

Quantum field measurements with(out) superluminal signalling

Jason Pye

work with José de Ramón, Eduardo Martín-Martínez

The Time Machine Factory, 24 Sep 2019



UNIVERSITY OF
WATERLOO



Theory of measurements/operations for QFT

How to do quantum operations in a relativistic spacetime?

- i.e. how to introduce agents
- finite times (push beyond S-matrix)

Contexts:

- Time machine paradoxes, BH entropy, ...
- QI experiments in relativistic regime (q. optics, cQED, ...)
- Foundational question: What can be measured in principle?

QM operations \rightarrow QFT?

Naively, can use standard idealised measurements from QM

- PVMs: projectors $\rho \rightarrow P_a \rho P_a$

e.g., from self-adjoint $A = \sum_a a P_a \rightarrow \langle A \rangle$

- POVMs: $\rho \rightarrow M_i \rho M_i^\dagger$

- Channels (CPTP maps):

$$\rho \rightarrow \sum_i M_i \rho M_i^\dagger$$

$$\text{s.t. } \sum_i M_i^\dagger M_i = 1$$

Does relativity impose further requirements?

Impossible Measurements on Quantum Fields*

RAFAEL D. SORKIN

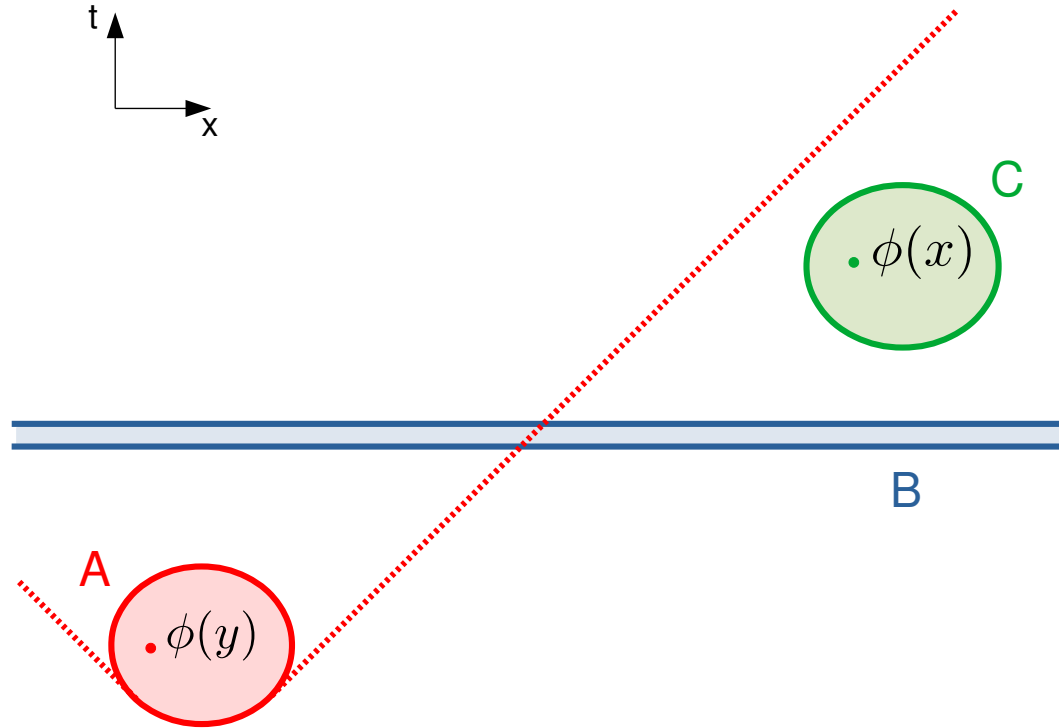
Department of Physics, Syracuse University, Syracuse NY 13244-1130

20 Feb 1993

Abstract

It is shown that the attempt to extend the notion of ideal measurement to quantum field theory leads to a conflict with locality, because (for most observables) the state vector reduction associated with an ideal measurement acts to transmit information faster than light. Two examples of such information-transfer are given, first in the quantum mechanics of a pair of coupled subsystems, and then for

Sorkin's Impossible Measurement



- Initial state $\rho_0 = |0\rangle\langle 0|$
- **Operation in A**

$$U = e^{i\lambda\phi(y)} \quad (\lambda \in \mathbb{R})$$

$$\rho = U\rho_0U^\dagger$$
- **Measurement in B (unconditional)**

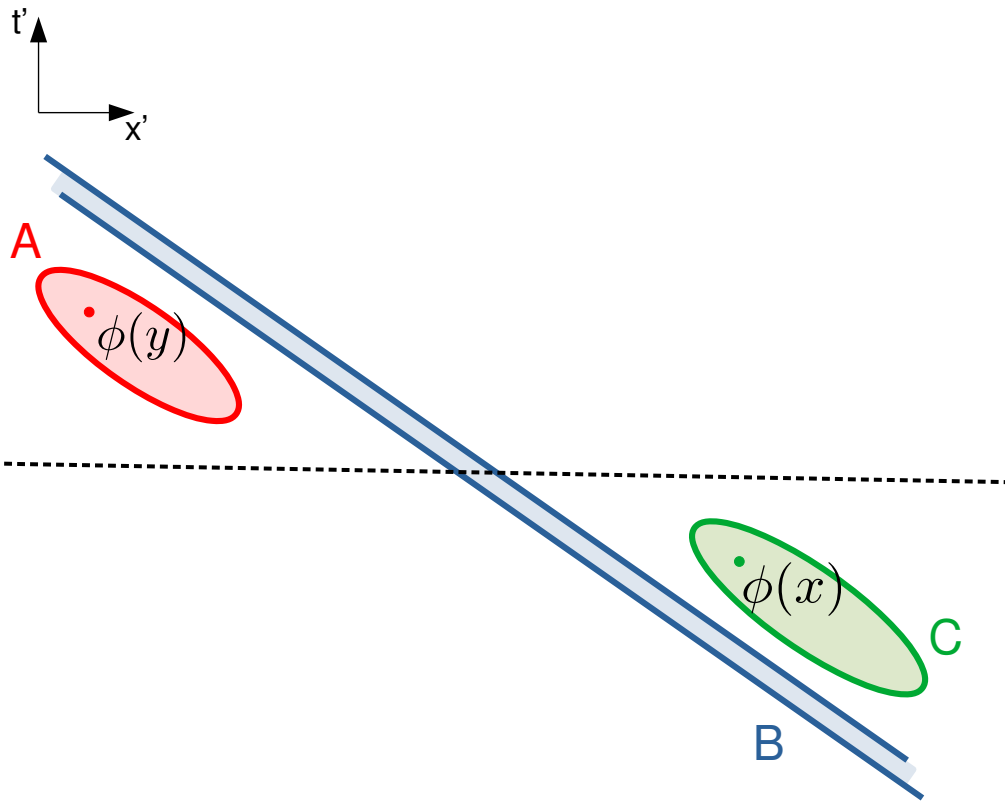
$$P = |\psi\rangle\langle\psi|$$

with $|\psi\rangle := \alpha|0\rangle + \beta|1\rangle \quad (|1\rangle \in \mathcal{H}_1)$

$$\rho' = P\rho P + (1 - P)\rho(1 - P)$$
- **Expectation value in C**

$$\langle\phi(x)\rangle_{\rho'} \text{ depends on } \lambda!$$

Sorkin's Impossible Measurement



- Initial state $\rho_0 = |0\rangle\langle 0|$

- **Operation in A**

$$U = e^{i\lambda\phi(y)} \quad (\lambda \in \mathbb{R})$$

$$\rho = U\rho_0U^\dagger$$

- **Measurement in B (unconditional)**

$$P = |\psi\rangle\langle\psi|$$

$$\text{with } |\psi\rangle := \alpha|0\rangle + \beta|1\rangle \quad (|1\rangle \in \mathcal{H}_1)$$

$$\rho' = P\rho P + (1 - P)\rho(1 - P)$$

- **Expectation value in C**

$$\langle\phi(x)\rangle_{\rho'} \text{ depends on } \lambda!$$

Implications

Assuming theory should not admit superluminal signalling of this kind...

- Sorkin's quantum measure (path integral); OR
- (Operator framework): not every idealized quantum operation is admissible
 - Measurement postulates vs. Relativistic spacetime structure

$$A = \sum_a a P_a$$

$$p_a = \text{tr}(\rho P_a)$$

$$\rho \rightarrow P_a \rho P_a$$

vs.

$$[\phi(x), \phi(y)] = 0,$$

for $(x - y)^2 > 0$ (spacelike)

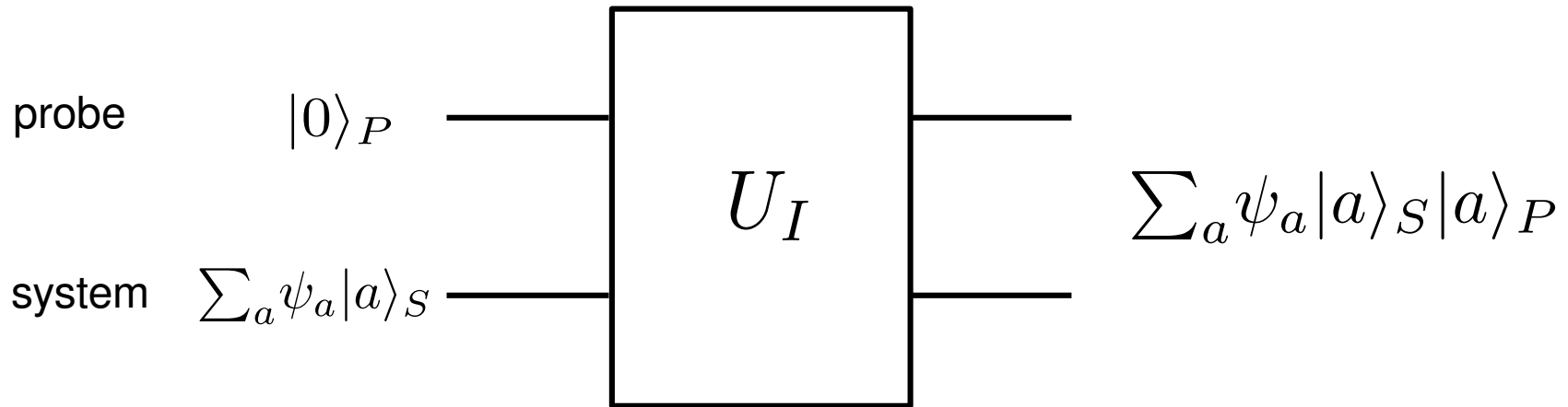
- Counterexample to: self-adjoint \implies observable

Aim

- What measurements are/are not allowed?
 - What precisely is the issue with Sorkin's measurement?
- Would like:
 - List of observables, e.g., $H, P, N, \phi(f), \phi(f)^2, \dots$?
 - Characterization of allowable PVMs/POVMs, channels

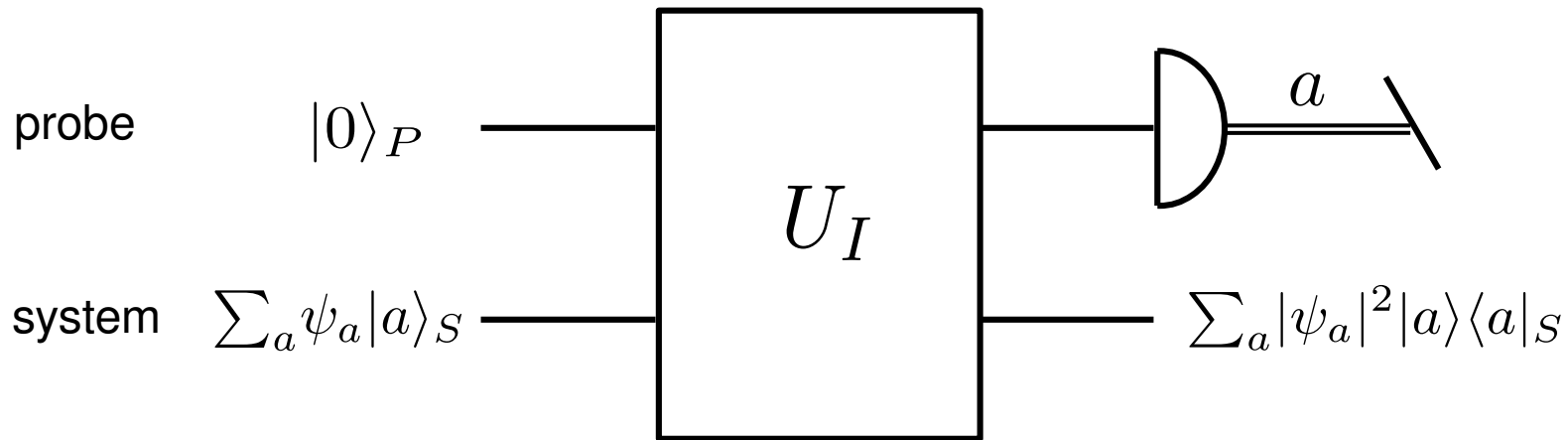
How would these measurements be performed?

- Measurements via interactions
 - Idealized von Neumann measurement



How would these measurements be performed?

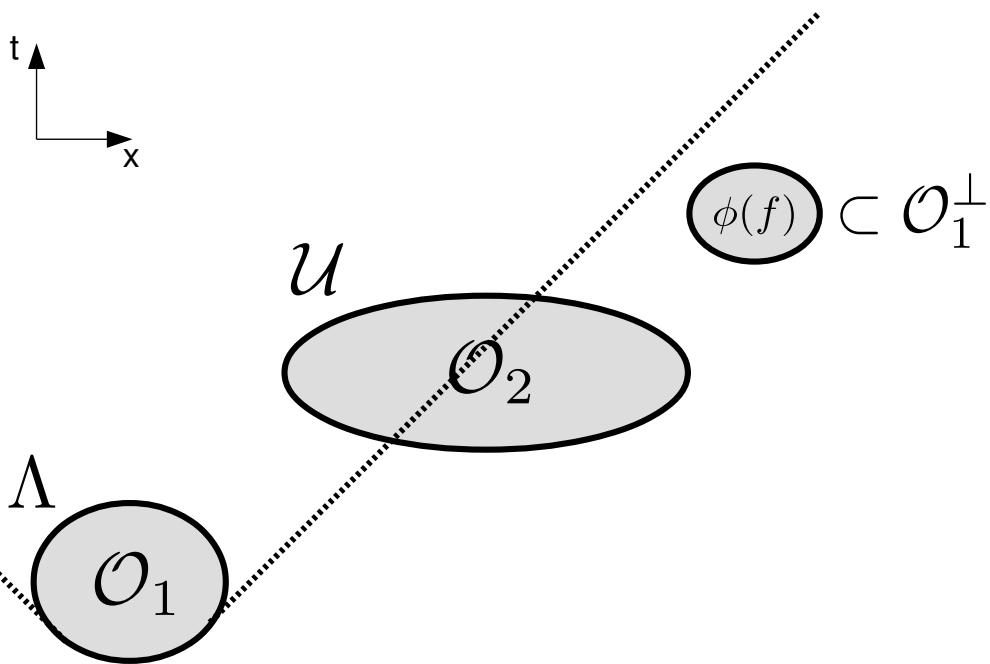
- Measurements via interactions
 - Idealized von Neumann measurement



channel on system $\mathcal{U} : \rho \mapsto \sum_a P_a \rho P_a$

Causality condition

Formulate condition that operations should satisfy:



- For any map in (any) \mathcal{O}_1
 $\Lambda : \mathcal{A}(\mathcal{O}_1) \rightarrow \mathcal{A}(\mathcal{O}_1)$
- Operation (CP, unital) in \mathcal{O}_2
 $\mathcal{U} : \mathcal{A}(\mathcal{O}_2) \rightarrow \mathcal{A}(\mathcal{O}_2)$

should be s.t.

- any state $\omega : \mathcal{A}(\mathcal{M}) \rightarrow \mathbb{C}$
- any $A \in \mathcal{A}(\mathcal{O}_1^\perp)$

$$\omega(\Lambda \circ \mathcal{U}(A)) \stackrel{!}{=} \omega(\mathcal{U}(A))$$

Measurements via interactions with probes

- Sorkin-type impossible measurement

$$\begin{aligned} H_I(t) &= \lambda \mu(t) \otimes \phi(f)^2 \\ &= \lambda \mu(t) \otimes \int d\mathbf{x}d\mathbf{y} f(t, \mathbf{x})f(t, \mathbf{y})\phi(t, \mathbf{x})\phi(t, \mathbf{y}) \end{aligned}$$

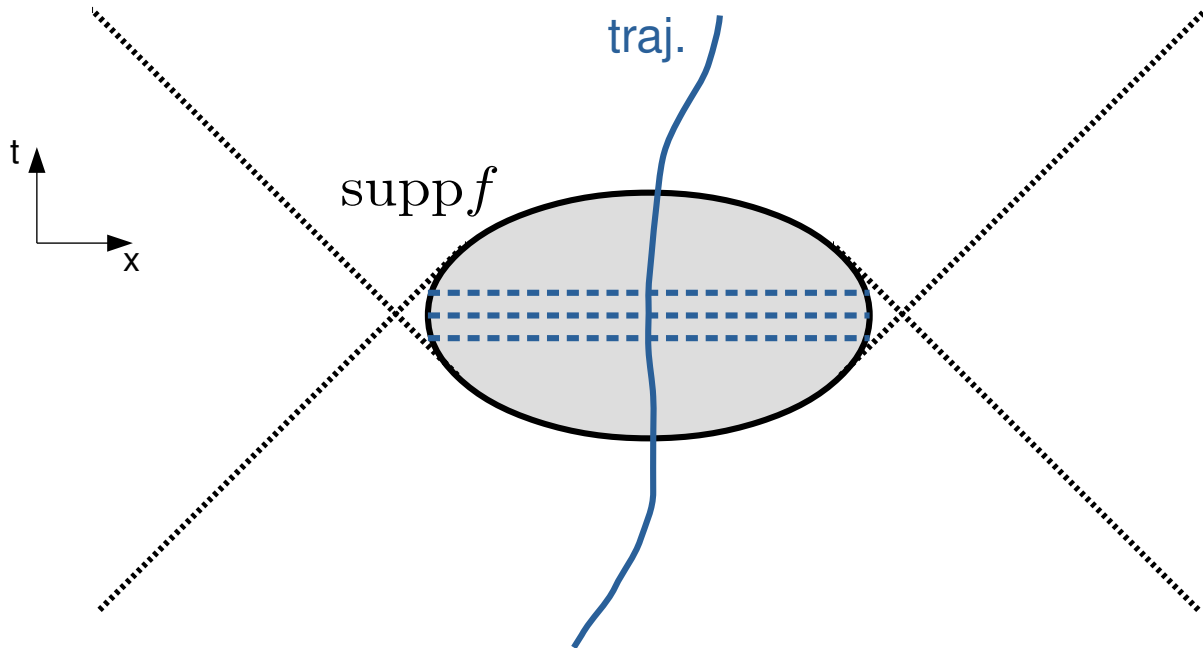
- Allowed, e.g., Weyl channels

$$\rho \rightarrow \sum_a p_a e^{-ia\phi(f)} \rho e^{ia\phi(f)}$$

Issues with internal dynamics

- Non-relativistic probes highly constrained

$$H = H_F \otimes 1 + 1 \otimes \cancel{H_D} + H_I$$



Unruh-deWitt-like model

$$H_I = \lambda p(t) \otimes \phi(f)$$

Towards a general characterization

- Non-relativistic probes generate examples of allowable measurements/operations
- May not exhaust set of allowable idealized operations
 - (If not, some kind of dilation theorem?)
- Source in the classical interaction?
- Eventually should take axiomatic approach to avoid proving existence of general classes of models

Summary

- Gained insight into issues causing superluminal signalling for idealized measurements
 - Non-local interaction terms, internal probe dynamics, ...
- Simple class (Weyl channels) of allowable operations generated by non-relativistic probes
- **Message**: be careful when performing quantum measurements/operations on relativistic spacetime!