

**News from theoretical AGB models:
are we going in the right direction?**

S. Cristallo & D. Vescovi



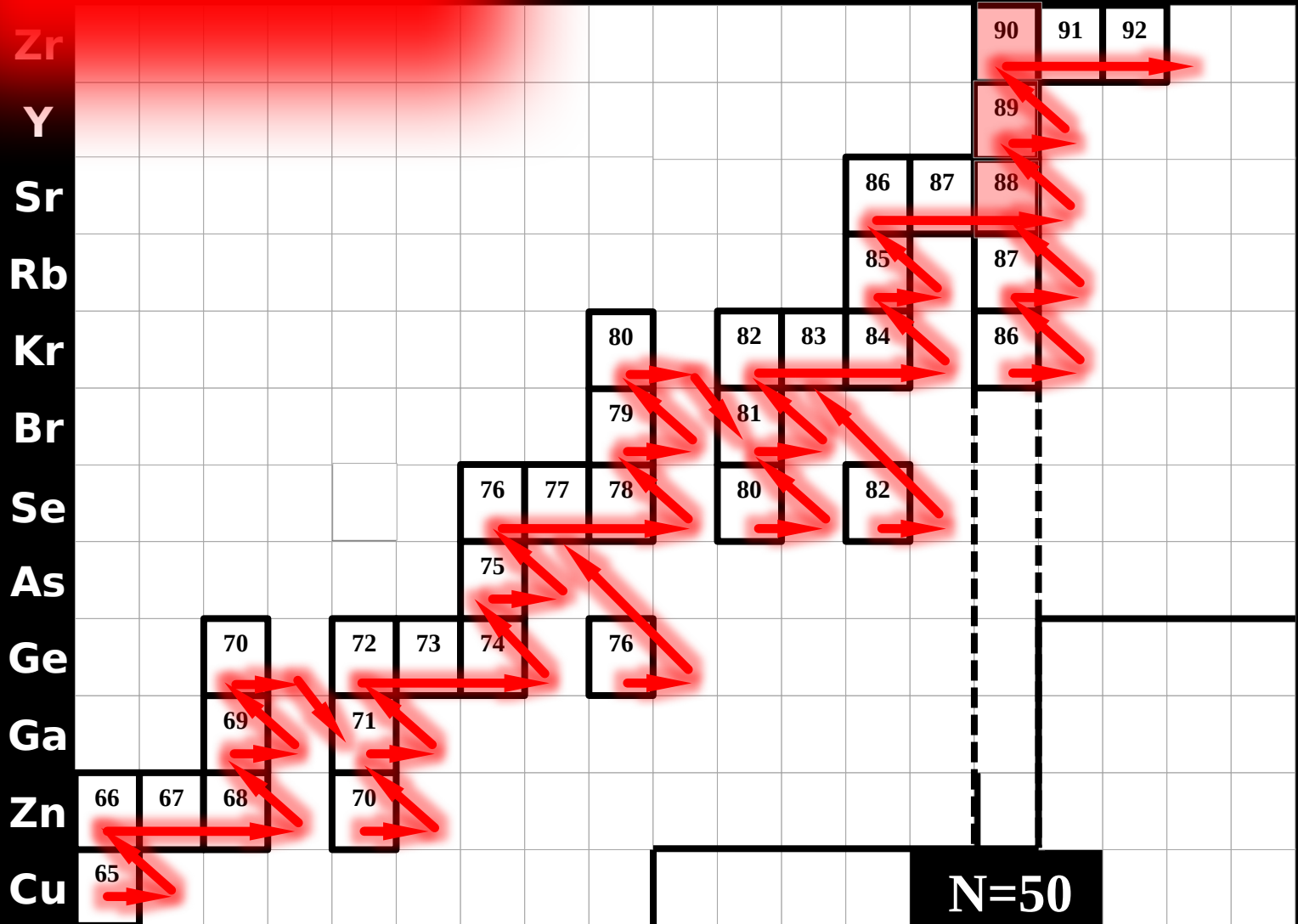
INAFA - Osservatorio Astronomico d'Abruzzo

OUTLINE

- Brief introduction to AGB stars and the **s-process**
- PAST & PRESENT: the **FRUITY** database
- FUTURE: **magnetic models**
- Comparison to **observations**

Proton number

S
process
 $N_n \sim 10^{10} \text{ n/cm}^3$



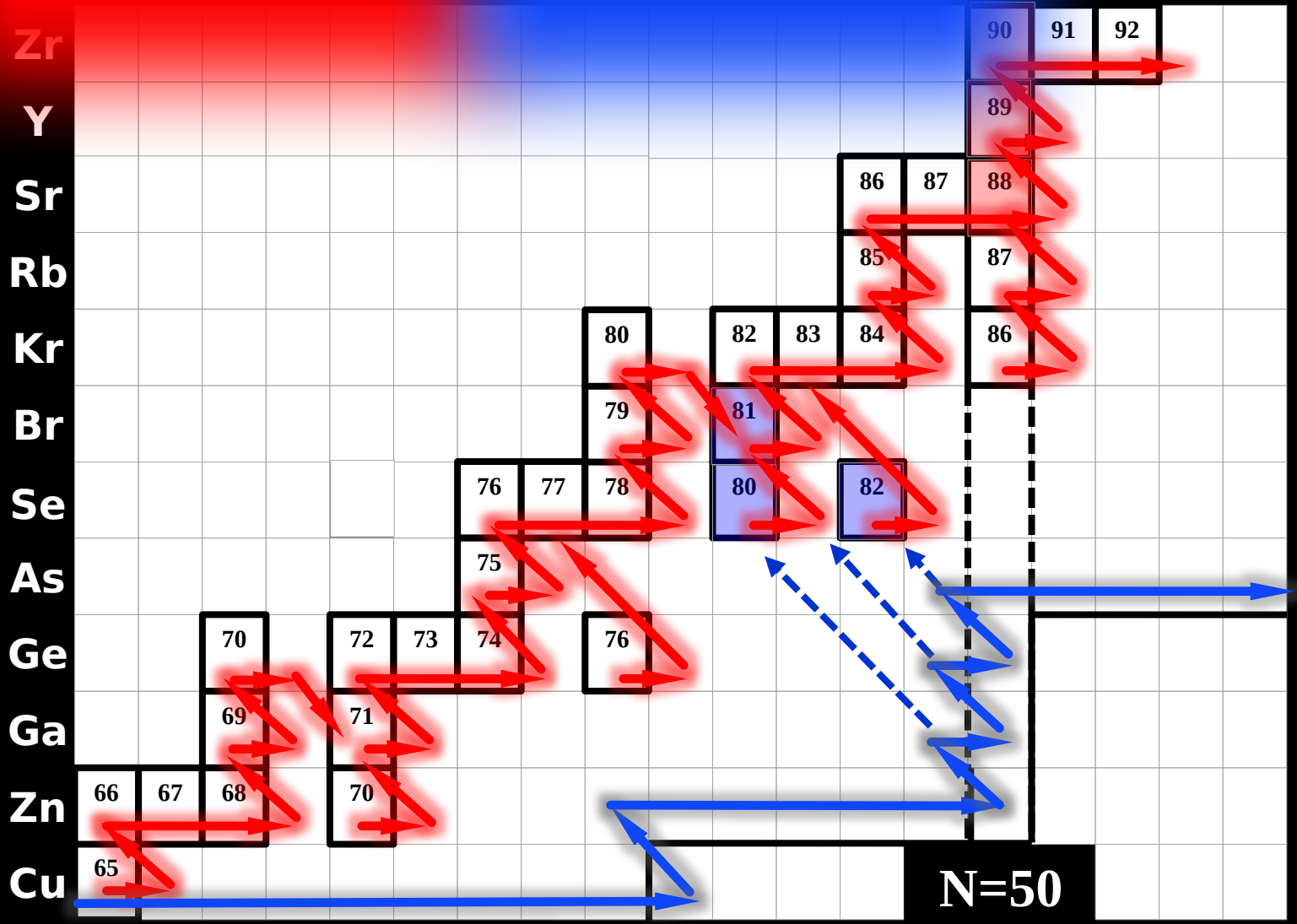
Neutron number

N=50

Proton number

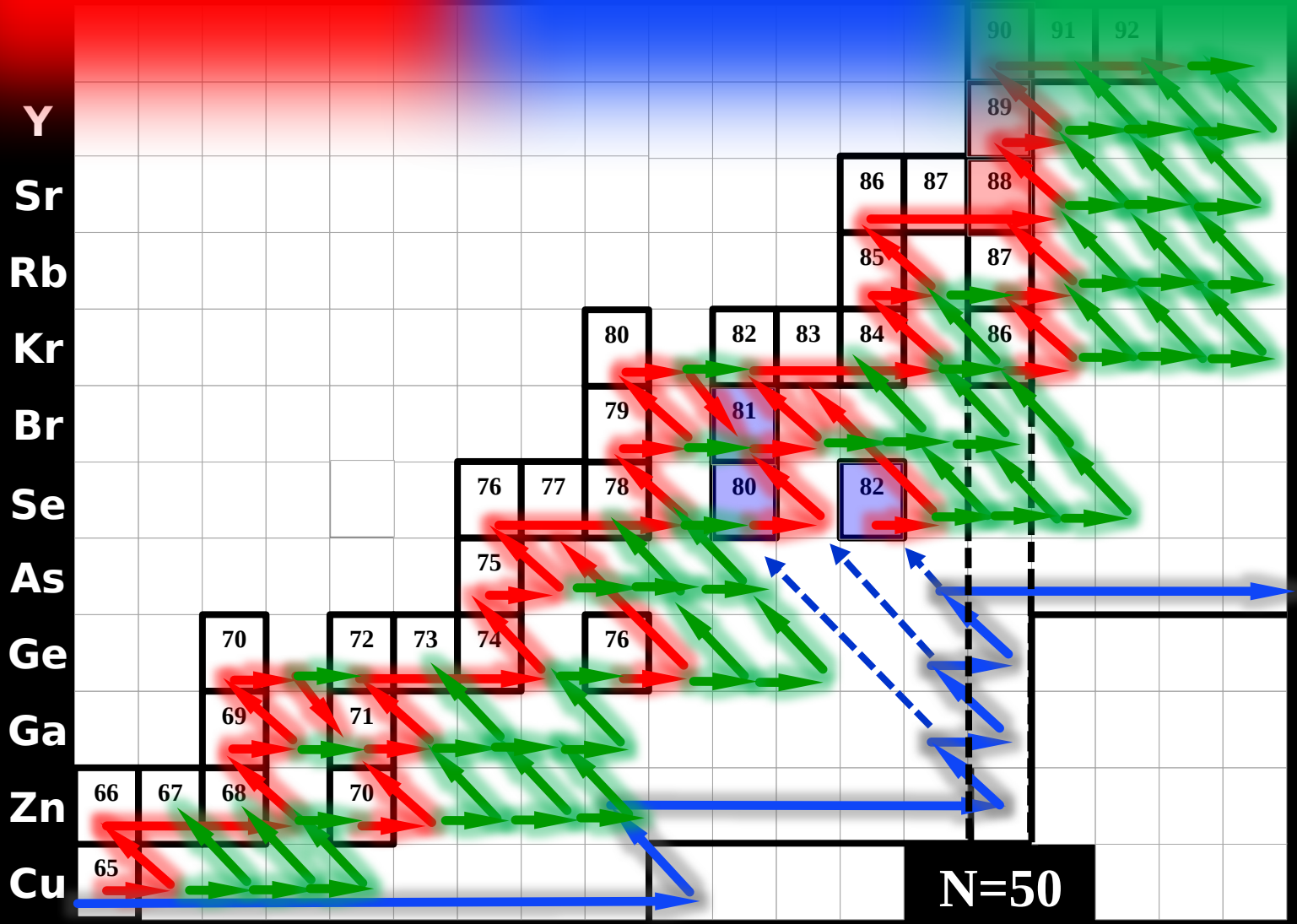
s
process
 $N_n \sim 10^8 \text{ n/cm}^3$

r
process
 $N_n > 10^{10} \text{ n/cm}^3$



Proton number

s process $N_n \sim 10^8 \text{ n/cm}^3$ **r** process $N_n > 10^8 \text{ n/cm}^3$ **i** process $N_n \sim 10^{17} \text{ n/cm}^3$



Neutron number

F.R.U.I.T.Y.

Full-Network Repository of Updated Isotopic Tables & Yields

F.R.U.I.T.Y.
(Full-Network Repository of Updated Isotopic Tables & Yields)

Select Data:

MODEL SELECTION	OUTPUT SELECTION	OUTPUT FORMAT	
Mass (M_{\odot}) ---	Nuclides Properties	Multiple Table format ⁽¹⁰⁾	Single Table format ⁽¹¹⁾
Metallicity (Z) ⁽¹⁾ ---		<input checked="" type="radio"/> Elements ^(3,4) Z: All	<input checked="" type="radio"/> All Dredge Up Episodes ⁽¹²⁾
Initial Rotational Velocity (IRV) ⁽²⁾ 0		<input type="radio"/> Isotopes ⁽⁵⁾ A: All Z: All	
¹³ C Pocket ⁽⁹⁾ Standard		<input type="radio"/> s-process ⁽⁶⁾ : [hs/ls], [Pb/hs], ...	<input type="radio"/> Final
	<input type="radio"/> Net ⁽⁸⁾ Yields ⁽⁷⁾ A: All Z: All		<input type="radio"/> Final
	<input type="radio"/> Total		

[NOTES ON THE MODELS \(pdf file\)](#)

SC+ 2011,2015

fruity.oa-abruzzo.inaf.it

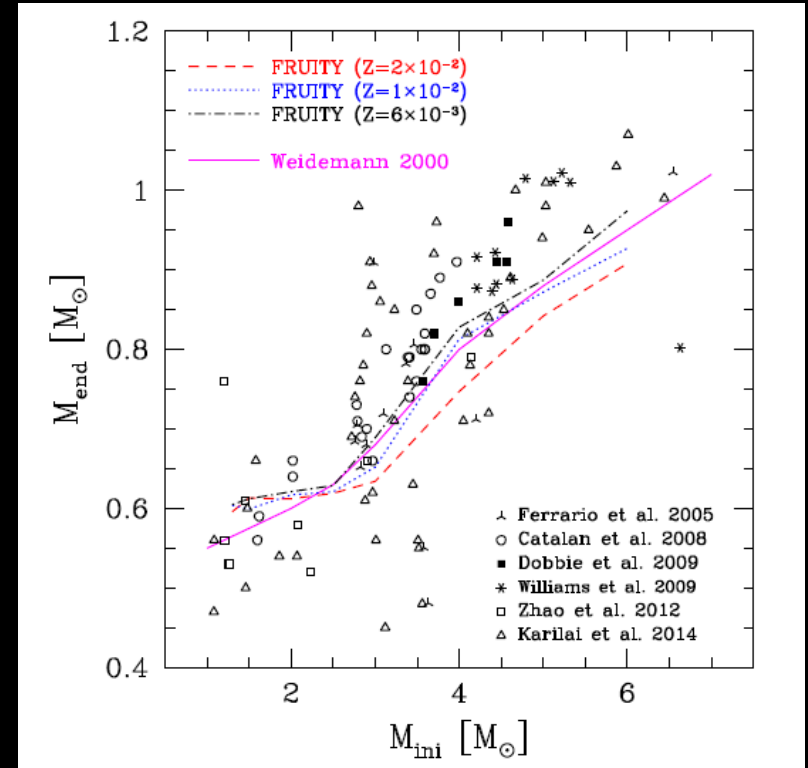
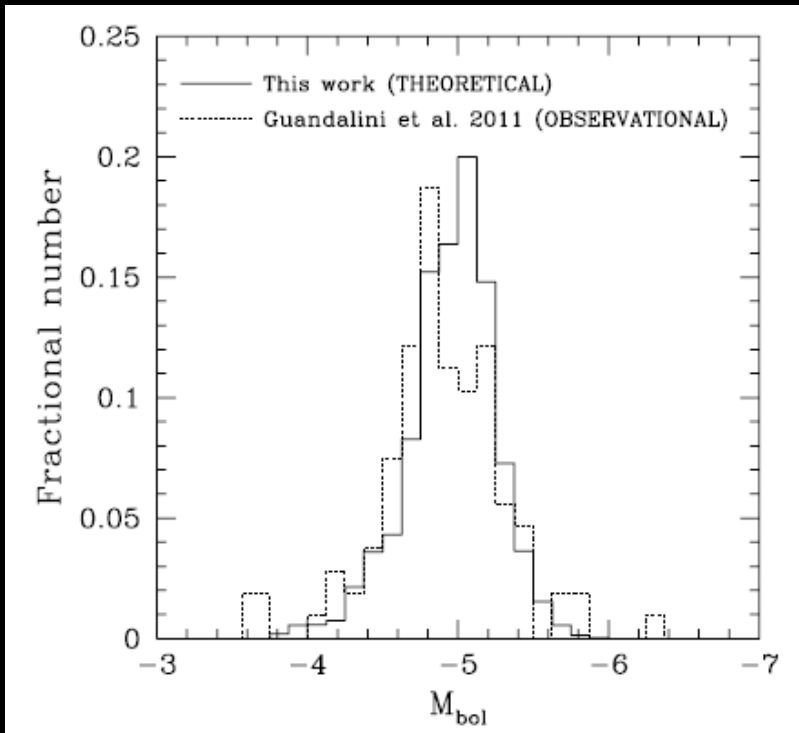
$$-2.85 \leq [Fe/H] \leq +0.3$$

$$1.3 \leq M/M_{sun} \leq 6.0$$

OBSERVATIONS (physics)

Initial-to-final mass relations

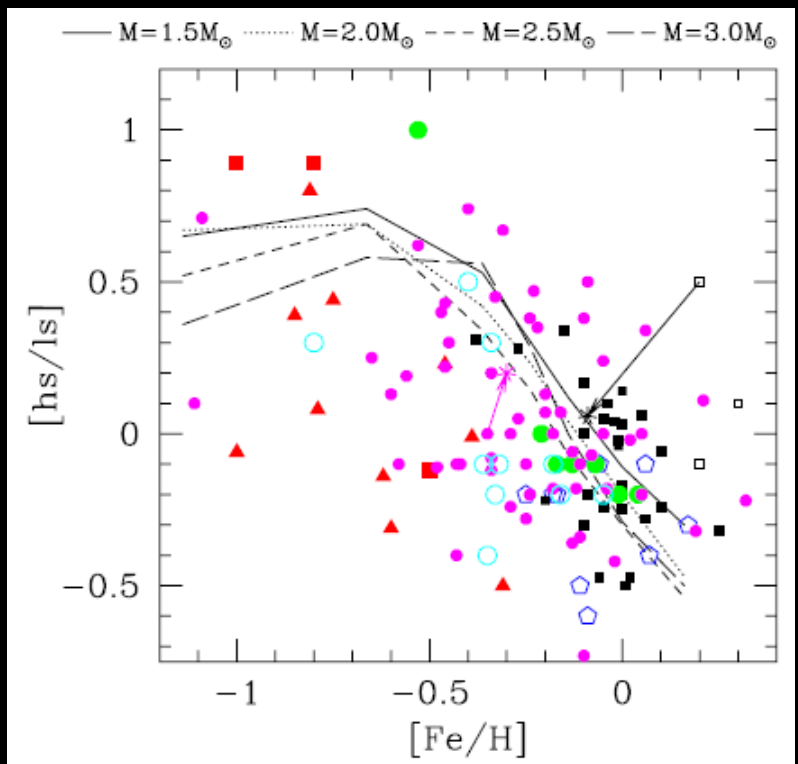
Luminosity Function of C-stars



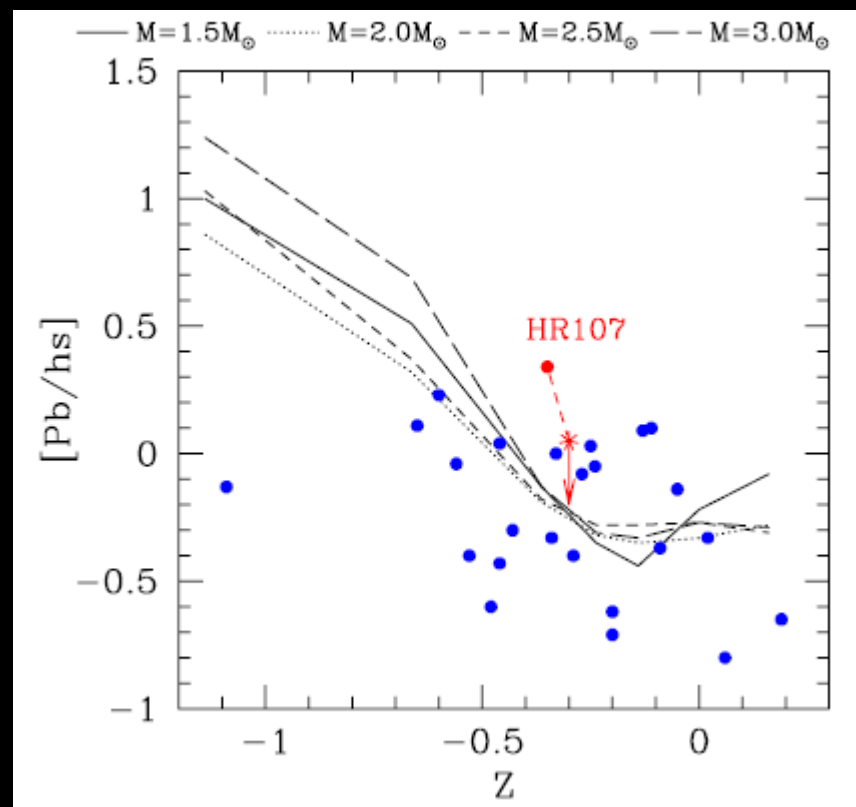
SC+ 2011

OBSERVATIONS (spectroscopy)

Second to first s-process peak



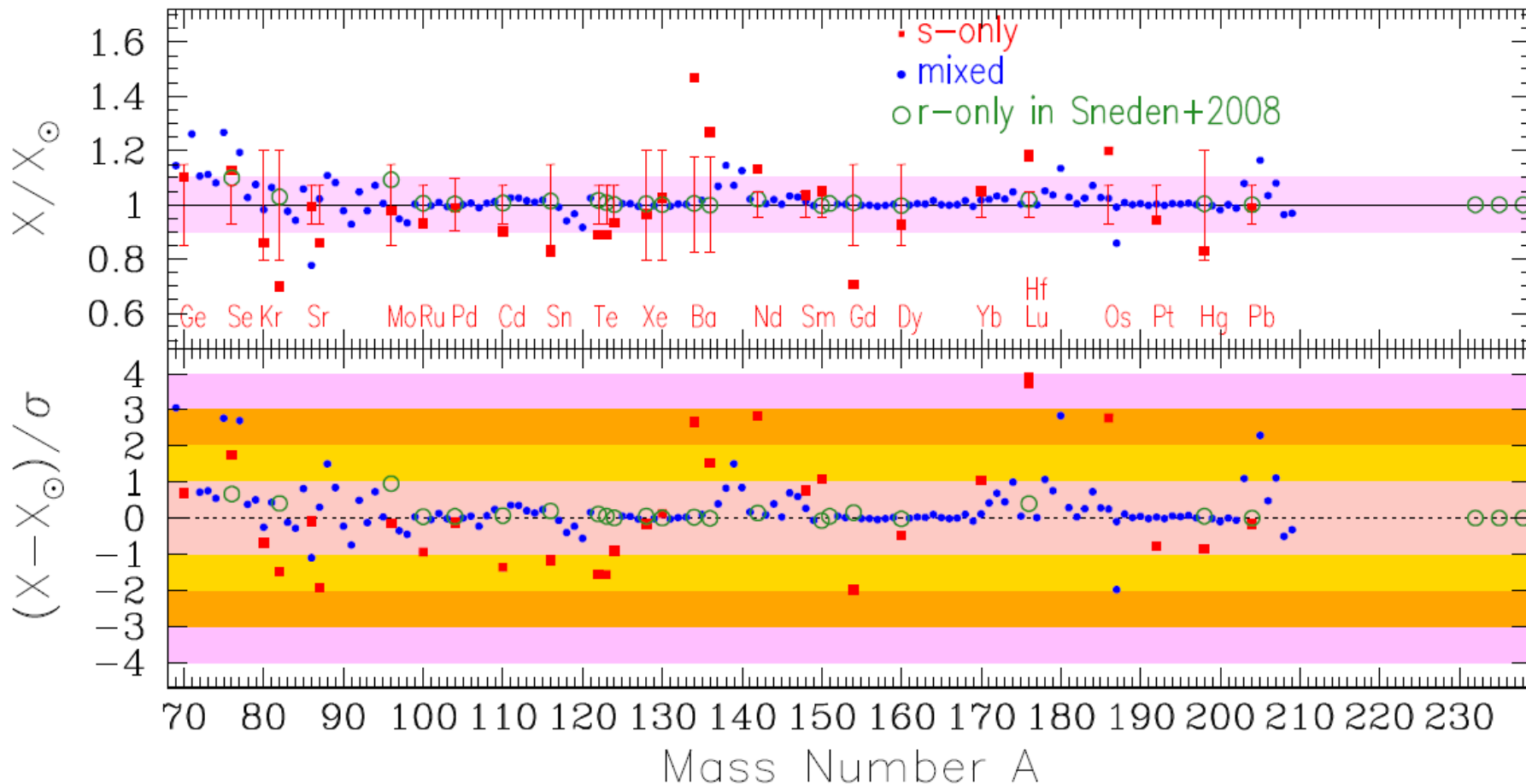
Third to second s-process peak



SC+ 2011

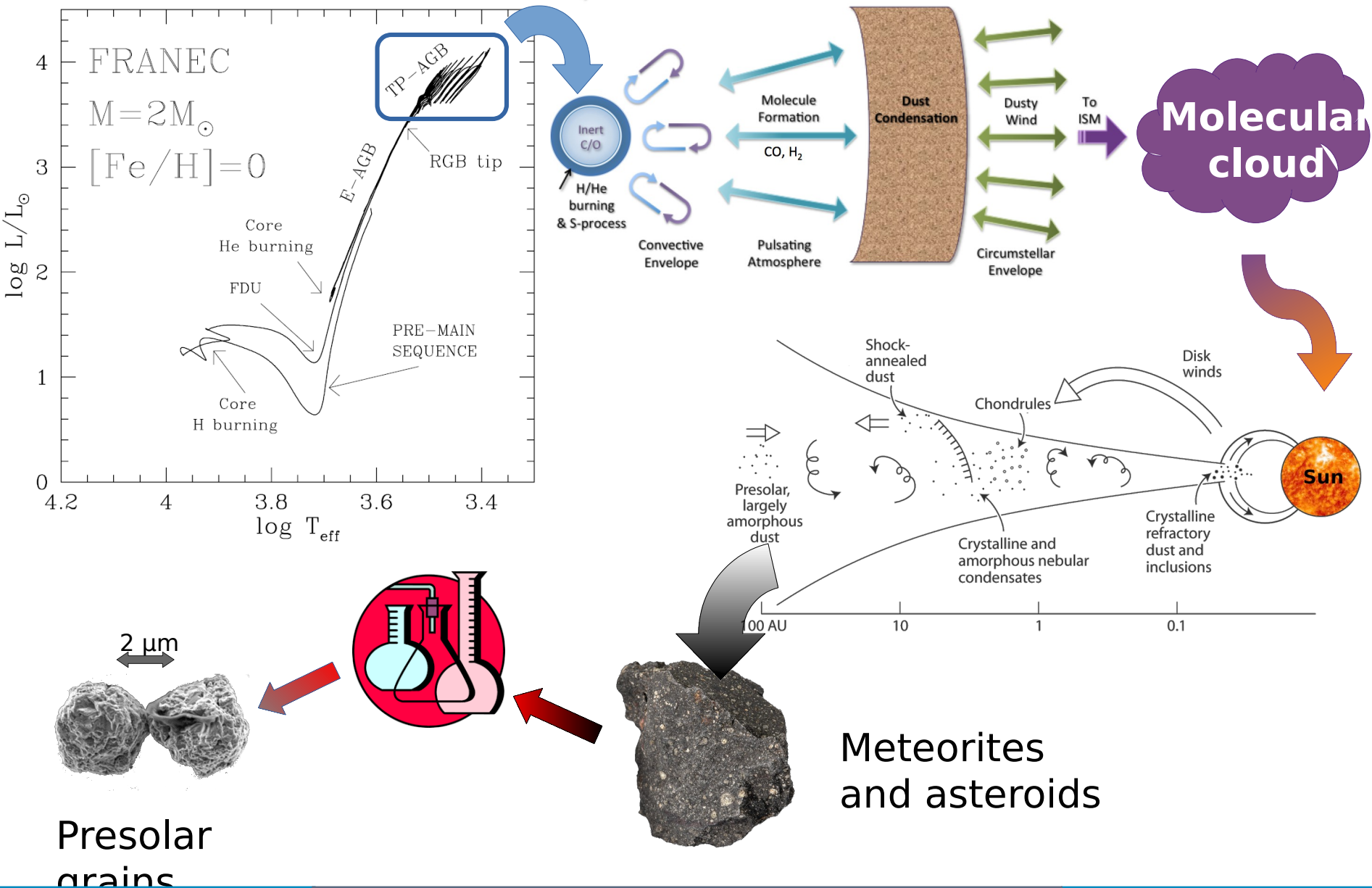
[ls/Fe] = [Sr, Y, Zr/Fe]
 [hs/Fe] = [Ba, La, Ce, Nd/Fe]
 [hs/ls] = [hs/Fe] - [ls/Fe]

Comparison to solar distribution



Prantzos+ 2020

AGB stars and presolar SiC grains




Presolar grains

Meteorites and asteroids

CONSTRAINTS TO STELLAR MODELS

Telescopes


Laboratories

 Large uncertainties
(≈ 2)


Small
uncertainties
($\approx 10\%$)


 No infos on isotopic
composition

Isotopic
abundances

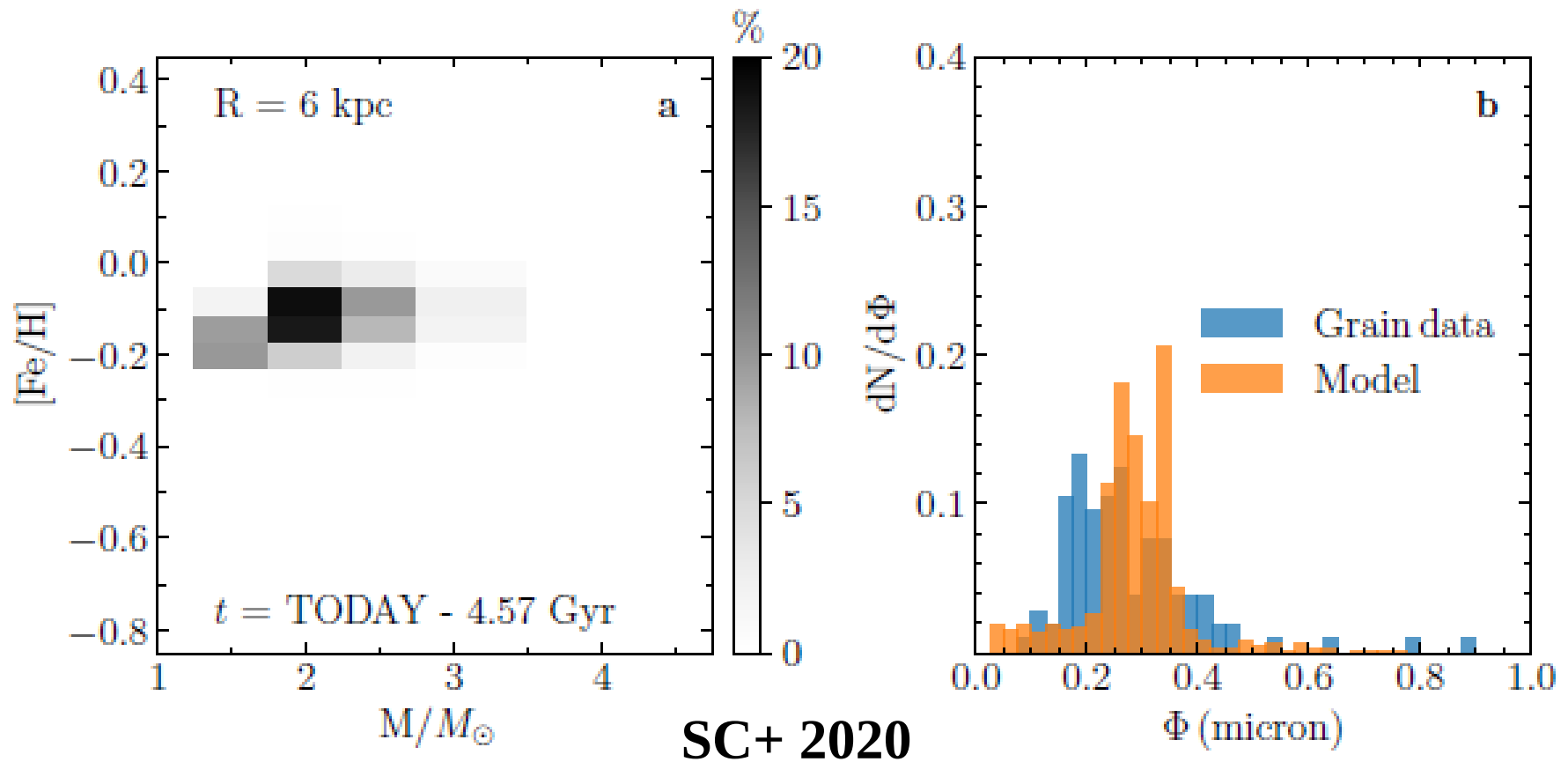
 Elemental
abundances

No infos on elemental
abundances 

 We know the origin:
able to define at least
the metallicity

We guess which
was the parent
star 

First important question is: which stars are progenitor of pre-solar SiC grains?

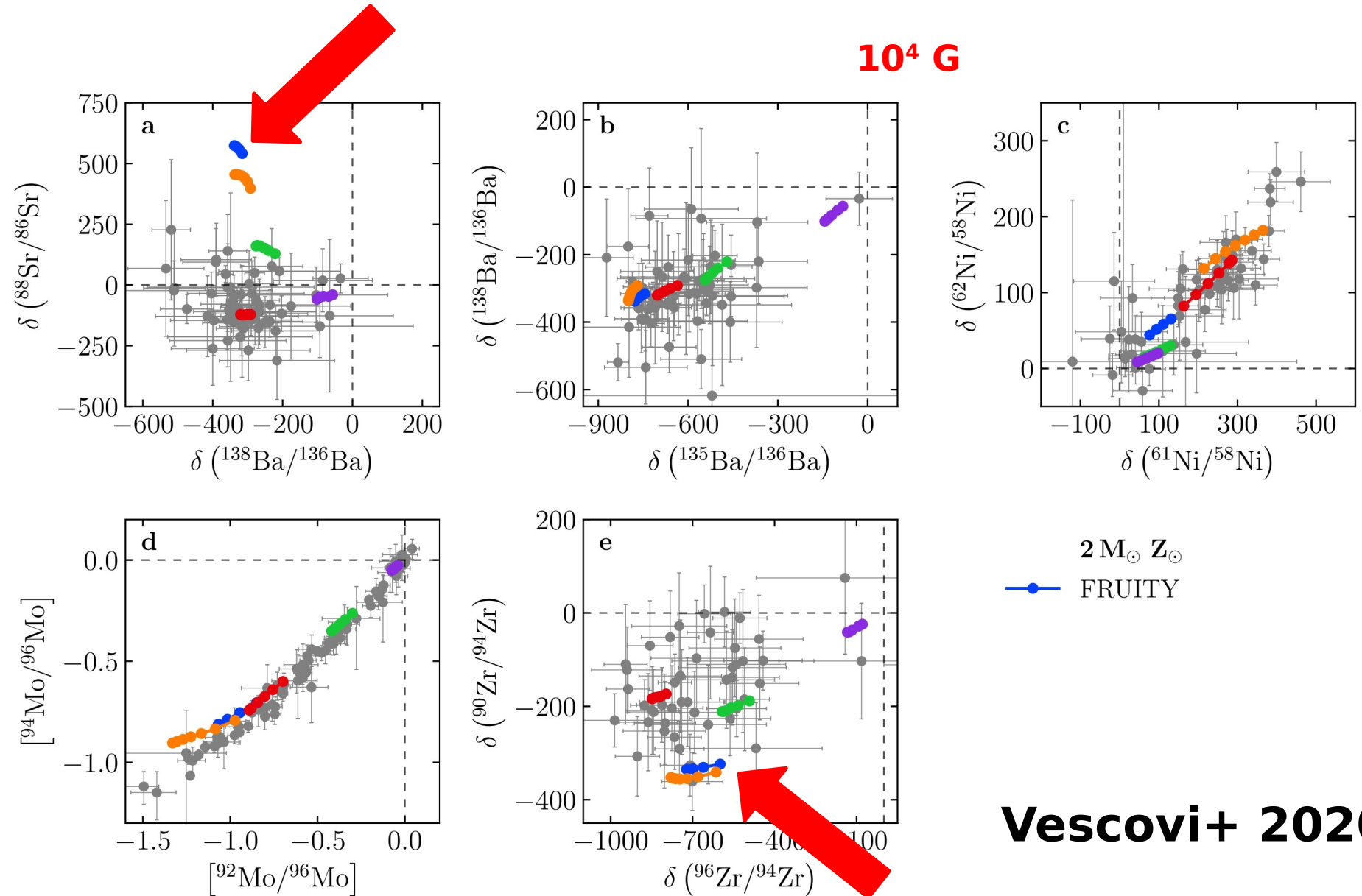


Our models predict that the SiC production at the epoch of the Solar System formation is dominated by contributions from AGB stars with $M \approx 2M_{\text{SUN}}$ and $Z \approx Z_{\text{SUN}}$, which are thus likely the parent stars of presolar SiC grains identified in extraterrestrial materials.

SiC Grains

$$\delta(^{88}\text{Sr}/^{86}\text{Sr}) = [(^{88}\text{Sr}/^{86}\text{Sr})_{\text{grain}} - (^{88}\text{Sr}/^{86}\text{Sr})_{\text{SS}} - 1] \times 1000$$

10⁴ G

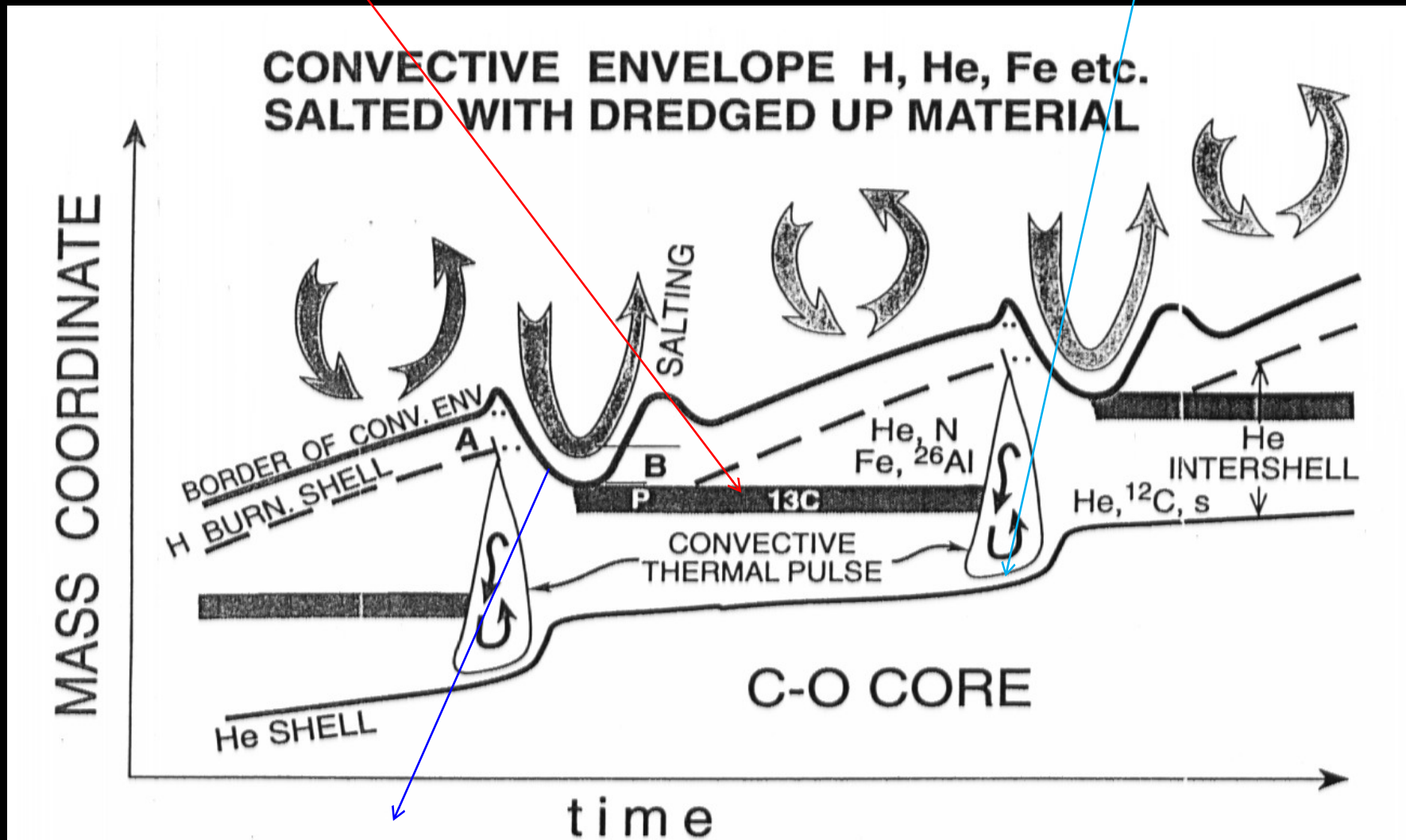


Vescovi+ 2020

The s-process in AGB stars

$^{13}\text{C}(\alpha, n)^{16}\text{O}$ reaction

$^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction



Third Dredge Up (TDU)

Busso et al. 1999

The ^{13}C pocket in stellar evolutionary models

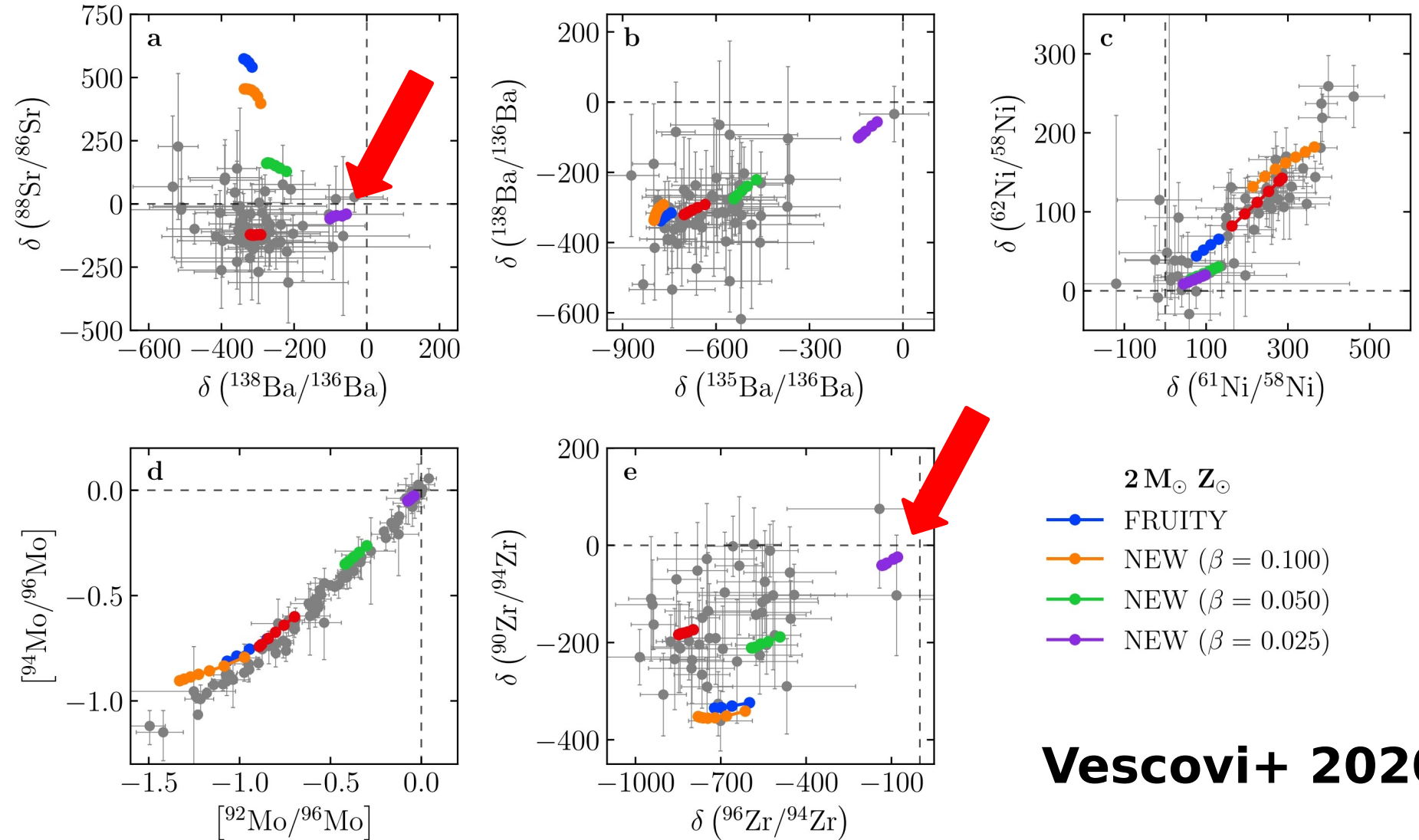
- ✓ **Opacity induced overshoot** (SC+2009)
- ✓ **Convective Boundary Mixing + Gravity Waves** (Battino+ 2017)
- ✓ **Magnetic-induced mixing** (Vescovi+2020)

How does the ^{13}C pocket change?

- ✓ **Rotation-induced mixing** (Herwig+ 2003; Siess+ 2004; Piersanti+ 2013)

SiC Grains

10⁴ G



Vescovi+ 2020

A new working hypothesis: magnetic induced mixing

Nucci & Busso 2014

$$v_r = \frac{dw(t)}{dt} r^{-(k+1)}$$

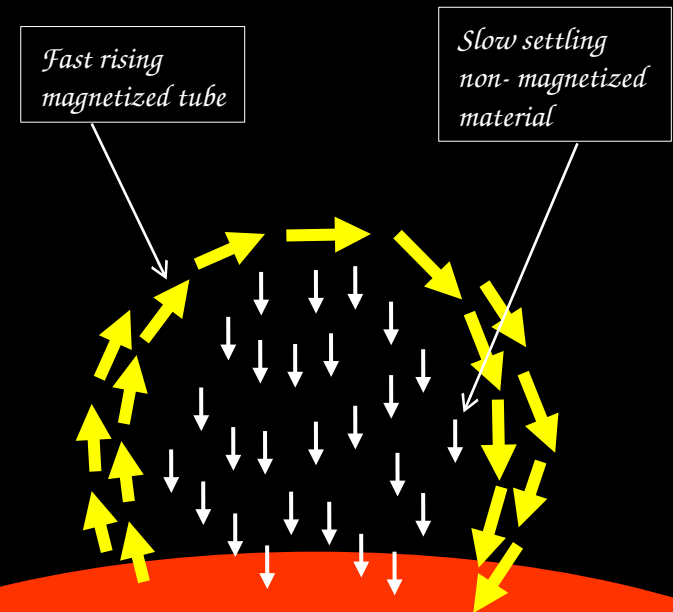
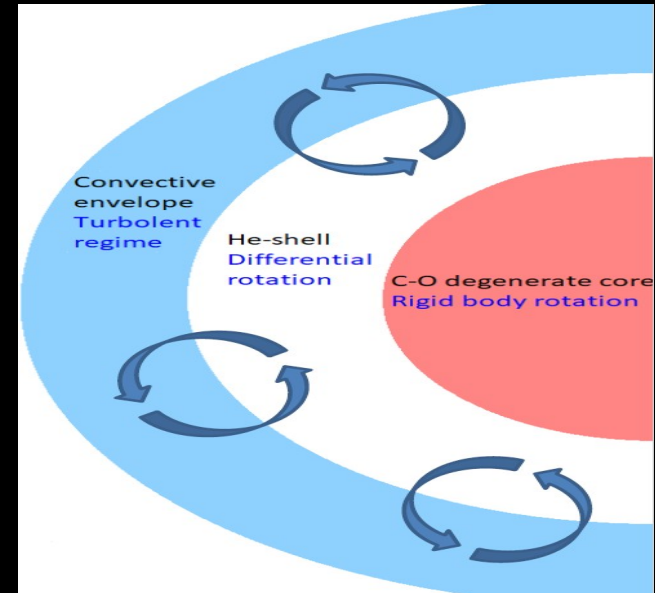
$$B_\phi = \Phi(\xi)r^{k+1}, \quad [\xi = -(k+2)w(t) + r^{k+2}].$$

A magnetized stellar plasma in the quasi-ideal MHD regime, with a density distribution closely following a power law as a function of the radius ($\rho \propto r^k$, with $k < -1$), reaches a dynamic equilibrium and is in radial expansion.

The result above is analytically exact.

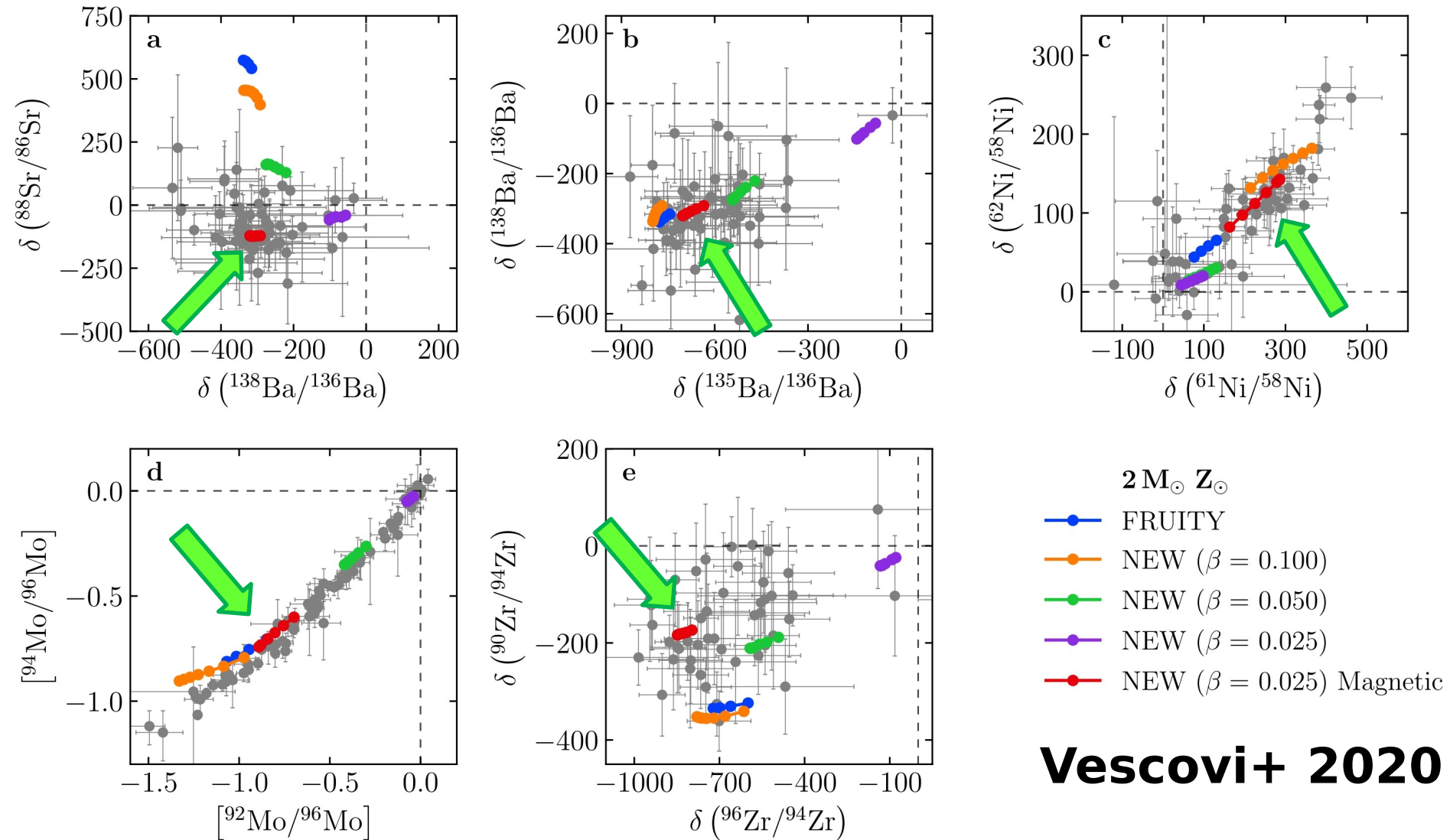
$$v_{down}(r) = v(r_p) \frac{\rho(r_p)}{\rho(r_{h+1})} \left(\frac{r_h}{r_p}\right)^{k+2} \left(\frac{r_h}{r}\right)^{k+1}$$

In strong field regimes, the magnetic field tends to concentrate in flux tubes. As a consequence of the magnetic extra-pressure, these tubes are buoyant (see, e.g., Parker 1955). Due to the effect of the magnetic buoyancy, a matter flow is pushed from the He-intershell toward the envelope. This, in turn, induces a downflow flux, in order to guarantee mass conservation.



SiC Grains

- **Magnetic** contribution account for SiC data!!
- Best fit for $u_p = 5 \times 10^{-5}$ cm/s and $B_\phi = 5 \times 10^4$ G

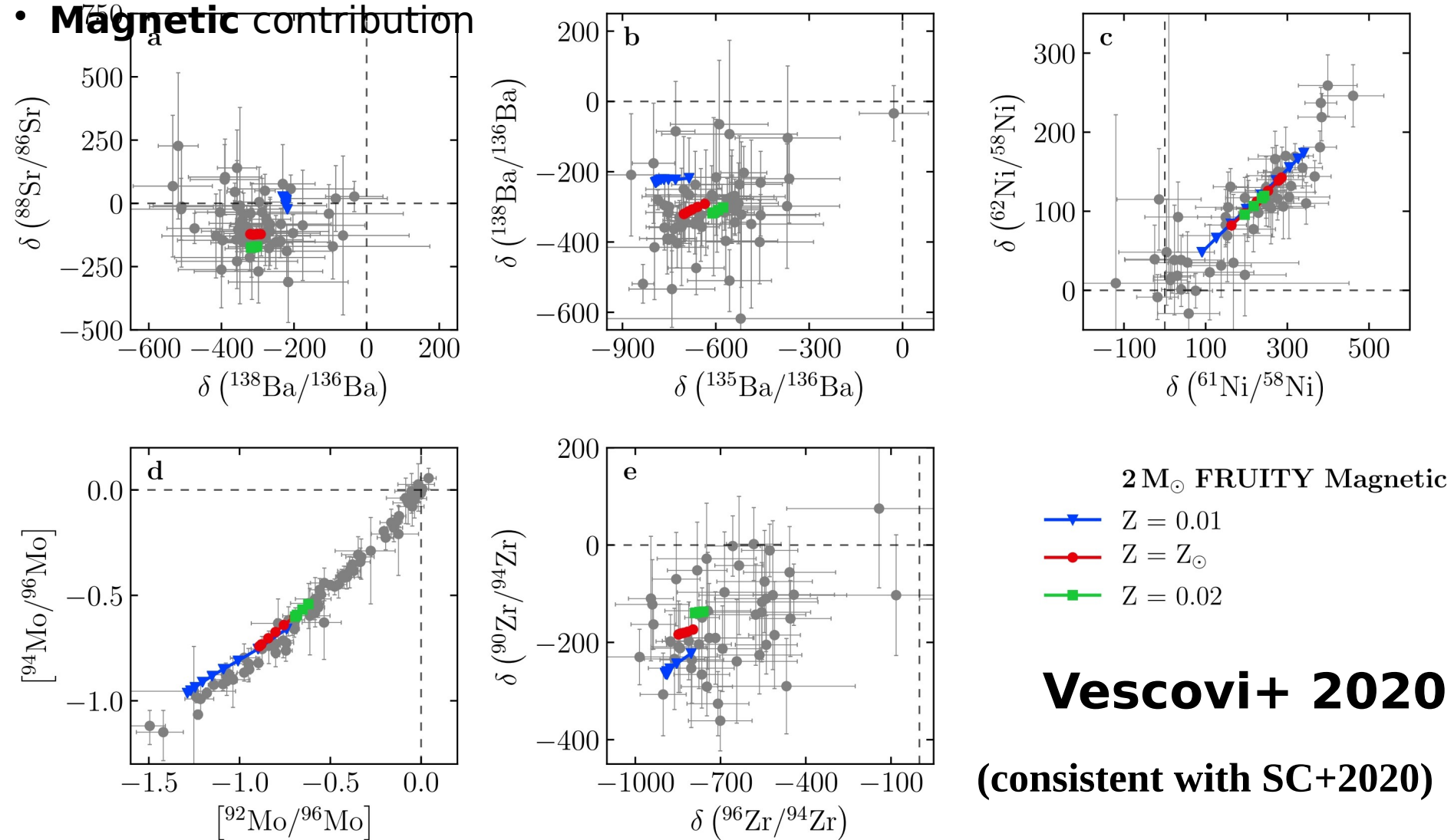


Vescovi+ 2020

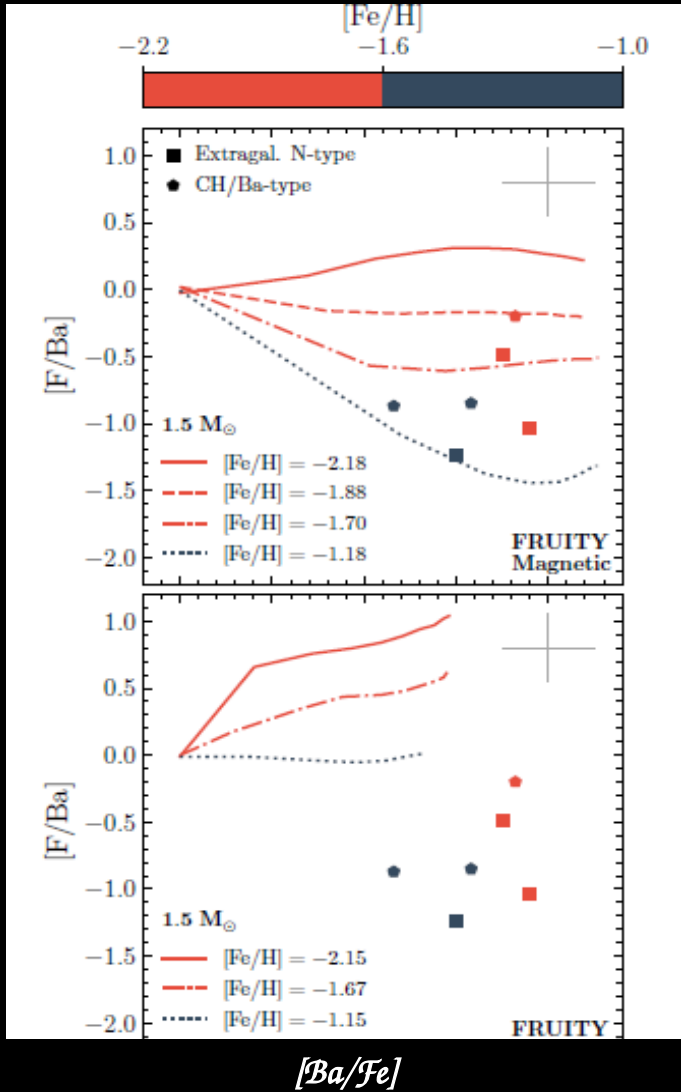
SiC Grains I

- Stellar models with **same initial mass** ($2 M_{\odot}$) and **close-to-solar metallicity**

- Magnetic contribution**



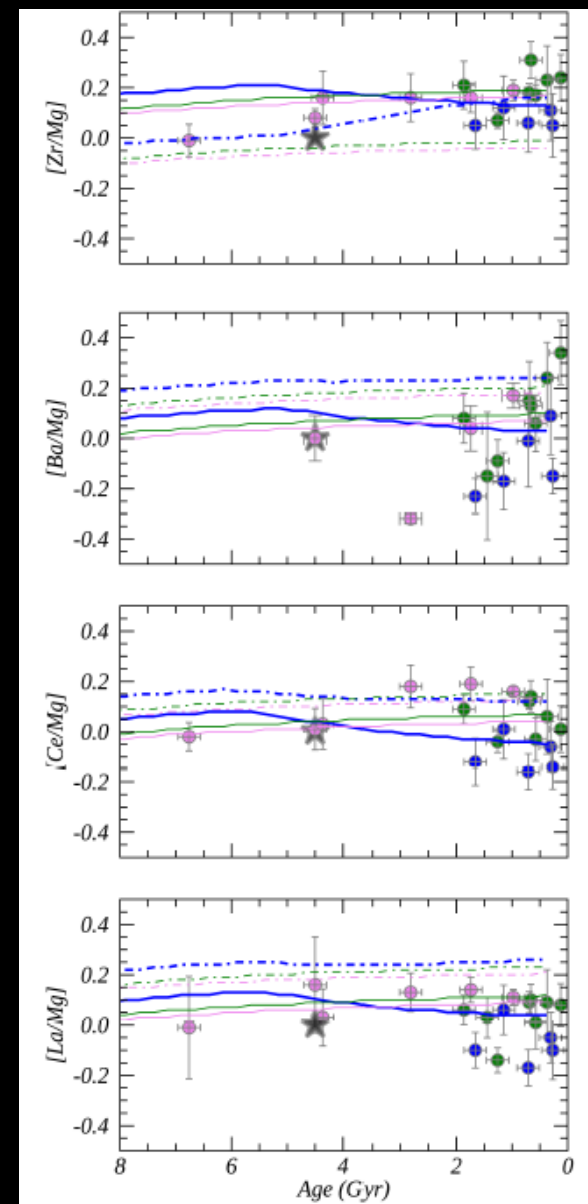
Magnetic induced mixing



Vescovi+ 2021

Better agreement

1. Young Open Clusters;
2. Fluorine @ low Z.

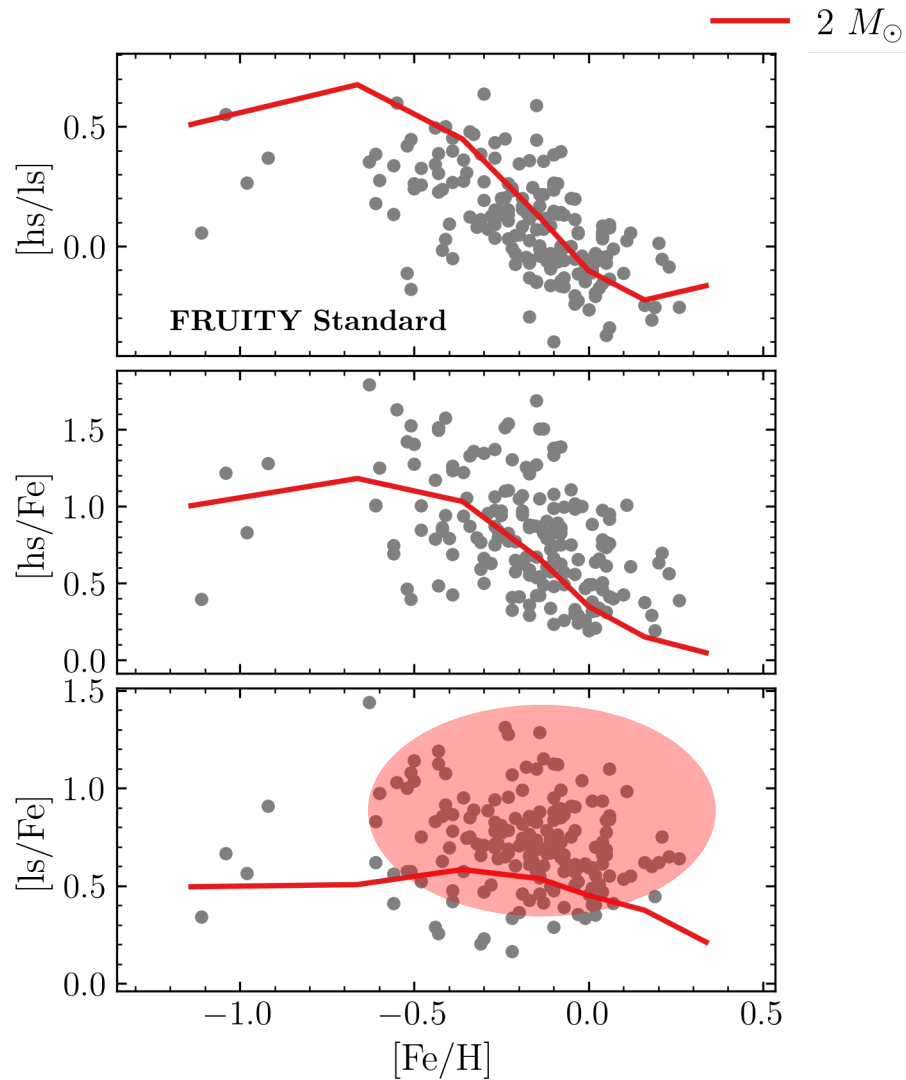


Magrini+ 2021

WHAT'S NEXT?

- **Ba stars**
- **C-stars**
- **S-stars**
- **Post-AGB stars**

Ba-stars



The origine of magnetic fields in stars

Still largely debated topic:

1. fossil relics in stably stratified radiative regions (inherited from previous evolutionary phases);
2. dynamo-generated in turbulent convective layers.

Since the time-scale for ohmic decay of a large-scale field is typically longer than stellar lifetime, the **radiative regions** may be regarded as perfectly conducting and the magnetic field is then "**frozen**" into the plasma.

During a Thermal Pulse, **turbulence** leads to rapid reconnection that dissipates any large-scale coherent field.

HOWEVER, convection, rotation, and shear within the convective region will regenerate the field through **dynamo action**: numerical simulations suggest that convective layers are site of very efficient small-scale dynamos.

BUT: we are interested in a **axisymmetric toroidal magnetic field**.

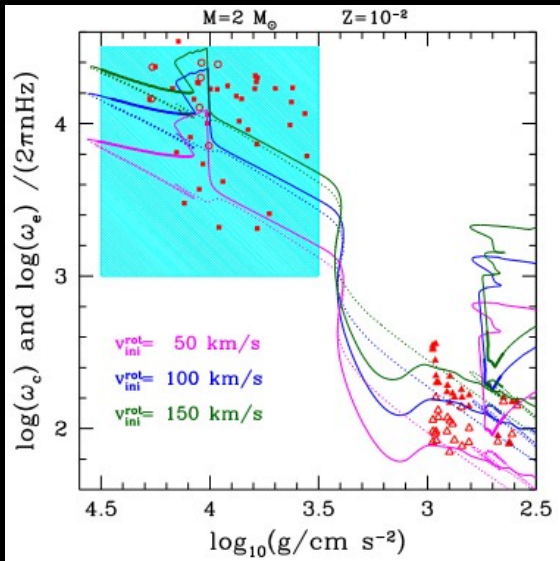
The origine of magnetic fields in stars

Such a field could be achieved through the stretching of a preexisting low-magnitude poloidal field in the radiative zone below the convective envelope after the quenching of a thermal pulse, via the action of differential rotation around the rotation axis

$$\frac{\partial B_\phi}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} (\Omega r^2 B_p) = \Omega q B_p$$

$$\frac{1}{2} (r\Omega(t))^2 + \frac{B_\phi^2}{8\pi\rho} = \frac{1}{2} (r\Omega_{ini})^2$$

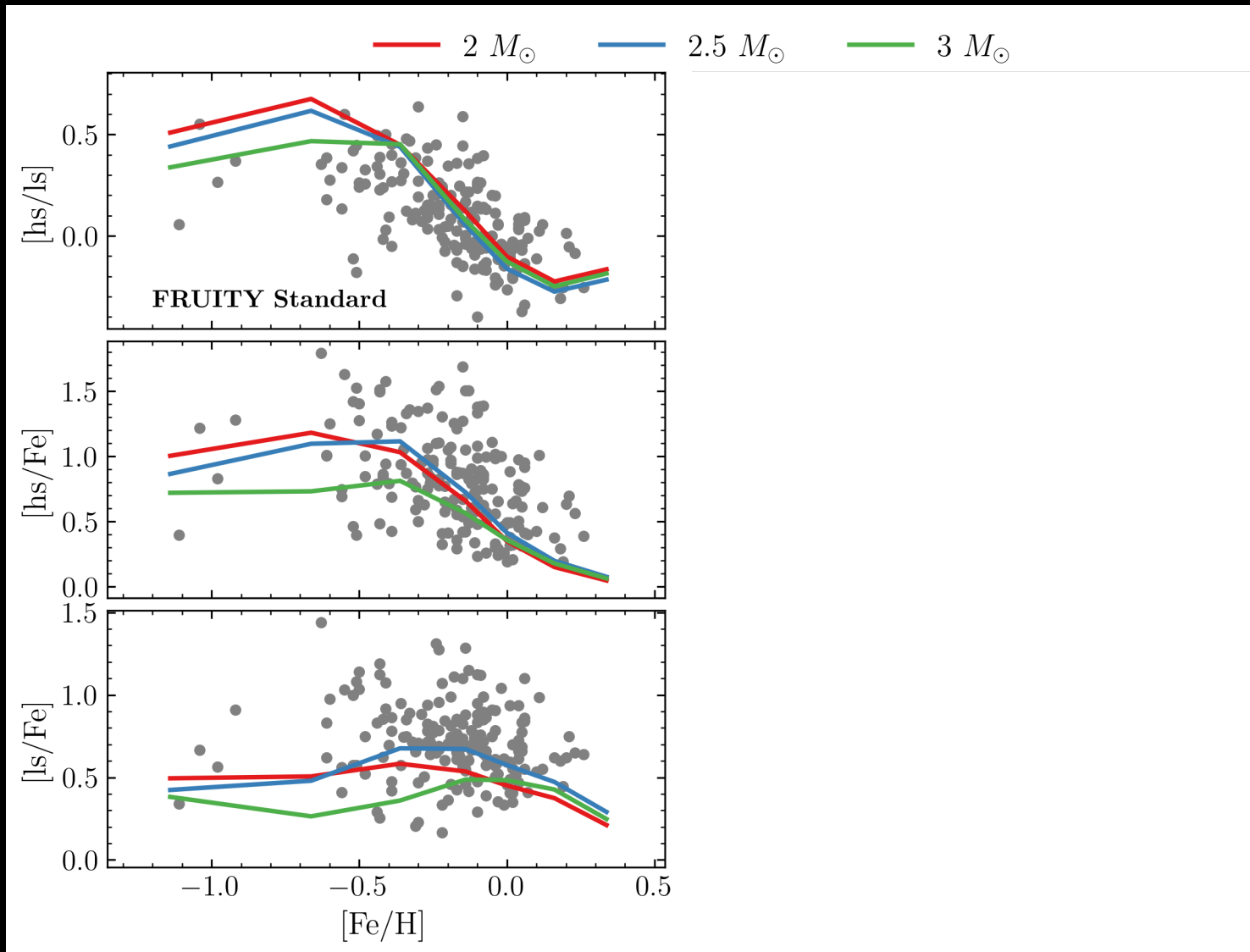
$$B_p \approx 10 \text{ G} \quad B_\phi \approx 10^5 \text{ G}$$



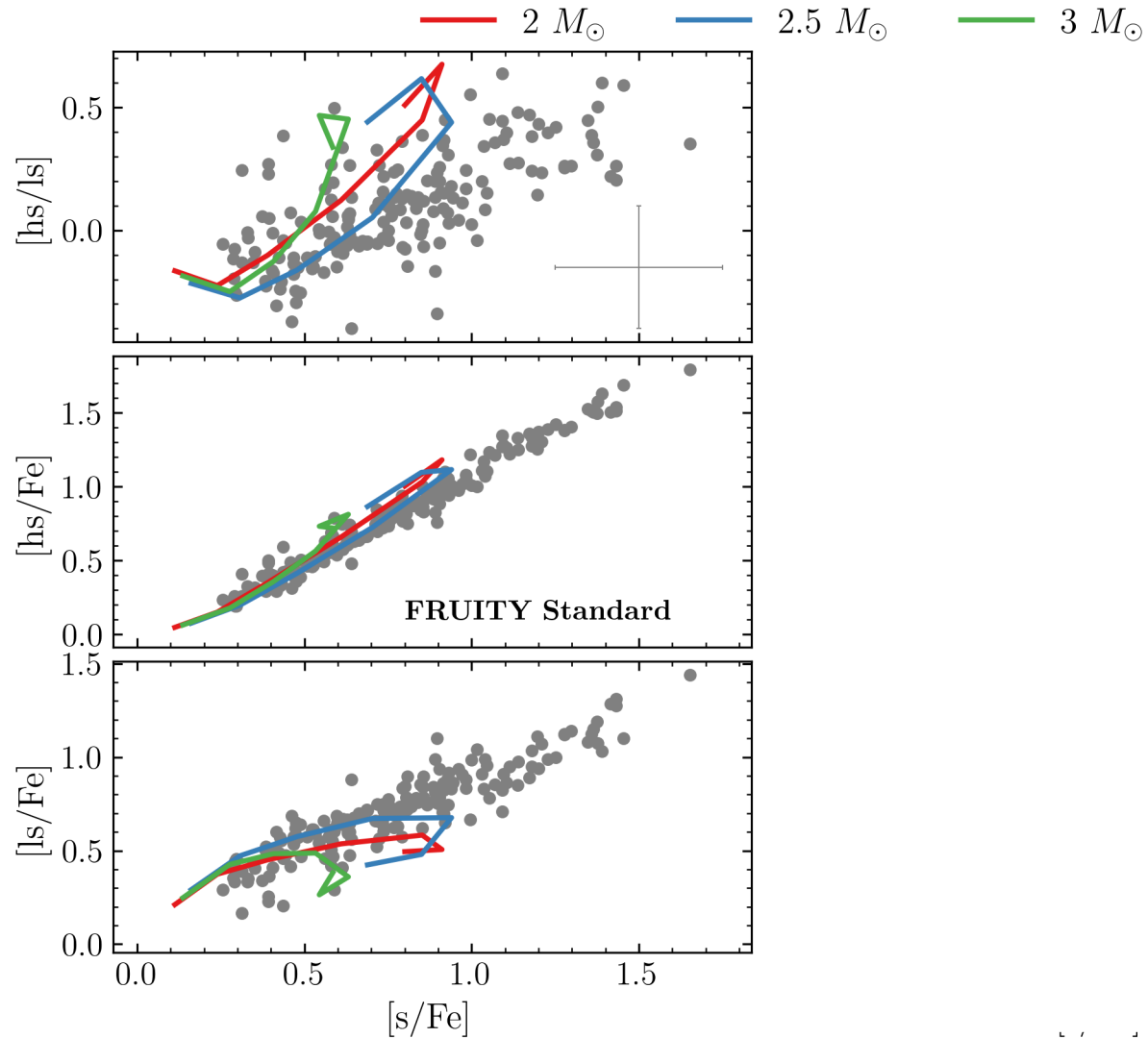
Artificial viscosity added to match asteroseismic data.

1. *Different initial rotation velocities;*
2. *Different TDU efficiencies.*

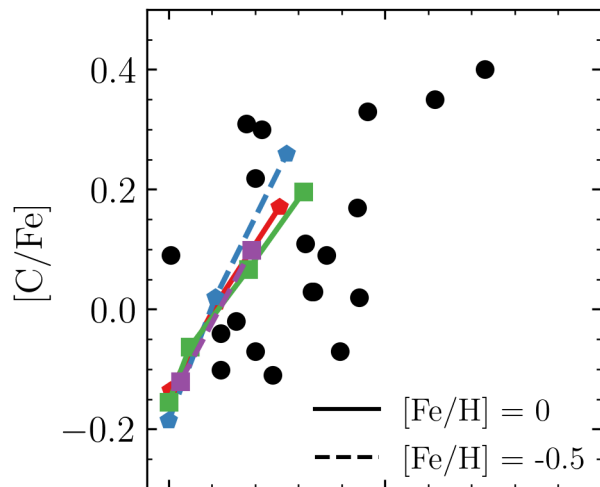
Ba-stars



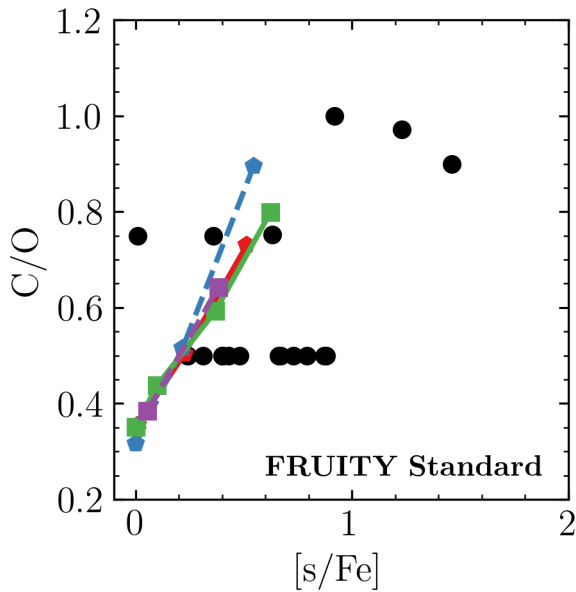
Ba-stars



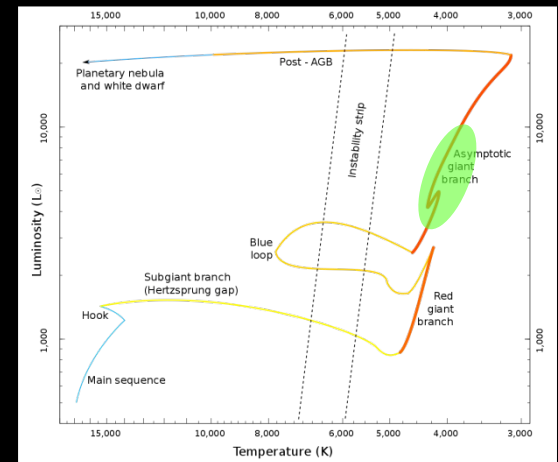
S-stars



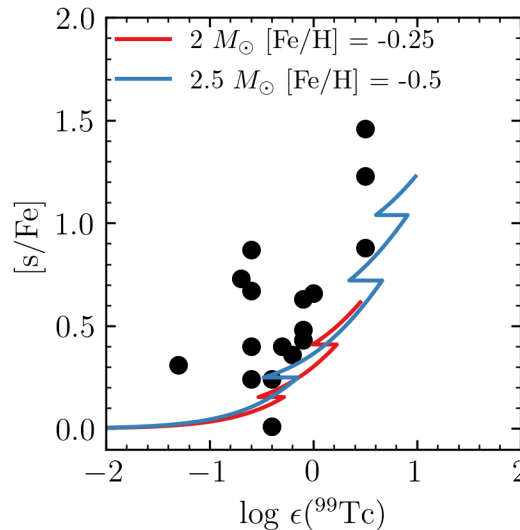
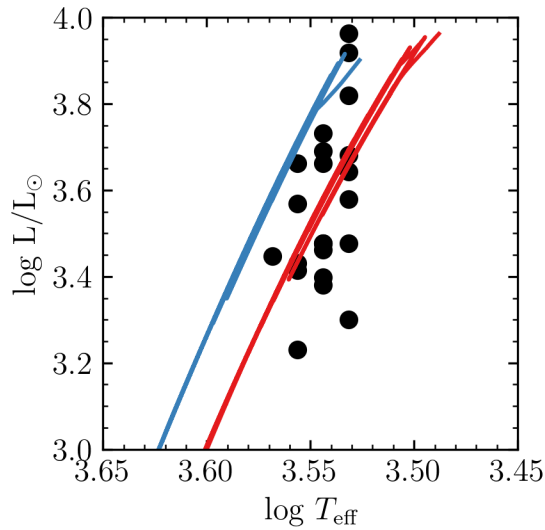
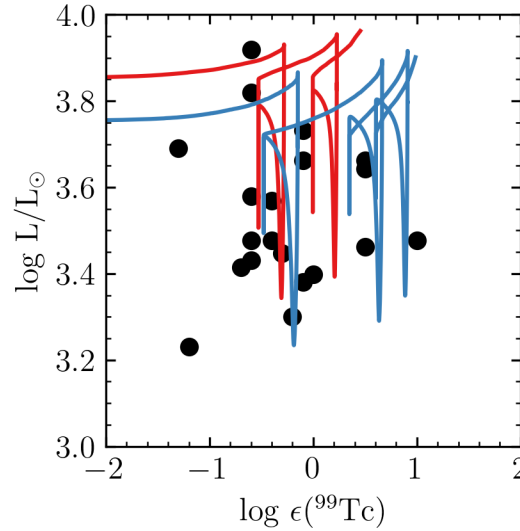
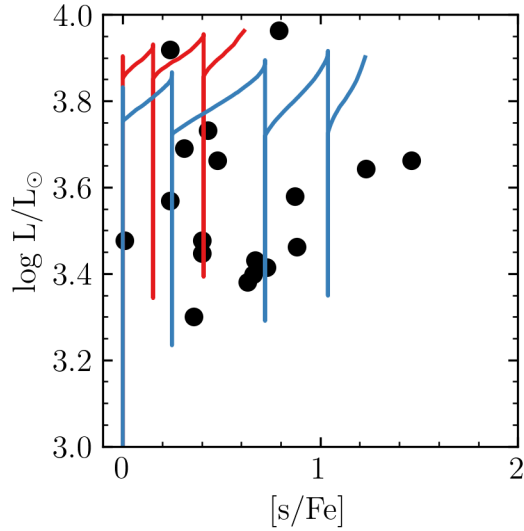
- ◆ $2 M_{\odot}$
- $2.5 M_{\odot}$



$$0.3 < C/O < 1.0$$

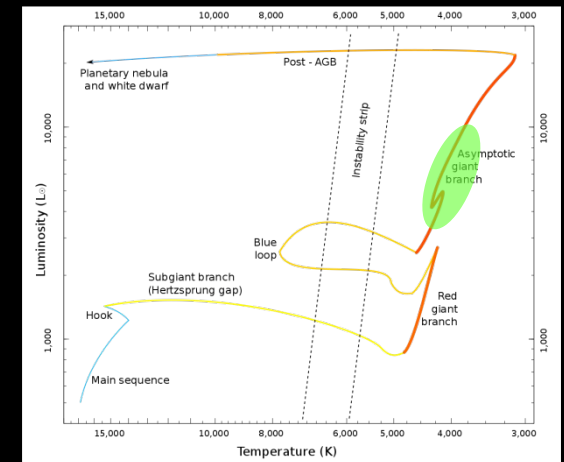


S-stars



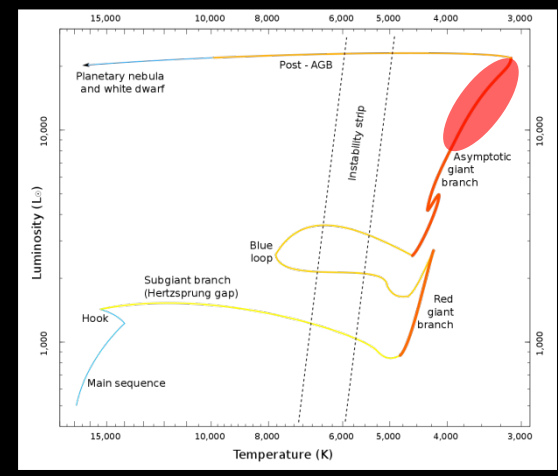
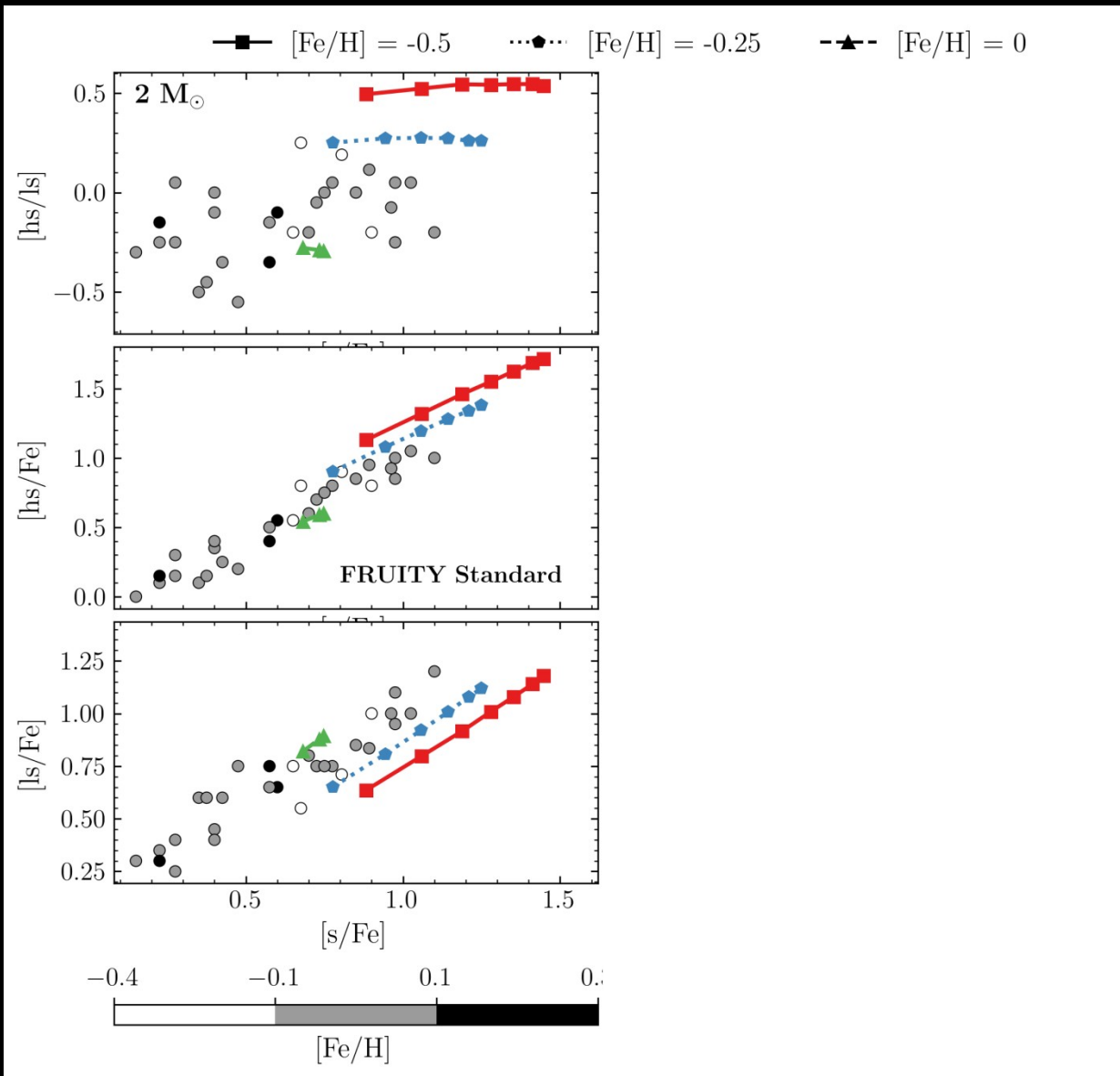
$$0.3 < \text{C/O} < 1.0$$

Tc is freshly produced by TDU [Merrill 1952]

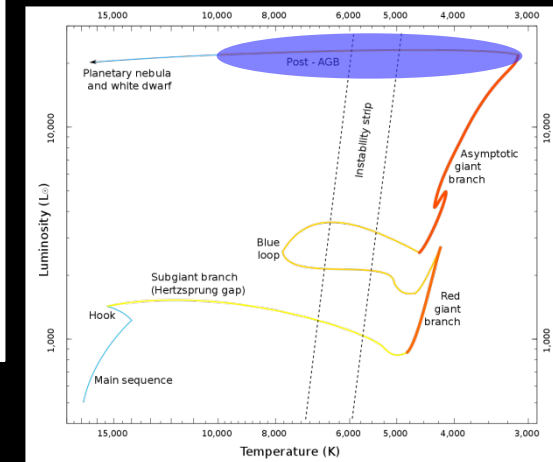
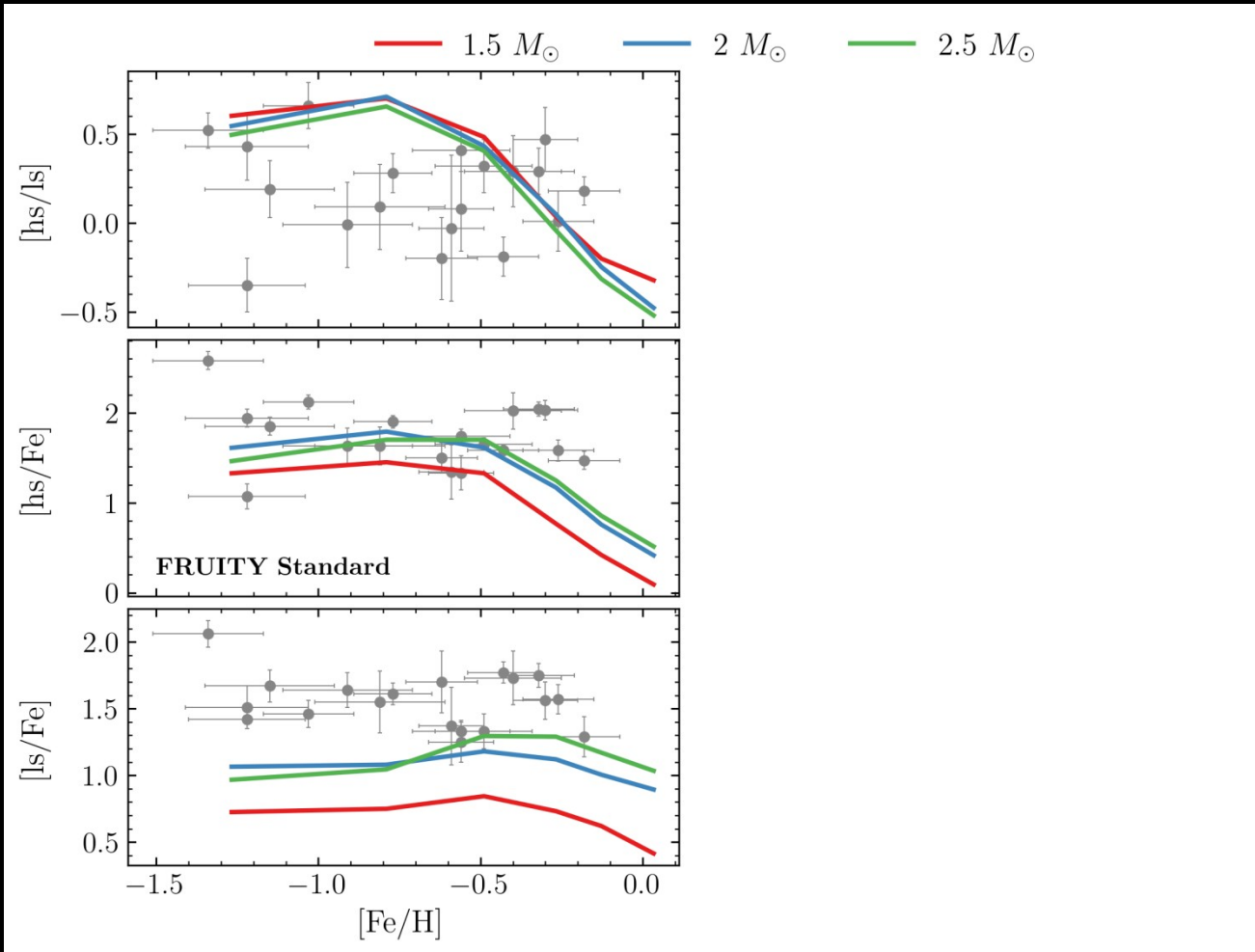


C-stars

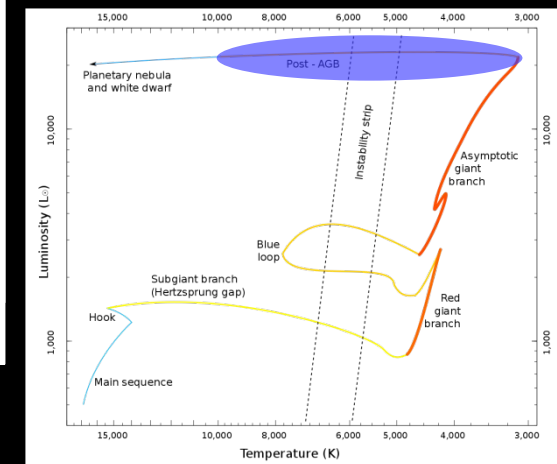
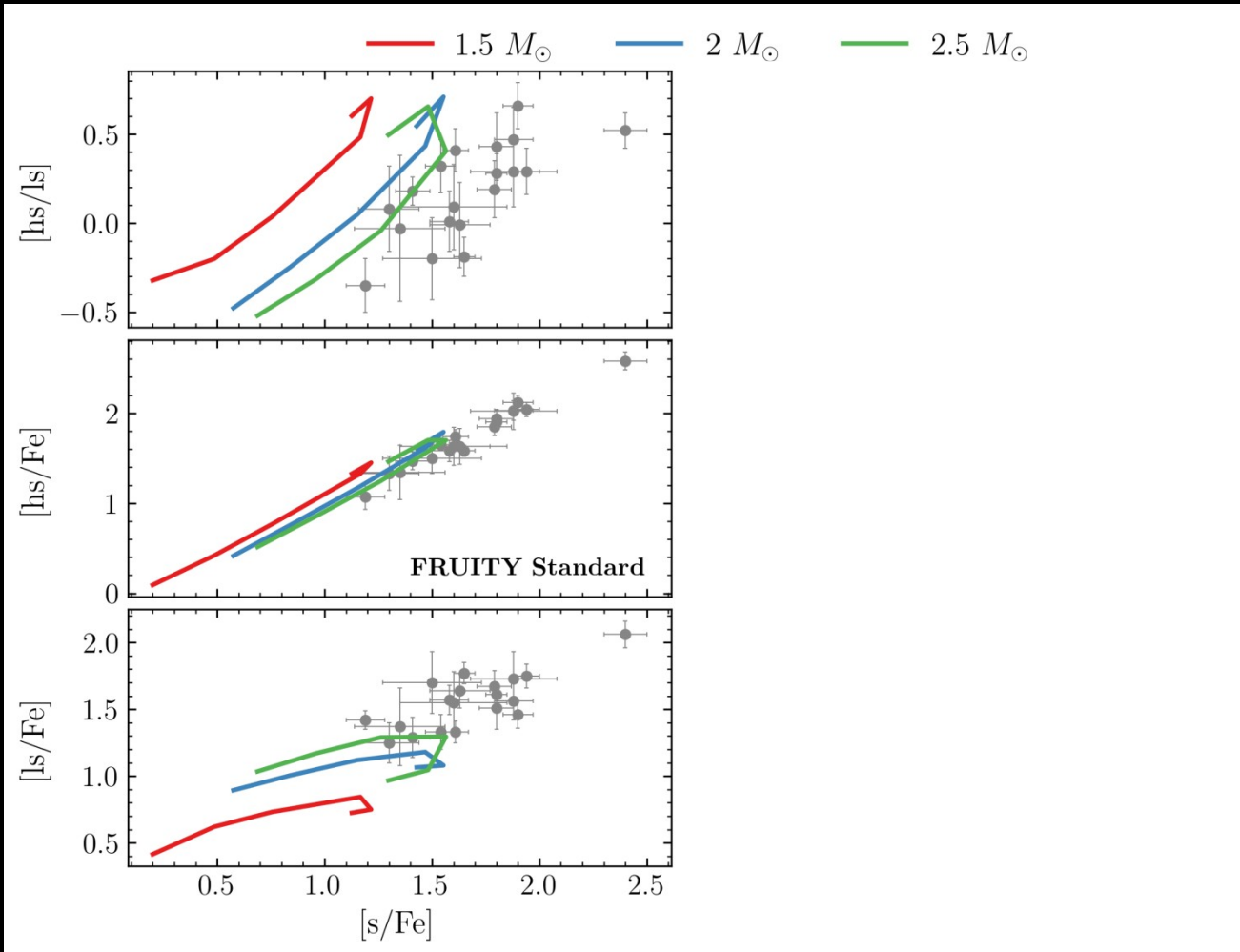
$C/O > 1.0$



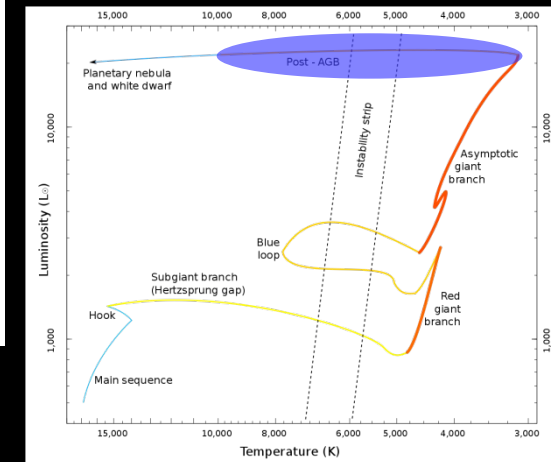
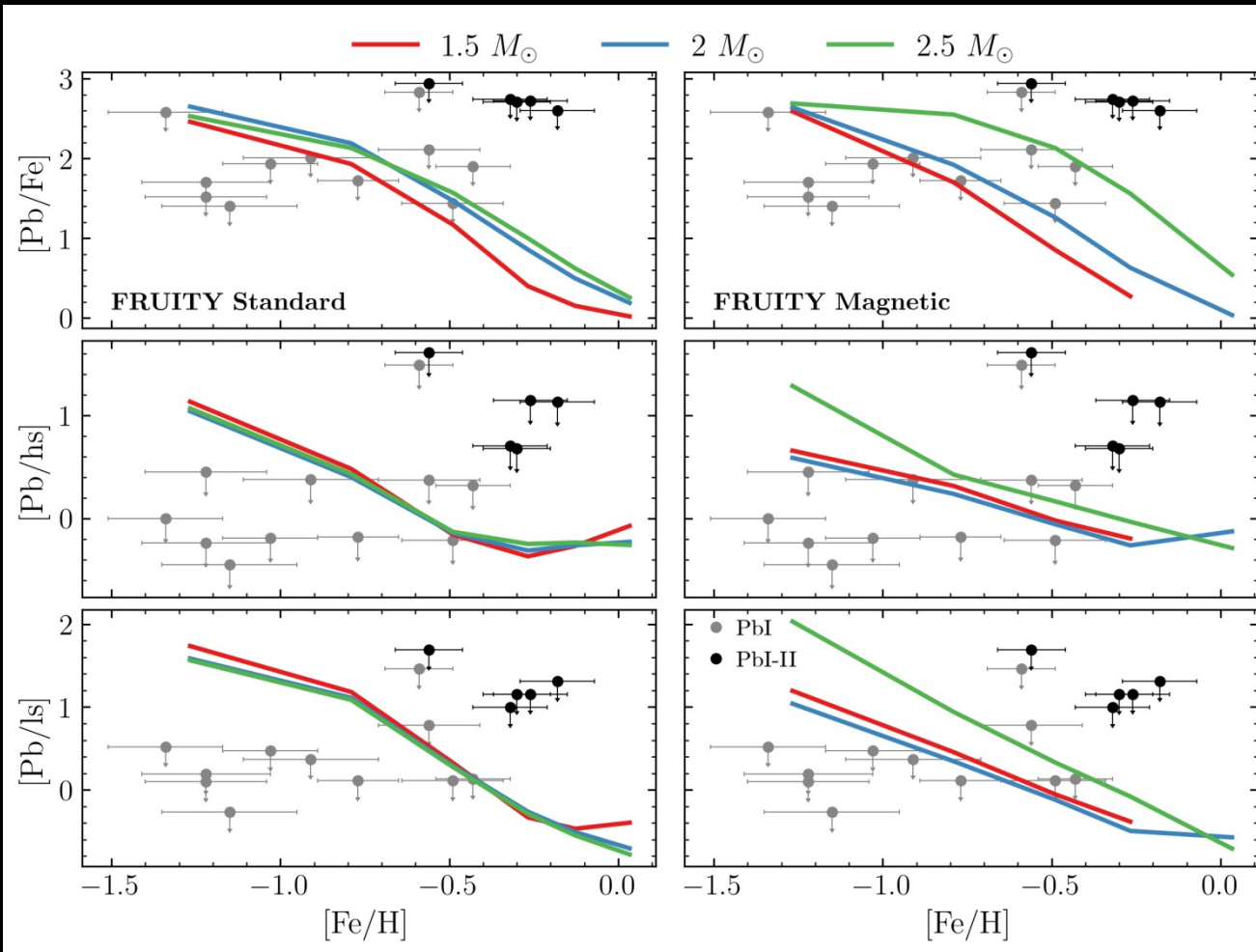
Post-AGB stars



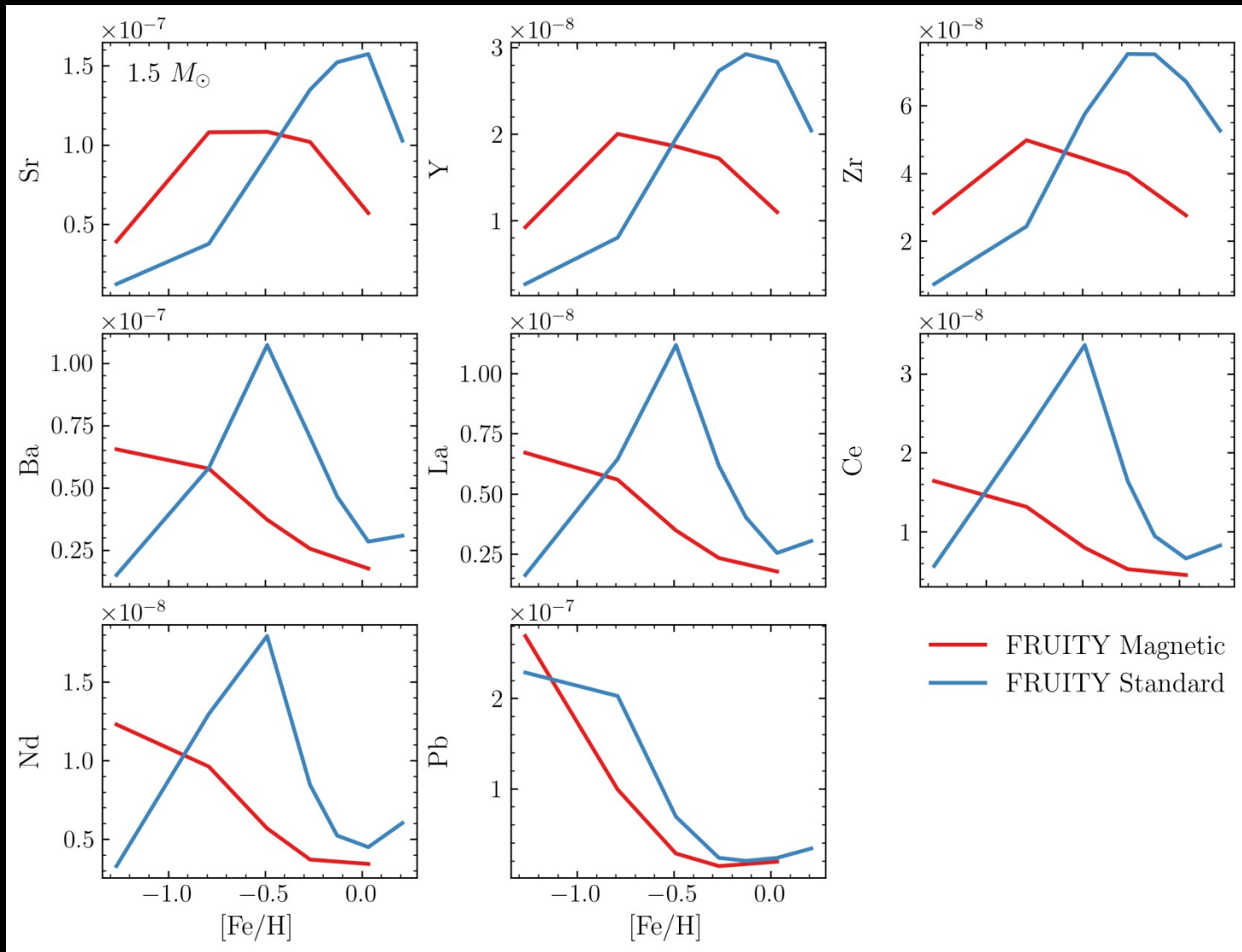
Post-AGB stars



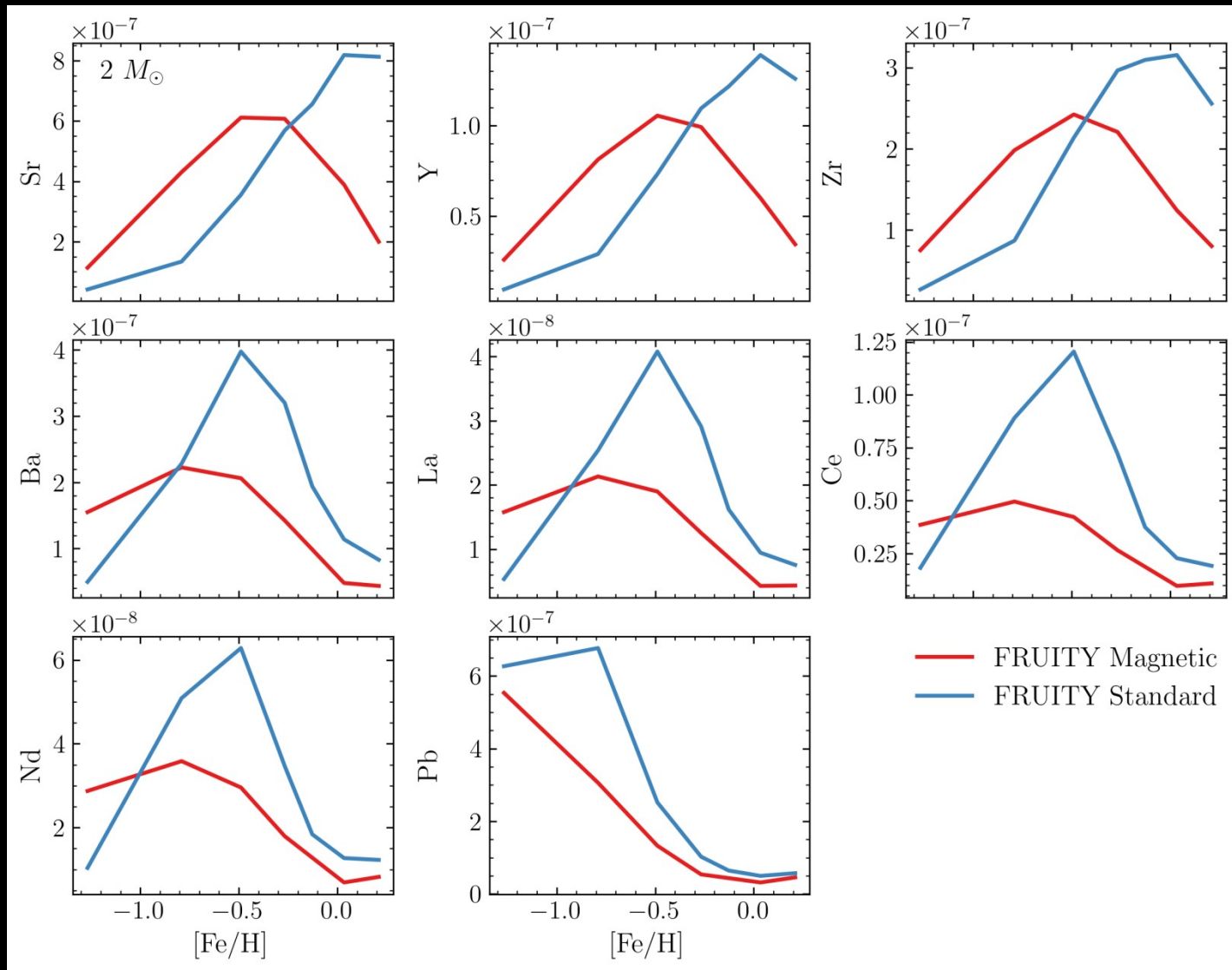
Post-AGB stars



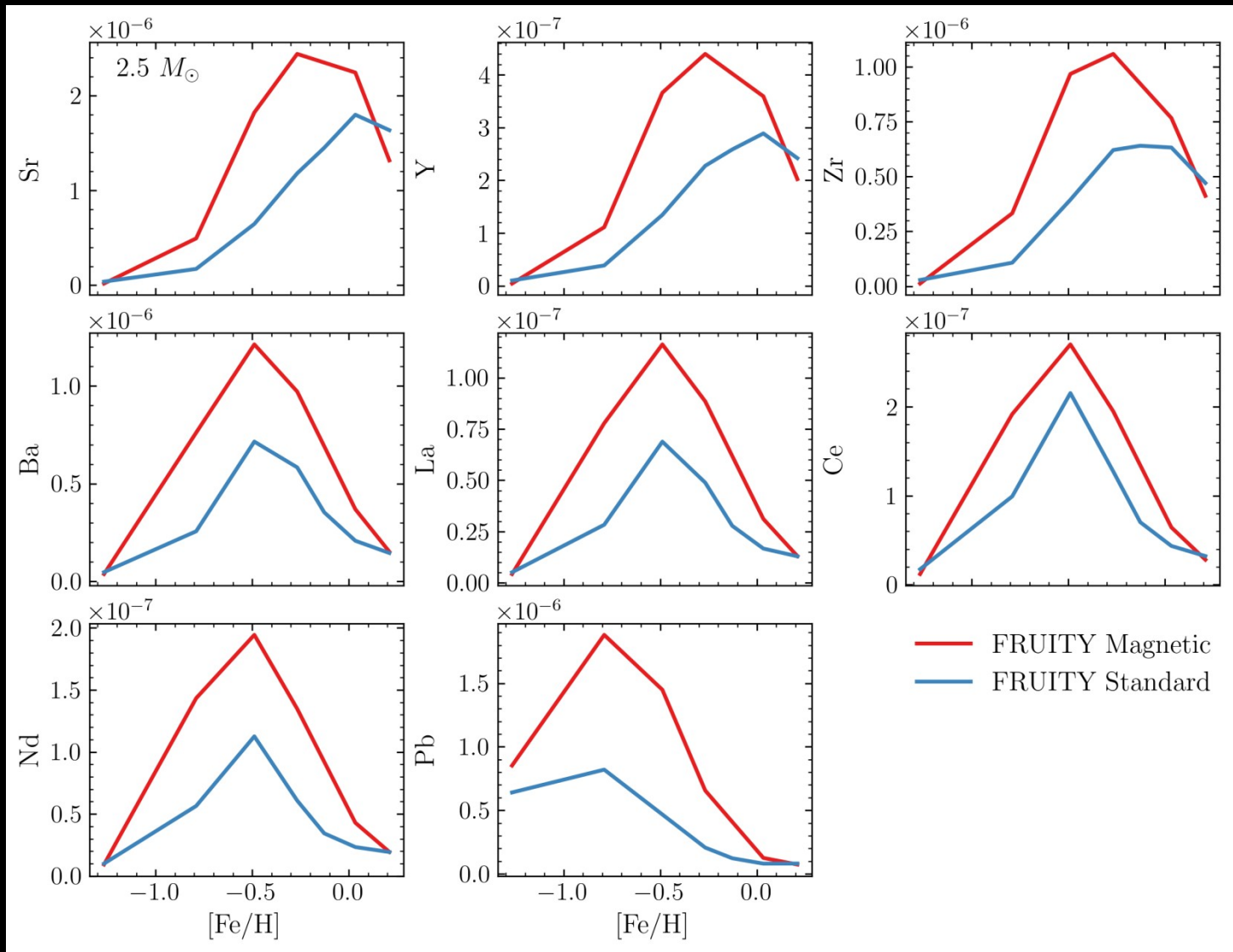
...and finally YIELDS!



...and finally YIELDS!



...and finally YIELDS!



...and finally YIELDS!

