

Digging for cosmological constraints out of high odds J-PLUS DR3 galaxies

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Outline

- The *r<21, odds>0.8* peculiar sample: dN/dz, and impact level from additive and multiplicative systematics
- Tomography from *z=0.05* up to *z=0.25* with 2D clustering (**ADF**) and angular redshift fluctuations (**ARF**)
- Modeling data in the deep non-linear regime with linear theory: the reference of the *MICE* mock catalogue
- \cdot Comparison with J-PLUS DR3. Tomographic constraints on the bias, peculiar velocities and lensing of the CMB

About \sim 2,800 sq.deg

APM = Andrés del Pino Molina's galaxy catalog RvM = Rodrigo von Marten's galaxy catalog

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r<21, odds>0.8, APM

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APM TOMOGRAPHY

$$
W(z) = \frac{dN}{dz} \times \exp{-(z - z_{\text{obs}})^2 / (2\sigma_z^2)}
$$

* 3 different Gaussian widths: $\sigma_z = 0.01, 0.03, 0.05$

$$
n_g^{\text{obs}}(\hat{n}) = (n_g(\hat{n}) + \vec{\epsilon} \cdot \vec{\mathbf{M}})(\vec{\beta} \cdot \delta \vec{\mathbf{M}})
$$

RvM TOMOGRAPHY

$$
W(z) = \frac{dN}{dz} \times \exp{-(z - z_{\text{obs}})^2 / (2\sigma_z^2)}
$$

* 3 different Gaussian widths: $\sigma_z = 0.01, 0.03, 0.05$

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$$

PRE ADF, $z=0.07,\,\sigma_z=0.05$

POST ADF $z=0.07,\:\sigma_z=0.05$

RvM

PRE - POST, $z=0.07,\,\sigma_z=0.05$

2D clustering, source counts in footprint, angular density fluctuations (ADF)

RvM

POST ADF $z = 0.07$, $\sigma_z = 0.05$

RvM Angular redshift fluctuations (ARF) (Under any given redshift shell, much more Gaussian observable)

PRE ARF, $z = 0.07$, $\sigma_z = 0.05$

 δ_z

ADF covariance matrices, APM

 \boldsymbol{z}

 $\rm 0.35$

 0.05

 0.15

0.25

 \overline{z}

 0.25

 \boldsymbol{z}

 $\,0.35\,$

 -0.50

 -0.75

 -1.00

 $\rm 0.05$

 $0.15\,$

 $\rm 0.35$

 0.35

ADF covariance matrices, RvM

0.35

0.35

ARF covariance matrices, APM

 \boldsymbol{z}

For **both** ADF and ARF the observed angular power spectra break like this:

 $\overline{175}$

200

75

100

Multipole ℓ

125

150

50

 $\overline{25}$

Unlike the ADF, **the ARF are sensitive to errors in photometric redshifts:**

$$
\delta z^{\text{photo}}(\hat{n}) = \frac{1}{N} \int dz \frac{d\bar{N}}{dz} (1 + \delta_g) \left(z_H + z_{\text{vel}} + z_{\text{error}} - \bar{z} \right) \exp \left[- (z_H + z_{\text{vel}} + z_{\text{error}} - z_{\text{center}_j})^2 / (2\sigma_z^2) \right] \Rightarrow
$$

$$
\langle (\delta z^{\text{photo}})^2(\hat{n}) \rangle^2 \simeq \exp - \left[(\sigma_{\text{Err}}/\sigma_z)^2 \right] \langle \delta z^2(\hat{n}) \rangle^2
$$

We shall be measuring the following set of parameters:

$$
\{\sigma_{\text{photo-z}}, b_{i=1,\text{nshell}}, A_{\text{vel}}\}\
$$

... or ...

$$
\{\sigma_{\text{photo}-z,i=1,\text{nshell}}, b_{i=1,\text{nshell}}, A_{\text{vel}}\}
$$

For **both** ADF and ARF the observed angular power spectra break like this:

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For **both** ADF and ARF the observed angular power spectra break like this:

Real J-PLUS DR3 **ADFs** (RvM)

J-PLUS DR3 is probing *J*-PLUS DR3 is probing
 deeply in the non-linear regime ...

J-PLUS DR3 is probing *deeply* in the non-linear regime ...

J-PLUS DR3 is probing *deeply* in the non-linear regime

MCMC on the parameter set

$$
\{\sigma_{\text{photo}-z}, b_{i=1,\text{nshell}}, A_{\text{vel}}\}
$$

obtained upon a "simplified" covariance matrix

MCMC on the parameter set

$$
\{\sigma_{\text{photo-z,i=1,nshell}}, b_{i=1,nshell}, A_{\text{vel}}\}
$$

obtained upon a "simplified" covariance matrix

Could the CMB help out here?

We can cross-correlate J-PLUS DR3 maps with maps of lensing convergence, that are sensitive to the projected gravitational potential. The *z-*window function for this cross-correlation peaks typically at *z~2,* but it is wide and there may be some signal with J-PLUS DR3 …

MCMC on the parameter set

$$
\{\sigma_{\text{photo-z}}, b_{i=1,\text{nshell}}, A_{\text{vel}}\}
$$

obtained upon a "simplified" covariance matrix

MCMC on the parameter set + CMB κ map

$$
\{\sigma_{\text{photo}-z}, b_{i=1,\text{nshell}}, A_{\text{vel}}\}
$$

obtained upon a "simplified" covariance matrix

Do we see a similar trend in **mocks**?

Let's look at the **MICE** simulation:

 $\Omega_m = 0.25; \ \Omega_\Lambda = 0.75; \ \Omega_b = 0.044; \ n_S = 0.95; \ h = 0.7; \ w_0 = -1; \ w_a = 0$

We impose the same *r<21* cut, but we cannot apply the same odds cut since we lack J-PLUS photometry. We impose the same *dN/dz:*

MICE original

 10°

The box is projected into an octant, 8,000 sq.deg

After fixing everything and neglecting photo-*z* errors, ADF typically provide **higher bias** values than ARF. The bias seems to decrease vs *z, contrary* to the expected behavior of the bias versus redshift, maybe hinting to probing scales that are *less* linear at higher *z-*s … To be confirmed with the real J-PLUS mock !

Conclusions:

- A full **pipeline** for conducting **ADF & ARF 2D tomography** on any LSS survey (J-PLUS, J-PAS, eBOSS, DES, *Euclid* …) is in place and working
- When applied on J-PLUS DR3, we find that the **linear model provides good fit to both ADF/ARF** observations up to *z~0.2.* Bias values of order unity, with a clear increasing trend in *z,* are found.
- The values of the bias are, however, clearly *discrepant.* This points to **different sensitivities** of ADF and ARF to *non-linear* effects, which itself is a good test for spotting non-linear contamination. This seems to be confirmed when looking at the MICE mock.
- J-PLUS DR3 is at best **mildly correlated to** *Planck'***s lensing** convergence map, which points to lower values of the bias than those inferred by ADF, pointing again to non-linear contribution.
- A **deeper** analysis of the **J-PLUS mock** (Izquierdo-Villalba et al. 2019) will be conducted before the submission of this work for publication (together with the systematics pipeline one — hopefully before the end of this 2023 year — *BTW, A&A or MNRAS?*