

# Concurrent spectral and light curve synchro-curvature modelling of pulsars

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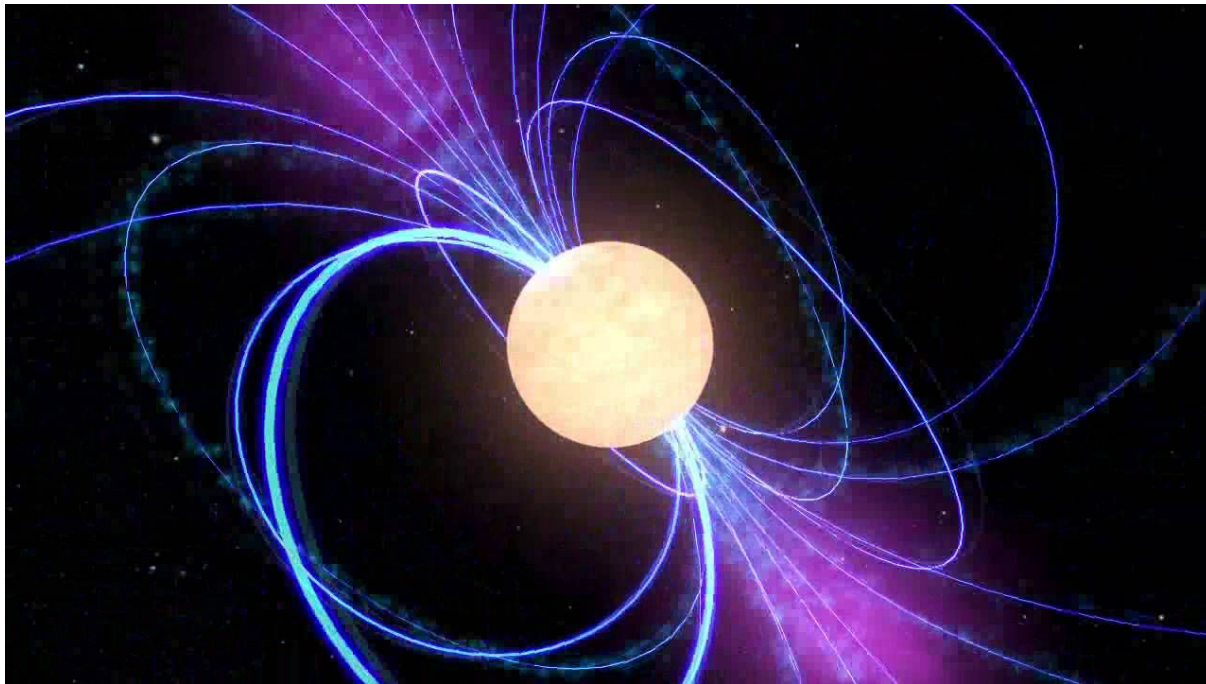
Co-authors: Diego F Torres, Daniele Viganò

8<sup>th</sup> Heidelberg International Symposium on High-Energy Gamma-ray Astronomy ( $\Gamma$ -2024)

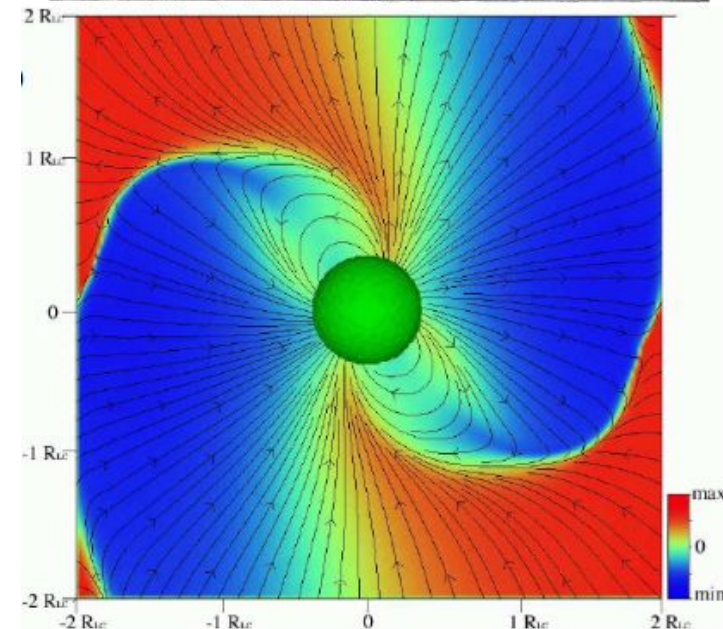
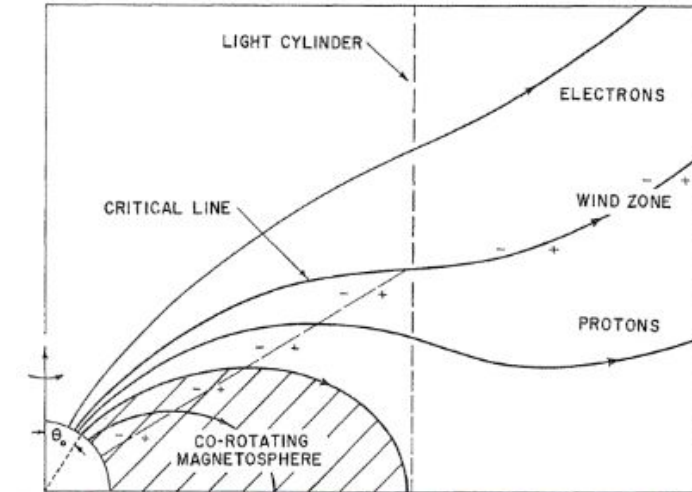
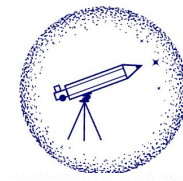
2-6 September 2024 - Milano, Italy

# Motivation of the project

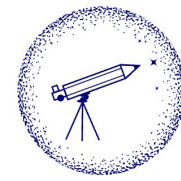
- High-energy pulsars are those with detected pulsated high-energy emission, gamma-rays and X-rays.



Credit: NASA/Fermi

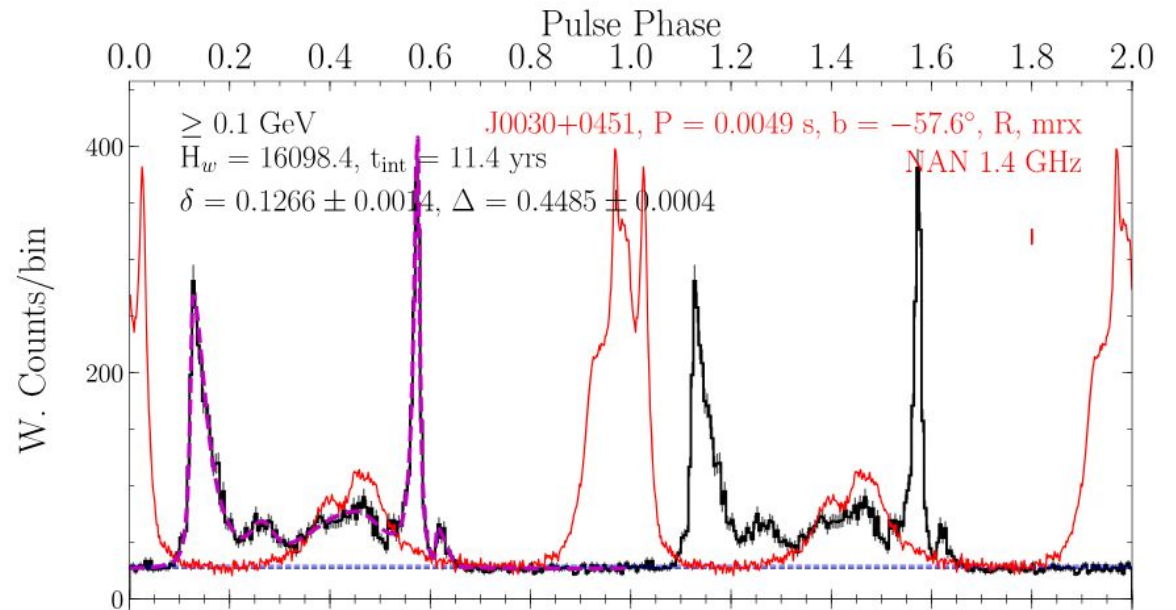
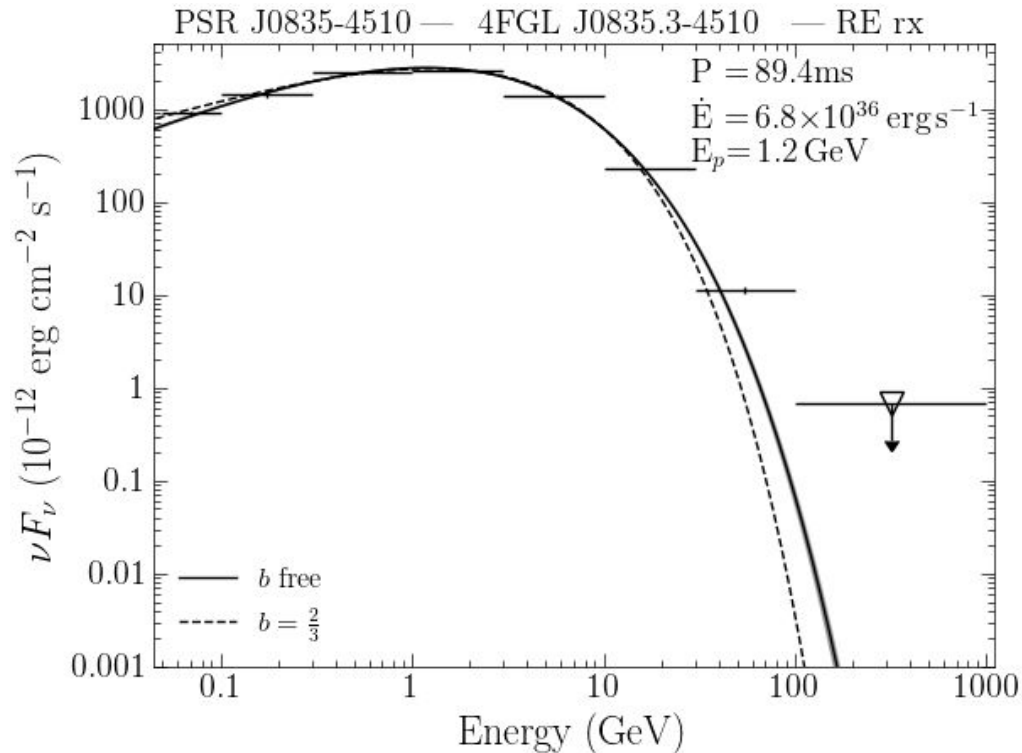


[Goldreich &  
Julian (1969);  
Spitkovsky  
(2006)]



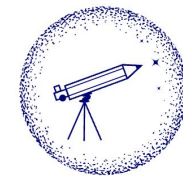
# Motivation of the project

- Models can be tested with e.g. the data of the gamma-ray emitting pulsars released in the 3PC (Smith et al. (2023))

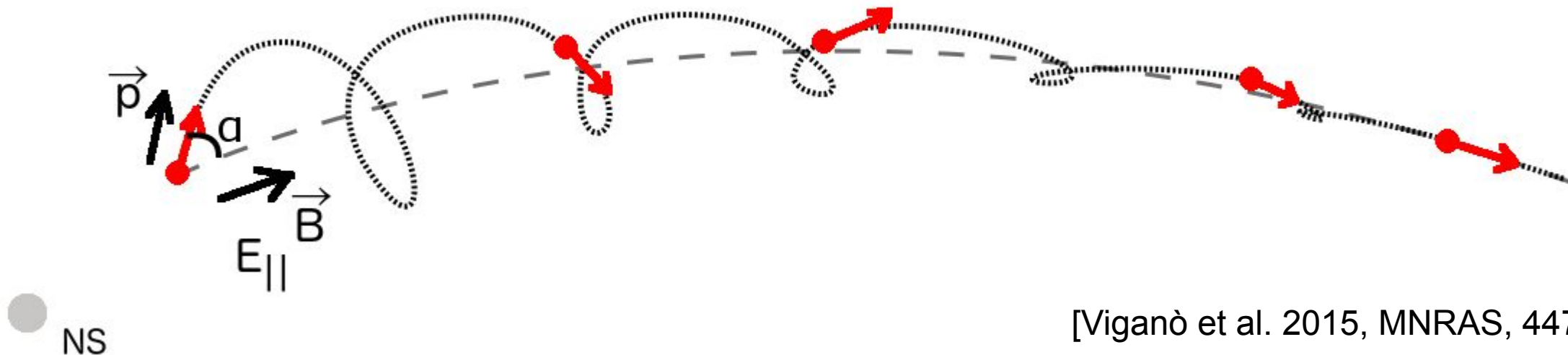


- Our goal is to reproduce observational data with a model that contains simple but realistic physics and is computationally affordable

# Spectral model: particle dynamics

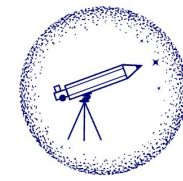


- We follow the dynamics of the emitting particles, ruled by electric acceleration and synchro-curvature losses and with two free parameters involved:  $E_{||}$ ,  $b$
- Solving the equation of motion gives the evolution of the relativistic momentum and of the Lorentz factor  $\Gamma$  and pitch angle  $\alpha$



[Viganò et al. 2015, MNRAS, 447, 1164]

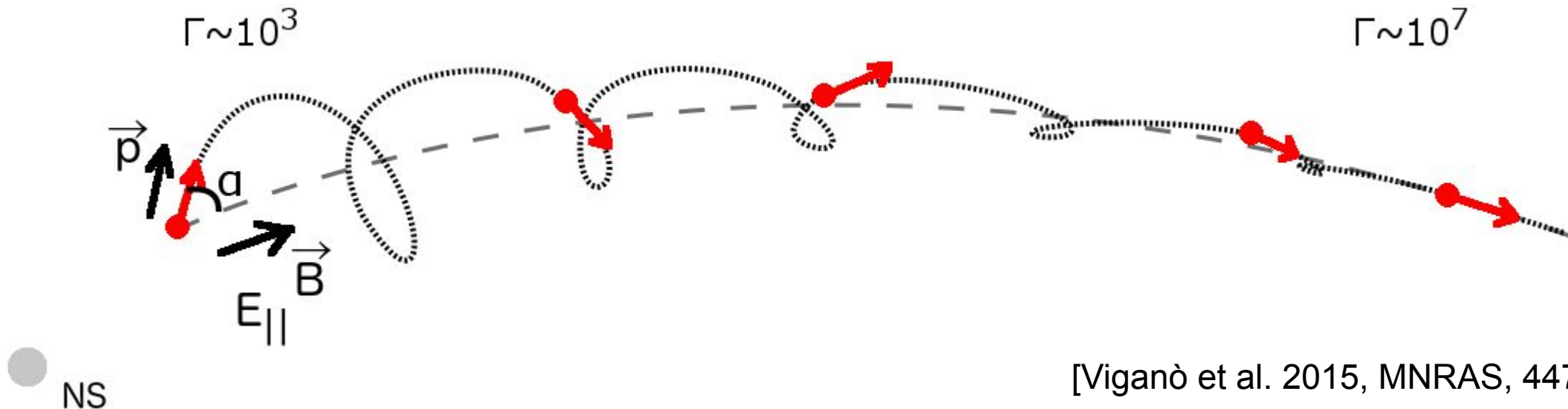
# Spectral model: particle dynamics



- Power of synchrotron and curvature radiations depend on the Lorentz factor  $\Gamma$  of particles differently

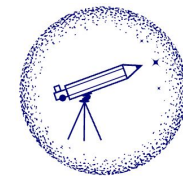
$$P_{syn} = \frac{2}{3} \frac{(Ze)^4 B^2 (\Gamma^2 - 1) \sin^2 \alpha}{m^2 c^3}$$

$$P_c = \frac{2}{3} \frac{(Ze)^2 c \Gamma^4}{r_c^2}$$

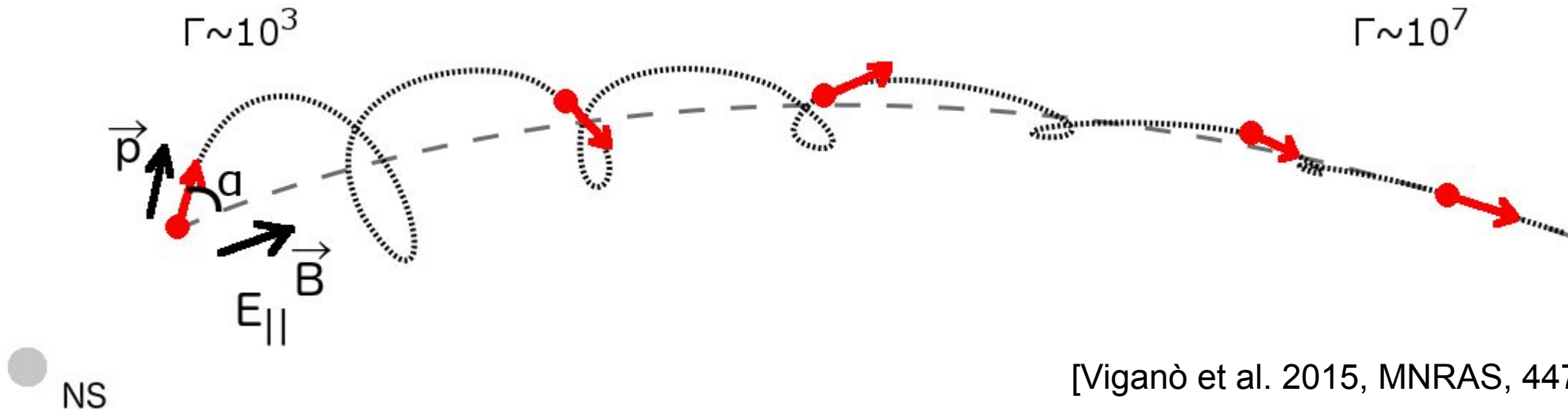


[Viganò et al. 2015, MNRAS, 447, 1164]

# Spectral model: emission

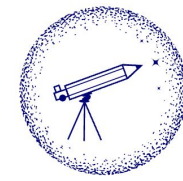


- Synchro-curvature formulae gives the emission of the particles all along the trajectory, which convolved with an effective particle distribution gives the total radiation from the emission region
- We produce theoretical spectra with just three free parameter ( $E_{||}$ ,  $b$ ,  $x_0$ ) and a normalization factor



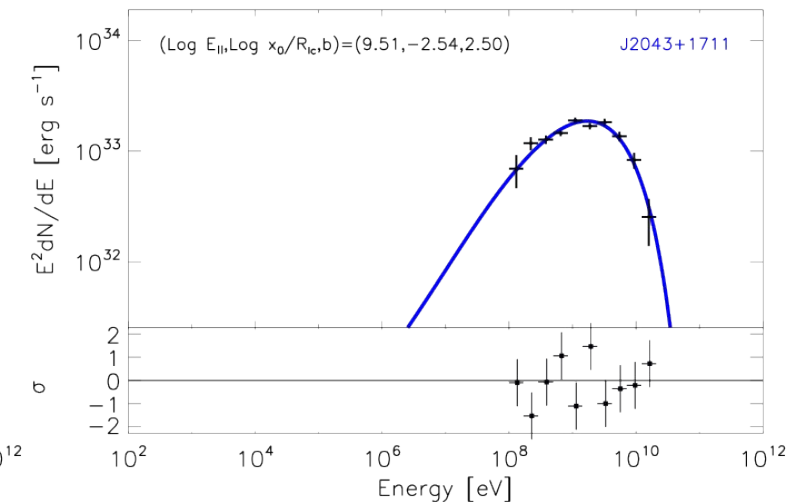
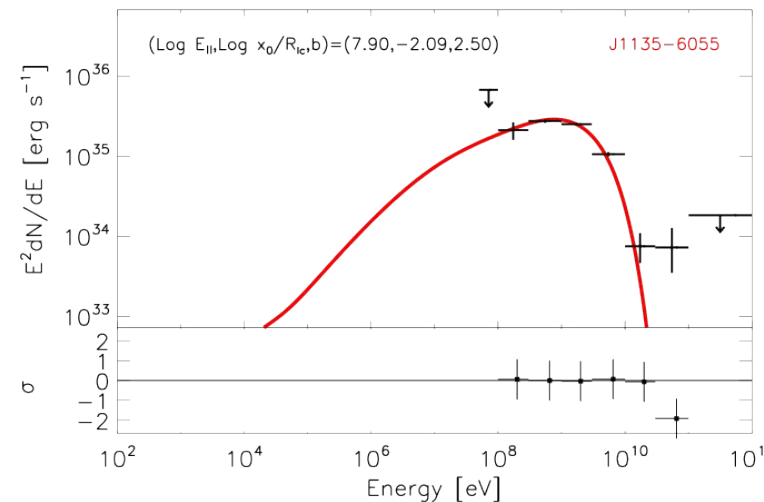
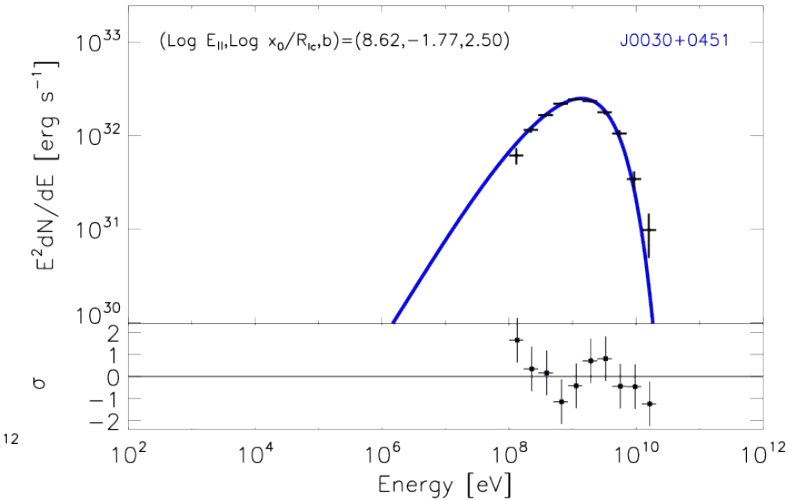
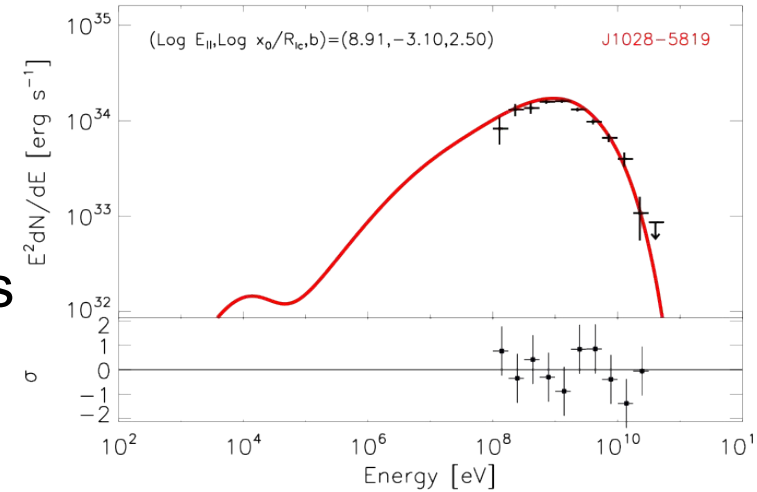
[Viganò et al. 2015, MNRAS, 447, 1164]

# Spectral fitting



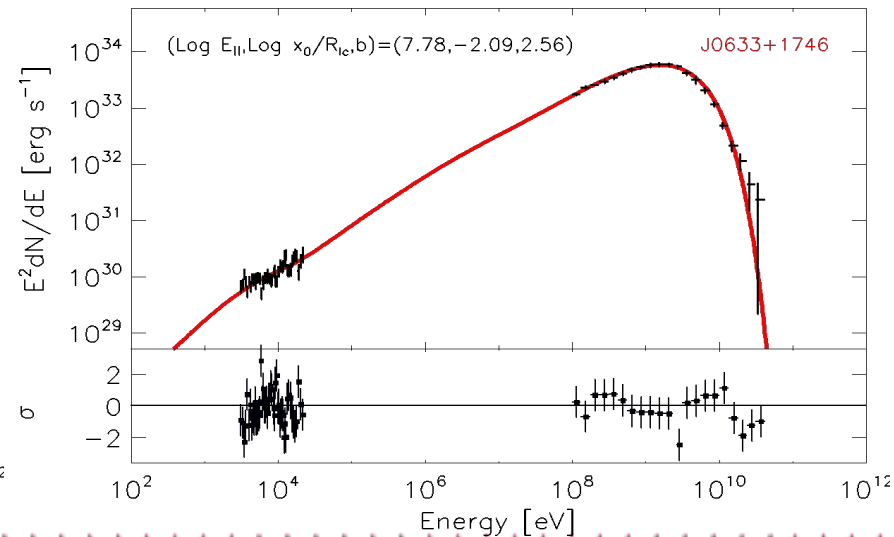
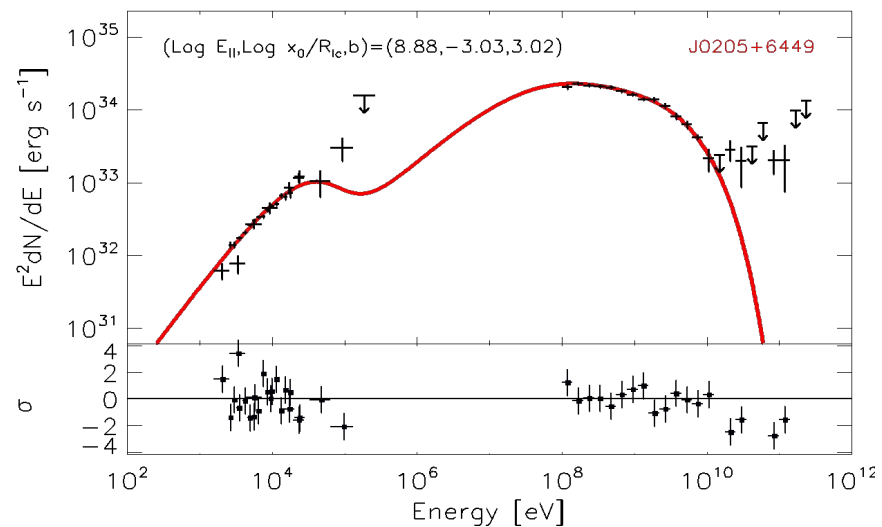
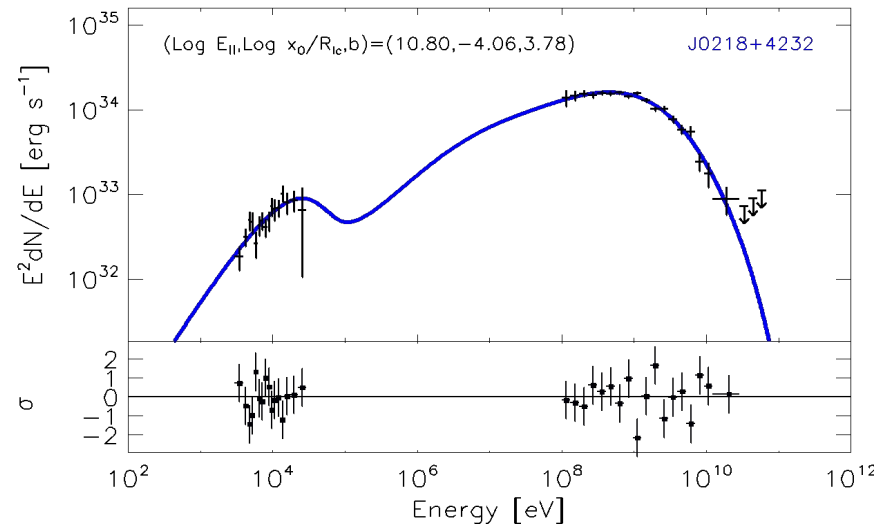
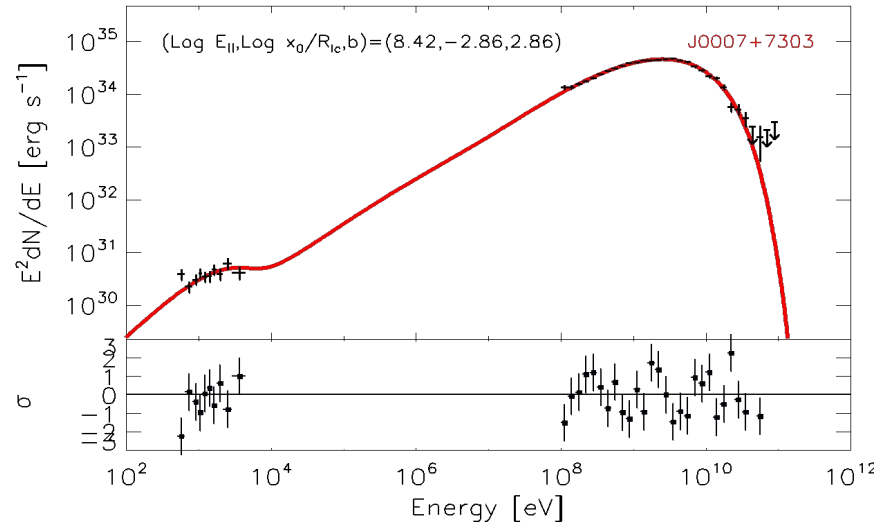
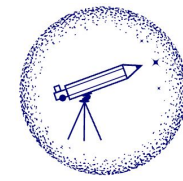
[Viganò et al. 2015, MNRAS, 453, 2599; Íñiguez-Pascual et al. (in prep.)]

- The model successfully fitted the 117 gamma ray-pulsars on the 2PC (Abdo et al. (2013)) and the ~300 gamma-ray pulsars on the 3PC (Smith et al. (2023))
- A relevant synchrotron contribution is needed to match gamma-ray spectra, implying that synchro-curvature radiation is an appropriate mechanism to explain the emission from these objects



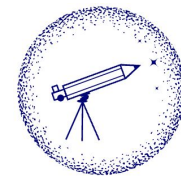
# Spectral fitting

[Torres 2018, Nat. Astron., 2, 247;  
Torres et al. 2019, MNRAS, 489, 5494;  
Íñiguez-Pascual, Viganò & Torres 2022, MNRAS, 516, 2475]



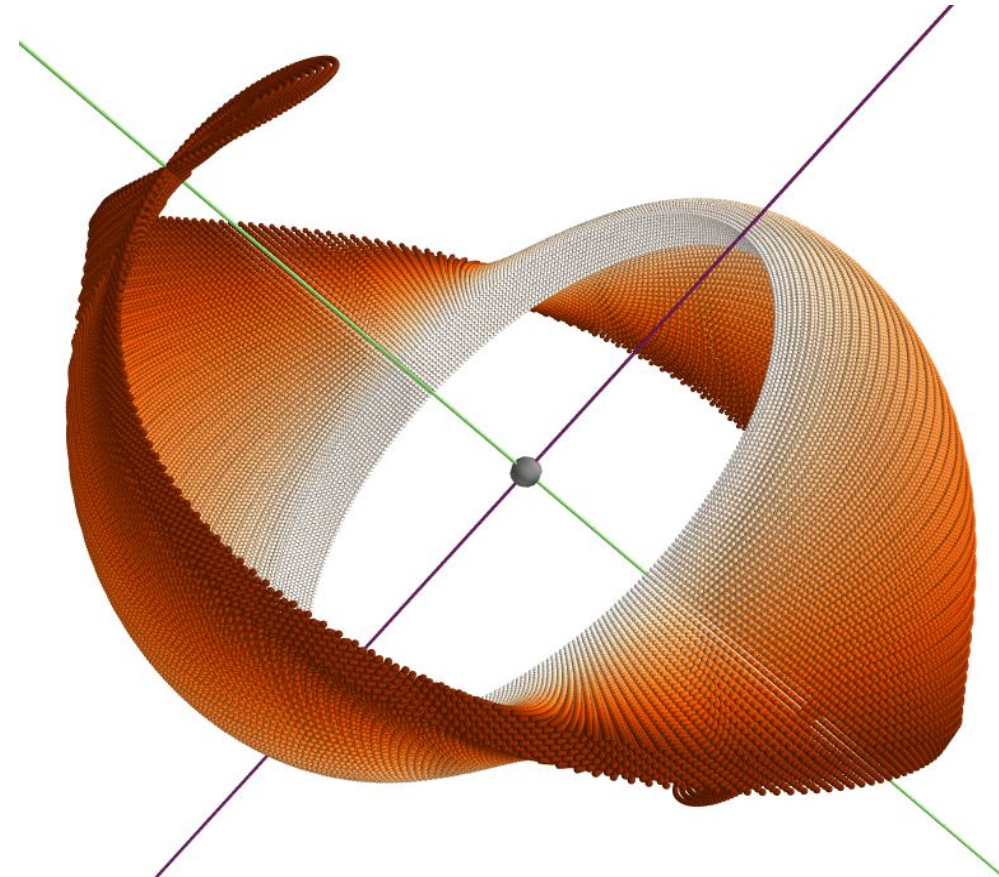
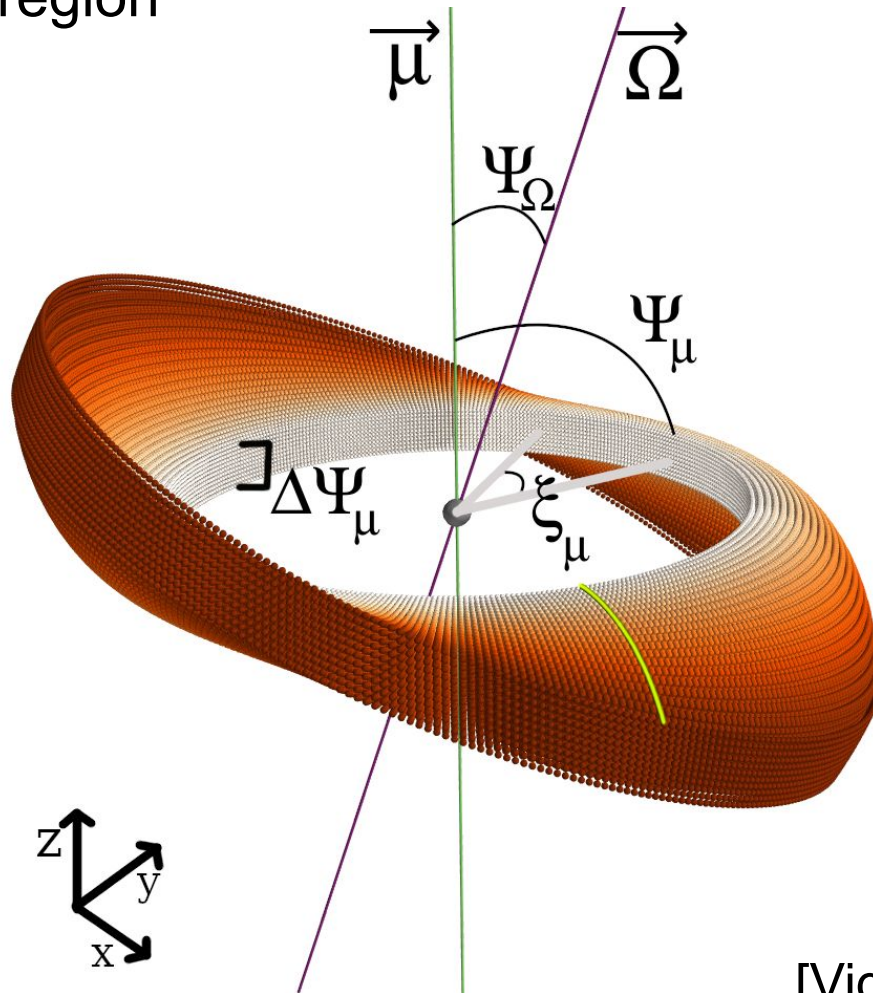
- The model can be extended to resemble the X-ray regime too, getting to do it in a majority of the ~40 high-energy pulsars





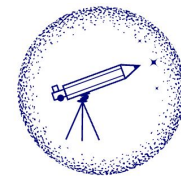
# Geometrical model

- The inclination angle and the meridional extent define the geometry of the emission region

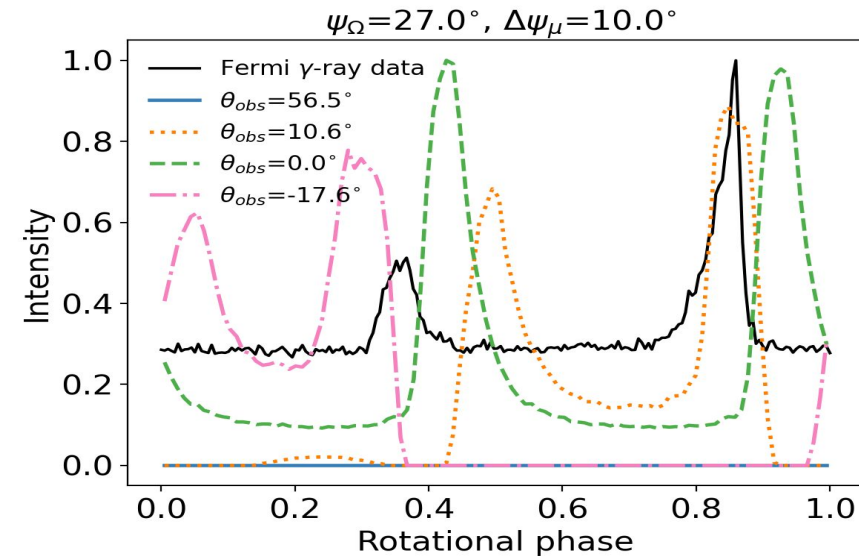
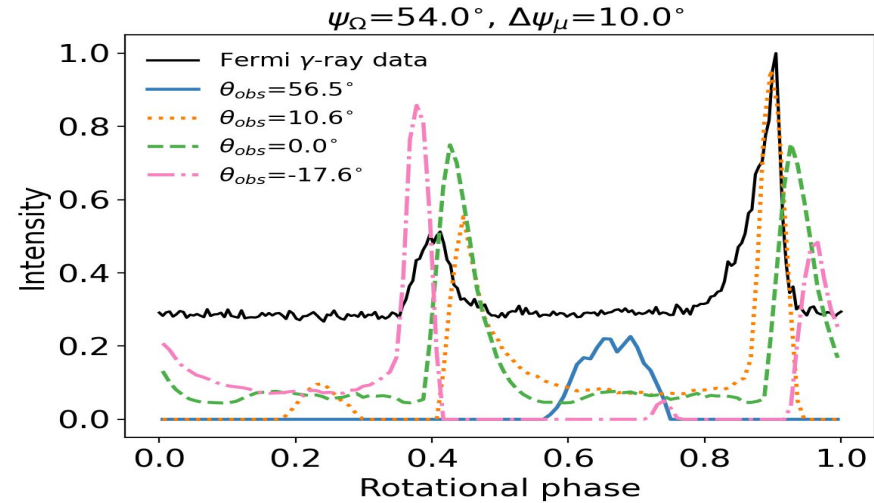
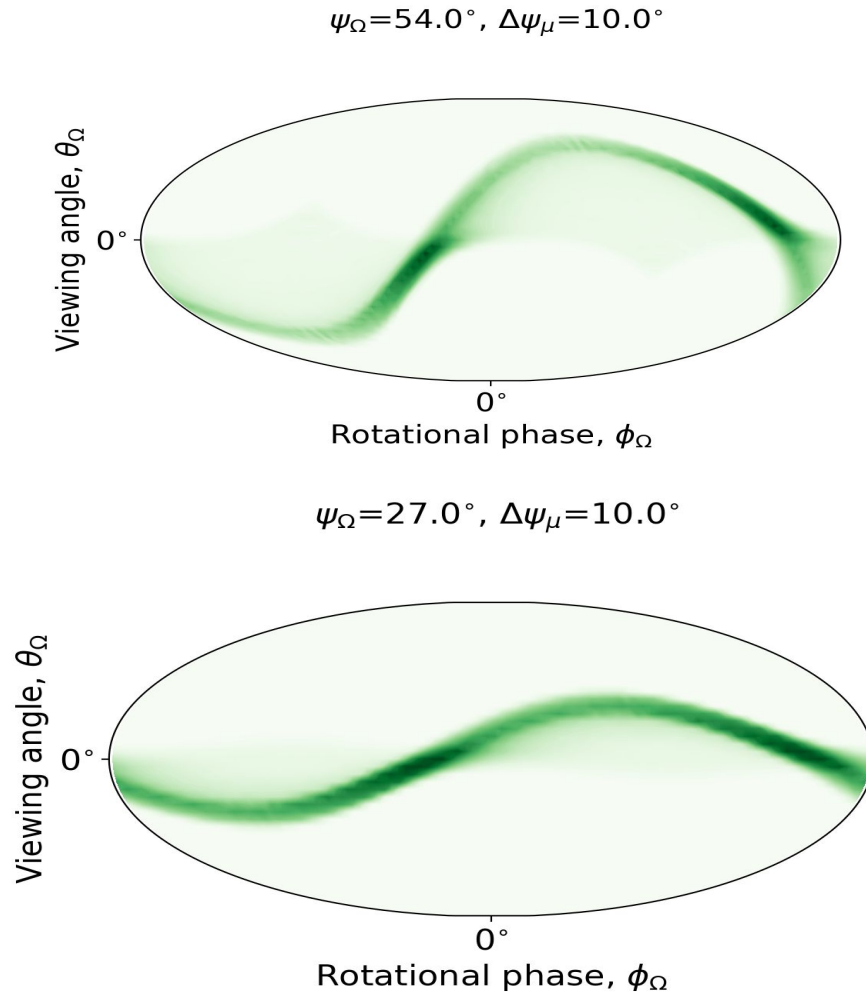


[Viganò & Torres (2019); Íñiguez-Pascual, Torres & Viganò (2024)]

# Emission maps and light curves



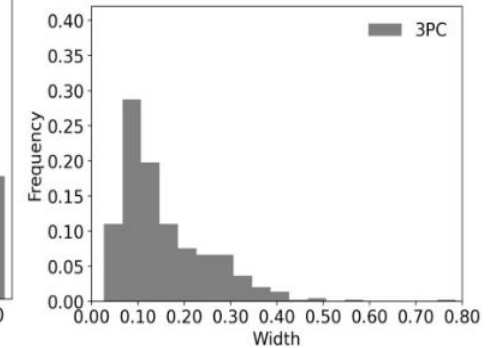
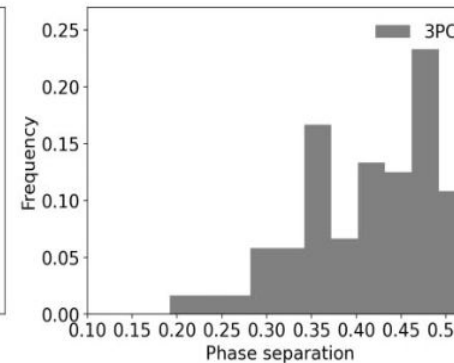
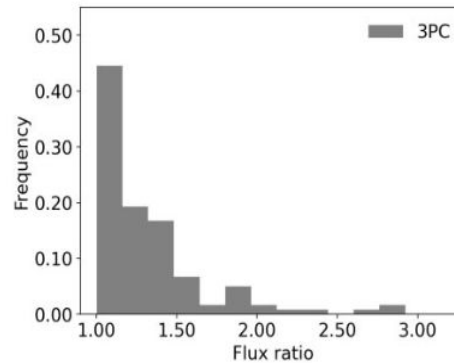
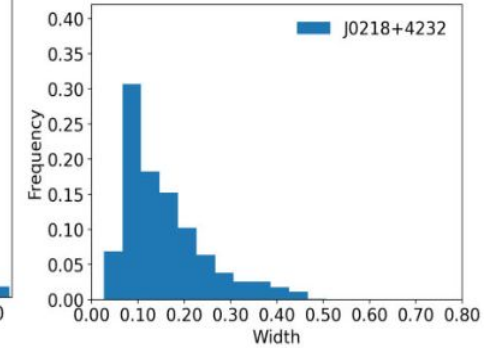
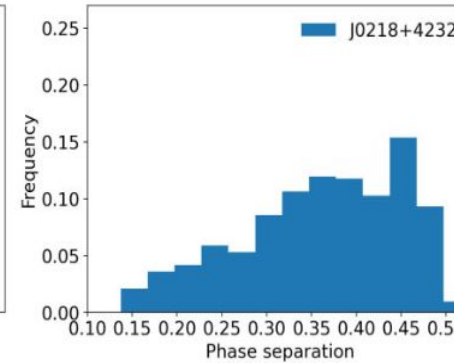
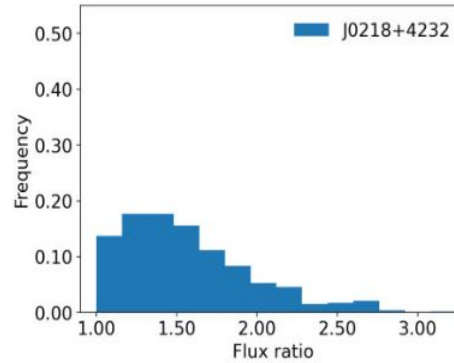
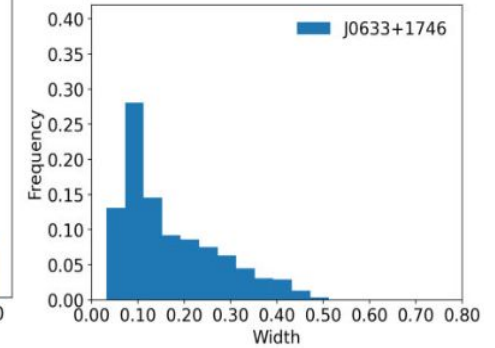
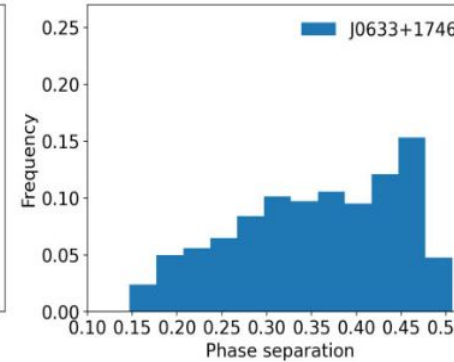
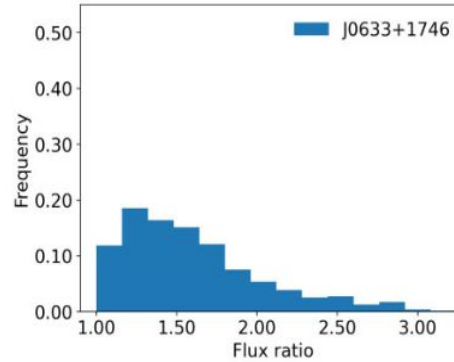
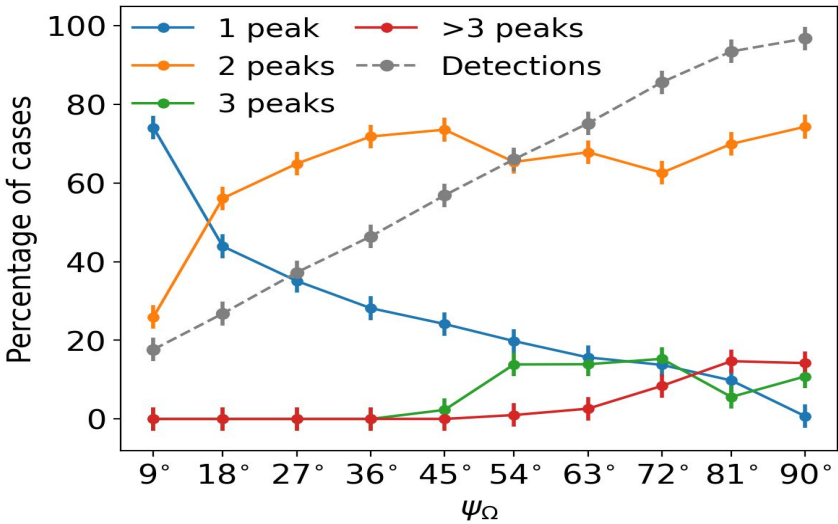
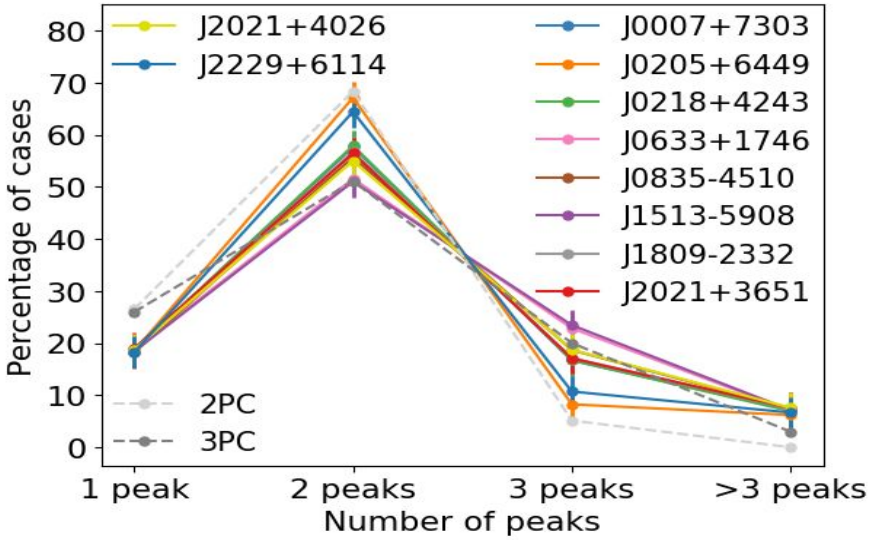
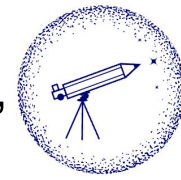
- We can build emission maps (skymaps), from which light curves are obtained



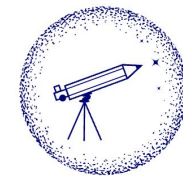
[Íñiguez-Pascual,  
Torres & Viganò  
2024, MNRAS,  
530, 1550]

# Light curves statistics

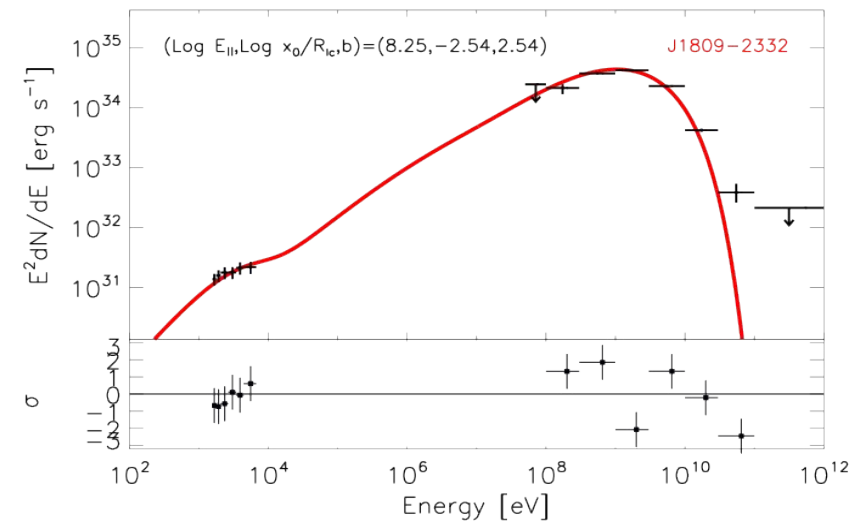
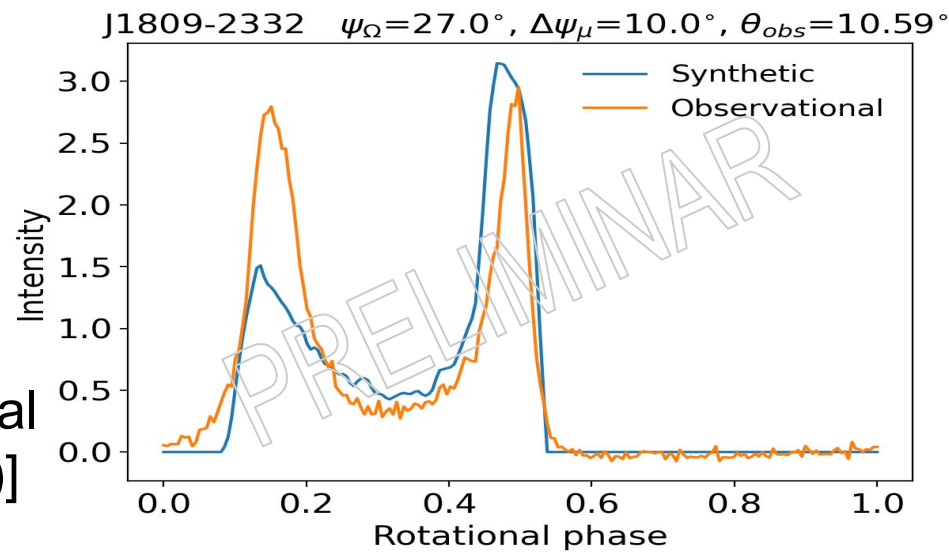
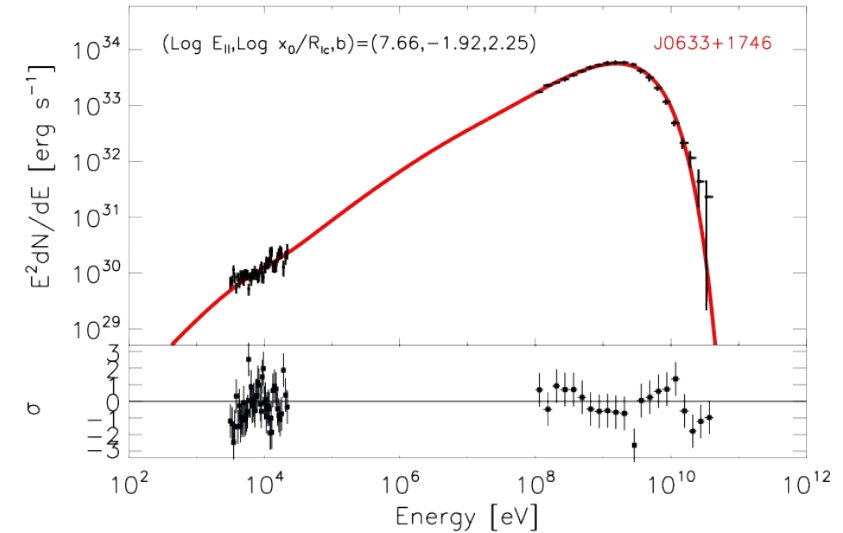
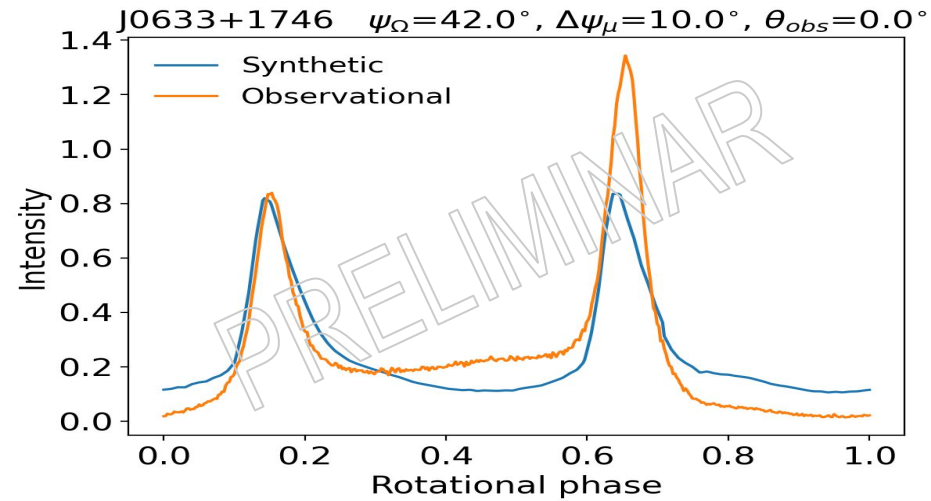
[Íñiguez-Pascual, Torres & Viganò 2024, MNRAS, 530, 1550]



# Light curves fitting



- Fitting synthetic light curves to observational ones, concurrently to the spectral fitting



[Íñiguez-Pascual  
et al. (in prep.)]

# Conclusions

- Our spectral model is able to properly fit the population of high-energy emitting pulsars, showing that synchro-curvature radiation is an appropriate mechanism to explain the emission from these objects

*[Íñiguez-Pascual D., Viganò D., Torres D. F., 2022, MNRAS, 516, 2475 (2208.05549)]*

- The geometrical model reproduces the variety of observational gamma-ray light curves but still cannot capture all their small scale features

*[Íñiguez-Pascual D., Torres D. F., Viganò D., 2024, MNRAS, 530, 1550 (2404.01926)]*

## Future prospects

- Improve the spectral and geometrical models by including more realistic physics while keeping our effective approach
- Possess a fully working concurrent spectral and light curves fitting

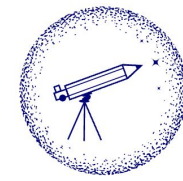
# Thank you



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# Back up slide: particle dynamics formulae



- The equation of motion of charged particles balances electric acceleration and synchro-curvature losses

$$\frac{d\vec{p}}{dt} = ZeE_{\parallel}\hat{b} - (P_{sc}/v)\hat{p}$$

$$\vec{p} = \Gamma m\vec{v}$$

[Viganò et al. (2014)]

- We solve it numerically, considering separately the components parallel and perpendicular to the trajectory

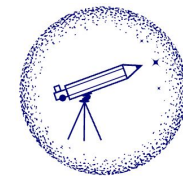
$$\frac{d(p \cos \alpha)}{dt} = ZeE_{\parallel} - \frac{P_{sc}}{v} \cos \alpha$$

$$\frac{d(p \sin \alpha)}{dt} = - \frac{P_{sc}}{v} \sin \alpha$$

- Local magnetic field strength and curvature radius are parametrized in an effective way:

$$B = B_{\star} \left( \frac{R_{\star}}{x} \right)^b \quad r_c = R_{lc} \left( \frac{x}{R_{lc}} \right)^{\eta}$$

# Back up slide: synchro-curvature formulae



- Single-particle synchro-curvature power spectra:

$$\frac{dP_{sc}}{dE} = \frac{\sqrt{3}(Ze)^2\Gamma y}{4\pi\hbar r_{eff}} [(1+z)F(y) - (1-z)K_{2/3}(y)]$$

- Synchro-curvature power:

$$P_{sc} = \frac{2(Ze)^2\Gamma^4 c}{3r_c^2} g_r$$

[Cheng & Zhang (1996),  
Viganò et al. (2014)]

$$r_{eff} = \frac{r_c}{\cos^2 \alpha} \left( 1 + \xi + \frac{r_{gyr}}{r_c} \right)^{-1}$$

$$z = (Q_2 r_{eff})^{-2}$$

$$y = \frac{E}{E_c}$$

$$Q_2^2 = \frac{\cos^4 \alpha}{r_c^2} \left[ 1 + 3\xi + \xi^2 + \frac{r_{gyr}}{r_c} \right]$$

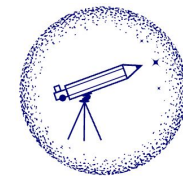
$$E_c = \frac{3}{2} \hbar c Q_2 \Gamma^3$$

$$F(y) = \int_y^\infty K_{5/3}(y') dy'$$

$$\xi = \frac{r_c}{r_{gyr}} \frac{\sin^2 \alpha}{\cos^2 \alpha}$$



# Back up slide: synchro-curvature formulae



- Convoluting the single-particle power spectra with an effective particle distribution,

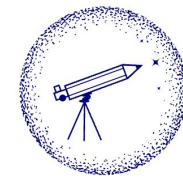
$$\frac{dN}{dx} = N_0 \frac{e^{-(x-x_{min})/x_0}}{x_0(1 - e^{-(x_{max}-x_{min})/x_0})}$$

We obtain the total emission from the region:

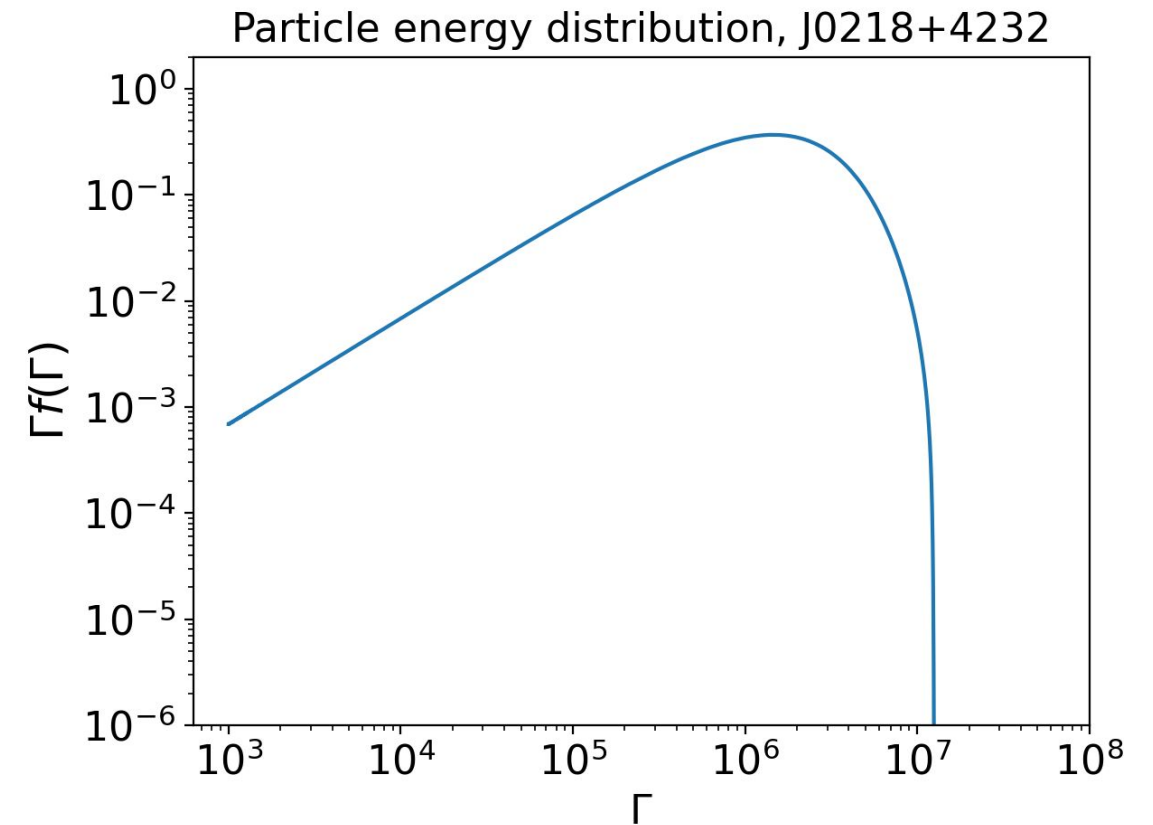
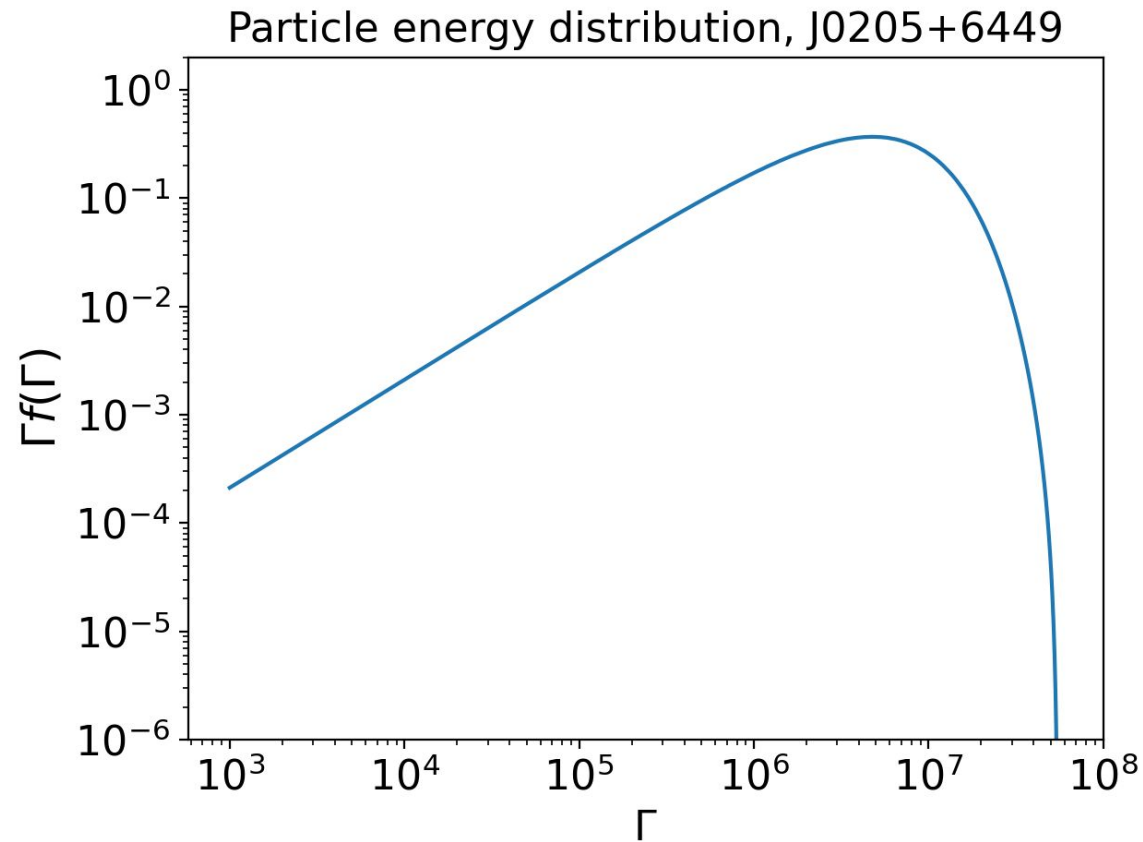
$$\frac{dP_{tot}}{dE} = \int_{x_{min}}^{x_{max}} \frac{dP_{sc}}{dE} \frac{dN}{dx} dx$$

[Cheng & Zhang (1996),  
Viganò et al. (2014)]

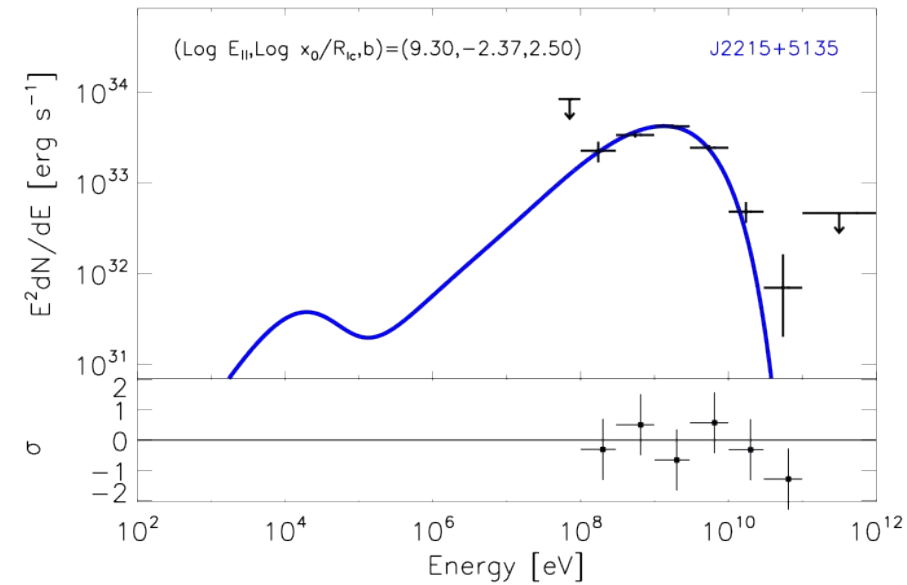
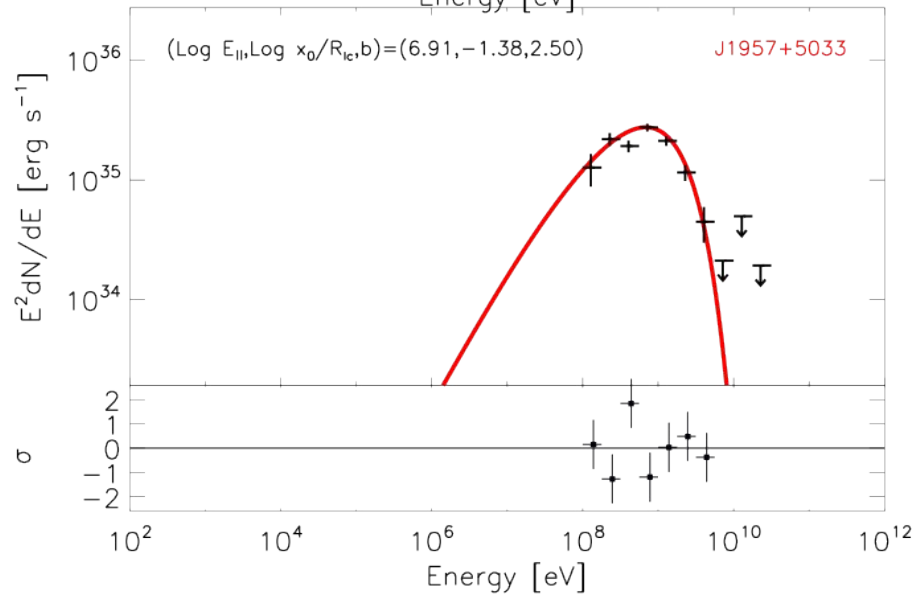
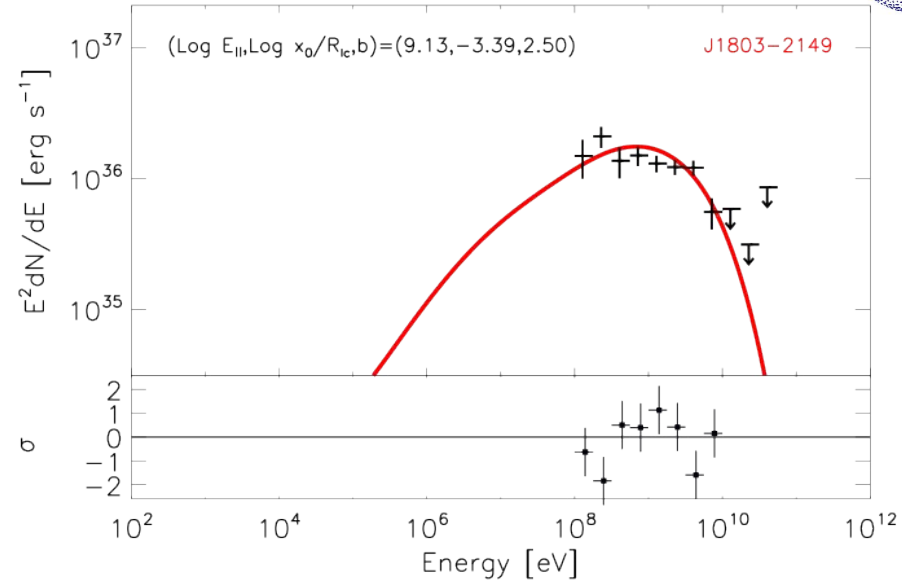
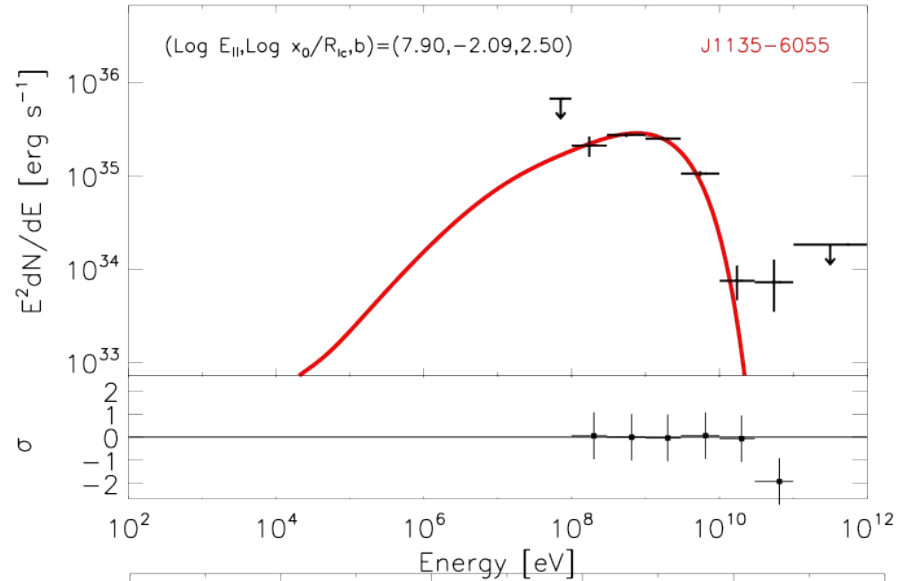
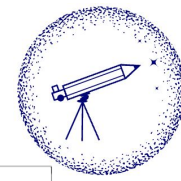
# Back up slide: particle energy distribution



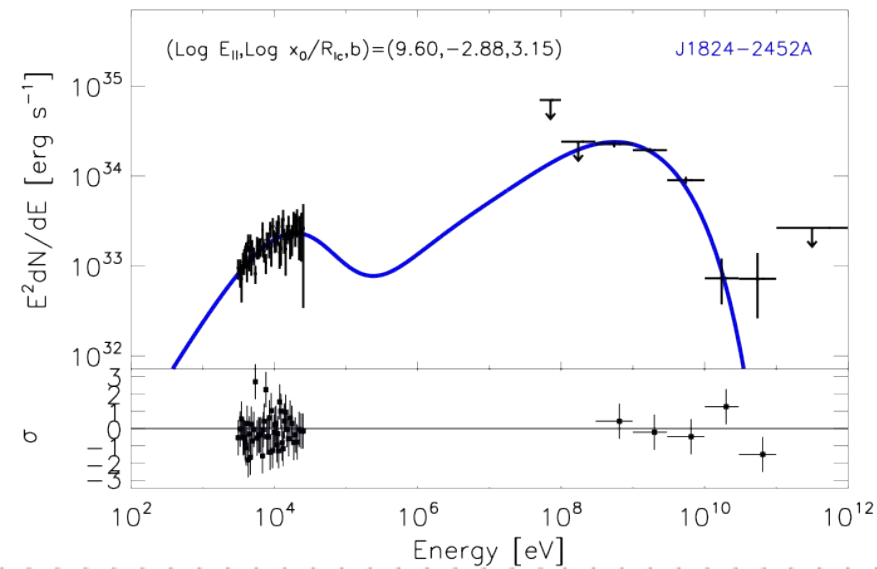
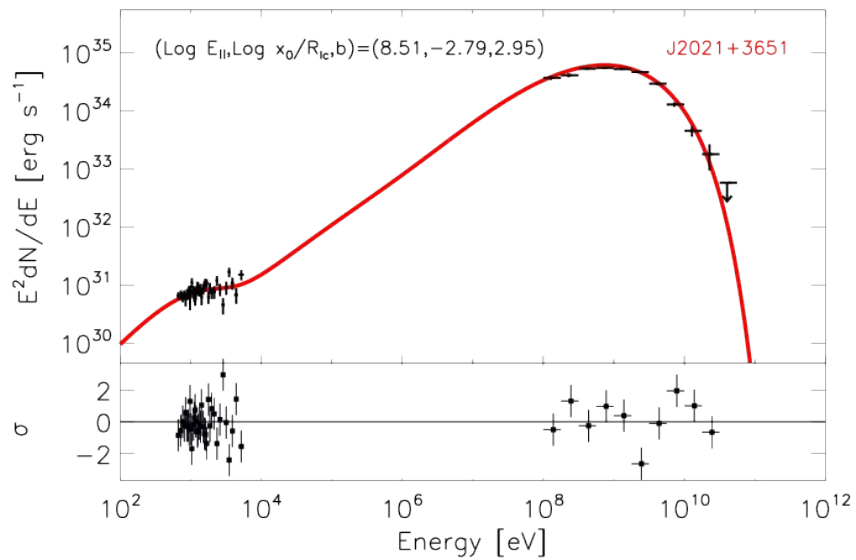
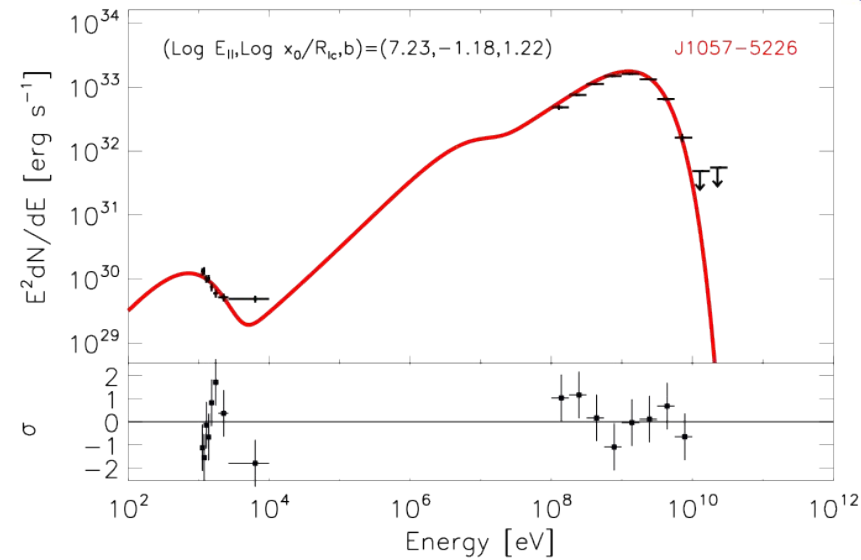
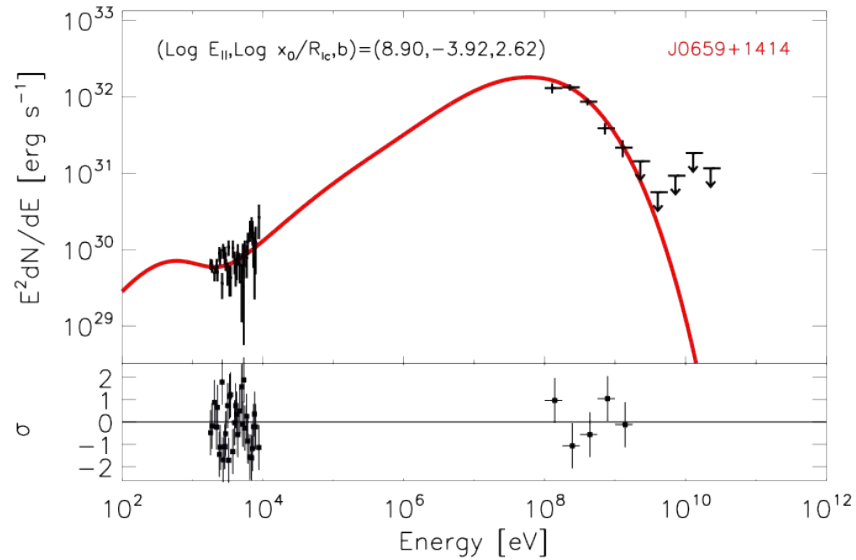
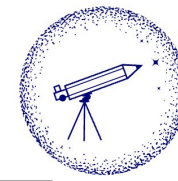
- Lorentz factors  $\Gamma$  typically range from  $10^3$  to  $10^7$ .



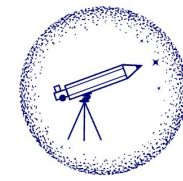
# Back up slide: more spectral fits



# Back up slide: more spectral fits



# Back up slide: geometrical model



- Frenet-Serret equations allow to geometrically describe a curved trajectory with torsion

$$\begin{aligned}\frac{d\hat{t}}{d\lambda} &= \frac{1}{r_c} \hat{n} , \\ \frac{d\hat{n}}{d\lambda} &= -\frac{1}{r_c} \hat{t} + \tau \hat{b} , \\ \frac{d\hat{b}}{d\lambda} &= -\tau \hat{n} .\end{aligned}$$

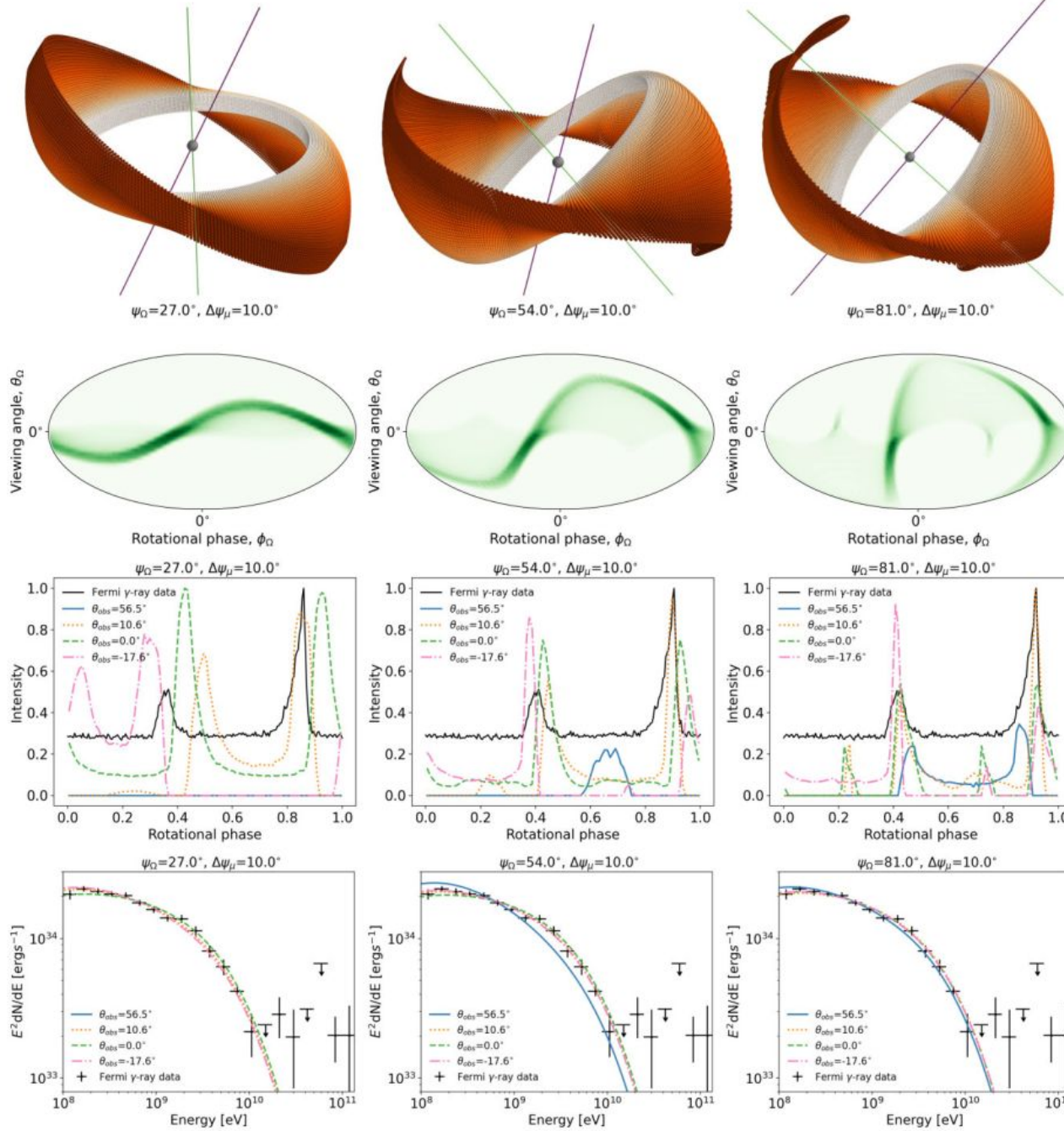
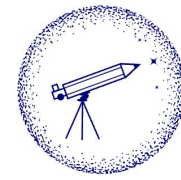
- The emission region is build around a centered value of the magnetic colatitude:

$$\Psi_\mu^c(\xi_\mu, R, \Psi_\Omega) = \frac{\pi}{2} + A(R, \Psi_\Omega) \sin(\xi_\mu - \pi/2)$$

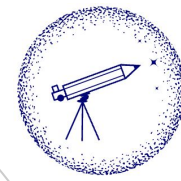
with

$$A(R, \psi_\Omega) = K \Psi_\Omega (R/R_{lc} - R^0/R_{lc})^2$$

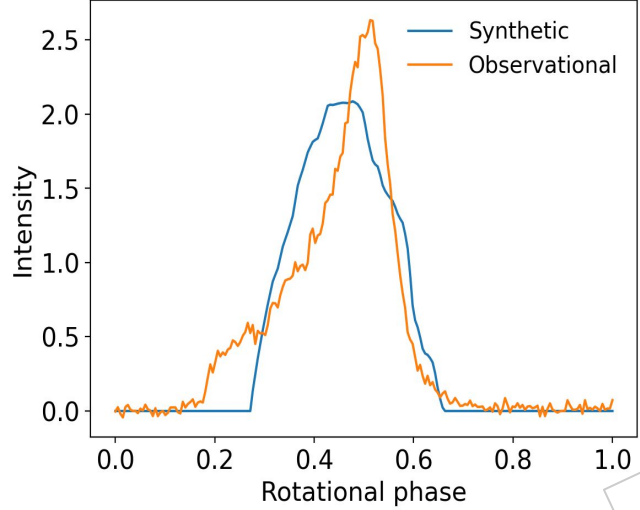
# Back up slide: geometrical model



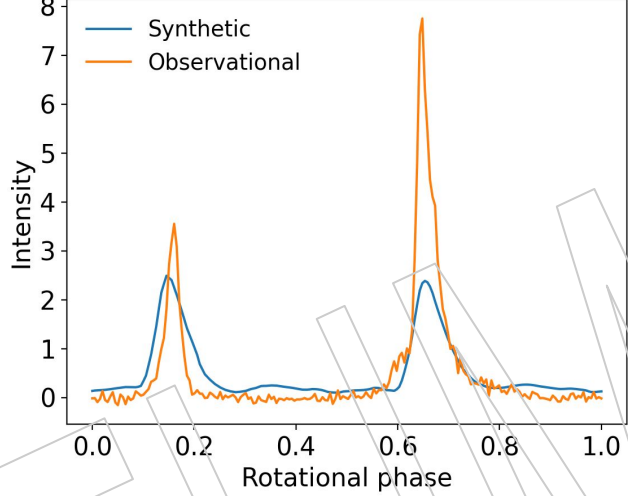
# Back up slide: light curve fits



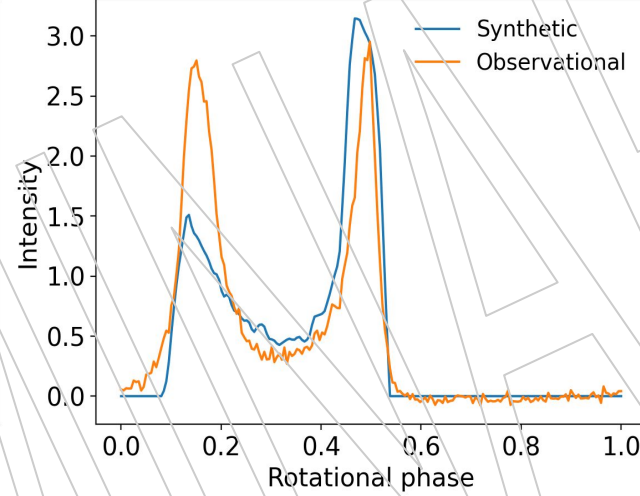
J2229+6114  $\psi_{\Omega}=3.0^{\circ}$ ,  $\Delta\psi_{\mu}=10.0^{\circ}$ ,  $\theta_{obs}=-7.06$



J2032+4127  $\psi_{\Omega}=51.0^{\circ}$ ,  $\Delta\psi_{\mu}=10.0^{\circ}$ ,  $\theta_{obs}=0.0$



J1809-2332  $\psi_{\Omega}=27.0^{\circ}$ ,  $\Delta\psi_{\mu}=10.0^{\circ}$ ,  $\theta_{obs}=10.59^{\circ}$



J2238+5903  $\psi_{\Omega}=57.0^{\circ}$ ,  $\Delta\psi_{\mu}=10.0^{\circ}$ ,  $\theta_{obs}=0.0^{\circ}$

