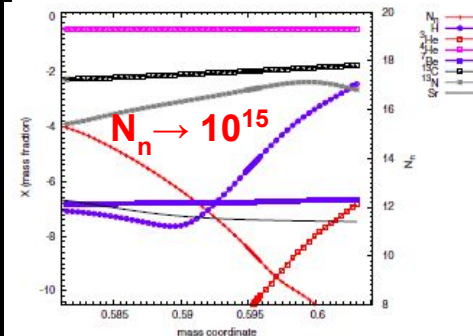
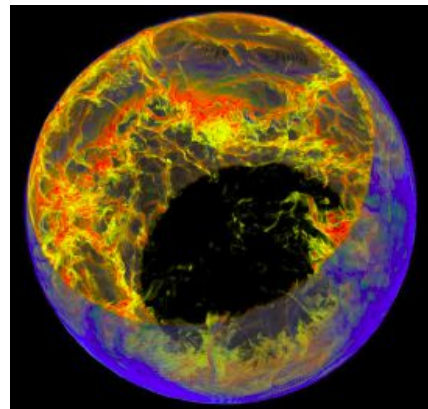


Figure 7. Abundance profiles snapshot (RUN103), just before the mixing split is imposed, demonstrates the simultaneous action of nucleosynthesis and mixing on similar timescales.

Low metallicity, low mass AGB

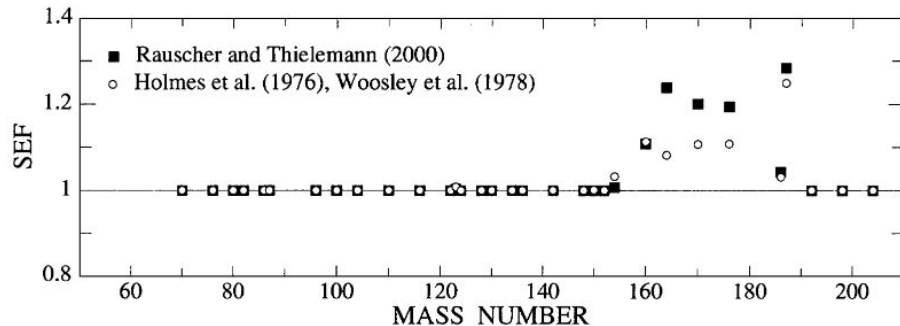
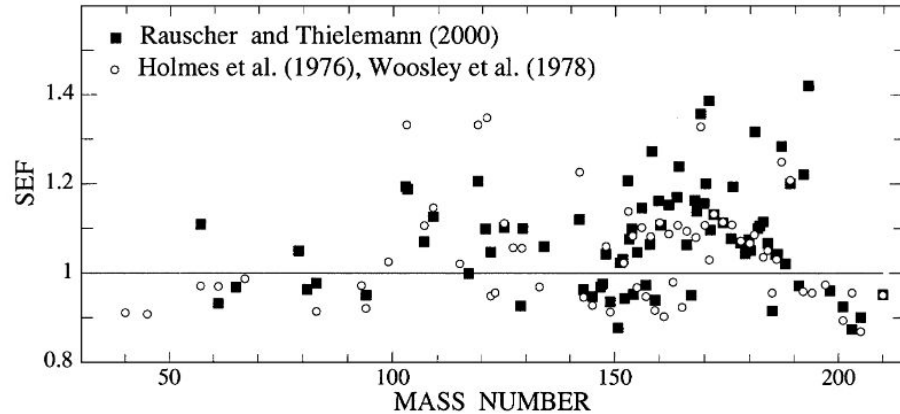
Intermediate (i)-process

$T > 10^8 \text{ K}$

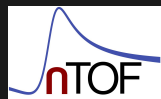
$N_n = 10^{13} - 10^{16} \text{ cm}^{-3}$

i-process: few units away from stability →
Direct neutron capture is challenging → Future
neutron beams & techniques

- 1) Early AGB phase of metal-poor low-mass stars, in which the entropy barrier between the hydrogen- and helium-rich zones can be surmounted by the energy released by a thermal pulse ([Fujimoto et al. 2000](#); [Chieffi et al. 2001](#); [Siess et al. 2002](#); [Iwamoto et al. 2004](#); [Cristallo et al. 2009, 2016](#); [Suda & Fujimoto 2010](#)). The penetration of the helium driven convection zone in the hydrogen-rich layers is favored in metal-poor stars as a result of the absence of heavy elements.
- 2) Core helium flash of very low-metallicity low-mass stars ([Fujimoto et al. 1990, 2000](#); [Schlattl et al. 2001](#); [Suda & Fujimoto 2010](#); [Campbell et al. 2010](#); [Cruz et al. 2013](#)). A
- 3) Very late thermal pulses of post-AGB stars ([Herwig et al. 2011](#)).
- 4) Rapidly accreting carbon-oxygen (C-O) or oxygen-neon (O-Ne) white dwarfs (RAWDs) in close binary systems could also host *i*-process nucleosynthesis ([Denissenkov et al. 2017, 2019](#)). In this scenario, the helium shell formed by the combustion of the accreted hydrogen material becomes thermally unstable. As a consequence of convection, protons from the accreted hydrogen-rich envelope may reach the helium-burning shell and lead to the production of an intense neutron burst.
- 5) Super-AGB stars ($7 M_{\odot} \lesssim M_{\text{ini}} \lesssim 10 M_{\odot}$). [Jones et al. \(2016\)](#) show that at low metallicity ($Z \leq 10^{-3}$), a $7 M_{\odot}$ model can experience a proton ingestion episode during the thermally pulsing super-AGB phase leading to *i*-process nucleosynthesis.
- 6) Helium shell of very low- or zero-metallicity massive stars ($M_{\text{ini}} > 10 M_{\odot}$). These stars can experience a proton ingestion.



- 25% of all isotopes involved in the s process show SEF corrections of more than 2%.
- Significant corrections are to be expected for odd and/or deformed nuclei with excited states well below 100 keV.
- s-only isotopes little affected, except for the mass region $160 < A < 190$

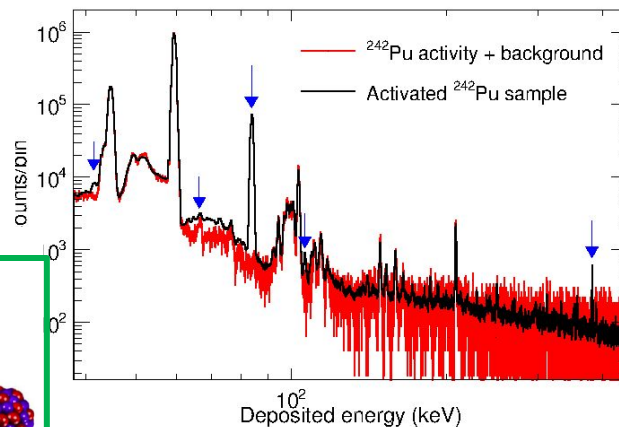


Intro techniques & recent upgrades

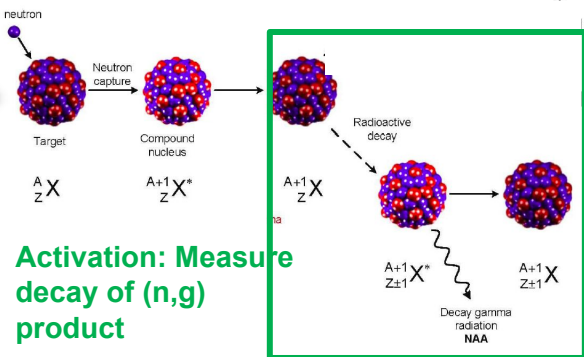




Sample inside the HPGe shielding



g-ray spectrum from the decay



Peak Fit area



$$N_0 = \frac{R_j(t_d = 0)}{F_{geom} F_{abs} \epsilon_j I_j}$$

Number of (n,g) reactions



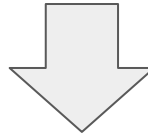
$$\frac{\text{MACS, Pu, kT}}{\text{MACS, Au, kT}} = \frac{F_c^{242\text{Pu}}}{F_c^{197\text{Au}}} \frac{N_0^{243\text{Pu}}}{N_0^{198\text{Au}}} \frac{n_{197\text{Au}}}{n_{242\text{Pu}}}$$

MACS relative to reference (Au-197)

$$Y(E_n) = (1 - e^{-n_C \sigma_{tot}(E_n)}) \frac{\sigma_\gamma(E_n)}{\sigma_{tot}(E_n)} + Y_{MS}(E_n),$$

Theory

Capture yield = Probability for an
incident neutron to undergo a
capture reaction



$$Y(E_n) = F_{norm}^{thr} \cdot \frac{C(E_n) - B(E_n)}{\Phi(E_n) \cdot \varepsilon_c}$$

Experimental

Capture yield = Probability for an incident neutron to undergo a capture reaction

TOF to neutron energy: $\text{TOF} - E_n$
Taking into account RF
(MC simulations)

Dead time:
correction vs. TOF

Background subtraction:
Direct + n/g scattering (MC)

$$Y(E_n) = F_{\text{norm}}^{\text{thr}} \cdot \frac{C(E_n) - B(E_n)}{\Phi(E_n) \varepsilon_c}$$

Efficiency: TED
 $\varepsilon_c \rightarrow E_c = S_n + E_n$

Flux: measured with different setups, unc. within 1-5%

Normalization:
Saturated Resonance
Method (^{197}Au)

$$F_{\text{norm}}^{\text{thr}} = \frac{\langle W^{\text{thr}} \rangle}{f_{\text{SRM}}^{\text{thr}}} \frac{F_{c,\text{Pu}}^{\text{thr}}}{F_{c,\text{Au}}^{\text{thr}}}$$

Efficiency correction factors :
low energy g-rays, EC
(cascade modelling + MC)

	TAC	C ₆ D ₆ TED
Advantages	Good background rejection	Low neutron sensitivity
	High efficiency	Simple set-up
	High resolution	Fast detectors
Drawbacks	Neutron sensitive	Poor background rejection
	Slow detectors: pile-up problems	Low efficiency
	Complex detector (42 crystals)	Software manipulation (PHWT) needed

Not able to measure in the keV range of astrophysical interest (dead time)

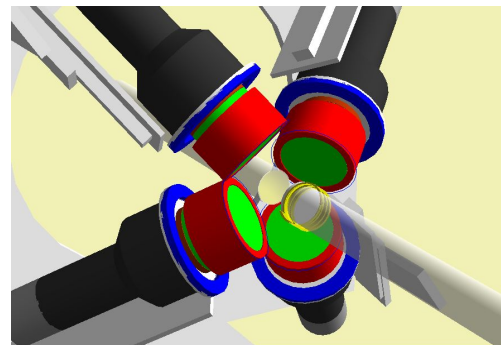
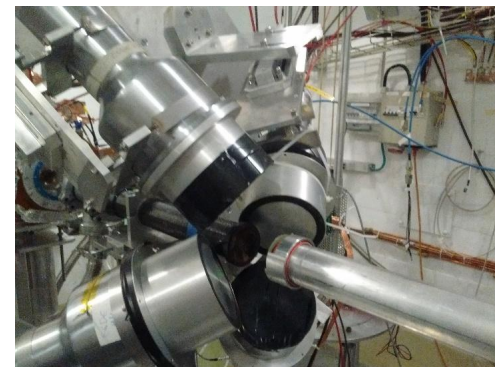
Usual setup for stellar (n,γ) measurements

State-of-the art for measurements in the astrophysical range of interest (1 - 100 keV):

- Fast response (10 ns width signals)
- Low neutron sensitivity

C₆D₆ TED

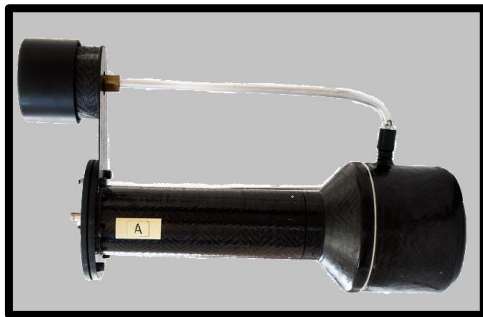
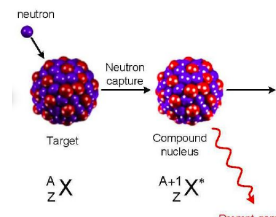
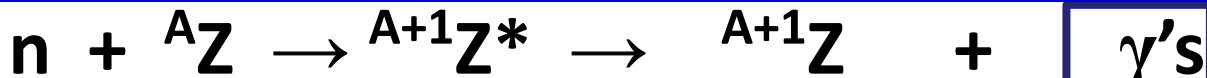
+ PHWT



Limitations:

- Background in the keV range due to dominant neutron scattering
- Limited by maximum count rate in high flux facilities such as n_TOF-EAR2

Solutions:
Recent progresses!



TOTAL ENERGY DETECTORS

Detect at-least/only ONE γ -ray



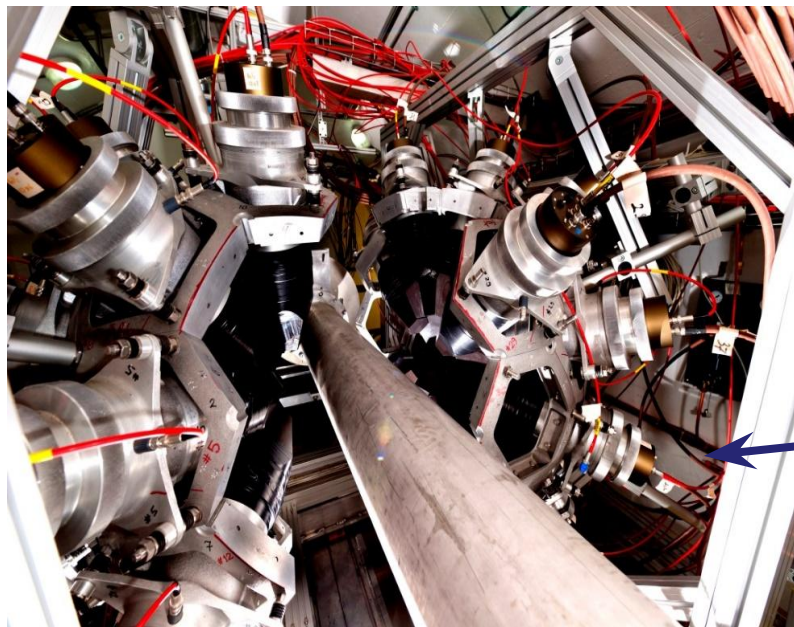
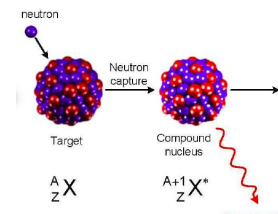
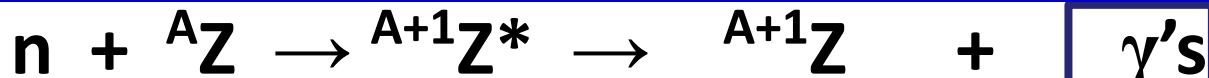
4 x D-benzene detectors (organic scintillator) →
Low neutron sensitivity, fast detectors, simple setup

TOF technique

Detect the γ -ray cascades
vs TOF (neutron energy)

TOTAL ABSORPTION

Detect ALL the γ -rays



The n_TOF Total Absorption Calorimeter (TAC)

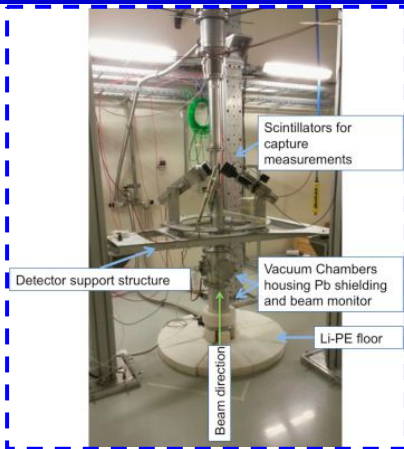
40 BaF_2 scintillator crystals $\rightarrow$ 95% solid angle

At n_TOF

Detect the γ -ray cascades vs TOF (neutron energy)

TOTAL ABSORPTION

Detect ALL the γ -rays

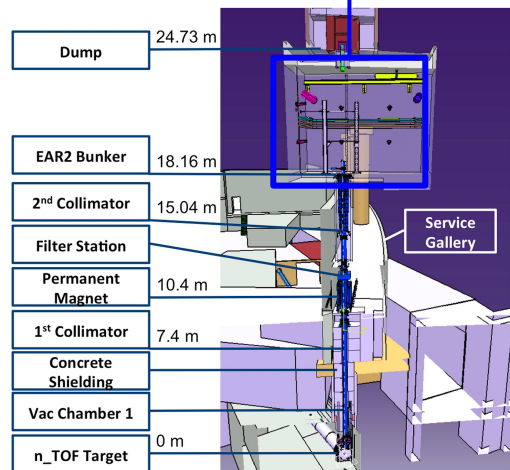
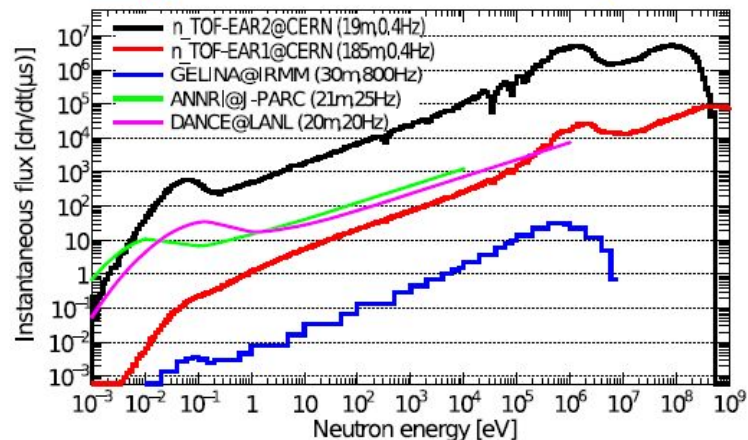


n_TOF EAR2 (20 m, vertical)

- Since **2015**: ~15 (n,γ) measurements

Key features:

- **x300** higher **instantaneous flux**
- Low mass / **radioactive samples**

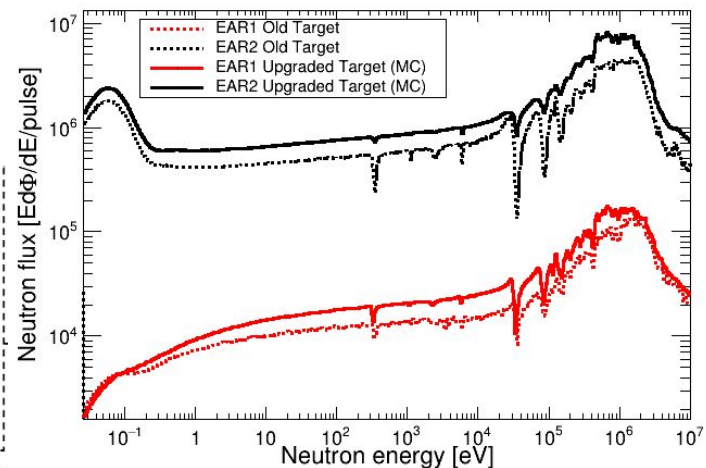
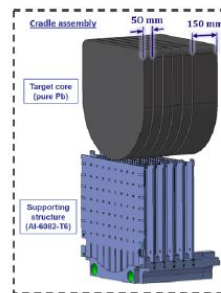


New spallation target

- Replaced in 2021

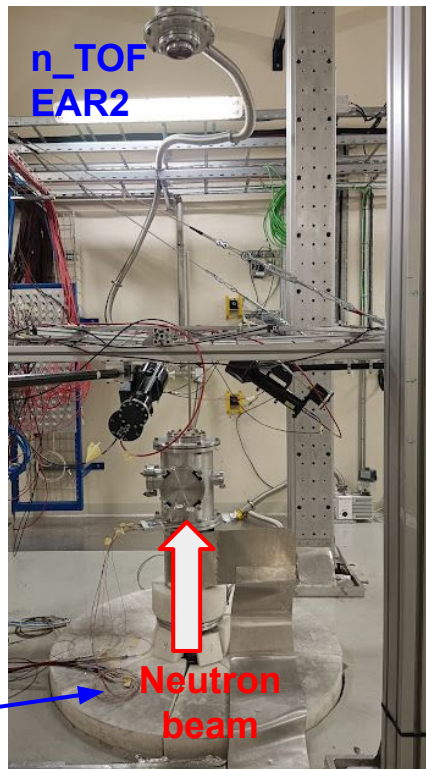
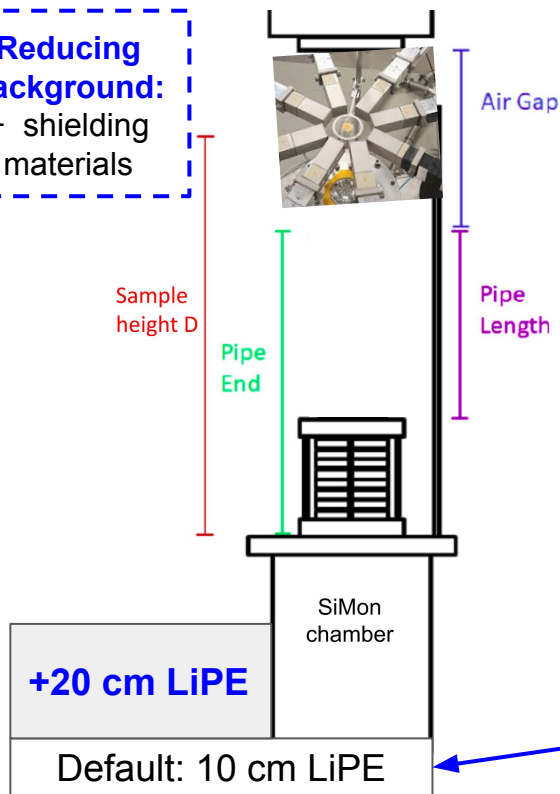
Key features:

- EAR2: 30-50% higher flux & improved energy resolution

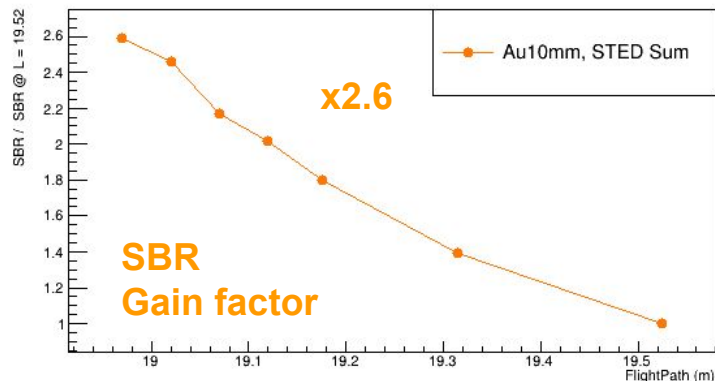


Short term: Optimization campaigns at EAR2 & new ideas to improve the SBR for future experiments

Reducing background:
+ shielding materials

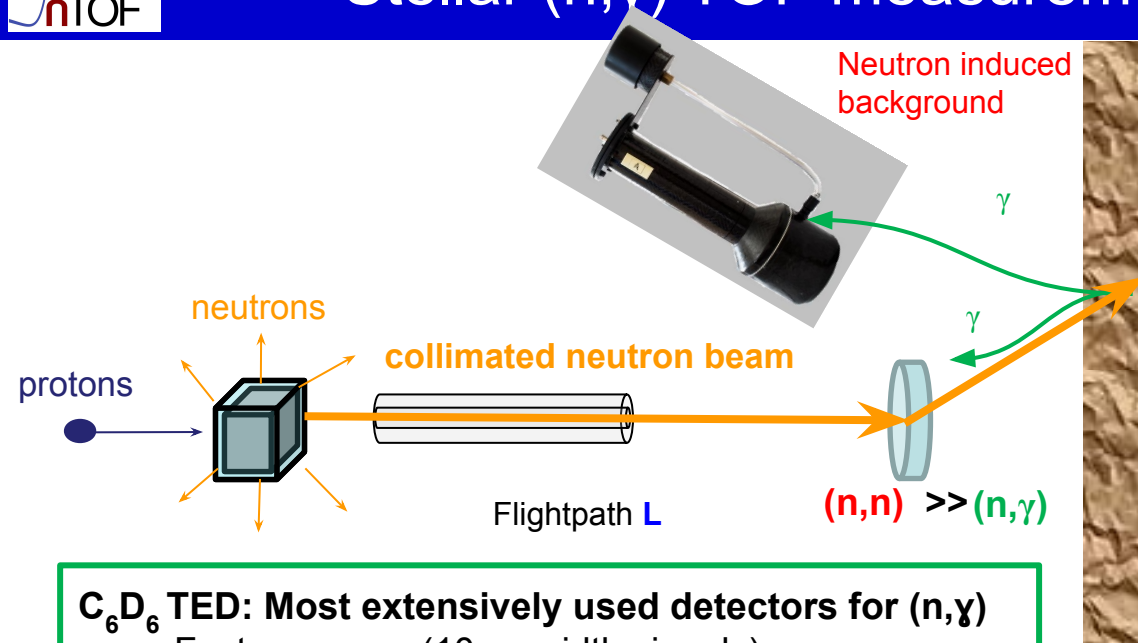


SBR Gain of a factor >2
by lowering 60cm the setup!



Long term:

- Change collimator design
- Upgraded experimental area (n-absorbers, enlarged: walls & floor further from beam)
- Avoid fastest neutrons: "look" only to moderator, new backwards TOF line

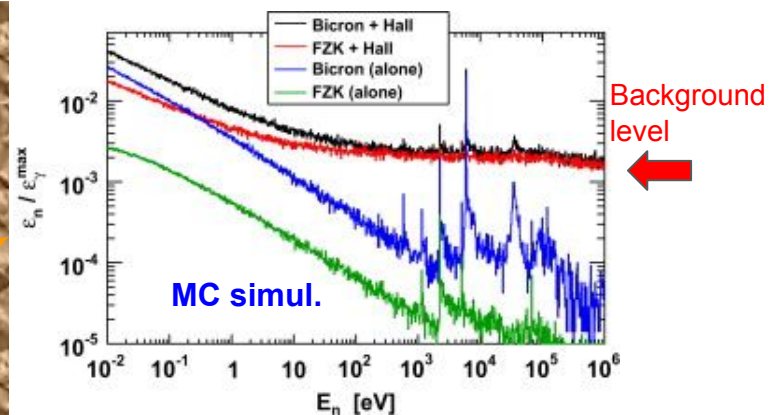


C_6D_6 TED: Most extensively used detectors for (n, γ)

- Fast response (10 ns width signals)
- Low neutron sensitivity

Limitation:

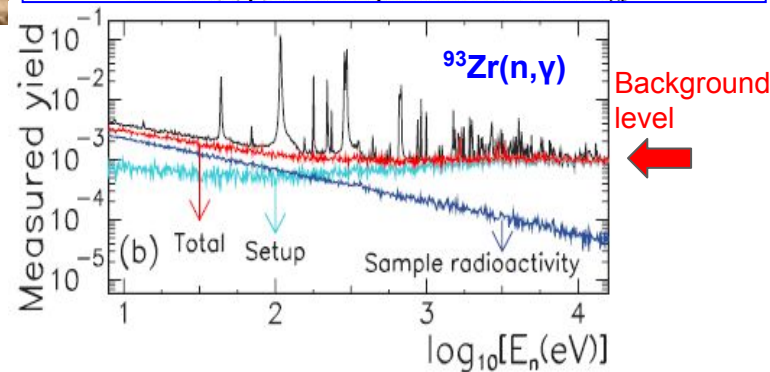
Poor background rejection capabilities.
In particular, background originated from neutron scattered in the sample



P. Zugec, et al., [Nucl. Instrum. Methods A 760, 57 \(2014\)](#);

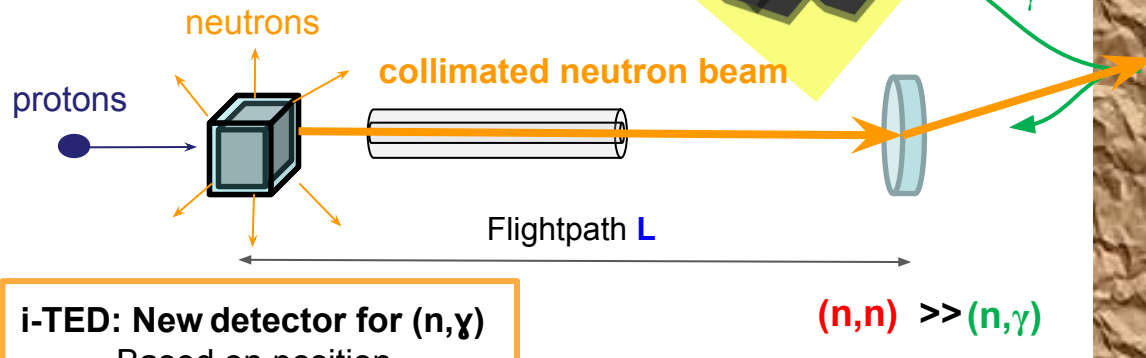
PHYSICAL REVIEW C 87, 014622 (2013)

The $^{93}\text{Zr}(n, \gamma)$ reaction up to 8 keV neutron energy



Contents lists available at ScienceDirect
Nuclear Instruments and Methods in Physics Research A
ELSEVIER
journal homepage: www.elsevier.com/locate/nima

i-TED: A novel concept for high-sensitivity (n,γ) cross-section measurements
C. Domingo-Pardo
Instituto de Física Corpuscular, CSIC-University of Valencia, Spain



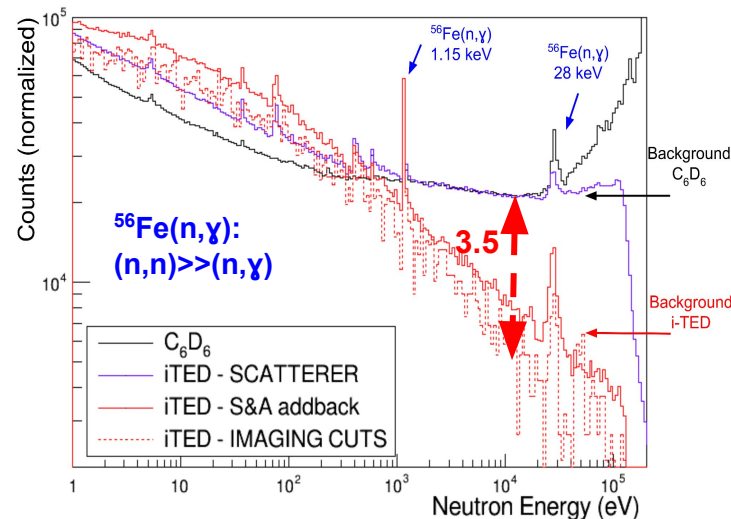
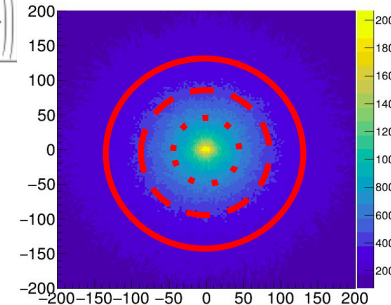
i-TED: New detector for (n,γ)

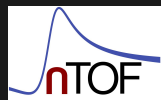
- Based on position sensitive detectors
- Imaging capabilities

Solution: i-TED → Compton Imaging technique to reduce the neutron background and enhance the detection sensitivity

$$\theta = \arccos \left(1 - m_e c^2 \left(\frac{1}{E_2} - \frac{1}{E_1 + E_2} \right) \right)$$

COMPTON IMAGING





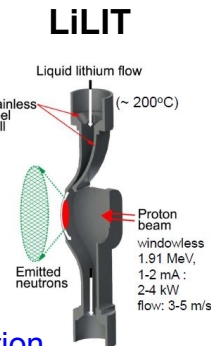
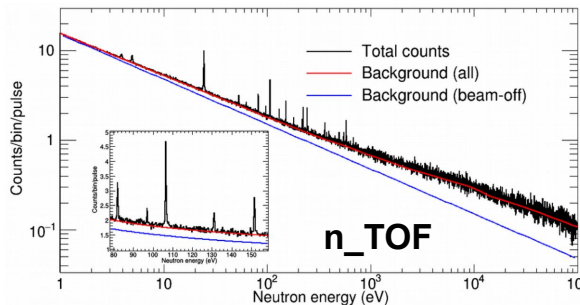
s-process branchings: + examples



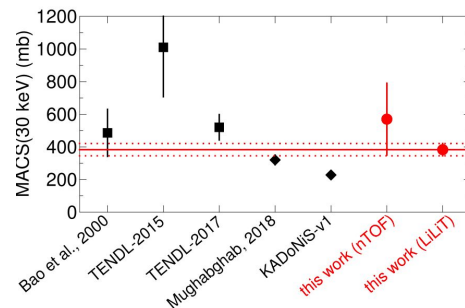
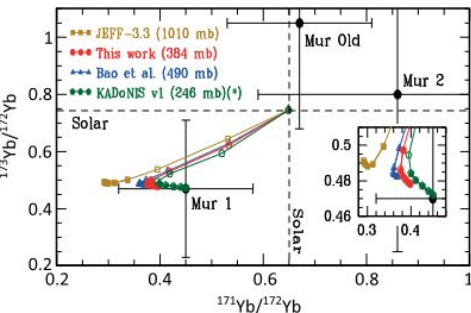
Tm-171(n,γ)

First measurement ever!
TOF @ n_TOF + activation @ LiLIT

m = 3.1 mg
97.9% pure
σ = 384(40) mb



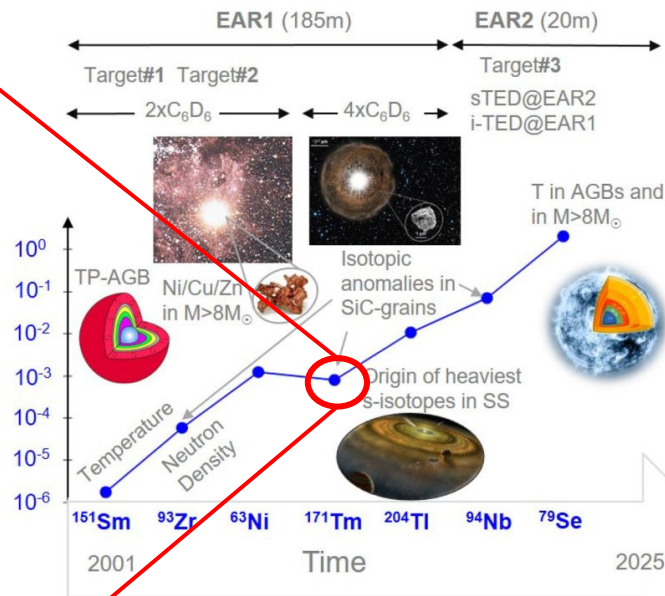
Consistent TOF & activation



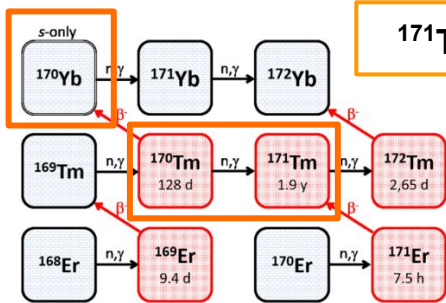
Yb-ratios SiC grains: Improved agreement vs JEFF-3.3

C. Guerrero et al., *Phys. Rev. Letters* **125**, 142701 (2020)

FOM = Difficulty Level
(mg·mb)⁻¹



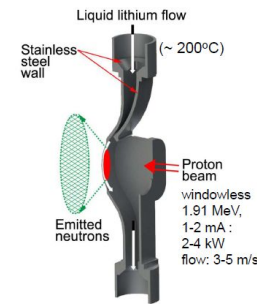
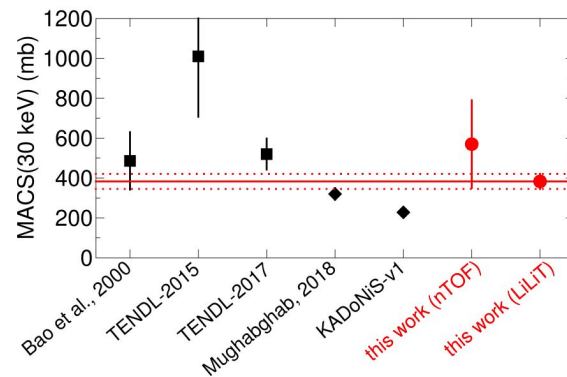
$^{171}\text{Tm}(n,g)$



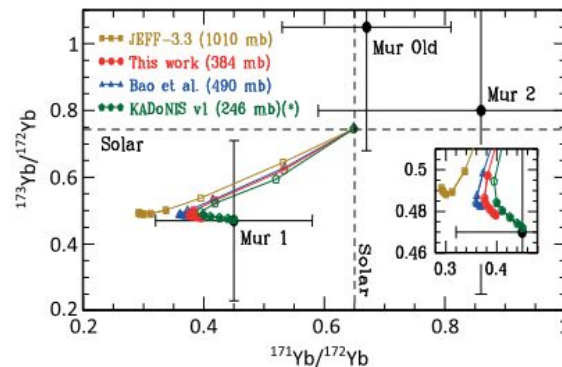
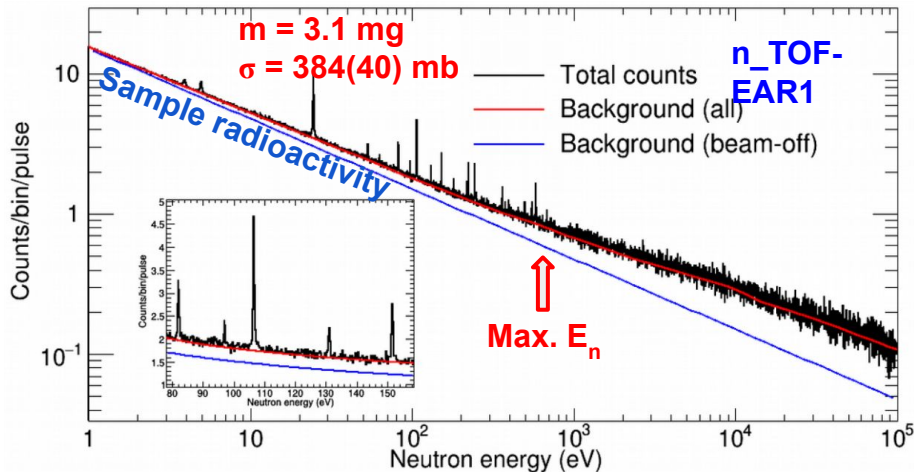
First TOF measurement ever! Combined TOF with activation MACS @ LiLIT

Neutron capture on the s-process branching point ^{171}Tm via time-of-flight and activation

C. Guerrero^{1,2}, J. Lerendegui-Marco¹, M. Paul³, M. Tessler⁴, S. Heinitz⁵, C. Domingo-Pardo⁶, S. Cristallo^{7,8}, R. Dressler⁵, S. Halfon⁴, N. Kivel⁵, U. Köster⁹, E. A. Mauger¹, T. Palchan-Hazan³, J. M. Quesada¹, D. Rochman⁵, D. Schumann⁵, L. Weissman⁴, O. Aberle¹⁰, S. Amaducci¹², J.



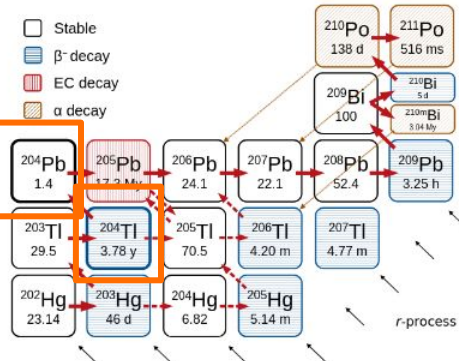
LiLIT



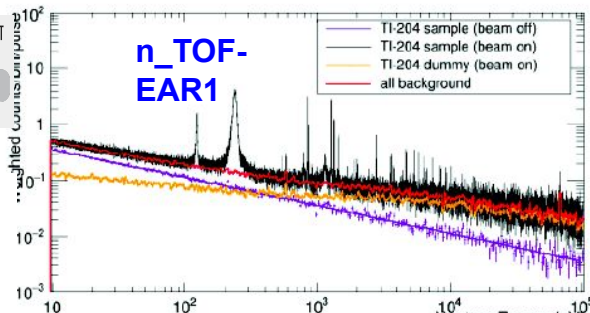
Yb-ratios from SiC grains:
JEFF-3.3 large discrepancies (2Mo, Z=0.01)

Tl-204(n,γ)

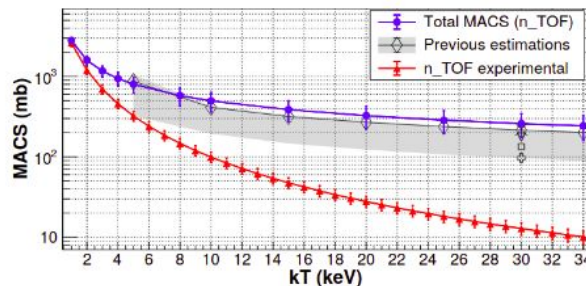
m = 9 mg
(+200 mg ²⁰³Tl)
σ = 260(90) mb



**First measurement of
 this cross section!**
 n_TOF Only



Relevant for the abundance of the
 heaviest s-only ²⁰⁴Pb

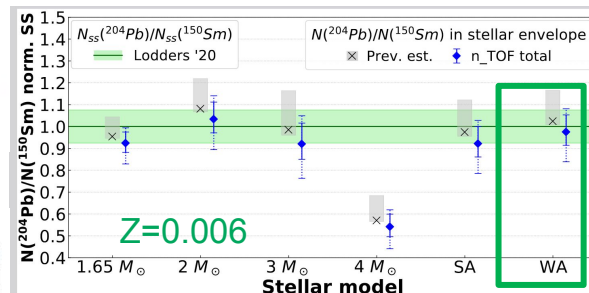


PHYSICAL REVIEW LETTERS 133, 052702 (2024)

Shedding Light on the Origin of ²⁰⁴Pb, the Heaviest s-Process-Only Isotope in the Solar System

A. Casanovas-Hoste^{1,2,3,*}, C. Domingo-Pardo,² J. Lerendegui-Marco,⁴ C. Guerrero,⁴ A. Tarifeño-Saldivia,² M. Krtićka,⁵

The neutron capture yield was measured by detecting the prompt de-excitation γ-rays emitted after each capture event, employing the standard setup at n_TOF of four C₆D₆ liquid scintillation detectors [42], which are optimized to minimize their neutron sensitivity [43]. Lead foils were placed on the detectors to reduce the impact of the γ-ray background arising from the ²⁰⁴Tl decay. By



Reproduces
 experim.
 abund.

Main limitations: Limited energy range due to mass, purity of the sample and background due to the sample activity →

Higher flux facilities and high purity samples are required

^{79}Se sample: a collaborative effort



~15x4mm

CERN: Encapsulated in
e-welded 6N Al+

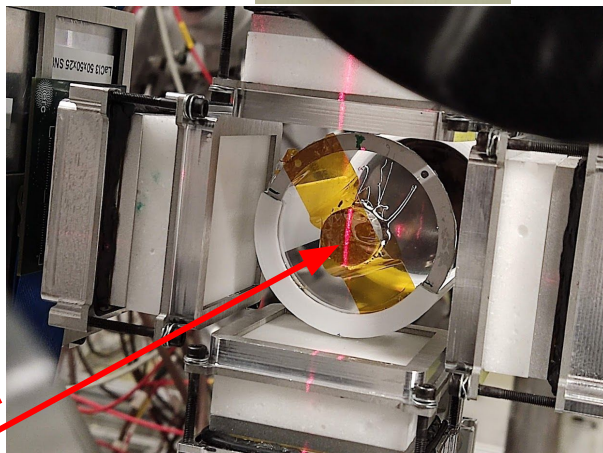
PAUL SCHERRER INSTITUT

PSI

PSI: $^{208}\text{Pb}^{78}\text{Se}$ alloy to
avoid low melting point



~17x6mm



Sample in the beam of n_TOF-EAR1

ILL: ~3 mg of ^{79}Se via $^{78}\text{Se}(n,\gamma)$



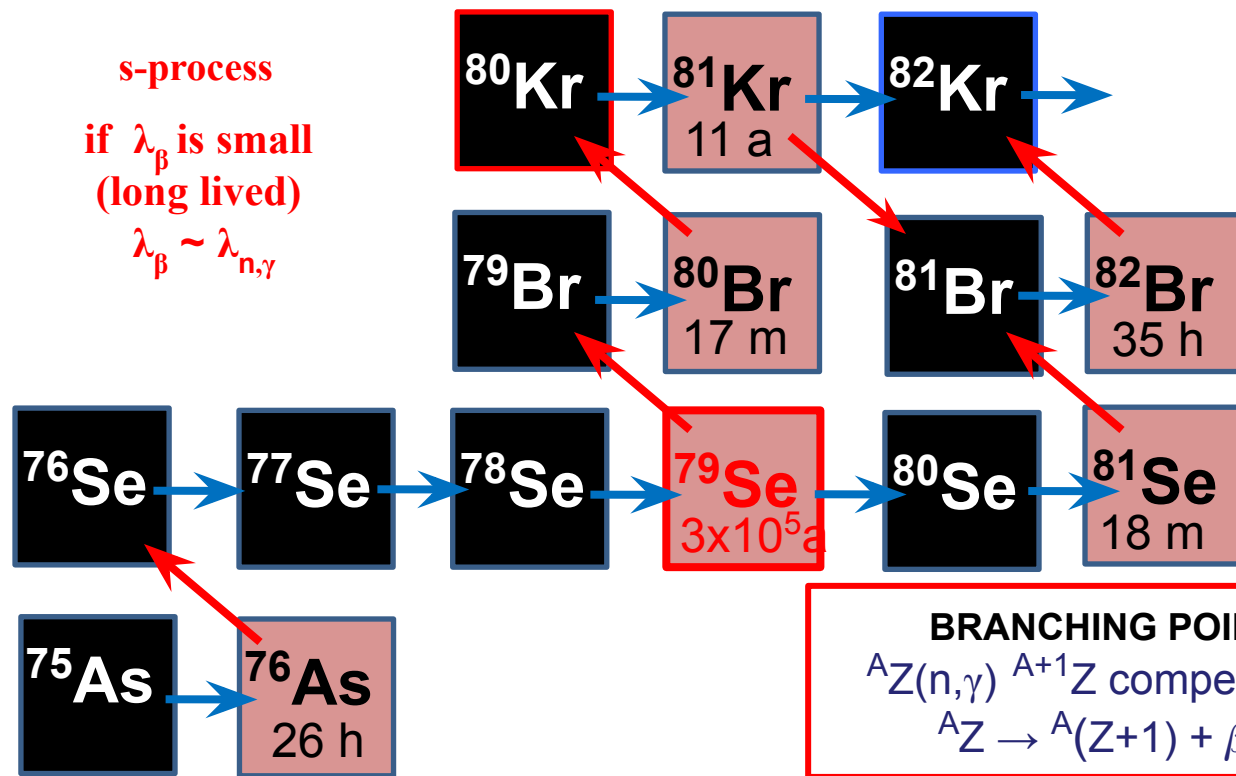
ILL
NEUTRONS
FOR SCIENCE

PSI: Characterized
active contaminants

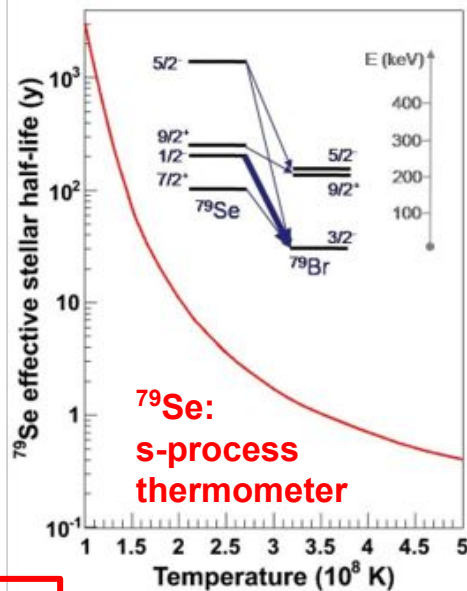
Isotope	Mass (g)
Se-78	1.064
Se-79	0.003
Pb-208	2.838
Al-27	1.0244

gamma
emitters

Isotope	Activity (Mbq)
Se-79	4.33E-01
Se-75	5.66E+00
Ag-110m	2.04E-01
Zn-65	2.33E-01
Co-60	1.40E+00



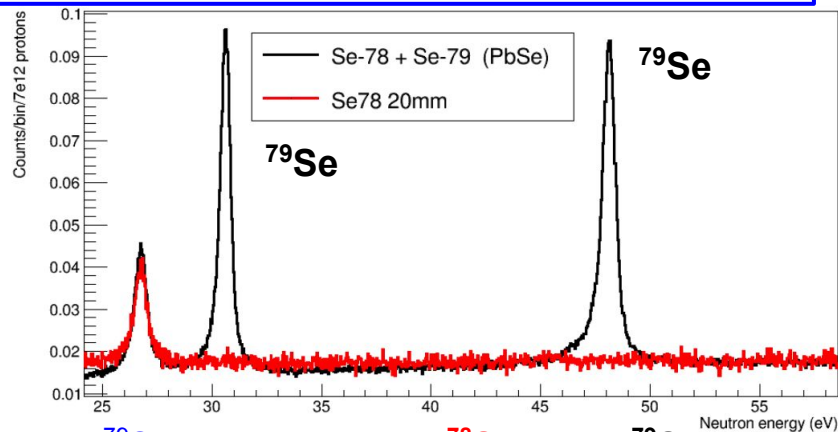
BRANCHING POINTS:
 $^A_Z(n,\gamma) ^{A+1}_Z$ competes with
 $^A_Z \rightarrow ^{A}_{(Z+1)} + \beta + \nu$



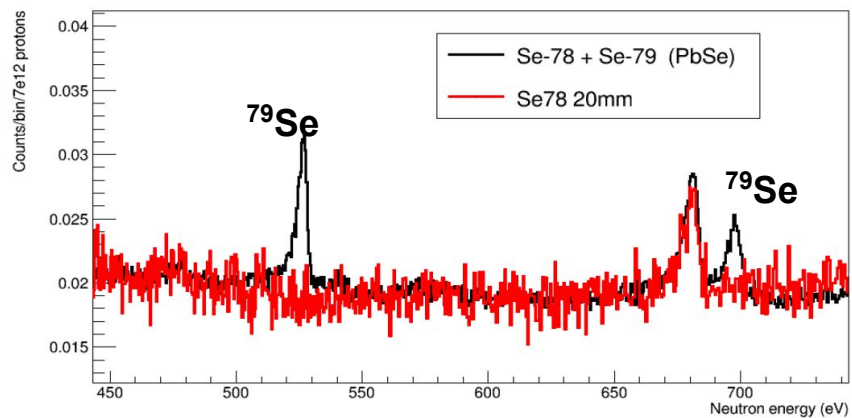
s-process: chain **n-capture** + **beta decays**
Heavier elements along stability valley

(n, γ) measurements + isotopic ratios provide tight constraints to stellar conditions (T, n-dens)

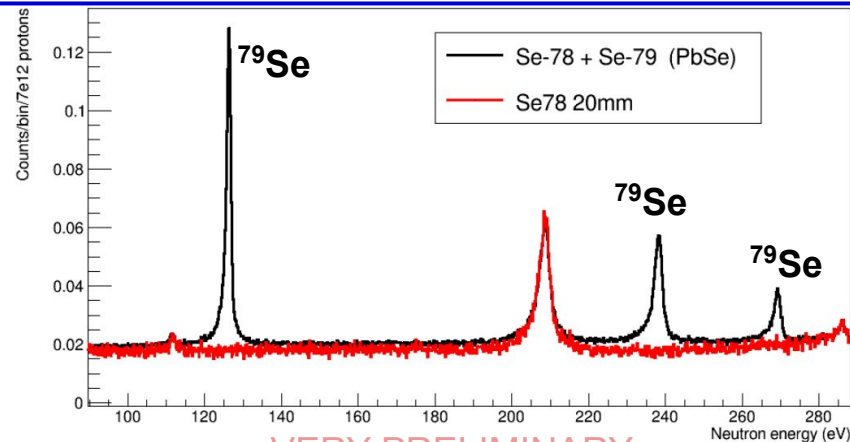
10-15 Resonances of Se-79 **First ever measured!**



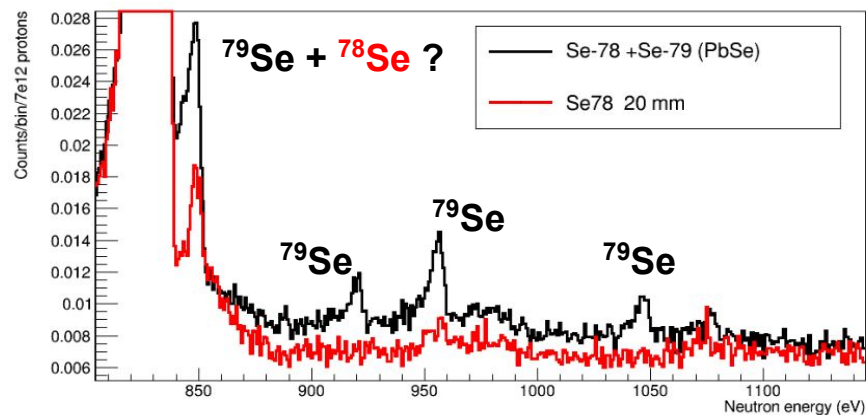
^{79}Se sample: 1060mg ^{78}Se + 3mg ^{79}Se



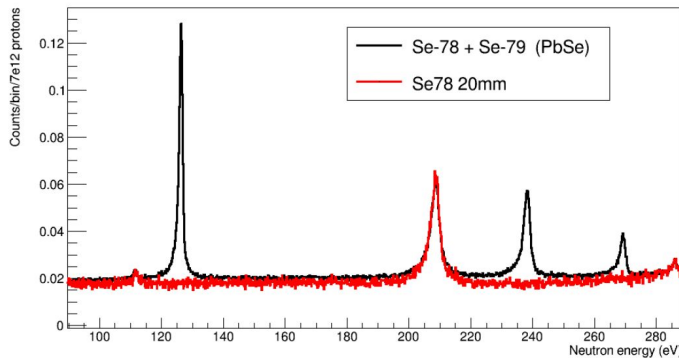
- **Se-79+Se-78 sample** vs **Se-78 sample**



VERY PRELIMINARY



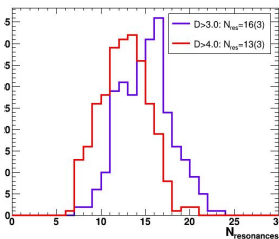
Experimental RRR: (s-wave Avg. Par.)



- MACS @30 keV: 20-26% unc (from avg. par. only)
- Stellar enhancement Factor = 1 at 8, 30 keV
- **Direct stringent constraint for the thermal conditions of AGB and MSs**

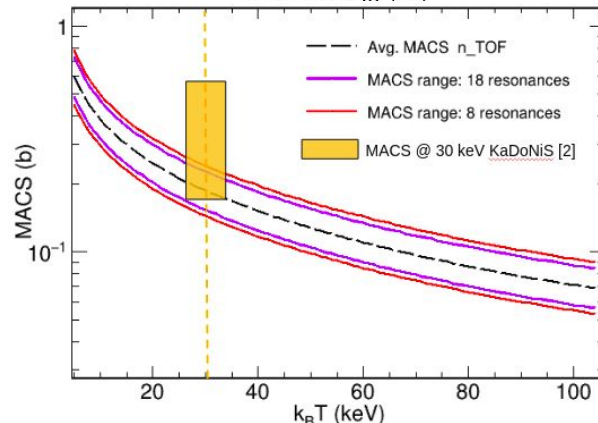
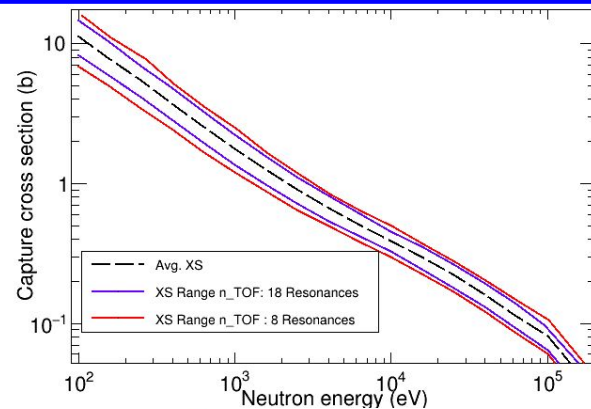
R-Matrix

Measured:
10-15
resonances



Uncertainty
 D_0, S_0
 $\sim \sqrt{N_{\text{res}}}$

Hauser-Feshbach (FITACS): CS in the URR



MACS @ different kT:
Constrain **theor. calculations**
in KaDoNiS (@ 30 keV)

REVIEW OF MODERN PHYSICS, VOLUME 83, JANUARY-MARCH 2011

Sample	Half-life (yr)	Q value (MeV)	Comment
⁶³ Ni	100.1	β^- , 0.066	C. Lederer <i>et al.</i> , Phys. Rev. Lett. 110, 022501 (2013)
⁷⁹ Se	2.95×10^5	β^- , 0.159	Lerendegui-Marco, J. et al., EPJ Web Conf. 279, 13001 (2023)
⁸¹ Kr	2.29×10^5	EC, 0.322	
⁸⁵ Kr	10.73	β^- , 0.687	Important branching, constrains neutron density in massive stars
⁹⁵ Zr	64.02 d	β^- , 1.125	Not feasible in near future, but important for neutron density low-mass AGB stars
¹³⁴ Cs	2.0652	β^- , 2.059	Important branching at $A = 134, 135$, sensitive to s -process temperature in low-mass AGB stars, measurement not feasible in near future
¹³⁵ Cs	2.3×10^6	β^- , 0.269	So far only activation measurement at $kT = 25$ keV by Patronis <i>et al.</i> (2004)
¹⁴⁷ Nd	10.981 d	β^- , 0.896	Important branching at $A = 147/148$, constrains neutron density in low-mass AGB stars
¹⁴⁷ Pm	2.6234	β^- , 0.225	C. Guerrero <i>et al.</i> , Phys. Letters B. 797, 134809 (2019)
¹⁴⁸ Pm	5.368 d	β^- , 2.464	Not feasible in the near future
¹⁵¹ Sm	90	β^- , 0.076	U. Abbondanno <i>et al.</i> , Phys. Rev. Lett. 93, 161103 (2004)
¹⁵⁴ Eu	8.593	β^- , 1.978	Complex branching at $A = 154, 155$, sensitive to temperature and neutron density
¹⁵⁵ Eu	4.753	β^- , 0.246	So far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1995)
¹⁵³ Gd	0.658	EC, 0.244	Part of branching at $A = 154, 155$
¹⁶⁰ Tb	0.198	β^- , 1.833	Weak temperature-sensitive branching, very challenging experiment
¹⁶³ Ho	4570	EC, 0.0026	Branching at $A = 163$ sensitive to mass density during s process, so far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1996b)
¹⁷⁰ Tm	0.352	β^- , 0.968	Important branching, constrains neutron density in low-mass AGB stars
¹⁷¹ Tm	1.921	β^- , 0.098	C. Guerrero, J. Lerendegui-Marco <i>et al.</i> , Phys. Rev. Letters 125, 142701 (2020)
¹⁷⁹ Ta	1.82	EC, 0.115	Crucial for s -process contribution to ¹⁷⁹ Ta, nature's rarest stable isotope
¹⁸⁵ W	0.206	β^- , 0.432	Important branching, sensitive to neutron density and s -process temperature in low-mass AGB stars
²⁰⁴ Tl	3.78	β^- , 0.763	A. Casanovas, C. Domingo <i>et al.</i> , Phys. Rev. Letters (submitted)

F. Käppeler *et al.*,
[Rev. Mod. Phys 83, 157 \(2011\)](#)

Significant progress in the last decade but still many measurements not feasible, in particular (n,g) via TOF
→ **WHY?**

Before 2015: only **2/21** of the key s -process isotopes measured by TOF
2015-2018: ¹⁷¹Tm, ²⁰⁴Tl, ¹⁴⁷Pm at CERN n_TOF and/or LiLIT (activation)
Very recent progresses: ⁷⁹Se, ⁹⁴Nb

Towards more unstable s-process branching point isotopes: **radioactive** → **more challenging**

1. Difficult to produce in sizable quantities → **Low capture/background ratio in TOF measurements**
2. Activity implies a considerable **radiation hazard**.
3. Activity represents an intense source of background for **standard measuring techniques**.

Example: $^{147}\text{Pm}(n,\gamma)$

m = 85 μg
 $\sigma = 826 \text{ mb}$

**n_TOF lowest mass
record at that point!**

**Current Limit
TOF (n,g):
5e17-1e18
atoms**

measurement features the smallest mass (85 μg) ever measured at the n_TOF facility. The mass was unexpectedly small due to the deviations in the assumed value for the thermal capture cross section of ^{146}Nd (seed irradiated at ILL). For this reason n_TOF-EAR2 was chosen over EAR1, featuring the highest instantaneous neutron flux among time-of-flight facilities worldwide [25]. The small mass just allowed to clearly measure the three largest resonances (see left panel of Figure 4) and identify ten

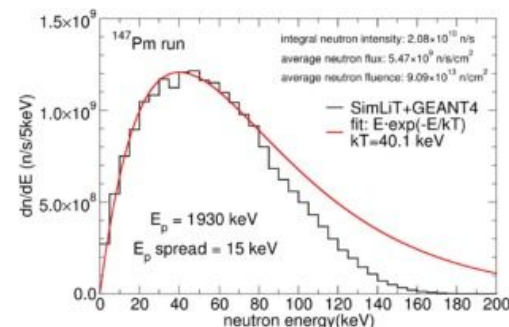
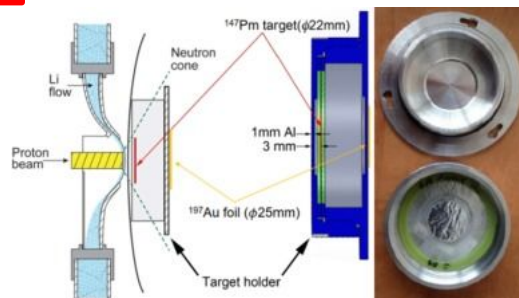
TOF measurement not fully successful due to limited mass

Only MACS via activation with a stellar spectra at LiLiT was feasible

C.Guerrero, M. Tessler, M. Paul, J. Lerendegui.,
[Phys. Letters B, 797, 134809 \(2019\)](#)

MACS @ 30 keV

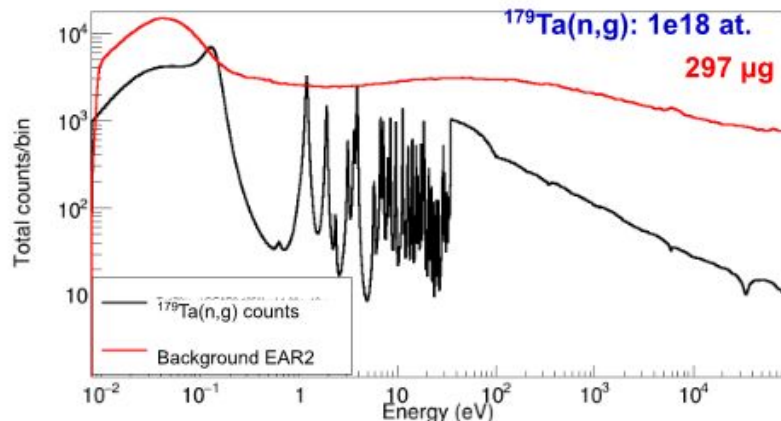
	$\langle \sigma(^{147} \rightarrow ^{148}\text{gPm}) \rangle_{30 \text{ keV}}$	$\langle \sigma(^{147} \rightarrow ^{148}\text{mPm}) \rangle_{30 \text{ keV}}$
This work	469(50) mb	357(27) mb
Reifarth et al.*	332(64) mb	419(58) mb



Towards more unstable s-process branching point isotopes: radioactive → challenges

1. Difficult to produce in sizable quantities → **Low capture/background ratio in TOF measurements**
2. Activity implies a considerable **radiation hazard**.
3. Activity represents an intense source of background for **standard measuring techniques**.

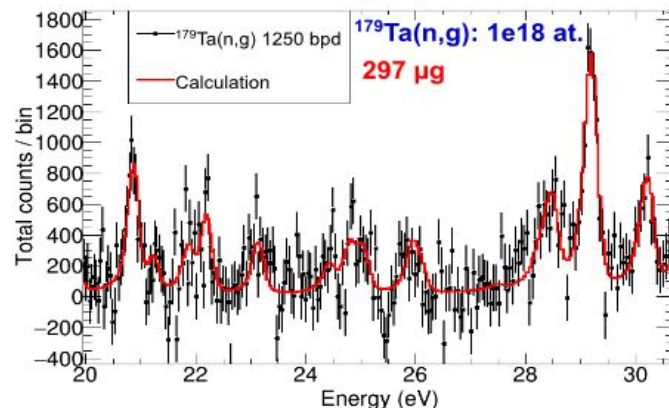
(n,g) vs background



**Current Limit
TOF (n,g):
5e17-1e18
atoms**

J. Lerendegui-Marco et al., EPJ-WOC (2023)
[arXiv:2310.15714](https://arxiv.org/abs/2310.15714)

Expected (n,g) yield with ~45 days of measurement

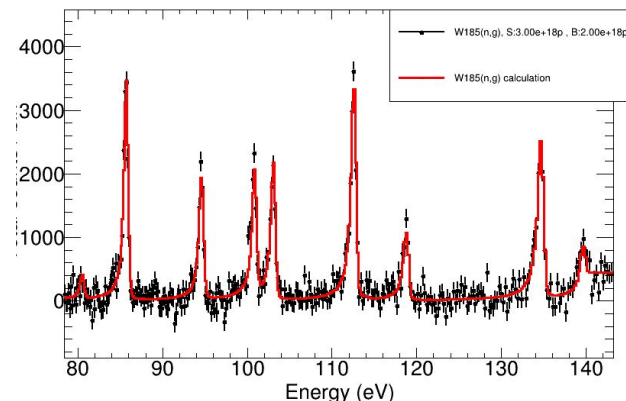
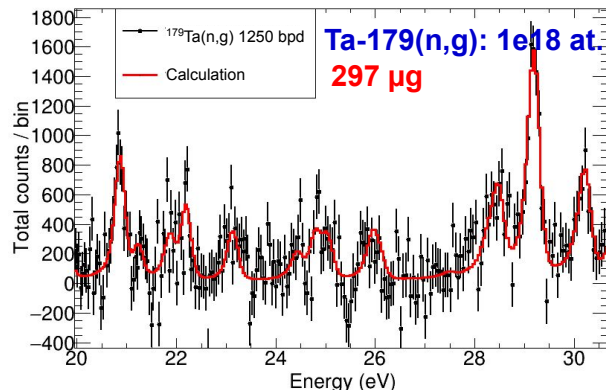
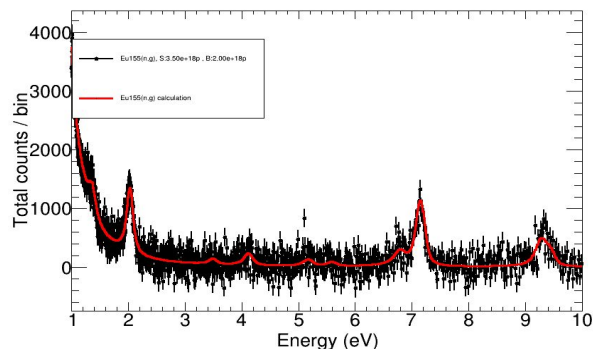
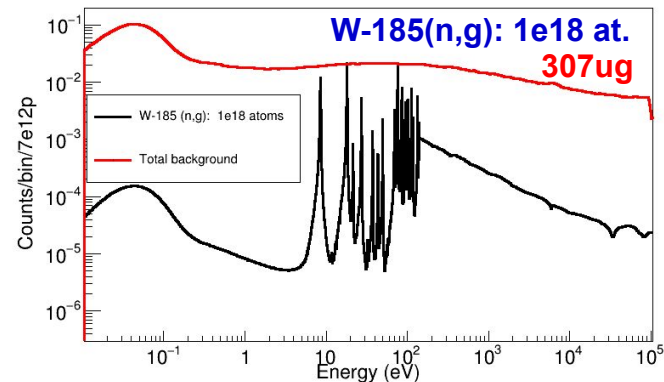
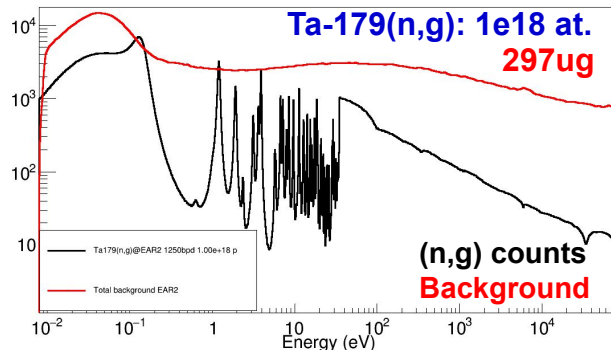
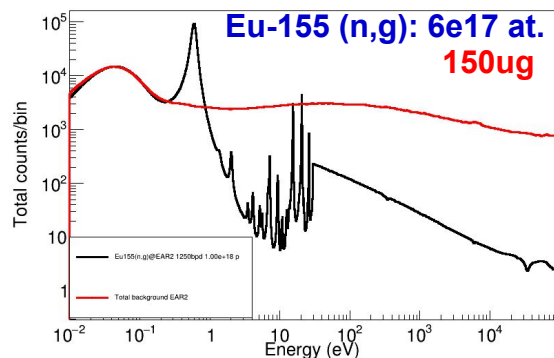


“First-time” key s-process branching that could be feasible:

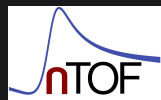
- Eu-155, Ta-179, W-185 (low g-ray background)
- **Mass required for TOF → Activity 5 - 100 GBq!**

Isotope	Mass(ug)	Natoms	Activity (Gbq)
Eu-155	257.390	1.00E+18	4.70E+00
Ta-179	297.243	1.00E+18	1.21E+01
W-185	307.207	1.00E+18	1.07E+02

Estimated results in the upgraded **n_TOF-EAR2** with the **best sensitivity achieved so far with the STEDs**



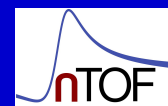
Limit: ~1e18 atoms (too high) currently required → We need complementary techniques → ACTIVATION!



s-process branchings: more details



Temperature and $T_{1/2}$ of ^{79}Se



PHYSICAL REVIEW C

VOLUME 38, NUMBER 1

JULY 1988

β -decay rate of ^{79m}Se and its consequences for the s -process temperature

N. Klay and F. Käppeler

Institut für Kernphysik III, Kernforschungszentrum Karlsruhe, D-7500 Karlsruhe, Federal Republic of Germany

(Received 21 December 1987)

The branching ratio between internal electromagnetic transitions and β^- decays of the isomer ^{79m}Se was determined experimentally. Extremely clean samples of ^{78}Se were activated with thermal neutrons at a high-flux reactor. A mini-orange-Si (Li) detection system was used to measure β^- particles and conversion electrons immediately after neutron irradiation. For the β^- decay we obtain $\log ft = 4.70^{+0.10}_{-0.09}$. Our present result was used to recalculate the temperature dependence of the effective β^- half-life of ^{79}Se in the stellar interior. In combination with the half-life deduced from a quantitative branching analysis, we obtain a possible temperature range between 182 and 295 million degrees for the weak component of the s process.

Nevertheless, the effective stellar half-life at temperatures above 200 million degrees is not changed by the assumed lower limit for the ground state half-life. This leads to the rather strange consequence that, at present, we know the stellar life time of ^{79}Se at a fixed temperature much better than its terrestrial one.

weak s -process component to range between 182 and 295 million degrees. **Need of improved cross sections**

The related uncertainty is mainly due to the uncertainty for the effective stellar half-life deduced from the σ/N systematics which can only be reduced by improved cross section measurements. The present experimental uncer-

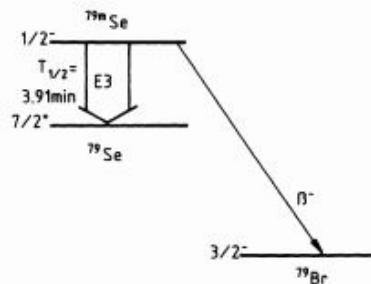
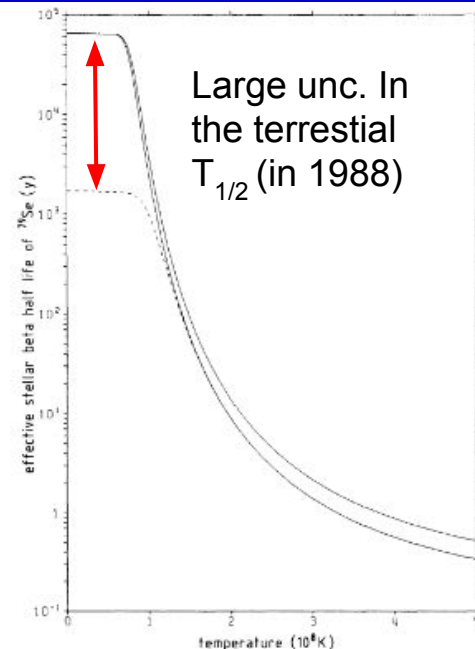
FIG. 1. Decay scheme of ^{79m}Se .

FIG. 7. Stellar half-life of ^{79}Se as a function of temperature. The error band reflects the uncertainty of our measurement. The dashed line refers to $\log ft = 8.5$ for the ground state decay.

Experimental measurement of the $\log ft$ of the beta decay of the ^{79m}Se state at 96 keV, thermally populated at stellar temperatures

Published: 22 November 1990

Meteoritic silicon carbide: pristine material from carbon stars

Roy S. Lewis, Sachiko Amari & Edward Anders

AGB Stars

All five noble gases in interstellar silicon carbide grains have grossly non-solar isotopic and elemental abundances that vary with grain size but are strikingly similar to calculated values for the helium-burning shell of low-mass carbon stars. Apparently these grains formed in carbon-star envelopes, and were impregnated with noble gas ions from a stellar wind. Meteoritic SiC provides a detailed record of nuclear and chemical processes in carbon stars.

The Kr isotopic ratios have been measured in bulk SiC acid residues providing details on AGB stars evolved prior to the formation of the Solar System.

Presolar grain measurements give the **most precise data** currently available on s-process nucleosynthesis (at least one order of magnitude better than spectroscopic observations)



Geochimica et Cosmochimica Acta

Volume 58, Issue 1, January 1994, Pages 471-494



Interstellar grains in meteorites: II. SiC and its noble gases

Roy S. Lewis, Sachiko Amari^{*}, Edward Anders[†]

interstellar SiC, isolated from the Murchison C2 chondrite. All are mixtures of a highly anomalous component bearing the isotopic signature of the astrophysical s-process and a more normal component, generally solar-like but with anomalies of up to 30% in the heavy isotopes. As these two components strikingly resemble predictions for the He-burning shells and envelopes of red giant carbon stars, it appears that the SiC grains are pristine circumstellar condensates from such stars. A number of elemental and isotopic ratios (such as $\text{K}^{80}\text{K}^{82}$ and $\text{K}^{86}\text{Kr}^{82}$) vary with grain size, suggesting that the SiC comes from carbon stars representing a range of masses, metallicities, temperatures, and neutron densities.

Published: 22 November 1990

Meteoritic silicon carbide: pristine material from carbon stars

Roy S. Lewis, Sachiko Amari & Edward Anders

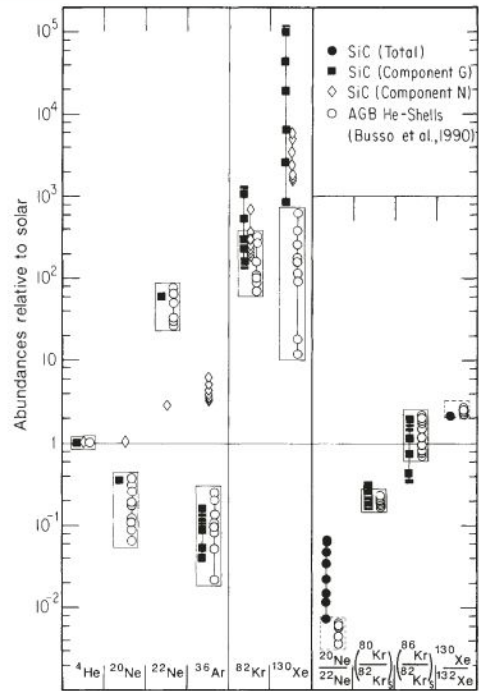


FIG. 3 Right, isotope ratios of SiC and AGB-star He shells, normalized to

AGB Stars

release temperature of the gas on heating or combustion^{7,8}. The variations of ⁸⁰Kr and ⁸⁶Kr reflect branching of the s-process at their radioactive progenitors, ⁷⁹Se and ⁸⁵Kr (ref. 7). These branchings depend sensitively on neutron density and temperature in the s-process region^{7,12}, and thus can provide clues about the stars in which the Kr was formed. What we would

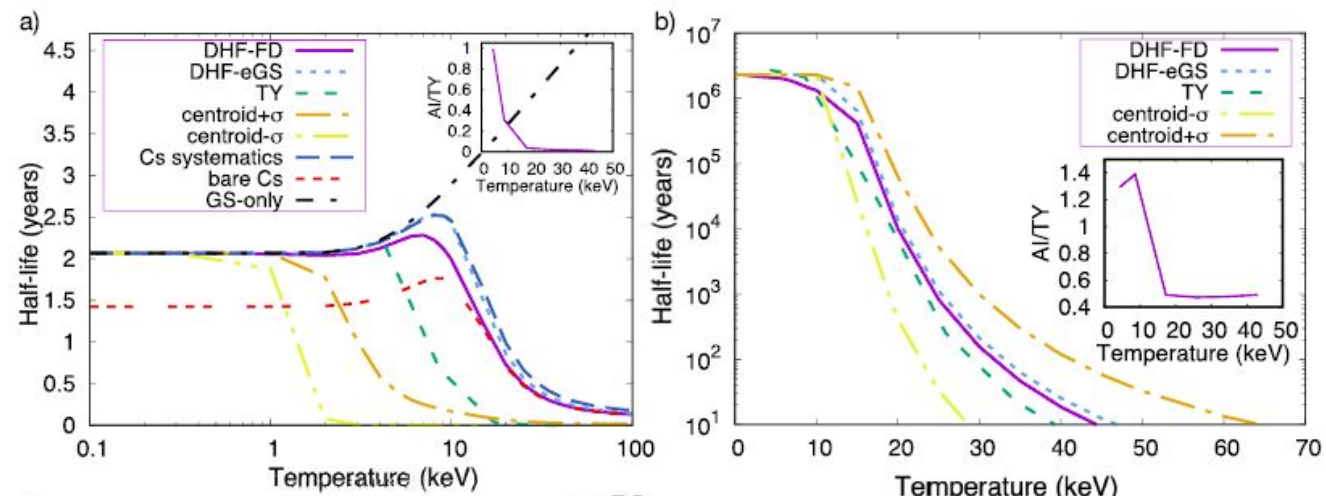
TABLE 1 Noble gases in SiC size fractions from Murchison chondrite

Sample	Size (μm)	³ He ⁴ He	²⁰ Ne	³⁸ Ar	⁸⁴ Kr	¹³⁰ Xe	^(80Kr) _s	^(86Kr) _s
		×10 ⁻⁴	²² Ne	³⁶ Ar	⁸² Kr	¹³² Xe		
KJA	0.05-0.1	1.32 (18)	0.9075 (10)	0.1927 (5)	3.702 (24)	0.3532 (10)	0.0597 (67)	0.621 (44)
KJB	0.1-0.2	2.00 (10)	0.8933 (3)	0.1919 (1)	3.603 (6)	0.3602 (4)	0.0586 (16)	0.695 (24)
KJC	0.2-0.3	1.48 (7)	0.6345 (3)	0.1907 (2)	3.693 (10)	0.3604 (6)	0.0526 (27)	1.145 (19)
KJD	0.3-0.5	1.16 (6)	0.4582 (5)	0.1935 (1)	3.769 (8)	0.3558 (5)	0.0490 (25)	1.733 (23)
KJE	0.5-0.8	0.81 (4)	0.2977 (5)	0.2051 (2)	3.756 (11)	0.3465 (10)	0.0451 (21)	2.363 (34)
KJF	0.8-1.5	0.56 (3)	0.2001 (2)	0.2210 (5)	3.776 (21)	0.3291 (22)	0.0365 (37)	2.710 (71)
KJG	1.5-3	0.53 (5)	0.1556 (2)	0.2313 (6)	3.915 (21)	0.2716 (28)	0.0336 (72)	2.872 (62)
KJH	3-5		0.0979 (19)	0.2090 (23)	4.500 (87)	0.1606 (52)	0.055 (58)	
Solar ²¹		1.42	13.7	0.188	4.988	0.1643		
He-Shell, Range*			0.05-0.084	0.57-1.25	2.2-2.6	0.39-0.44		
He-Shell, Typical		0	0.0808	0.74	2.55	0.485†		

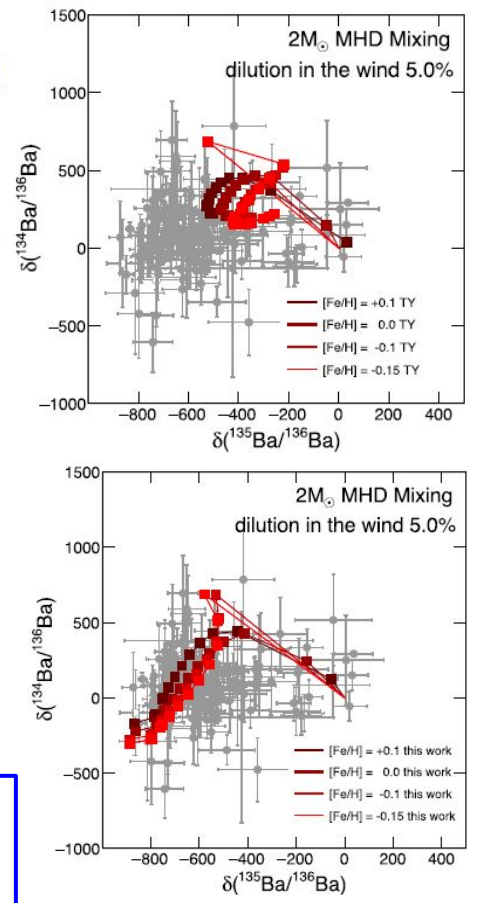
Very accurate 80Kr/82Kr abundances ratios

Theoretical Estimate of the Half-life for the Radioactive ^{134}Cs and ^{135}Cs in Astrophysical Scenarios

Simone Taioli^{1,2} ,
Diego Vescovi³ ,
Maurizio Busso^{4,5} ,
Sara Palmerini^{4,5,10} ,
Sergio Cristallo^{5,6} ,
Alberto Mengoni^{7,8} ,
and Stefano Simonucci^{5,9} 



Agreement with the data improves with
New beta decay rates of both $^{134,135}\text{Cs}$



Presolar Grain Isotopic Ratios as Constraints to Nuclear and Stellar Parameters of Asymptotic Giant Branch Star Nucleosynthesis

Sara Palmerini^{1,2}

Maurizio Busso^{1,2}

Diego Vescovi^{2,3}

Eugenia Naselli⁴

Angelo Pidatella⁴

Riccardo Mucciola^{1,2}

Sergio Cristallo^{2,5}

David Mescalì⁴

Alberto Mengoni^{6,7}

Stefano Simonucci^{2,8}

and Simone Taioli^{9,10,11}

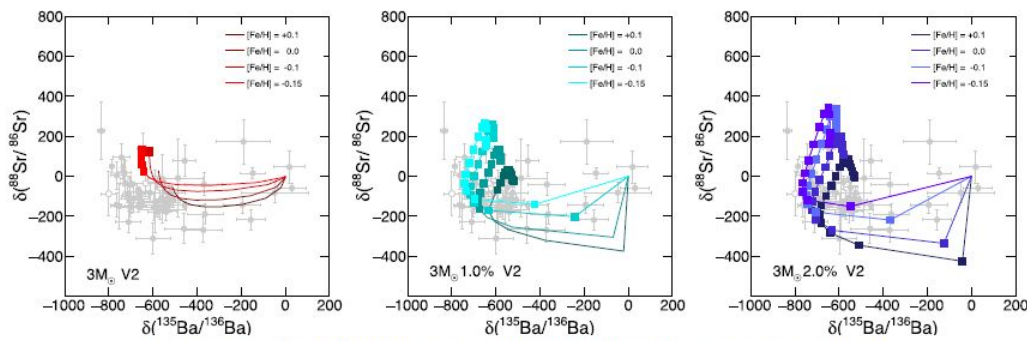


Figure 13. Same as Figure 12, but for the test models V2, with tentatively modified nuclear inputs (see Section 3).

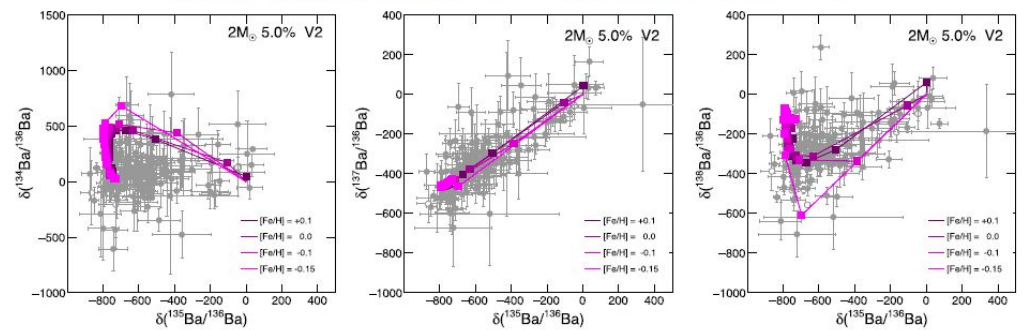


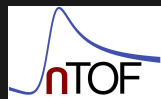
Figure 19. A comparison of model predictions from a representative case of $2 M_{\odot}$ models (full lines with heavy dots) with SiC data for various Ba isotopes and with the choice V2 for nuclear parameters, including revisions for the ^{134}Cs decay, as illustrated in Section 4. The meaning of the symbols is the same as in previous figures. The three panels represent the composition of winds with 5% of He-shell material added.

Sr/Ba ratios:

A longer half-live of ^{135}Cs at 30 keV or variations in the $^{135}\text{Cs}(n,g)$ cross section will improve the agreement with the data

Ba-ratios separately:

Revision of the capture cross sections of Cs-isotopes would improve the agreement

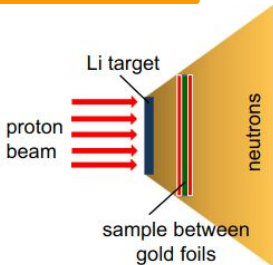


On-going & future developments

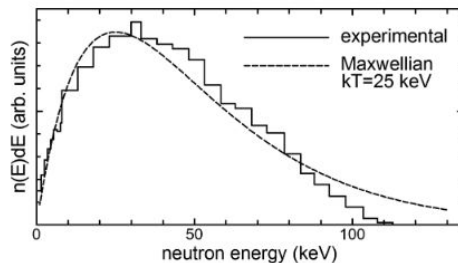


Low E (MeV) proton accelerators

${}^7\text{Li} (p,n) {}^7\text{Be}$



Quasi-maxwellian @ 25 keV

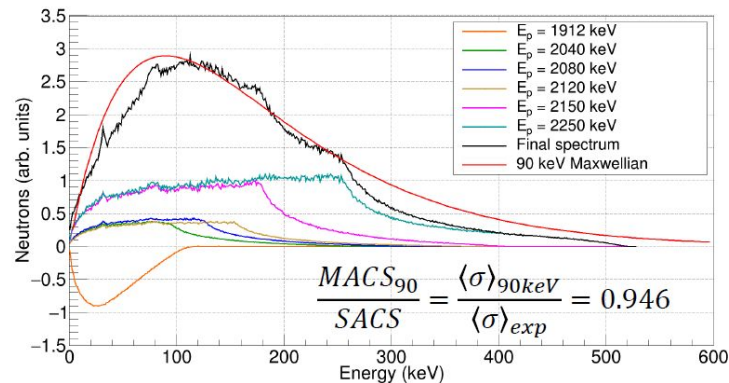


$E_p = 1912 \text{ keV}$ PHYSICAL REVIEW C **93**, 045803 (2016)

STANDARD METHOD



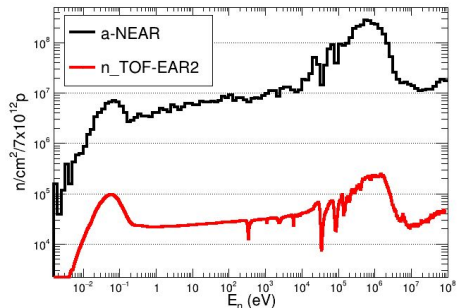
Different E_p combined $\rightarrow$ MACS @ Higher kTs



Med/High energy accelerator (~10MeV - ~GeV)

Target + moderator

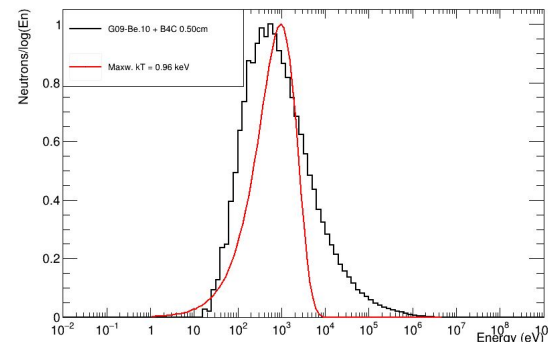
White neutron beam



n_TOF-NEAR APPROACH

Neutron energy shaping $\rightarrow$ MACS @ different kTs

Moderating
+
Filtering
assembly

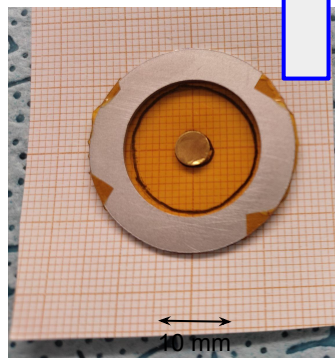


MACS measuring capability benchmark:

E. Stamati et al.,

[CERN-INTC-2022-008:](#)

[INTC-P-623 \(2022\)](#)

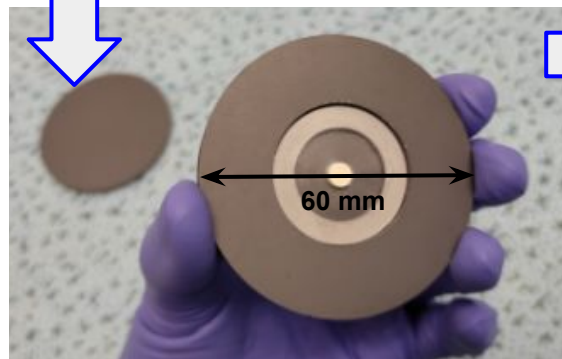


1) Sample preparation:
5 mm sample
on 30mm diam Al ring

2) Irradiation at NEAR:
One long irradiation (~1 week)



3) Access to NEAR (6h) + manual transport to decay station (~h)



4) Decay measurement at GEAR: HPGe decay station

Eff @ 3 cm:
1.6% @818keV



Main limitations:

- Not able to measure activation with short lived (s, min) (n,g) products
- Impossible to activate short lived (hours-days) targets

CYCLING @ NEAR: CYCLIC activation for (N,G) measurements

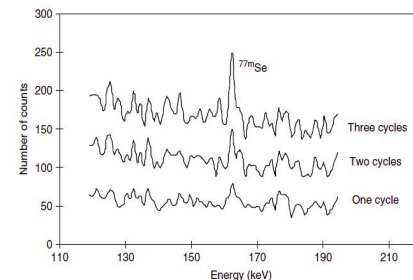
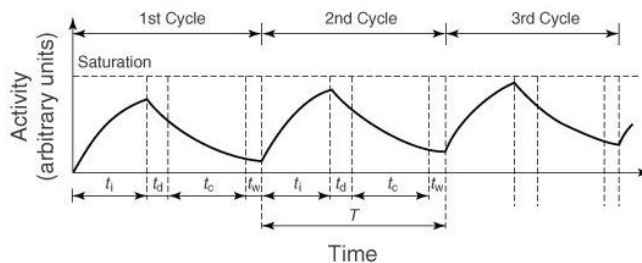
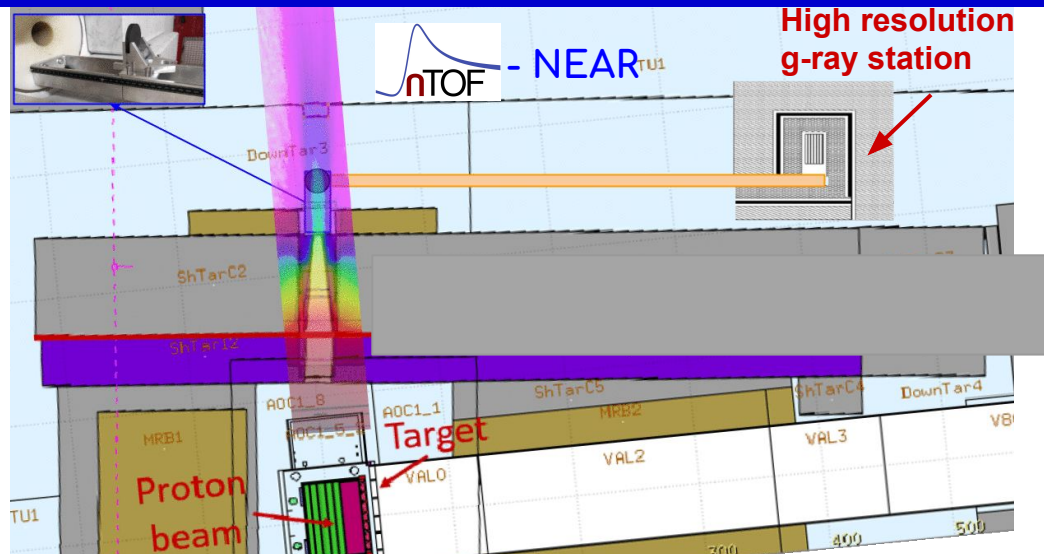
- Allows: activation with short lived (s, min) (n,g) products
- Especially interesting for unstable isotopes.
- Well suited to the rep. Rate ($>0.8\text{Hz}$) of n_TOF

Interesting cases for astrophysics:

- s-process/AGB: $^{107,109}\text{Ag}(n,\gamma)$, $^{26}\text{Mg}(n,\gamma)$, $^{50}\text{Ti}(n,\gamma)$, $^{19}\text{F}(n,\gamma)$, $^{60}\text{Fe}(n,\gamma)$
- i-process: $^{137}\text{Cs}(n,\gamma)$, $^{144}\text{Ce}(n,\gamma)$, $^{132}\text{Te}(n,\gamma)$,...

C. Domingo-Pardo, et al., [Eur. Phys. J. A 59, 8 \(2023\)](#)

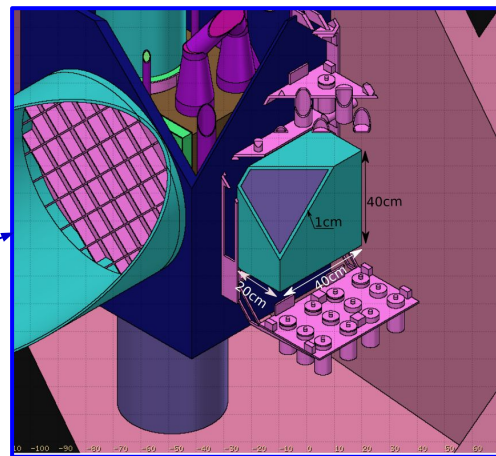
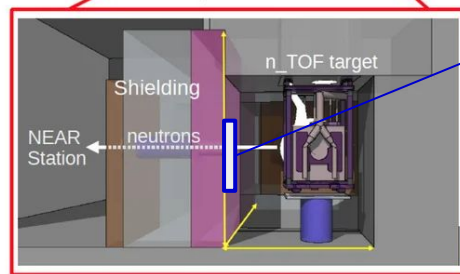
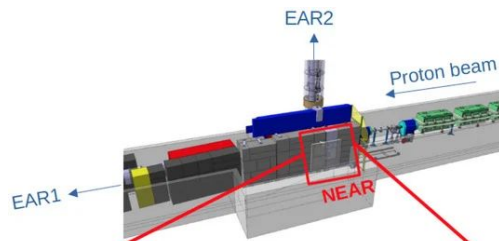
J. Lerendegui, M. Bacak [CERN-INTC-2022-018 ; INTC-I-241.](#)



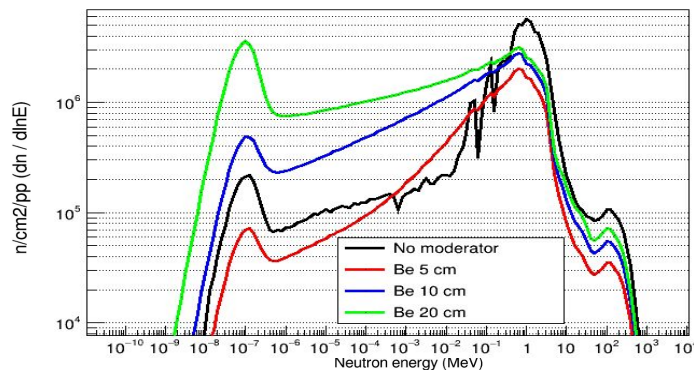
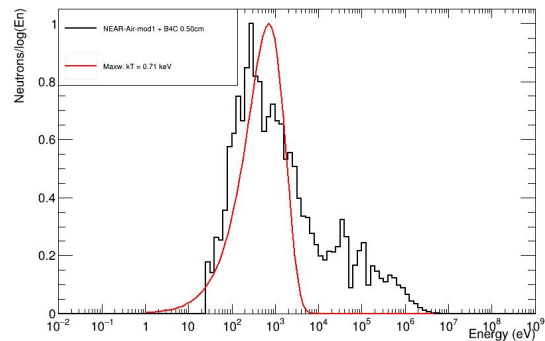
H. Beer, et al., [Nucl. Inst. Meths. 337, 2-3 \(1994\)](#)

Solution:

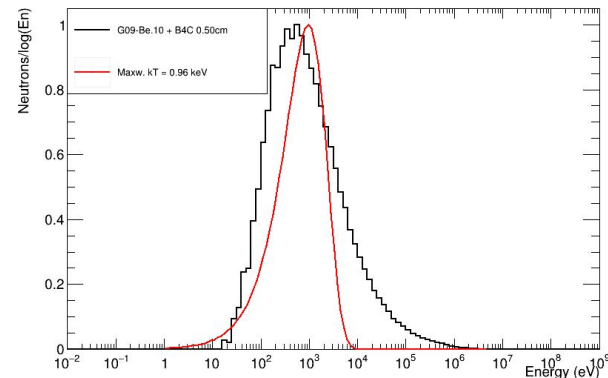
- Include a moderator in the target to improve the match between spectra and a Maxwellian spectra



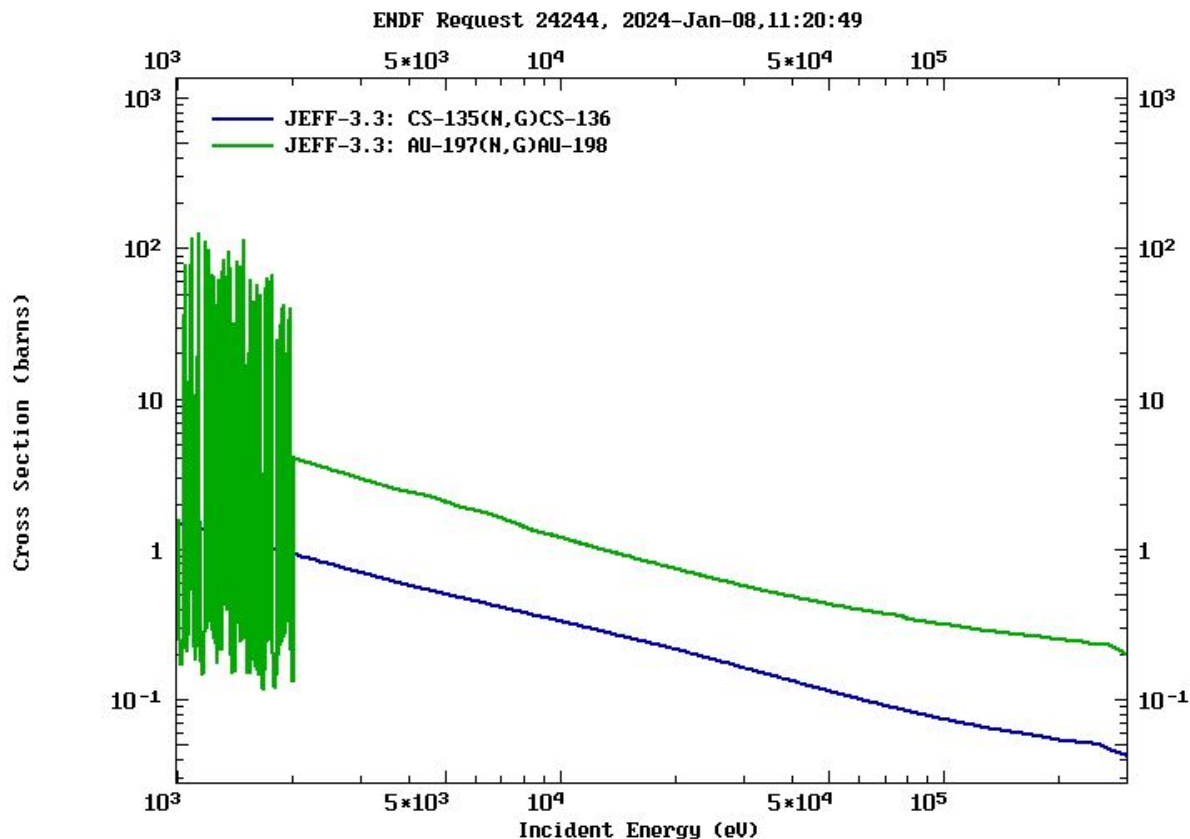
Current: No moderator



10 cm Be moderator



- + Very similar shape of the cross section
- + No large resonances in any of the isotopes
- + MACS/SACS expected to be not very different.
- + Shape of the cross sections can be used to improve the accuracy



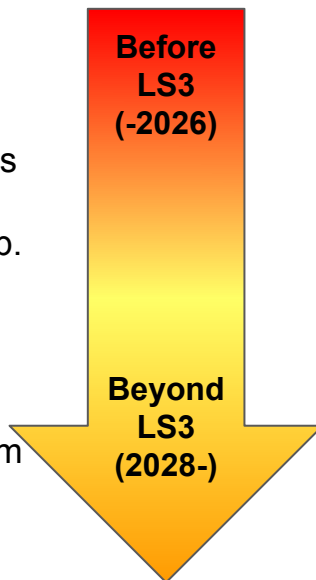
**GENERAL AIM: Enhancing the capabilities for new physics
Focusing on (n,g) measurements on unstable nuclei relevant for astrophysics:**

TOF Measurements on

lower mass unstable (s-process):

- Improve CR capability, efficiency & sensitivity (higher segmentation, larger arrays, solid scintillators, ...)
- Improve signal-to-background: New shieldings and collimator for EAR2 or modification of EAR2 with walls and ceiling further from setup.
- High purity samples (RITU project, PSI)
- Beamline looking to moderator and/or new beamlines (backwards): less background from high energy neutrons
- Higher intensity per pulse & higher rep. rate with similar t-resolution.

**Current
spallation target**



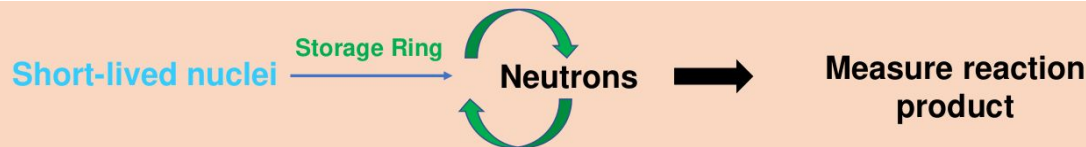
**New spallation
target**

NEAR: activations on short-lived isotopes (s/i-process)

- Installation of NEAR moderator (target) → Improve SACS/MACS
- Cycling: Rabbit system NEAR-ISR and (if feasible) decay station at NEAR.
- Activations inside the target shielding (factor x100 flux) → Rabbit system.
- Strengthen Synergy ISOLDE-NEAR for sample production and combined direct - surrogate reactions.
- Average intensity x10 higher with worse t-resol (suitable only for NEAR).

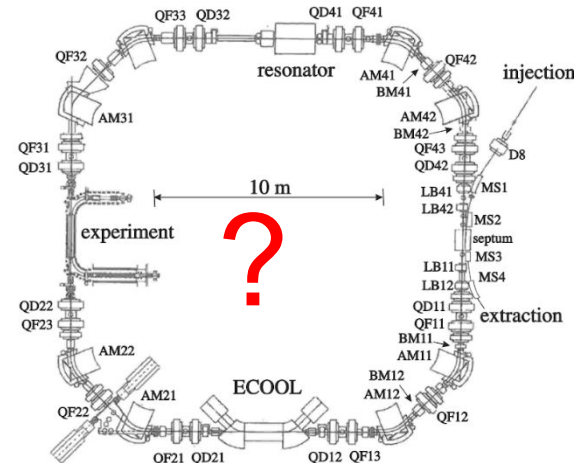
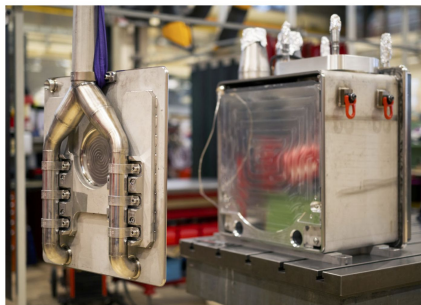
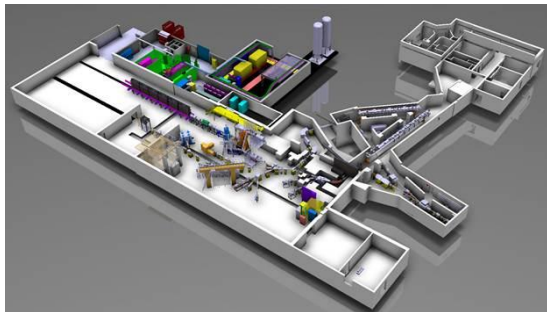
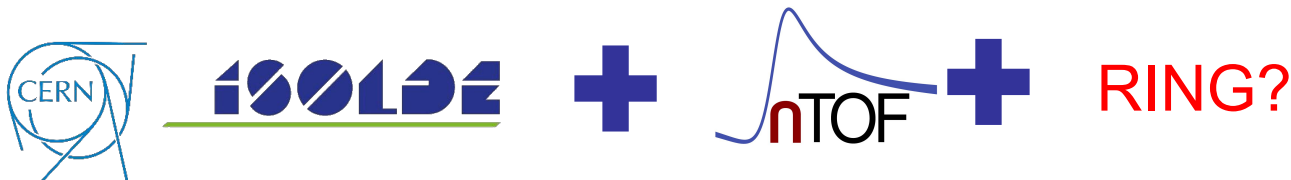
DIRECT (N,G) IN INVERSE KINEMATICS

Aim: short-lived (s- and i-process)



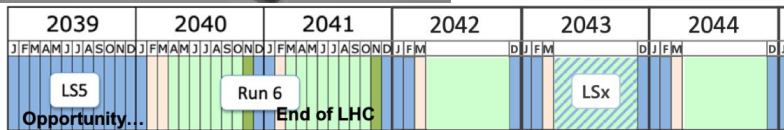
- Possible at CERN n_TOF beam & ISOLDE?

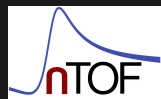
- *s*-process: ^{134}Cs ($t_{1/2} = 2.06$ y), ^{147}Nd ($t_{1/2} = 11$ d), ^{148}Pm ($t_{1/2} = 5.37$ d), ^{153}Gd ($t_{1/2} = 240.4$ d), ^{160}Tb ($t_{1/2} = 72.3$ d), ^{170}Tm ($t_{1/2} = 128.6$ d), and ^{204}Tl ($t_{1/2} = 3.78$ y).
- *i*-process: ^{66}Ni ($t_{1/2} = 54.6$ h), ^{75}Ga ($t_{1/2} = 126$ s), ^{135}I ($t_{1/2} = 6.58$ h), ^{141}Ba ($t_{1/2} = 18.27$ m), and ^{141}La ($t_{1/2} = 3.92$ h).



Project: Storage ring at HIE-ISOLDE

M. Grieser et al., Eur. Phys. J Spec. Top. 207, 1 (2012)

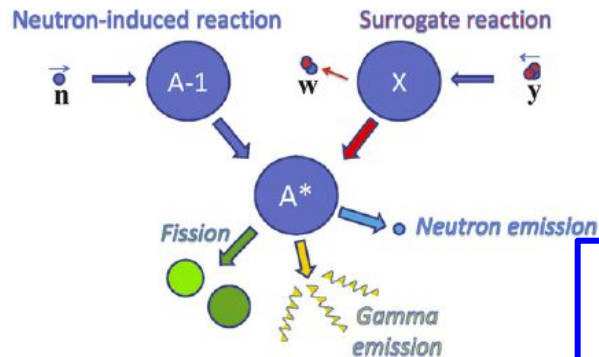




Other reaction channels / methods



SURROGATE REACTIONS



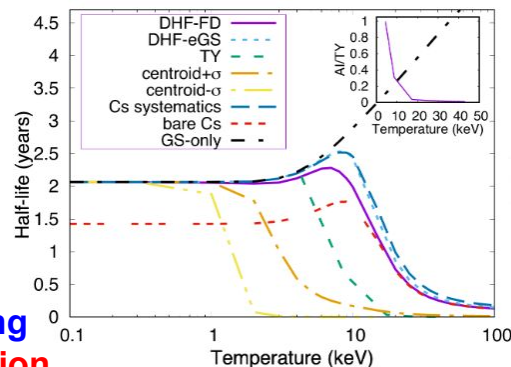
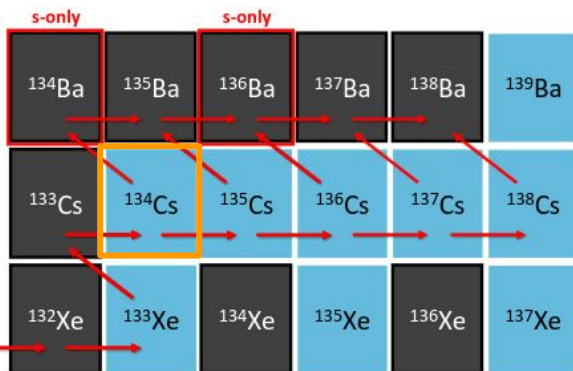
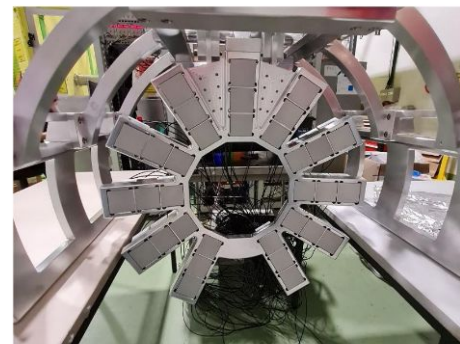
Towards Neutron Capture on Exotic Nuclei: Demonstrating ($d,p\gamma$) as a Surrogate Reaction for (n,γ)

A. Ratkiewicz,^{1,2,*} J. A. Cizewski,² J. E. Escher,¹ G. Potel,^{3,4} J. T. Harke,¹ R. J. Casperson,¹

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH
Proposal to the ISOLDE and Neutron Time-of-Flight Committee

Measurement of neutron capture cross section on ^{134}Cs through surrogate reaction ($d,p\gamma$) at ISS

January 10, 2024



Temperature-sensitive s-branching
No TOF (High activity) + No activation

Surrogate reaction Inverse kinematics
 ^{134}Cs beam @ ISOLDE + CD_2 target
ISOLDE Solenoid

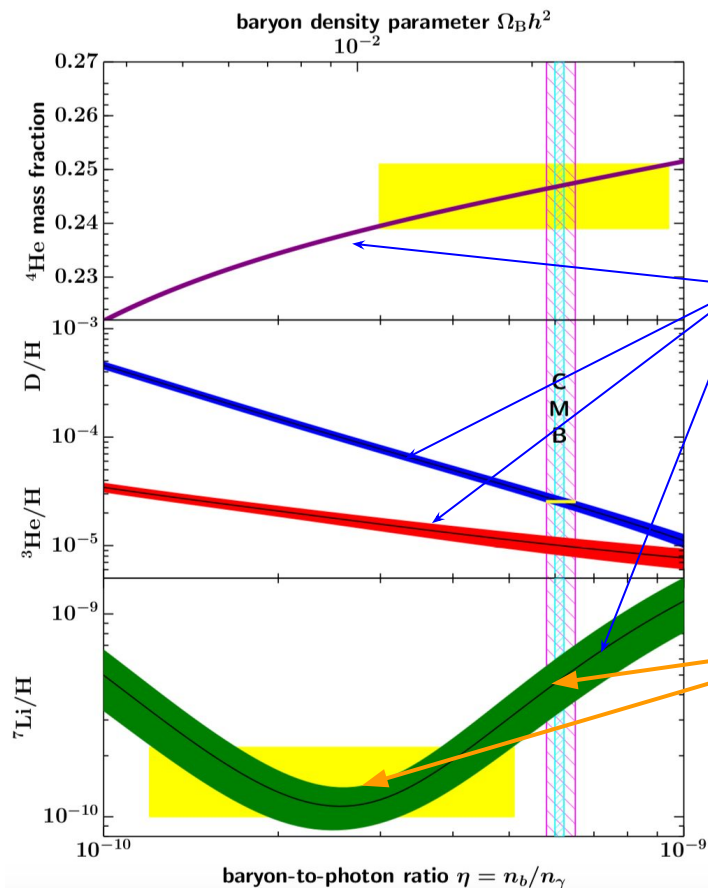
Future idea: complement (n,g) at NEAR (low E_n) + ($d,p\gamma$) at ISS-ISOLDE

s-process: production of the majority of elements heavier than Fe.

BBN and the Li-problem: Neutron cross sections are key to clarify the gross overestimate in BBN models of the primordial abundance of Lithium.

Neutron-induced reactions
in Nuclear Astrophysics

Neutron sources: neutron cross sections of light elements acting as neutron poison, or linked to stellar neutron sources.



BBN predictions

BBN and the Li-problem: Neutron cross sections are key to clarify the gross overestimate in BBN models of the primordial abundance of Lithium.

Production of ^7Li : dominated by the ^7Be EC decay after the BBN.

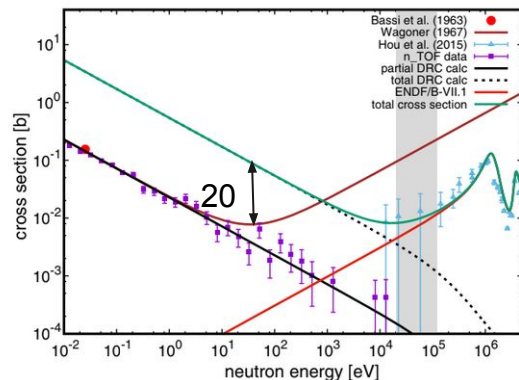
Lower ^7Li abundance? A higher rate for **neutron-induced** reactions on ^7Be ?

PRL 117, 152701 (2016)

PHYSICAL REVIEW LETTERS

week ending
7 OCTOBER 2016

$^7\text{Be}(n,\alpha)^4\text{He}$ Reaction and the Cosmological Lithium Problem: Measurement of the Cross Section in a Wide Energy Range at n_TOF at CERN

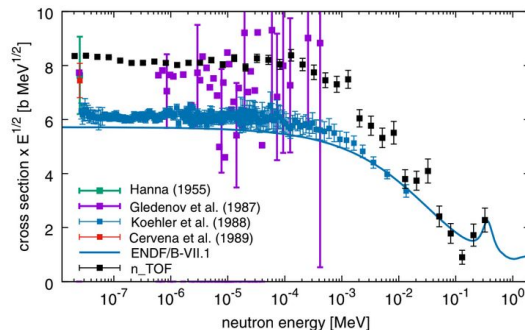


- $\text{Be-7}(n,\alpha)$: x20 vs calculations of BBN below 100 eV
- $\text{Be-7}(n,p)$: +40% vs previous at low En

BBN and the Li-problem: Neutron cross sections are key to clarify the gross overestimate in BBN models of the primordial abundance of Lithium.

PHYSICAL REVIEW LETTERS **121**, 042701 (2018)

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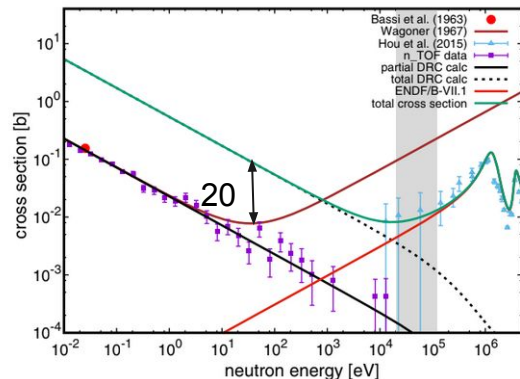
$^7\text{Be}(n, p)^7\text{Li}$ rate	η_{10}	Li/H yield
Cyburt (2004) rate [22]	6.09	5.46
This work [Eq. (3)]	6.09	5.26 ± 0.40
	5.8–6.6	4.73–6.23
Observations [1]		1.6 ± 0.3

PRL 117, 152701 (2016)

PHYSICAL REVIEW LETTERS

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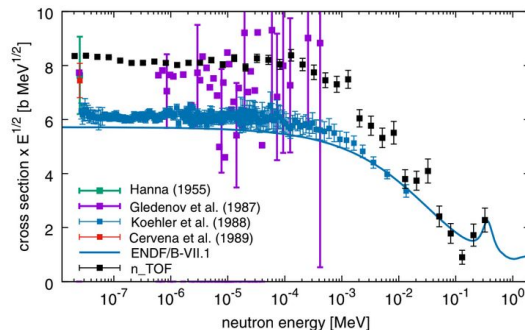
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- $\text{Be-7}(n,p)$: +40% vs previous at low E_n

Neutrons are not enough to solve the CLiP → Alternative physics or astromical scenarios

BBN and the Li-problem: Neutron cross sections are key to clarify the gross overestimate in BBN models of the primordial abundance of Lithium.

PHYSICAL REVIEW LETTERS **121**, 042701 (2018)

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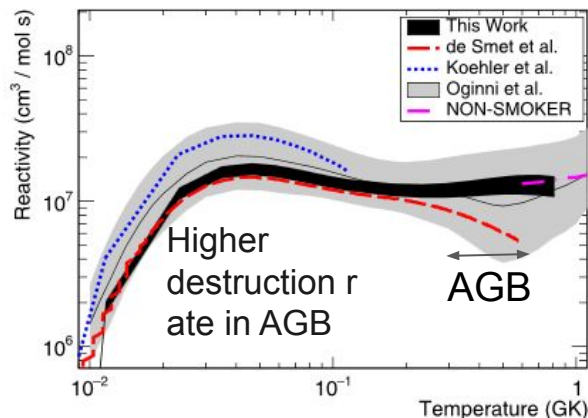
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PHYSICAL REVIEW C **104**, L032803 (2021)

Letter

Destruction of the cosmic γ -ray emitter ^{26}Al in massive stars: Study of the key $^{26}\text{Al}(n, \alpha)$ reaction

C. Lederer-Woods,^{1,*} P. J. Woods,¹ T. Davinson,¹ A. Estrade,^{1,*} J. Heyse,² D. Kahl,^{1,‡} S. J. Lonsdale,¹ C. Paradelo,²



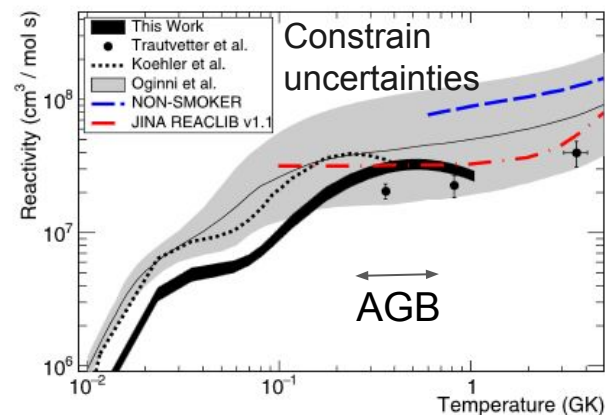
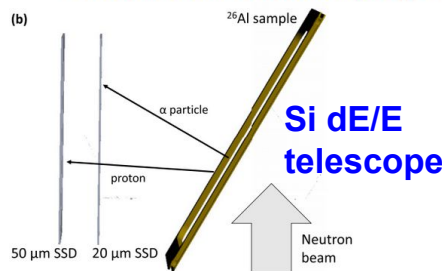
- Significant discrepancies in the Al-26 destruction rates.
- Impact in the Al-26 production in supernovae and AGB origin of our early solar system

Neutron sources: neutron cross sections of light elements acting as neutron poison, or linked to stellar neutron sources.

PHYSICAL REVIEW C **104**, L022803 (2021)

Letter

Destruction of the cosmic γ -ray emitter ^{26}Al in massive stars: Study of the key $^{26}\text{Al}(n, p)$ reaction

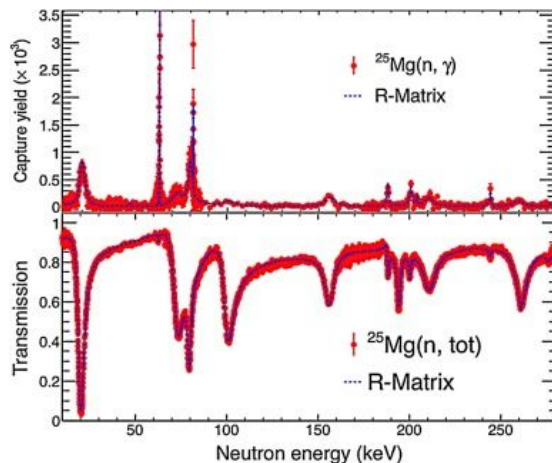
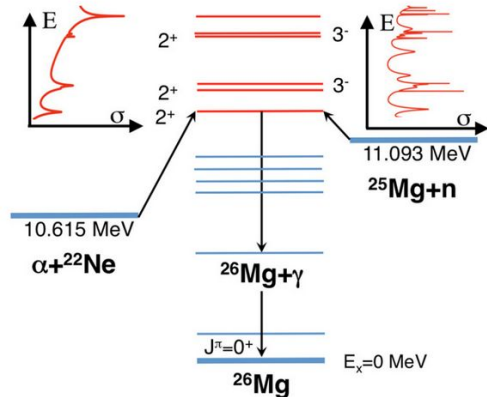


Next: extend >1 GK (Massive stars)

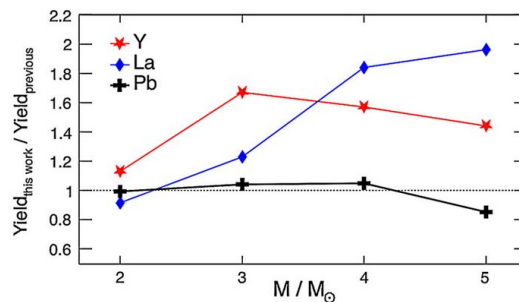
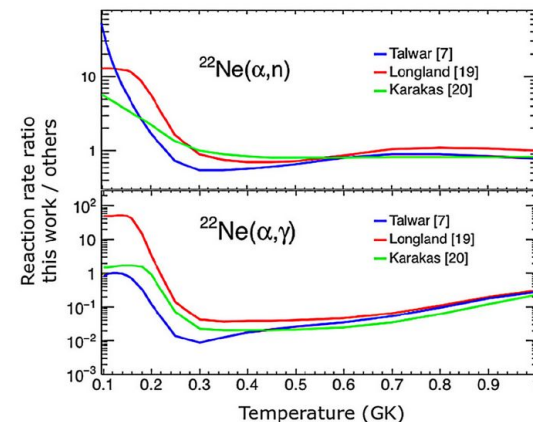
Neutron spectroscopy of ^{26}Mg states: Constraining the stellar neutron source $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$



C. Massimi^{a,b,*}, S. Altstadt^c, J. Andrzejewski^d, L. Audouin^e, M. Barbagallo^f, V. Bécarrés^g,



Neutron-spectroscopy to experimentally constrain the neutron source reaction $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$



- R-matrix on (n,g) → parity assignment of the excited states in ^{26}Mg
- $^{22}\text{Ne}(\alpha, n)$ rates enhanced at low temperature.
- Large impact in the AGB s-process abundances (2-5 $M_{\odot}$).

Neutron sources: neutron cross sections of light elements acting as neutron poison, or linked to stellar neutron sources.