

$$G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

$$S_B = \frac{k_B 4\pi G}{\hbar c} M^2$$

$$\Psi(x) = \frac{1}{\sqrt{k}} (A_+ e^{ikx} + A_- e^{-ikx}) \quad x < 0$$

$$k = \sqrt{2mE/\hbar^2}$$

$$\sigma = \frac{24\pi^3 L^2}{T^2 c^2 (1-e^2)}$$

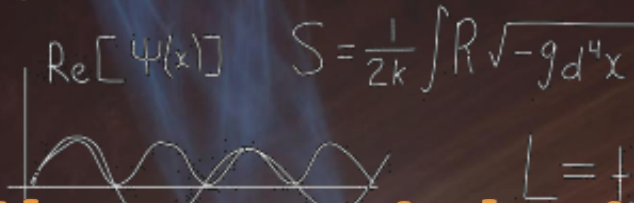
$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

Giorgio GALANTI

INAF - IASF-Milano

$$H = \frac{P^2}{2m} + V(r)$$

$$P = -i\hbar\nabla$$



$$S = \frac{1}{2k} \int R \sqrt{-g} d^4x$$

$$L = \frac{1}{4\pi} \left[\frac{1}{g^2} F_{IJ} F^{IJ} - i\lambda \Gamma^I D_I \lambda \right]$$

Axion-like particle impact in high-energy astrophysics

$$H|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

$$\delta(k_1 + k_2)$$

$$E = mc^2$$

$$E^2 = (pc)^2 + (mc^2)^2$$

$$r = \frac{\theta}{2\pi} + \frac{4\pi}{g^2}$$

$$I = \int e^{-ax^2/2} dx = \sqrt{\frac{2\pi}{a}}$$

$$E^2 = p^2 c^2 + m^2 c^4$$

$$\frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2} - \nabla^2 \psi + \frac{m^2 c^2}{\hbar^2} \psi = 0$$

$$p = \hbar k = \frac{h\nu}{c} = \frac{h}{\lambda}$$

8th Heidelberg International Symposium on High-Energy Gamma-Ray Astronomy

$$A_{ij} = \frac{8\pi h \nu^3}{c^3} B_{ij}$$

Milano, 2 September 2024

$$\frac{d}{dt} \langle A \rangle = \frac{1}{i\hbar} \langle [\hat{A}, \hat{H}] \rangle + \left\langle \frac{\partial \hat{A}}{\partial t} \right\rangle$$

$$i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2} \sum_{n=1}^N \frac{1}{m_n} \nabla_n^2 \psi + V\psi$$

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$dV = e^{-\int_t^s V(X(r)) dr} \left(X, s \right) \frac{\partial u}{\partial X} dW$$

$$\Omega_m = 10$$

Outline

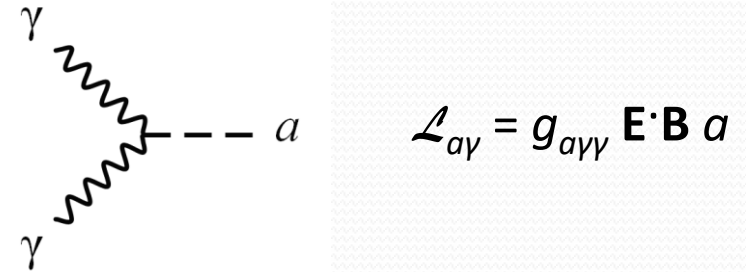
- Axion-like particles
- ALP impact in high-energy astrophysics
- Conclusions

Axion-like particles

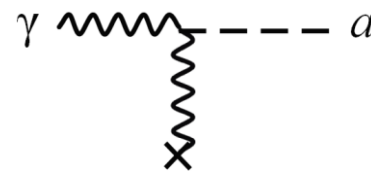
Axion-like particles (ALPs)

- Predicted by **String Theory**
- Very light particles ($m_a < 10^{-8}$ eV)
- Spin 0
- **Interaction with two photons** (coupling $g_{a\gamma\gamma}$)
- Subdominant interactions with other particles
- Possible candidate for dark matter
- Produce **spectral** and **polarization effects**

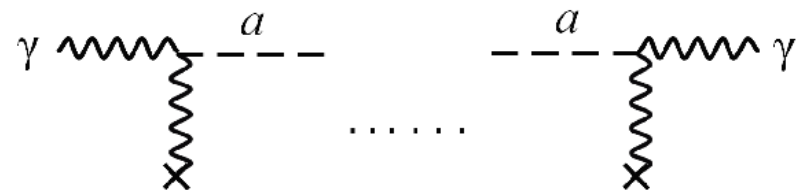
Two photons



In an external B field



Photon-ALP oscillations



ALPs in astrophysical context

- ALPs very **elusive** in laboratory experiments (low coupling) → **astrophysical environment** is the **best opportunity** to study ALPs and ALP effects (*for free*)
- Photon/ALP beam with energy $E \gg m_a$
- ALPs induce **modifications to astrophysical spectra**
 - For $E < 10$ GeV → negligible photon absorption due to EBL
 - **Photon-ALP interaction** produces effective **photon absorption**
 - For $E > 10$ GeV → photons absorbed by EBL ($\gamma\gamma \rightarrow e^+e^-$), **ALPs are not absorbed**
 - **Photon-ALP oscillations increase medium transparency**
- ALPs induce **modifications to photon polarization** (*birefringence, dichroism*)
- **HINTS** at ALP existence:
 - Explain how flat spectrum radio quasars (FSRQs) can emit up to 400 GeV
F. Tavecchio, M. Roncadelli, G. Galanti and G. Bonnoli, Phys. Rev. D, 86, 085036 (2012).
 - Solve the anomalous redshift dependence of blazar spectra
G. Galanti, M. Roncadelli, A. De Angelis and G. F. Bignami, MNRAS 493, 1553 (2020).
 - GRB 221009A
G. Galanti, L. Nava, M. Roncadelli, F. Tavecchio and G. Bonnoli, Phys. Rev. Lett. 131, 251001 (2023).

γ : photon

a : ALP

absorption: $\gamma + \gamma_{\text{Soft}} \rightarrow e^+ + e^-$

γ_{Soft} : EBL, BLR

$g_{a\gamma\gamma}$: $\gamma\gamma a$ coupling

E : γ electric field

B : external magnetic field

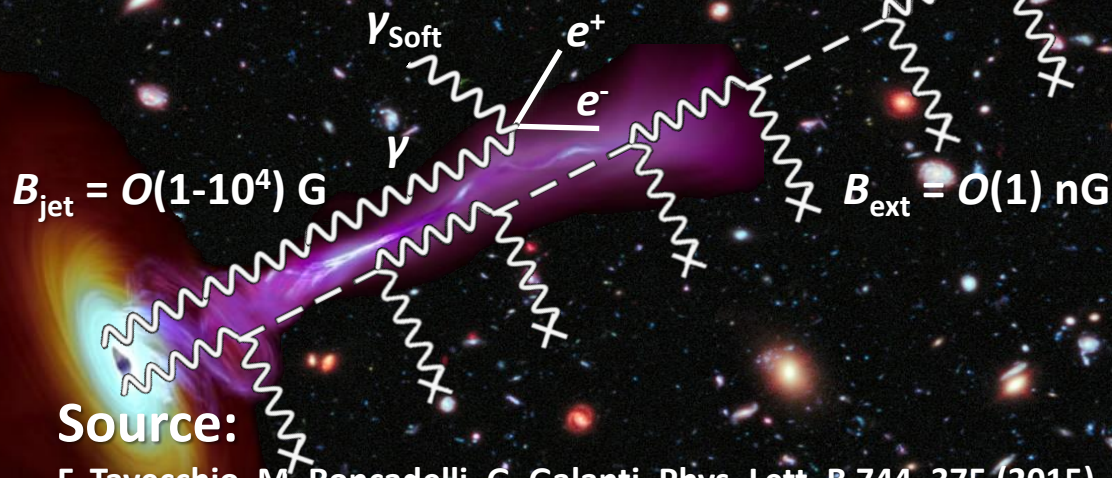
$$\mathcal{L}_{a\gamma} = g_{a\gamma\gamma} \mathbf{E} \cdot \mathbf{B} a$$

$$B_{\text{clu}} = O(10) \mu\text{G}$$

Galaxy cluster:

M. Meyer, D. Montanino, J. Conrad, JCAP 09, 003 (2014).

G. Galanti, M. Roncadelli, F. Tavecchio, E. Costa, Phys. Rev. D 107, 103007 (2023).



Source:

F. Tavecchio, M. Roncadelli, G. Galanti, Phys. Lett. B 744, 375 (2015).

G. Galanti, L. Nava, M. Roncadelli, F. Tavecchio, G. Bonnoli, Phys. Rev. Lett. 131, 251001 (2023).

Milky Way:

D. Horns, L. Maccione, M. Meyer et al., Phys. Rev. D, 86, 075024 (2012).

G. Galanti, F. Tavecchio, M. Roncadelli, C. Evoli, MNRAS 487, 123 (2019).

$$B_{\text{MW}} = O(1) \mu\text{G}$$



Extragalactic space:

A. Mirizzi and D. Montanino, JCAP 12, 004 (2009).

G. Galanti and M. Roncadelli, Phys. Rev. D 98, 043018 (2018).

G. Galanti and M. Roncadelli, JHEAp, 20 1-17 (2018).

γ : photon

a : ALP

absorption: $\gamma + e^-$ **SPECTRA**

γ_{Soft} : EBL, BIR

• **Standard physics**

- photon absorption

• **With ALPs**

- lower photon absorption
- spectral irregularities

$g_{a\gamma\gamma}$: $\gamma\gamma a$ coupling

E : γ electric field

B : external magnetic field

$\alpha_{\text{eff}} = g_{a\gamma\gamma} E \cdot B$

$B_{\text{clu}} = O(10) \mu\text{G}$

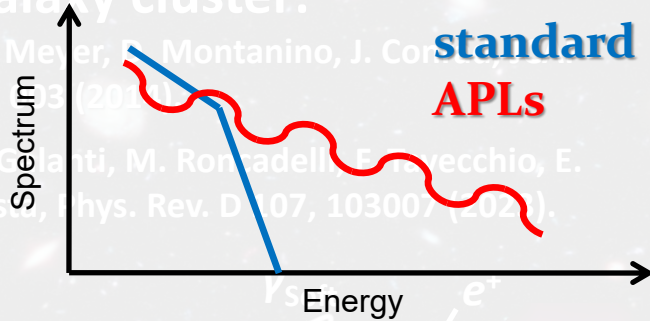
Galaxy cluster:

M. Meyer, D. Montanino, J. Cosmi, **standard**

09, 03 (2021)

APLs

G. Galanti, M. Roncadelli, F. Tavecchio, E. Cosmi, Phys. Rev. D 107, 103007 (2023)



Milky Way:

D. Horns, L. Maccione, M. Meyer et al., Phys. Rev. D, 86, 075024 (2012).

G. Galanti, F. Tavecchio, M. Roncadelli, C. Evoli, MNRAS 487, 123 (2019).

$B_{\text{MW}} = O(1) \mu\text{G}$

POLARIZATION

• **Standard physics**

- unpolarized photon case

• **With ALPs**

- modified polarization
- partially polarized photons

$B_{\text{jet}} = O(1-10^4) \text{G}$

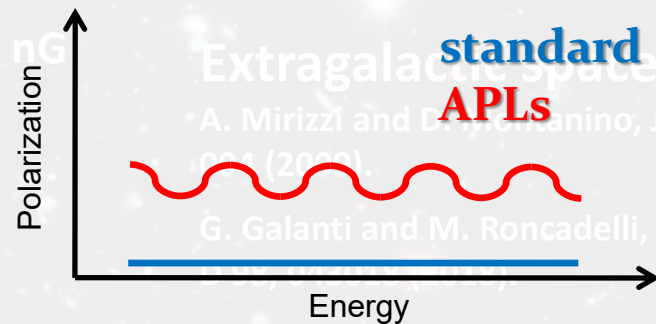
$B_{\text{ext}} = O(1) \text{nG}$

Extragalactic space:

standard
APLs

A. Mirizzi and D. Montanino, JCAP 12, 031 (2019).

G. Galanti and M. Roncadelli, Phys. Rev.



G. Galanti and M. Roncadelli, JHEAp, 20 1-17 (2018).

Source:

F. Tavecchio, M. Roncadelli, G. Galanti, Phys. Lett. B 744, 375 (2015).

G. Galanti, L. Nava, M. Roncadelli, F. Tavecchio, G. Bonnoli, Phys. Rev. Lett. 131, 251001 (2023).

ALP impact in high-energy astrophysics

Active Galactic Nuclei (AGN)

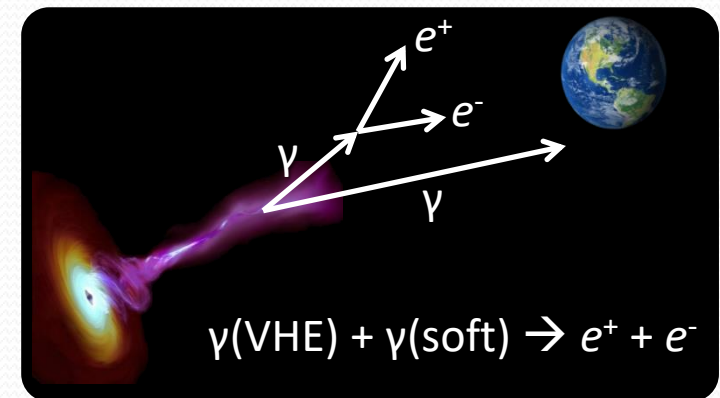
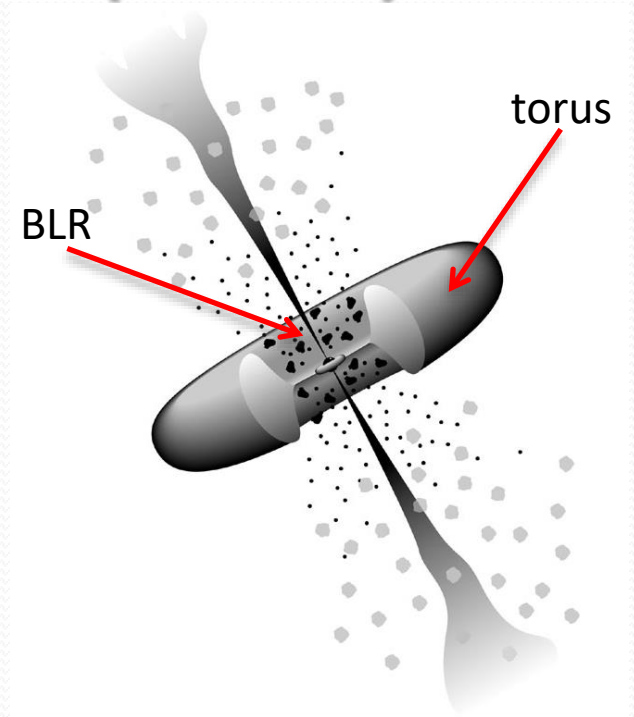
- Super massive black holes ($10^6 - 10^9 M_{\odot}$)
- Accretion disk
- Collimated jets (if towards us -> **blazars**)
- Photons produced at the jet base

BL Lacs:

- No broad line region (BLR)
- No dusty torus
- Absorption due to the extragalactic background light (EBL) for $E > 100$ GeV

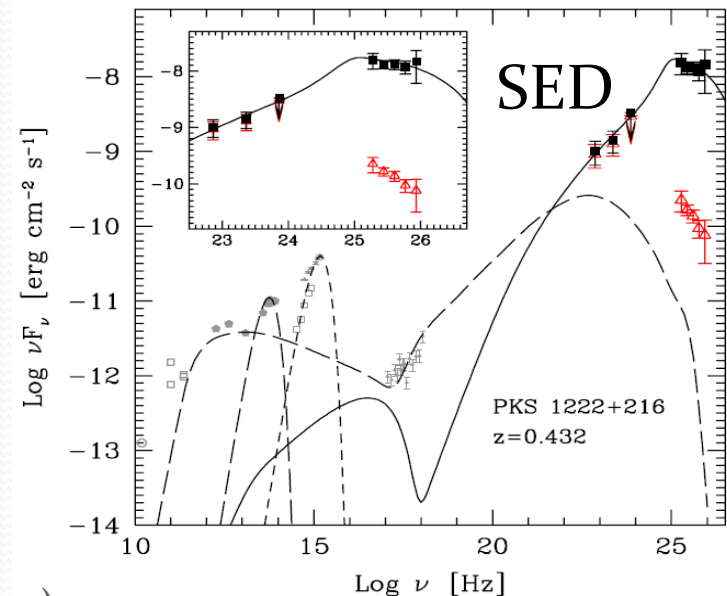
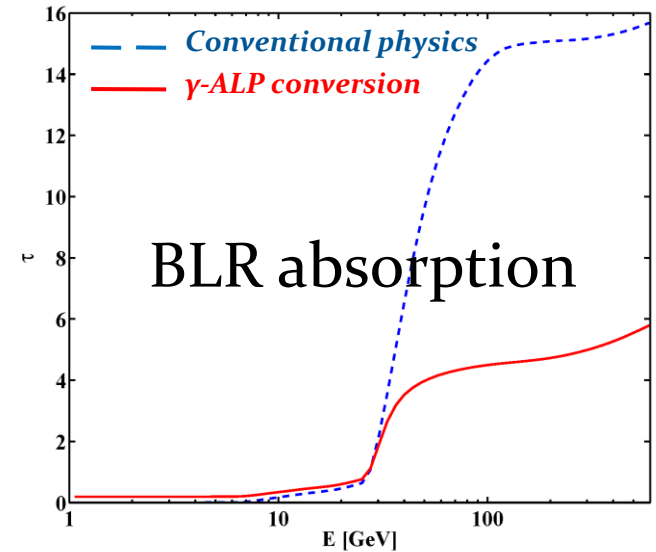
Flat spectrum radio quasars (FSRQs):

- Absorption due to the BLR for $E > 20$ GeV
- Absorption due to the dusty torus for $E > 300$ GeV
- Absorption due to the EBL for $E > 100$ GeV



ALPs in FSRQs

- **High BLR absorption** → no photons with $E > 20$ GeV predicted **BUT**
- **Photons observed up to 400 GeV**
- **Why?** Photon/ALP conversions?
 - $B_{\text{jet}} = 0.2$ G and scales as $1/\gamma$
 - $g_{a\gamma\gamma} = O(10^{-11})$ GeV $^{-1}$, $m_a < O(10^{-10})$ eV
 - BLR $n_{e,\text{BLR}} = 10^{10}$ cm $^{-3}$
- Photon-ALP **conversion** before the BLR
– **reconversion** outside the BLR
- -> BLR absorption **REDUCED**
- Physically motivated flux (SED)
- **First hint** at **ALP existence**



Anomalous z dependence of Blazars

- We consider all BL Lacs (HBL and IBL) with strong VHE spectrum:
 - In flare
 - $E > 100$ GeV
 - redshift up to $z = 0.6$

- Emitted spectra \rightarrow **power law**

$$\Phi_{\text{em}}(E) = \hat{K}_{\text{em}} E^{-\Gamma_{\text{em}}}$$

- Observed spectrum \rightarrow **power law**

$$\Phi_{\text{obs}}(E_0, z) = \hat{K}_{\text{obs}}(z) E_0^{-\Gamma_{\text{obs}}(z)}$$

- Emitted – observed spectrum relation

$$\Phi_{\text{obs}}(E_0, z) = P_{\gamma \rightarrow \gamma}(E_0, z) \Phi_{\text{em}}(E_0(1+z))$$

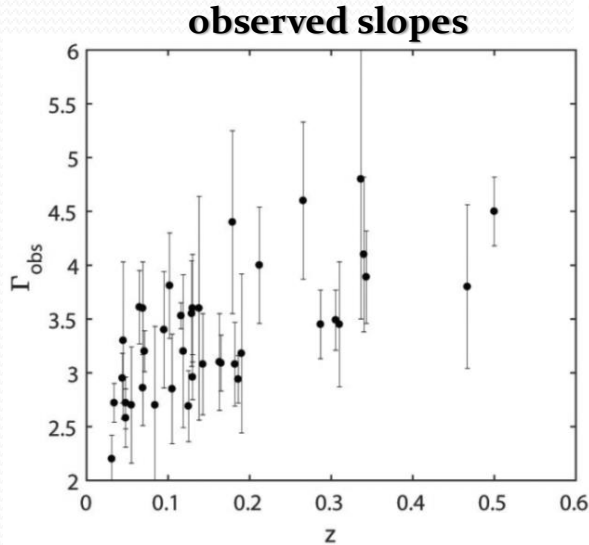
- We **deabsorb** the **observed spectrum**:

- if no ALPs \rightarrow EBL absorption only
- **with ALPs** \rightarrow EBL absorption and photon-ALP oscillations

Anomalous z dependence of Blazars (2)

Conventional Physics (CP):

- Anomalous redshift dependence of blazar spectra

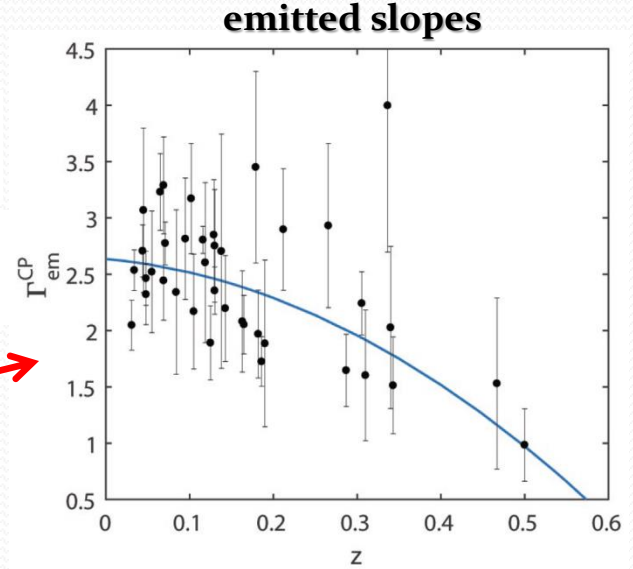


$$\Phi_{\text{em}}^{\text{CP}}(E_0(1+z)) = e^{\gamma_{\text{FR}}(E_0, z)} K_{\text{obs}}(z) \left(\frac{E_0}{E_{0,*}} \right)^{-\Gamma_{\text{obs}}(z)}$$

$$\Phi_{\text{em}}^{\text{CP,BF}}(E_0(1+z)) = K_{\text{em}}^{\text{CP}}(z) \left(\frac{E_0(1+z)}{E_{0,*}} \right)^{-\Gamma_{\text{em}}^{\text{CP}}(z)}$$

CP

ALP

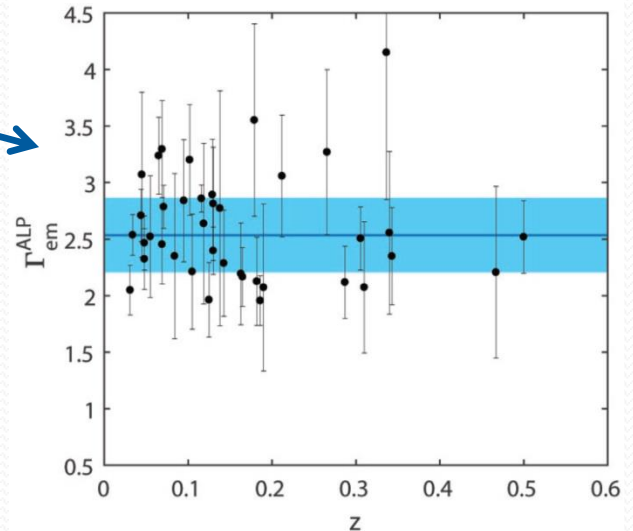


With ALPs:

- Anomaly **SOLVED**

$$\Phi_{\text{em}}^{\text{ALP}}(E_0(1+z)) = \left(P_{\gamma \rightarrow \gamma}^{\text{ALP}}(E_0, z) \right)^{-1} K_{\text{obs}}(z) \left(\frac{E_0}{E_{0,*}} \right)^{-\Gamma_{\text{obs}}(z)}$$

$$\Phi_{\text{em}}^{\text{ALP,BF}}(E_0(1+z)) = K_{\text{em}}^{\text{ALP}}(z) \left(\frac{E_0(1+z)}{E_{0,*}} \right)^{-\Gamma_{\text{em}}^{\text{ALP}}(z)}$$



Second hint at ALP existence (a new study gets similar results, see Dong+2023)

GRB 221009A

- Extremely luminous Gamma Ray Burst (GRB) at redshift $z = 0.151$
- Observed by:
 - Fermi-GBM, Fermi-LAT (Fermi 2023), Swift (Williams+2023)
 - **LHAASO** at $E \simeq (13-18)$ TeV within 2000 s after the initial burst (LHAASO 2022, 2023a,b)

BUT strong EBL absorption for $E \gtrsim 10$ TeV at $z = 0.151$ in Conventional Physics (CP)

EBL	15 TeV		18 TeV		100 TeV		251 TeV	
	τ_{CP}	P_{CP}	τ_{CP}	P_{CP}	τ_{CP}	P_{CP}	τ_{CP}	P_{CP}
D	12.7	3×10^{-6}	19.4	4×10^{-9}	350	2×10^{-152}	9654	~ 0
G	9.4	8×10^{-5}	13.1	2×10^{-6}	246	2×10^{-107}	9502	~ 0
FR	10.1	4×10^{-5}	14.1	7×10^{-7}	333	2×10^{-145}	15411	~ 0
SL	12.8	3×10^{-6}	18.3	10^{-8}	220	3×10^{-96}	>9251	~ 0

τ_{CP} -> optical depth; P_{CP} -> photon survival probability

D -> EBL model by Domínguez et al., 2011

G -> EBL model by Gilmore et al. 2012

FR -> EBL model by Franceschini & Rodighiero 2017

SL -> EBL model by Saldana-Lopez et al. 2021

QUESTION:

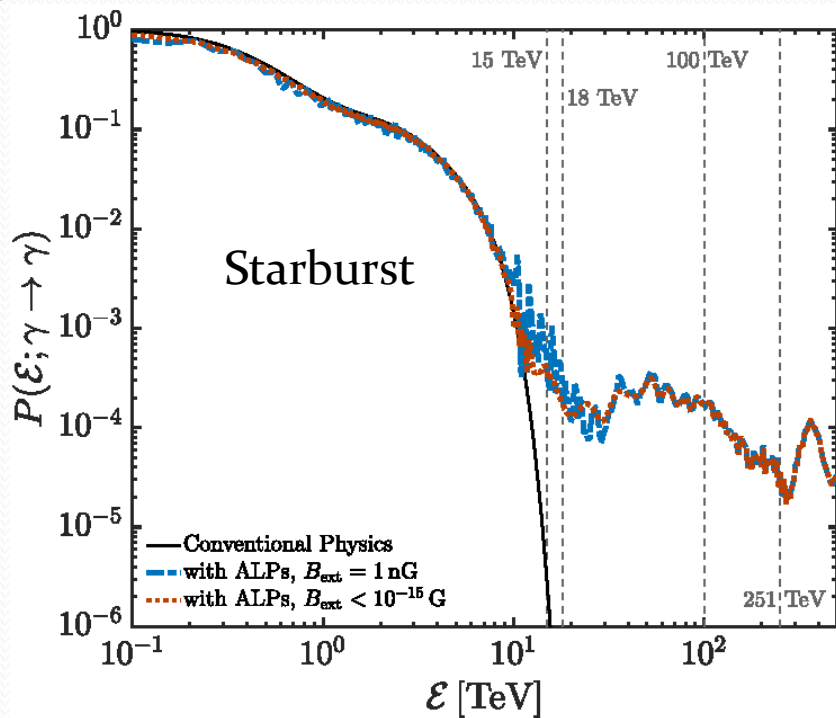
How can we have detected this GRB up to $E \simeq (13-18)$ TeV?



ANSWER:

with **axion-like particles (ALPs) !!!**

Photon survival probability



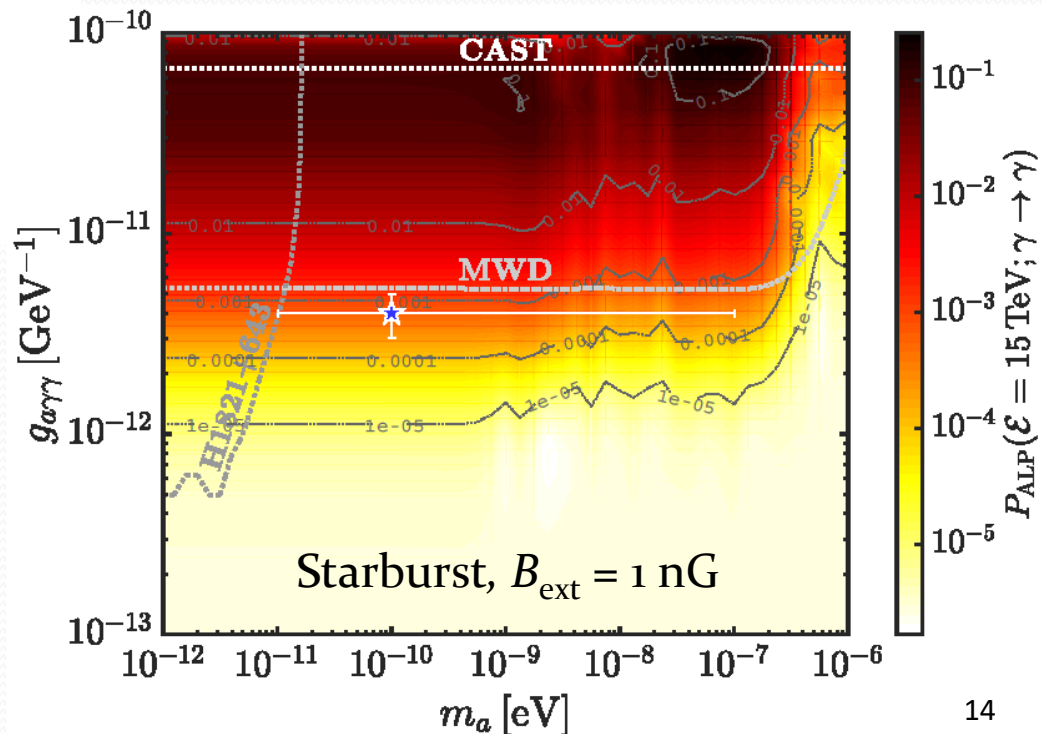
- We take $g_{a\gamma\gamma} = (3 - 5) \times 10^{-12} \text{ GeV}^{-1}$;
 $m_a = (10^{-11} - 10^{-7}) \text{ eV}$

- Within all most stringent ALP bounds (Sisk-Reynés+2022; Dessert+2022)

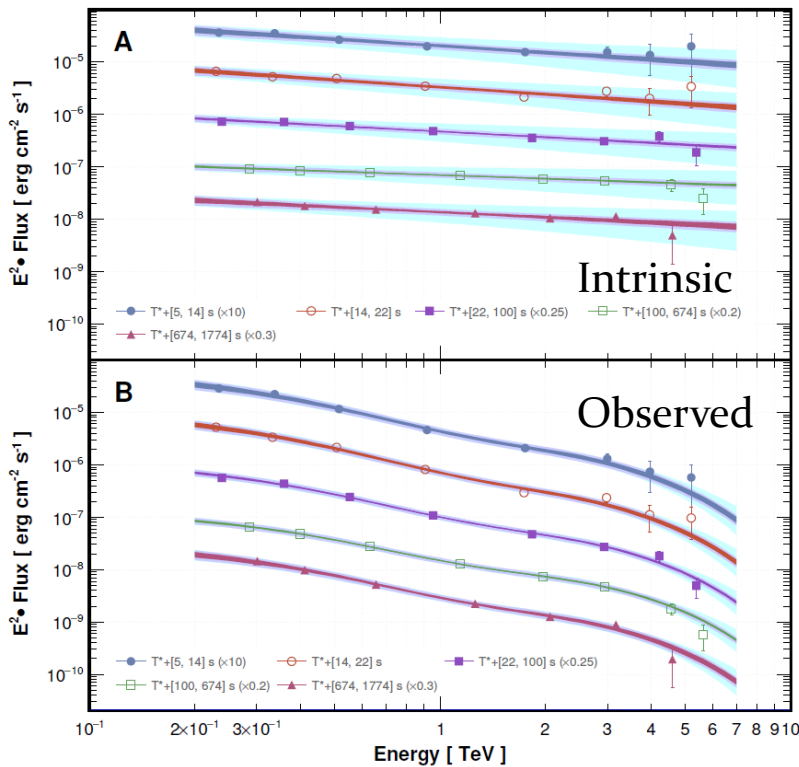
- Explain LHAASO**

- Compatible with other ALP hints**

- $P_{\text{CP}}(E, \gamma \rightarrow \gamma)$:
 - @ 15 TeV $\rightarrow \sim 3 \times 10^{-6}$
 - @ 18 TeV $\rightarrow \sim 1 \times 10^{-8}$
 - @ 100 TeV $\rightarrow \sim 3 \times 10^{-96}$
 - @ 251 TeV $\rightarrow \sim 0$
- $P_{\text{ALP}}(E, \gamma \rightarrow \gamma)$:
 - @ 15 TeV $\rightarrow \sim 6 \times 10^{-4}$
 - @ 18 TeV $\rightarrow \sim 3 \times 10^{-4}$
 - @ 100 TeV $\rightarrow \sim 2 \times 10^{-4}$
 - @ 251 TeV $\rightarrow \sim 3 \times 10^{-5}$



GRB 221009A spectrum

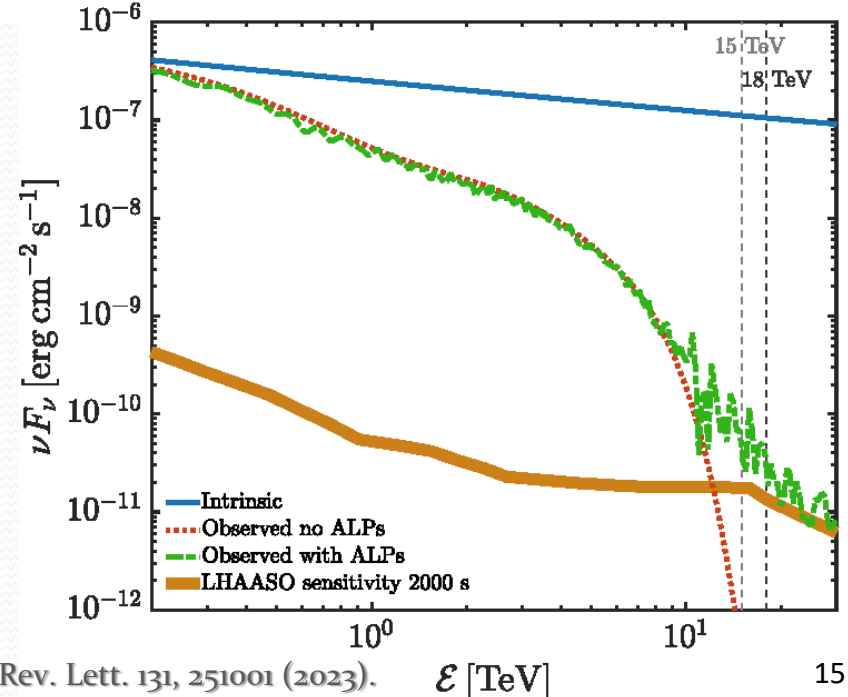


Adapted from LHAASO Collaboration, Science 380, 1390 (2023).

- **LHAASO spectrum** (only **WCDA**):
 - $E = (0.2 - 7) \text{ TeV}$
 - In 5 time intervals Δt_i
- **Intrinsic** spectrum \rightarrow power law with **no cutoff**:

$$\mathcal{F}_i(\mathcal{E}) \equiv \frac{dN}{d\mathcal{E}} = A \left(\frac{\mathcal{E}}{\text{TeV}} \right)^{-\gamma}$$
- **Average** intrinsic spectrum

$$\mathcal{F}_{\text{int}} \equiv \langle \mathcal{F}_i \rangle = \frac{1}{T} \sum_i^5 \Delta t_i \mathcal{F}_i \quad T \rightarrow \text{total duration}$$



- **Observed** spectrum \rightarrow extended to $\sim 20 \text{ TeV}$:

$$\mathcal{F}_{\text{obs}}(\mathcal{E}) = P(\mathcal{E}; \gamma \rightarrow \gamma) \mathcal{F}_{\text{int}}(\mathcal{E}) \quad \text{with } P = P_{\text{CP}} \text{ or } P_{\text{ALP}}$$

$$\nu F_\nu = \mathcal{E}^2 \mathcal{F}_{\text{obs}}$$
- **CP does NOT work** for E above $\sim 10 \text{ TeV}$
- **ALPs solve the problem** \rightarrow **Third hint**

Conclusions

DO ALPs EXIST?

- **Three hints** from astrophysical spectra with **same ALP parameters**
 - **Two indications** from **blazars**
 - The most recent and **strongest one** from **GRB 221009A**
- Additional hints from **photon polarization** expected
 - Photons diffusely emitted by galaxy clusters are **unpolarized** within standard physics
 - ALPs make them **partially polarized** (see Galanti+2023): if detected -> **ALP hint**

FINAL ANSWER:

- Within few years
- **Confirmed** or **disproved**:
 - From new spectral data by ASTRI, CTA, LHAASO
 - Possible polarization data from IXPE, NGXP, COSI, AMEGO
 - From laboratory experiments such as ALPSII, IAXO

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

$$S_B = \frac{k_B 4\pi G}{\hbar c} M^2$$

$$\psi(x) = \frac{1}{\sqrt{k}} (A_+ e^{ikx} + A_- e^{-ikx}) \quad x < 0$$

$$k = \sqrt{2mE/\hbar^2}$$

$$\sigma = \frac{24\pi^3 L^2}{T^2 c^2 (1-e^2)}$$

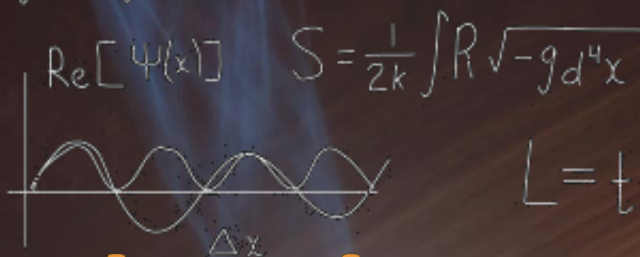
$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$



$$S = \frac{c^2 k A}{4 \hbar G}$$

$$H = \frac{p p}{2m} + V(r)$$

$$p = -i\hbar \nabla$$



$$L = \text{tr} \left[\frac{1}{g} F_{\mu\nu} F^{\mu\nu} - i\lambda \Gamma^i D_i \lambda \right]$$

$$H|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

Thank you

$$r = \frac{\theta}{2\pi} + \frac{4\pi}{g^2}$$

$$I = \int e^{-ax^2/2} dx = \sqrt{\frac{2\pi}{a}}$$

$$E^2 = p^2 c^2 + m^2 c^4$$

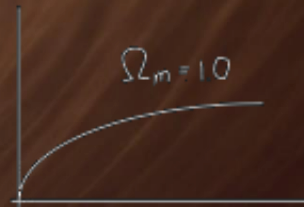
$$\frac{1}{c^2} \frac{\partial^2}{\partial t^2} \psi - \nabla^2 \psi + \frac{m^2 c^2}{\hbar^2} \psi = 0$$

$$p = \hbar k = \frac{h\nu}{c} = \frac{h}{\lambda}$$

$$A_{ij} = \frac{8\pi \hbar \nu^3}{c^3} B_{ij}$$

$$S_{fi} = \langle f | S | i \rangle$$

$$S = \frac{1}{2} \int d^4x \left(R + \frac{R^2}{6M^2} \right)$$



$$d\psi = e^{-\int_t^s V(X(r)) dr} (X, s) \frac{\partial u}{\partial X} dW$$

$$\frac{d}{dt} \langle A \rangle = \frac{1}{i\hbar} \langle [\hat{A}, \hat{H}] \rangle + \left\langle \frac{\partial \hat{A}}{\partial t} \right\rangle$$

$$i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2} \sum_{n=1}^N \frac{1}{m_n} \nabla_n^2 \psi + V\psi$$

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

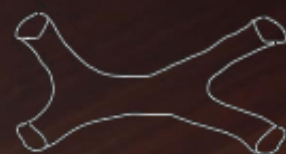
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$$k = \sqrt{2mE/\hbar^2}$$

$$\sigma = \frac{24\pi^3 L^2}{T^2 c^2 (1-e^2)}$$

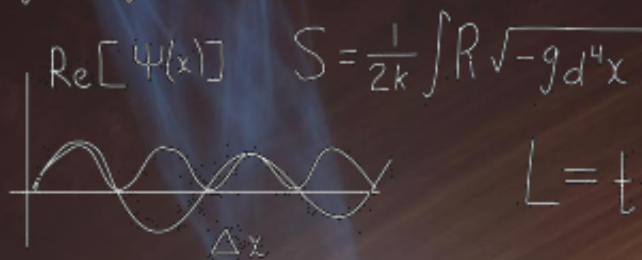
$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$



$$S = \frac{c^2 k A}{4 \hbar G}$$

$$H = \frac{P P}{2m} + V(r)$$

$$P = -i\hbar \nabla$$



$$S = \frac{1}{2k} \int R \sqrt{-g} d^4x$$

$$L = \text{tr} \left[\frac{1}{g} F_{\mu\nu} F^{\mu\nu} - i\lambda \Gamma^i D_i \lambda \right]$$

$$H|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

$$\frac{\delta(k_+ k_-)}{k^2}$$

$$E = mc^2$$

$$E^2 = (pc)^2 + (mc^2)^2$$

$$r = \frac{\theta}{2\pi} + \frac{4\pi}{g^2}$$

$$I = \int e^{-ax^2/2} dx = \sqrt{\frac{2\pi}{a}}$$

$$E^2 = p^2 c^2 + m^2 c^4$$

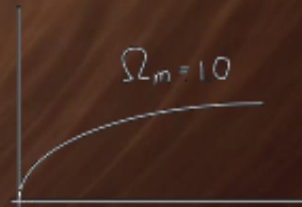
$$\frac{1}{c^2} \frac{\partial^2}{\partial t^2} \psi - \nabla^2 \psi + \frac{m^2 c^2}{\hbar^2} \psi = 0$$

$$p = \hbar k = \frac{h\nu}{c} = \frac{h}{\lambda}$$

$$A_{ij} = \frac{8\pi \hbar \nu^3}{c^3} B_{ij}$$

$$S_{fi} = \langle f | S | i \rangle$$

$$S = \frac{1}{2} \int d^4x \left(R + \frac{R^2}{6M^2} \right)$$



$$d\psi = e^{-\int_t^s V(X(r)) dr} \circ (X, s) \frac{\partial u}{\partial X} dW$$

$$\frac{d}{dt} \langle A \rangle = \frac{1}{i\hbar} \langle [\hat{A}, \hat{H}] \rangle + \left\langle \frac{\partial \hat{A}}{\partial t} \right\rangle$$

$$i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2} \sum_{n=1}^N \frac{1}{m_n} \nabla_n^2 \psi + V\psi$$

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$