

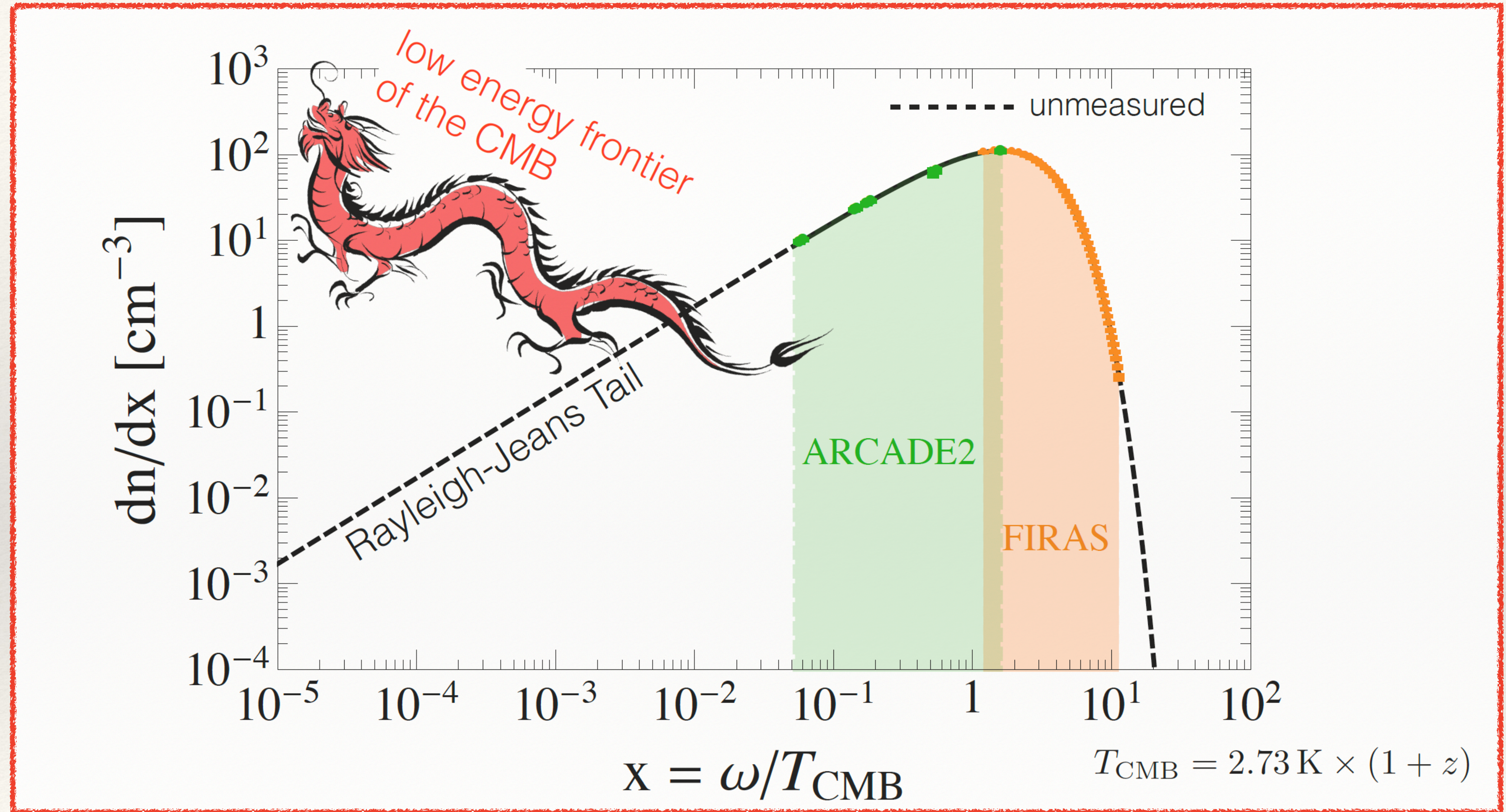
21cm and the RJ tail of the CMB

Andrea Caputo

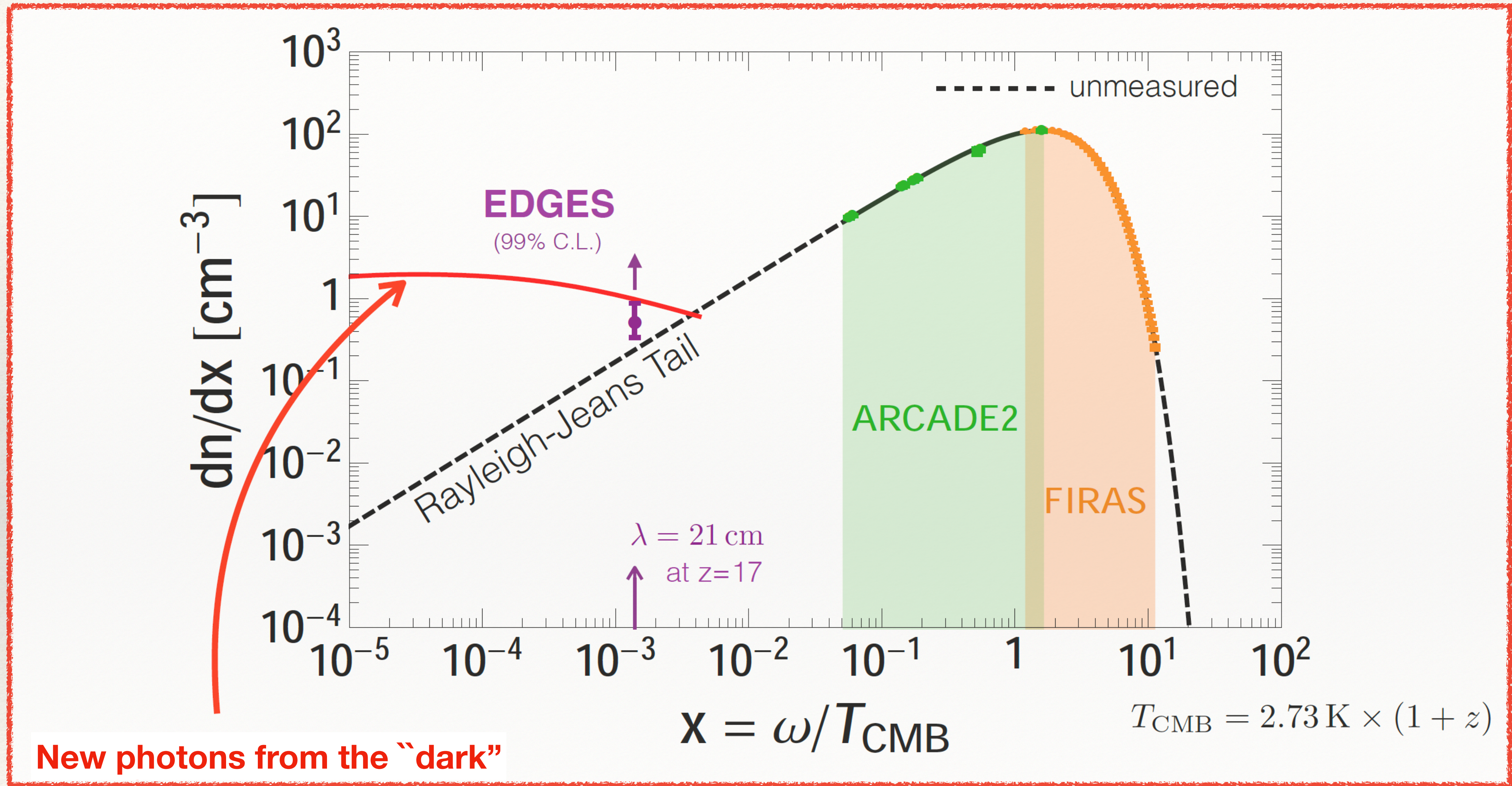
Trieste, 14 September 2023



Cosmic Microwave Background Spectrum



Cosmic Microwave Background Spectrum



21 cm, briefly

Brightness temperature

$$\Delta T_b^{21} \propto x_{HI} \left(1 - \frac{T_\gamma}{T_S} \right)$$

21 cm, briefly

Brightness temperature

$$\Delta T_b^{21} \propto x_{HI} \left(1 - \frac{T_\gamma}{T_S} \right)$$

$$T_S > T_\gamma \rightarrow \Delta T_b^{21} > 0 \text{ (emission)}$$

$$T_S < T_\gamma \rightarrow \Delta T_b^{21} < 0 \text{ (absorption)}$$

21 cm, briefly

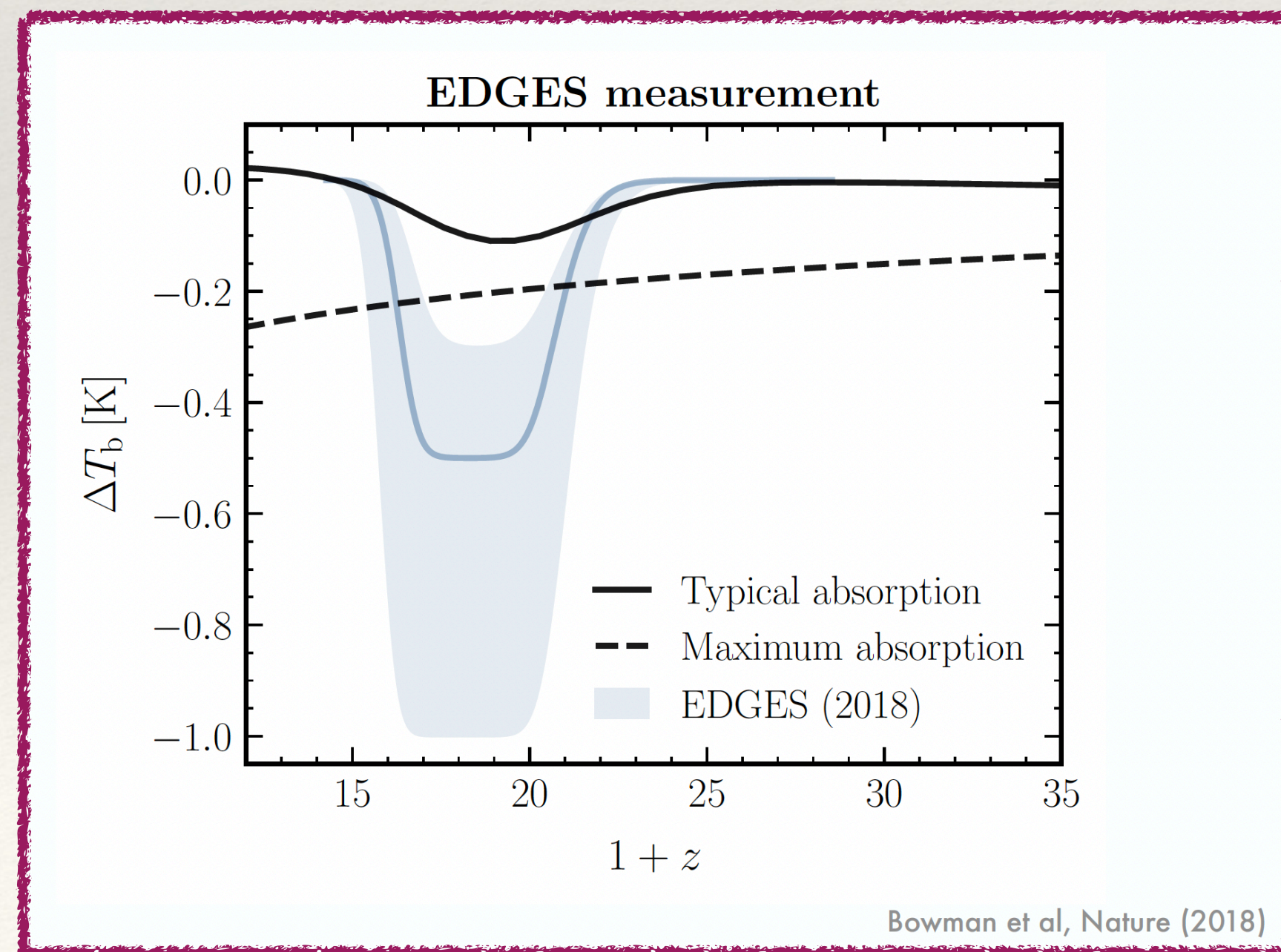
Brightness temperature

$$\Delta T_b^{21} \propto x_{HI} \left(1 - \frac{T_\gamma}{T_S} \right)$$

$$T_S > T_\gamma \rightarrow \Delta T_b^{21} > 0 \text{ (emission)}$$

$$T_S < T_\gamma \rightarrow \Delta T_b^{21} < 0 \text{ (absorption)}$$

More absorption than we thought?



21 cm, briefly

Brightness temperature

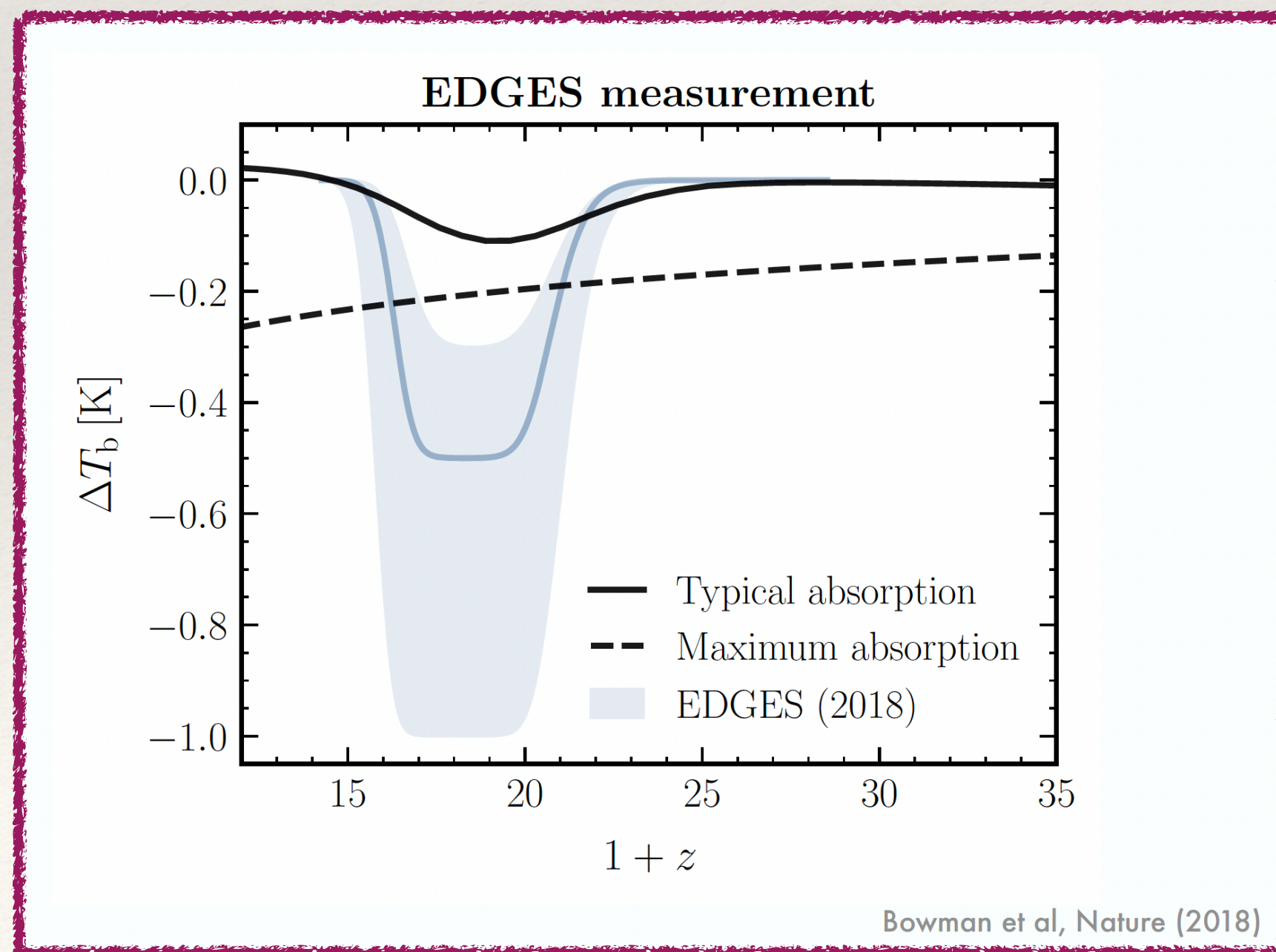
$$\Delta T_b^{21} \propto x_{HI} \left(1 - \frac{T_\gamma}{T_S} \right)$$

$$T_S > T_\gamma \rightarrow \Delta T_b^{21} > 0 \text{ (emission)}$$

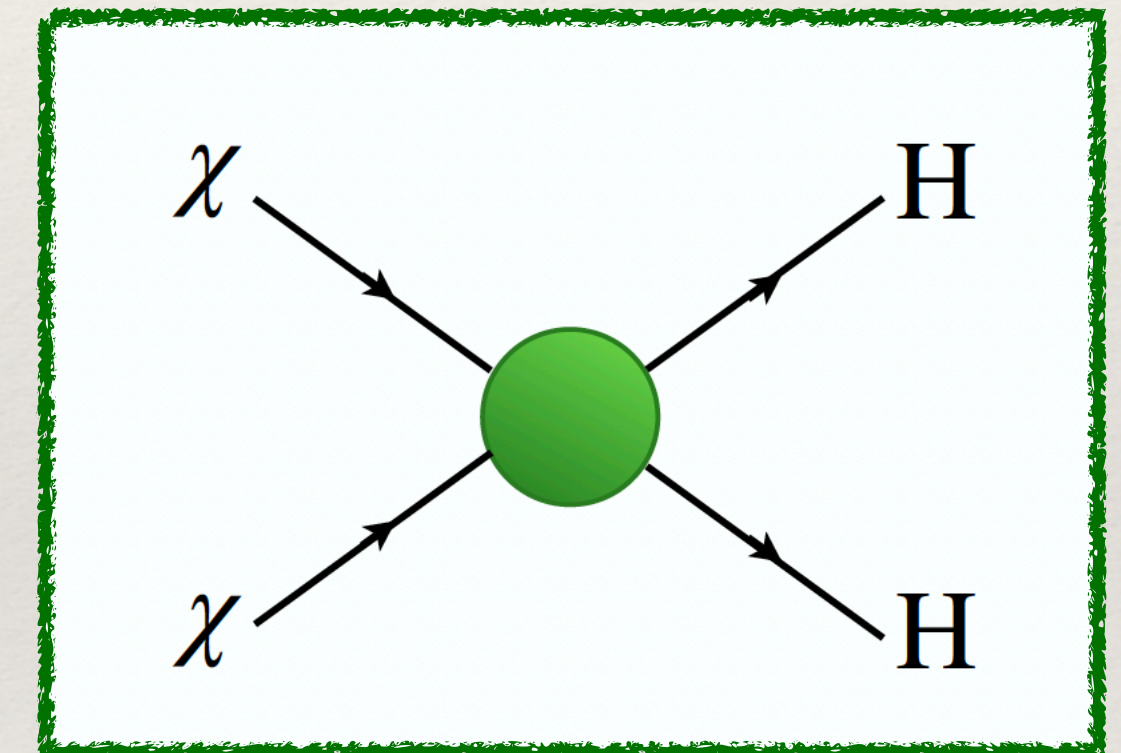
$$T_S < T_\gamma \rightarrow \Delta T_b^{21} < 0 \text{ (absorption)}$$

One can cool baryons

More absorption than we thought?



Muñoz & Loeb [1802.10094]
Falkowski & Petraki [1803.10096]
Barkana [1803.06698]
Barkana et al [1803.03091]
Berlin et al [1803.02804]
Liu et al [1908.06986]



21 cm, briefly

Brightness temperature

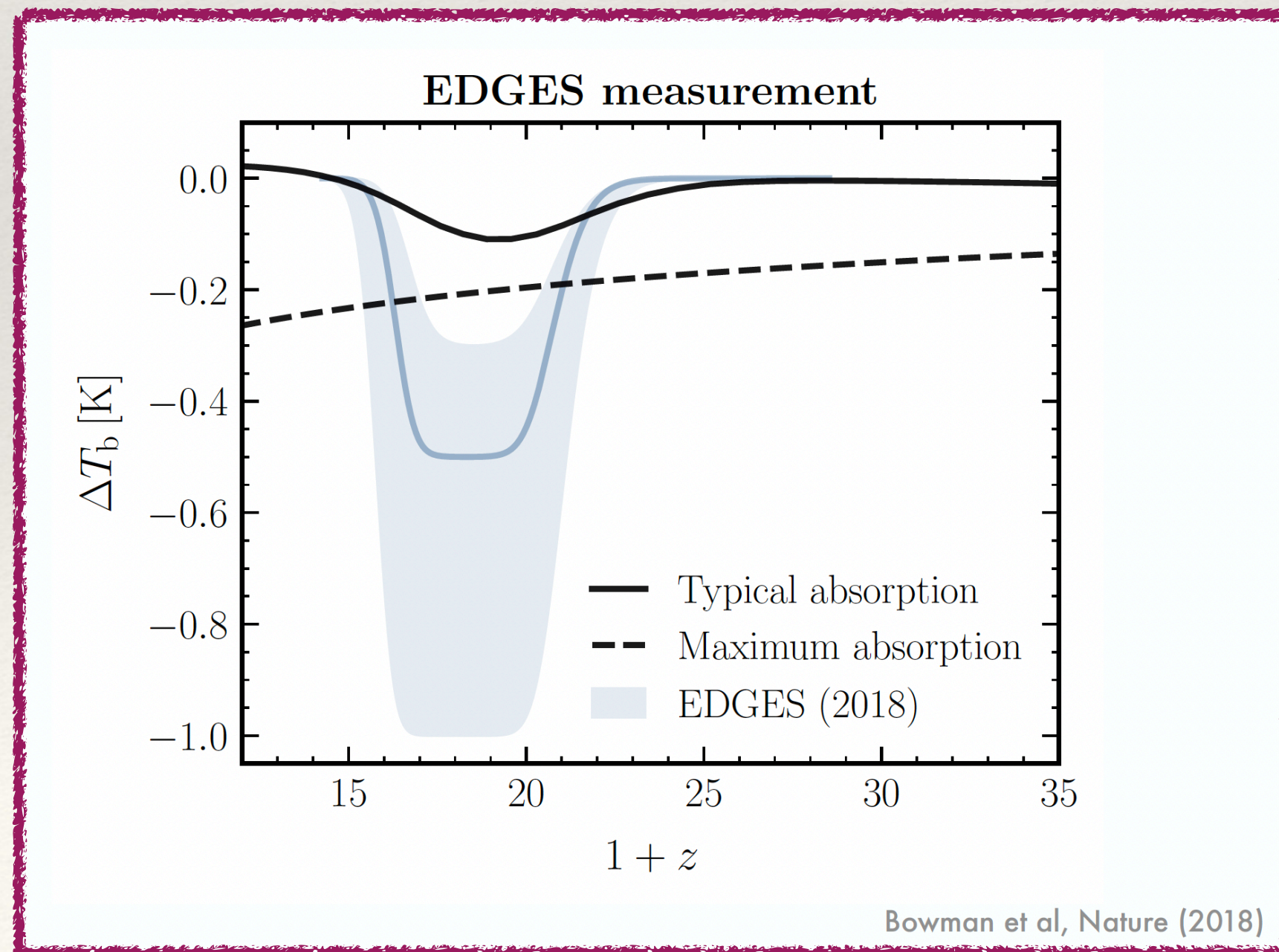
$$\Delta T_b^{21} \propto x_{HI} \left(1 - \frac{T_\gamma}{T_S} \right)$$

$$T_S > T_\gamma \rightarrow \Delta T_b^{21} > 0 \text{ (emission)}$$

$$T_S < T_\gamma \rightarrow \Delta T_b^{21} < 0 \text{ (absorption)}$$

One can cool baryons

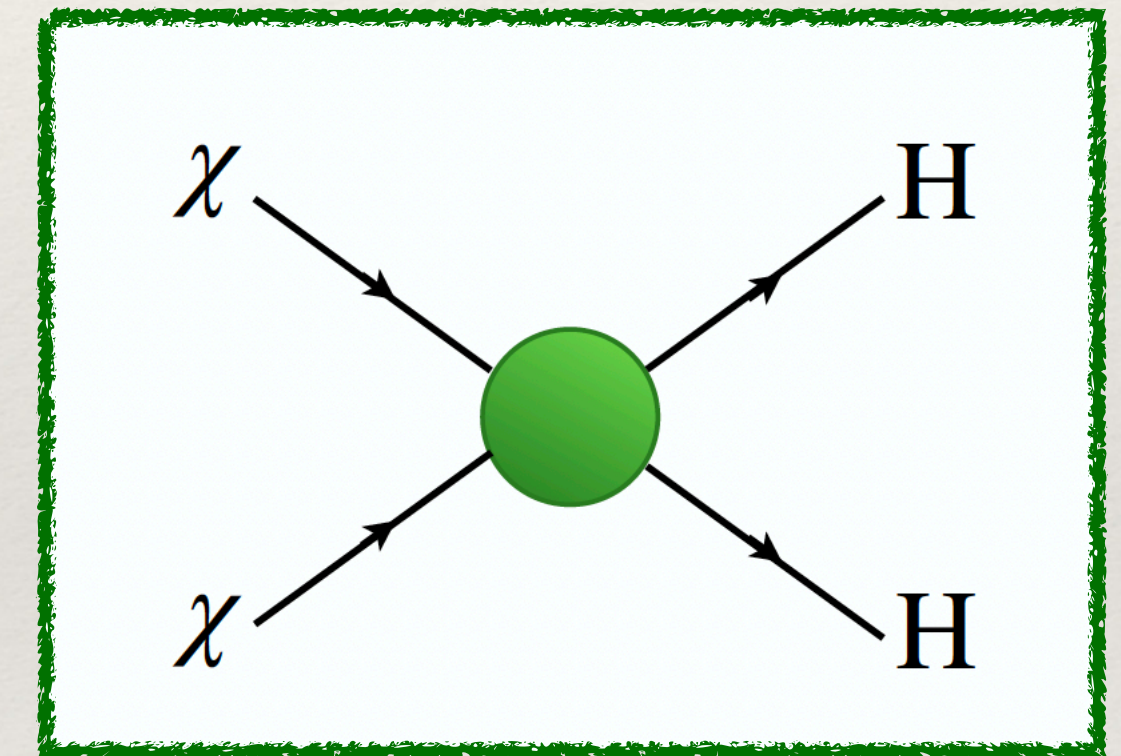
More absorption than we thought?



- Muñoz & Loeb [1802.10094]
- Falkowski & Petraki [1803.10096]
- Barkana [1803.06698]
- Barkana et al [1803.03091]
- Berlin et al [1803.02804]
- Liu et al [1908.06986]

One can add photons!

- Pospelov et al [1803.07048]
- Moroi, Nakayama, Tang [1804.10378]
- Choi, Seong, Yun [1911.00532]



$$n_\gamma^{\text{inj}} > n_{\text{CMB}} \quad (\lambda = 21\text{cm})$$

21 cm, briefly

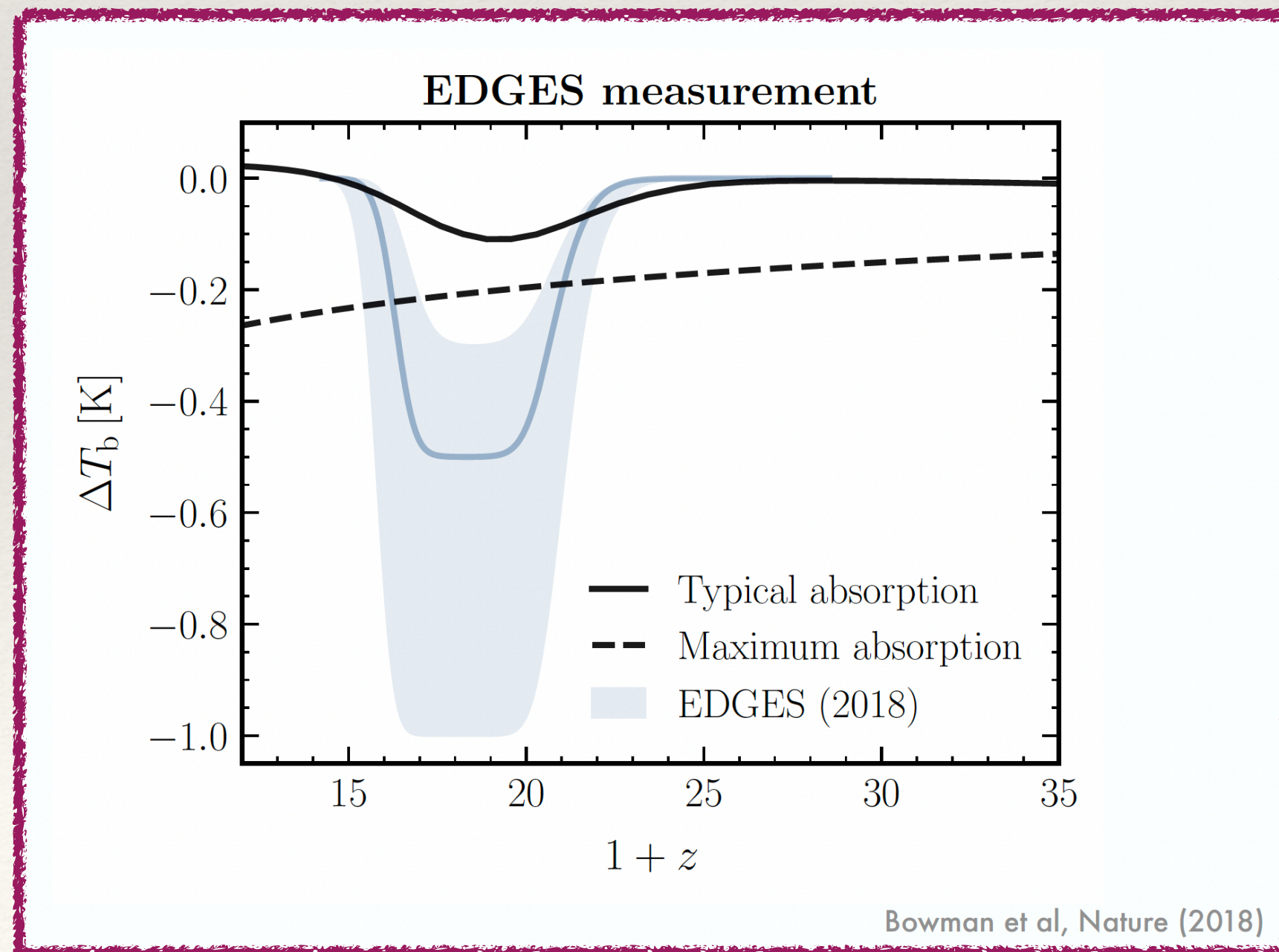
Brightness temperature

$$\Delta T_b^{21} \propto x_{HI} \left(1 - \frac{T_\gamma}{T_S} \right)$$

$$T_S > T_\gamma \rightarrow \Delta T_b^{21} > 0 \text{ (emission)}$$

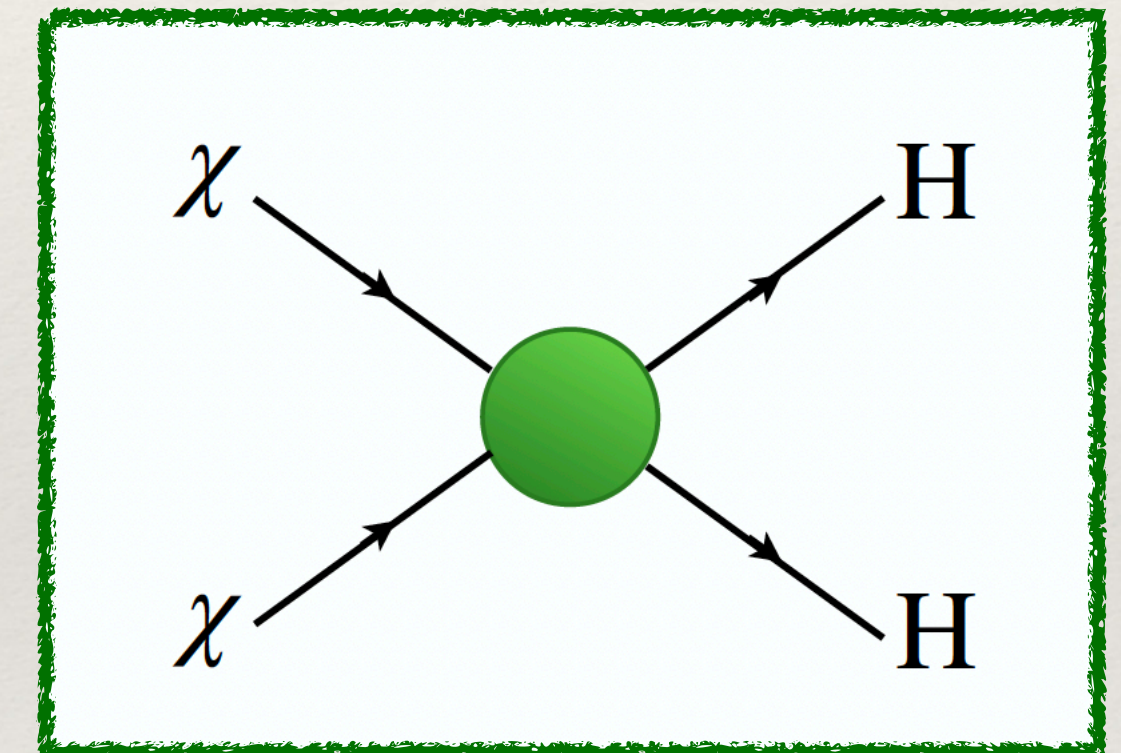
$$T_S < T_\gamma \rightarrow \Delta T_b^{21} < 0 \text{ (absorption)}$$

More absorption than we thought?



One can cool baryons

- Muñoz & Loeb [1802.10094]
- Falkowski & Petraki [1803.10096]
- Barkana [1803.06698]
- Barkana et al [1803.03091]
- Berlin et al [1803.02804]
- Liu et al [1908.06986]



One can add photons!

$$n_\gamma^{\text{inj}} > n_{\text{CMB}} \quad (\lambda = 21\text{cm})$$

$$n_{\text{RJ}} = \frac{1}{\pi^2} \int_0^{\omega_{\text{max}}} \frac{\omega^2 d\omega}{\exp[\omega/T] - 1} \simeq \frac{T\omega_{\text{max}}^2}{2\pi^2}$$

$$\simeq 0.21 x_{\text{max}}^2 n_{\text{CMB}}, \quad \hbar = c = k = 1 \text{ units}$$

$$n_{\text{RJ}} \sim 10^{-6} n_{\text{CMB}}$$

21 cm, briefly

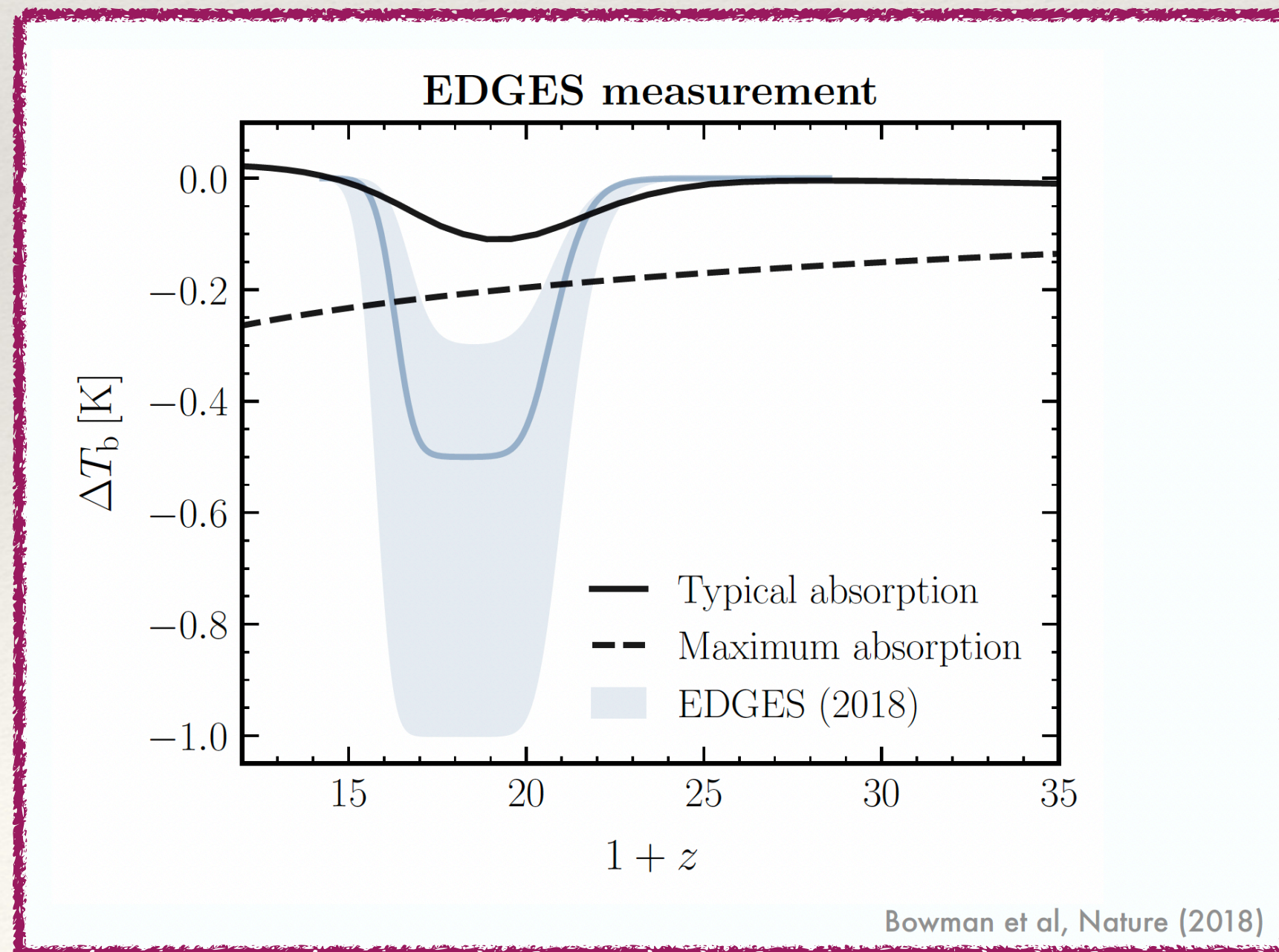
Brightness temperature

$$\Delta T_b^{21} \propto x_{HI} \left(1 - \frac{T_\gamma}{T_S} \right)$$

$$T_S > T_\gamma \rightarrow \Delta T_b^{21} > 0 \text{ (emission)}$$

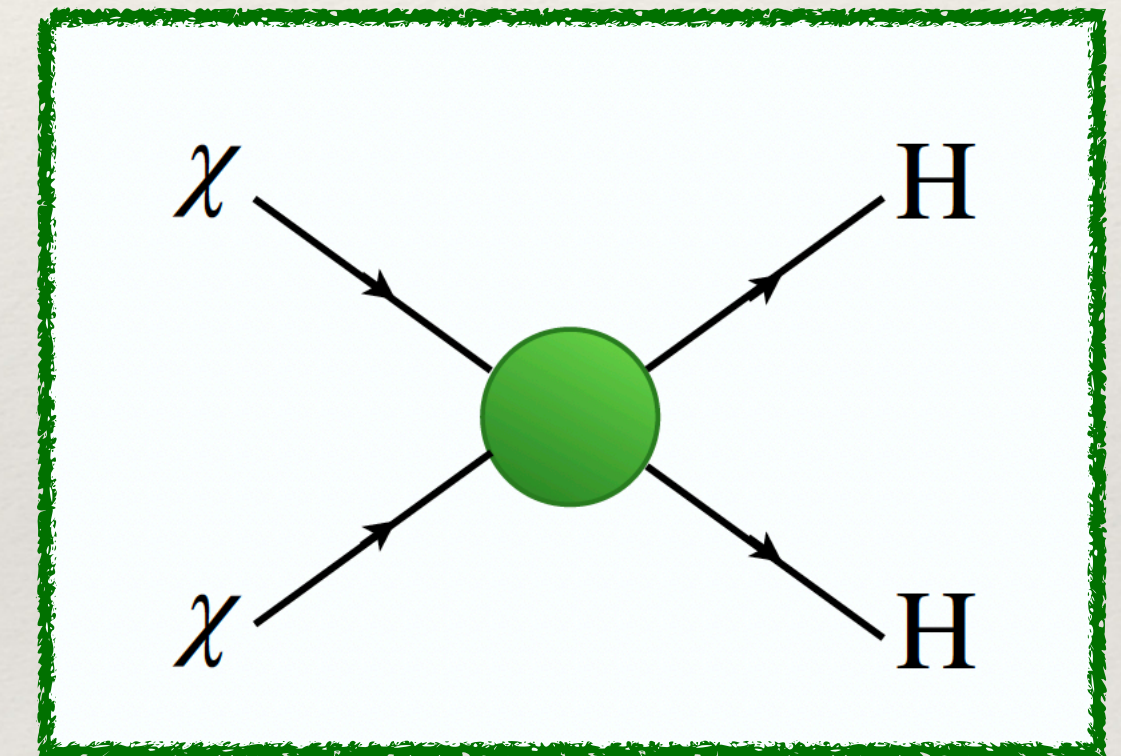
$$T_S < T_\gamma \rightarrow \Delta T_b^{21} < 0 \text{ (absorption)}$$

More absorption than we thought?



One can cool baryons

- Muñoz & Loeb [1802.10094]
- Falkowski & Petraki [1803.10096]
- Barkana [1803.06698]
- Barkana et al [1803.03091]
- Berlin et al [1803.02804]
- Liu et al [1908.06986]



One can add photons!

$$n_\gamma^{\text{inj}} > n_{\text{CMB}} \quad (\lambda = 21\text{cm})$$

$$n_{\text{RJ}} = \frac{1}{\pi^2} \int_0^{\omega_{\text{max}}} \frac{\omega^2 d\omega}{\exp[\omega/T] - 1} \simeq \frac{T\omega_{\text{max}}^2}{2\pi^2}$$

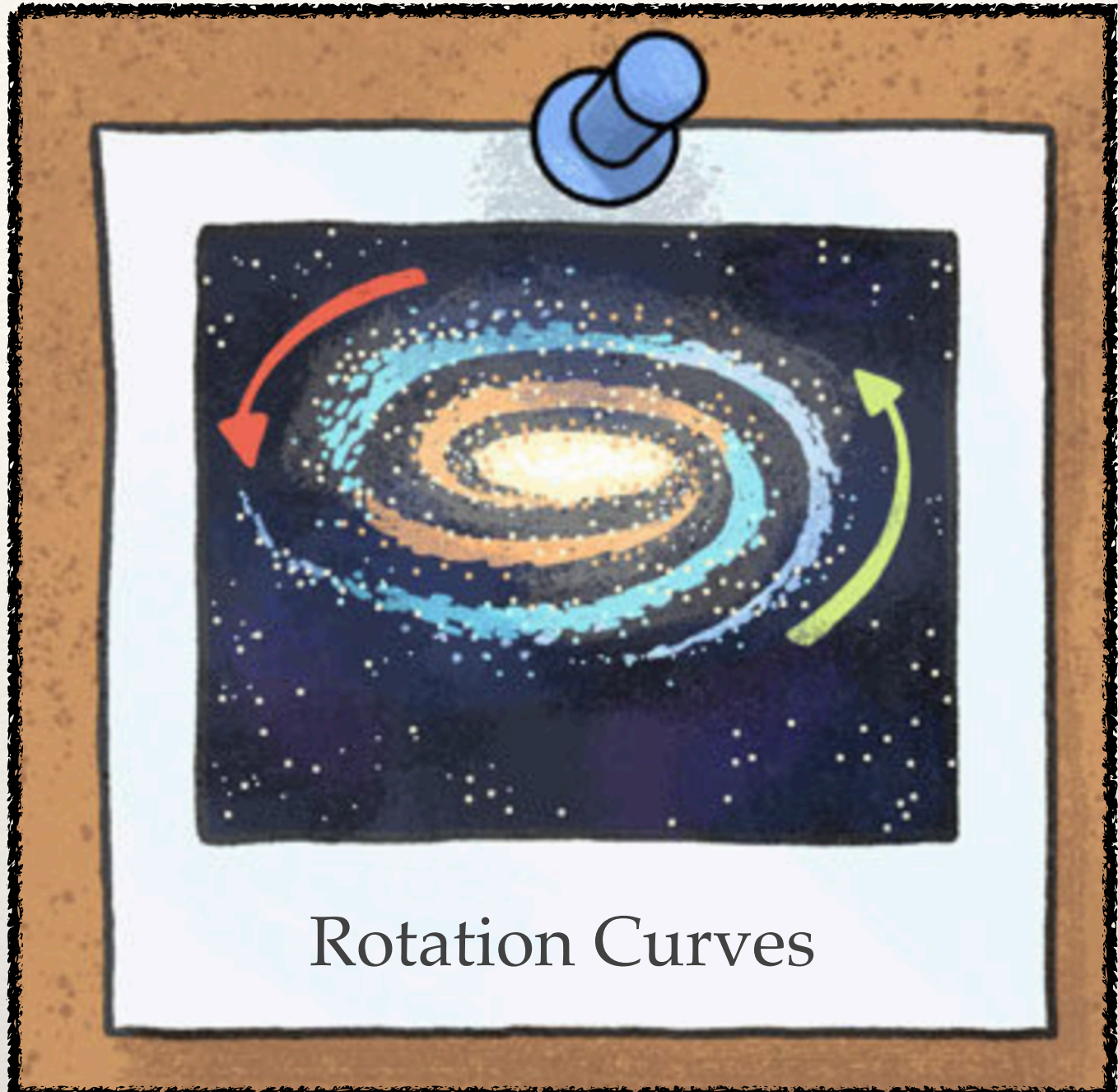
$$\simeq 0.21 x_{\text{max}}^2 n_{\text{CMB}}, \quad \hbar = c = k = 1 \text{ units}$$

$$n_{\text{RJ}} \sim 10^{-6} n_{\text{CMB}}$$

Add numerous, soft quanta

Particle physics motivation

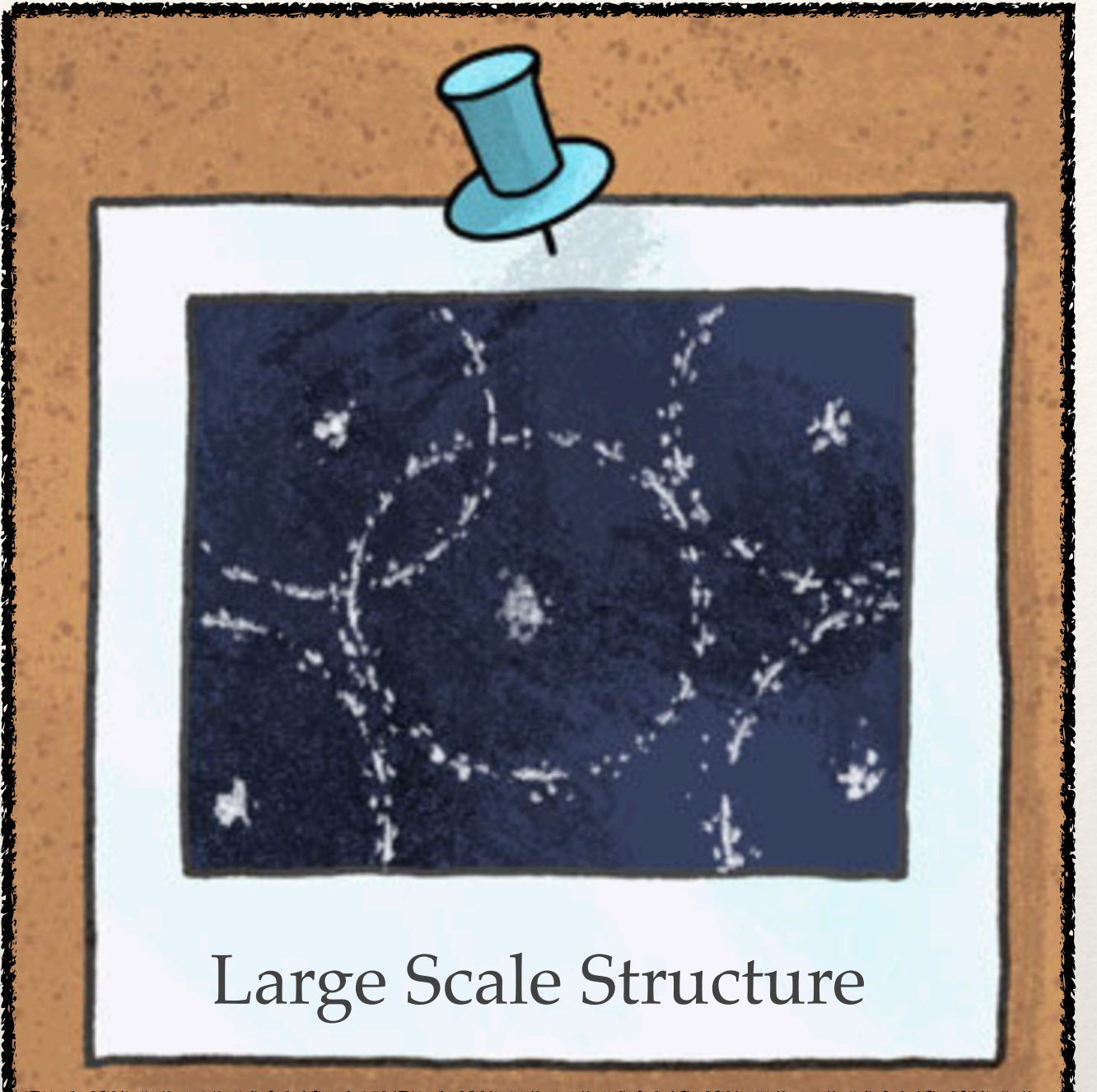
Dark Matter Mystery



Rotation Curves



Bullet Cluster



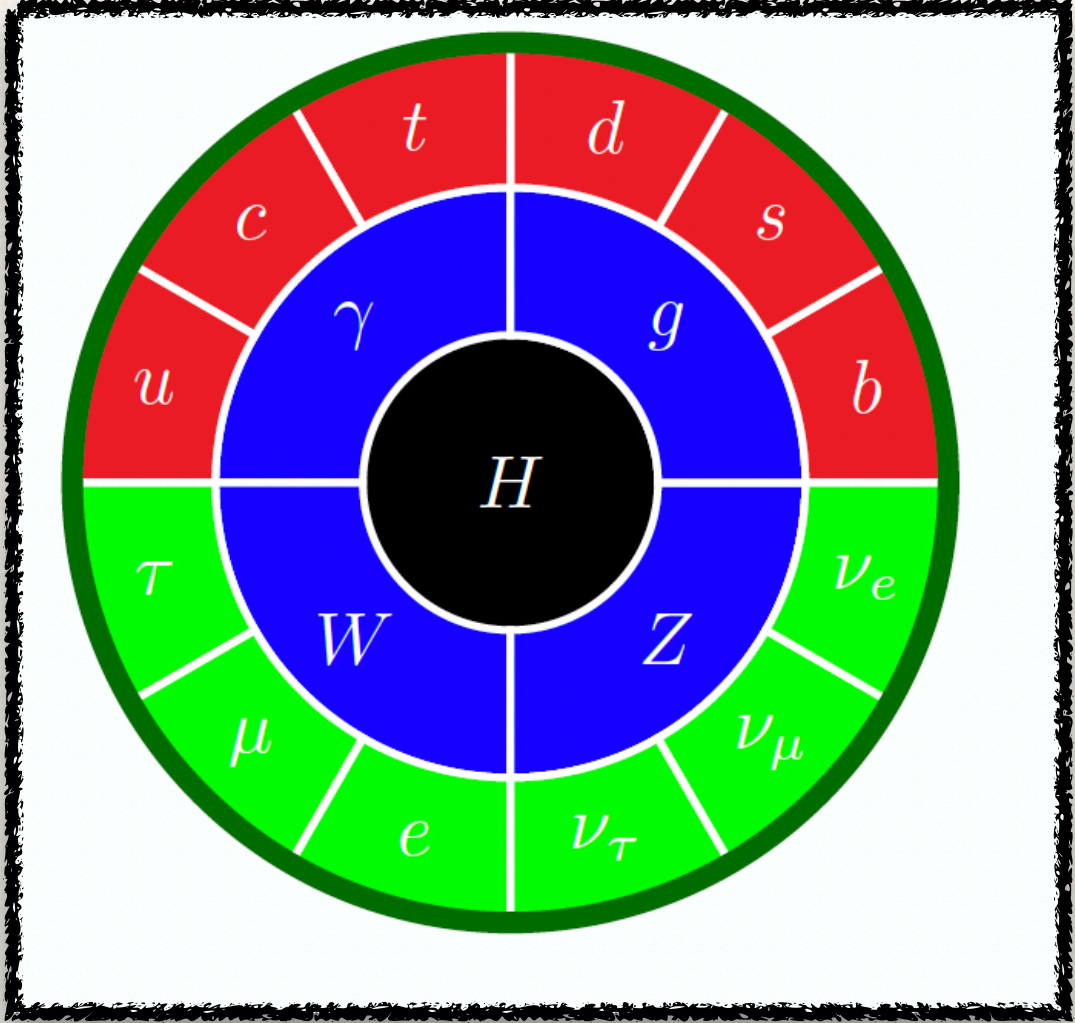
Large Scale Structure



Evidence at all scales!!



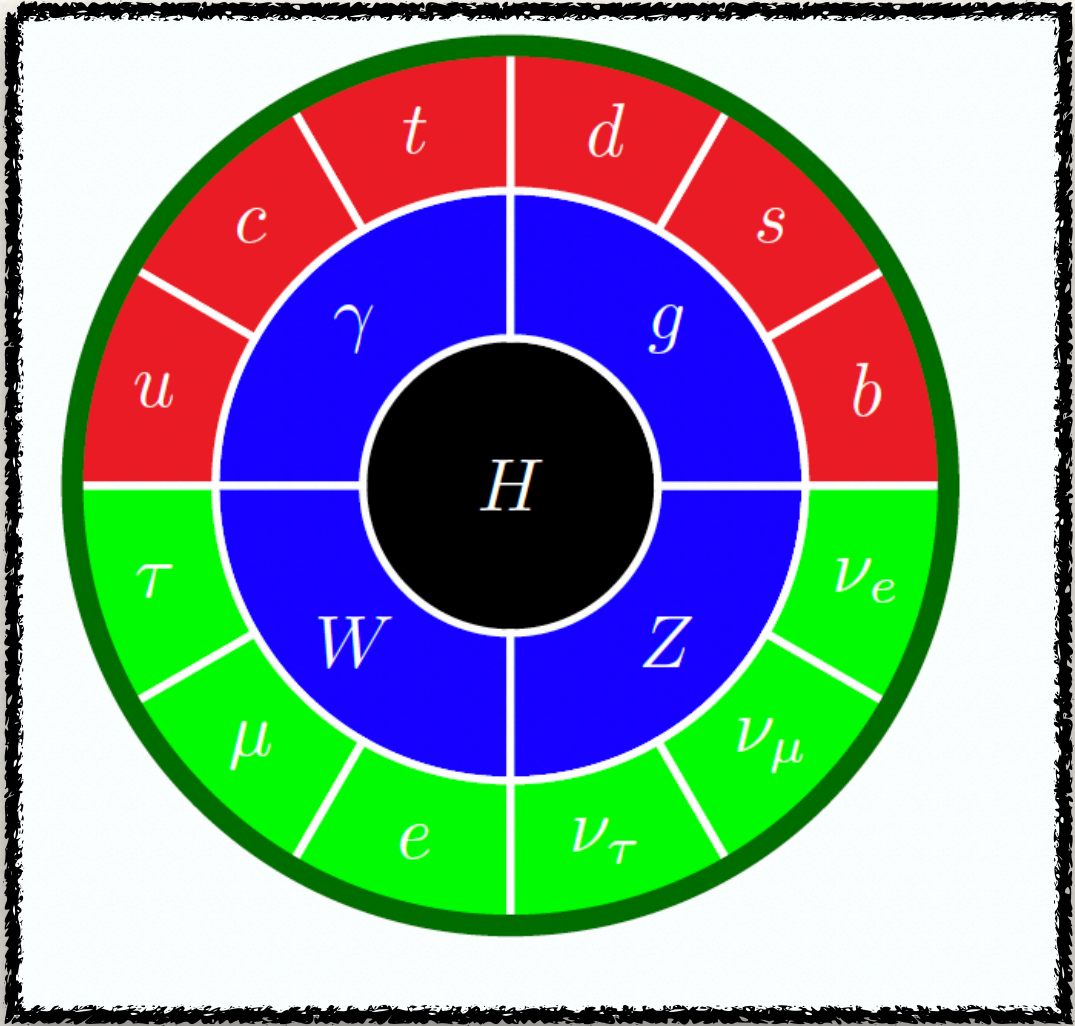
Standard Model



Dark Sector



Standard Model



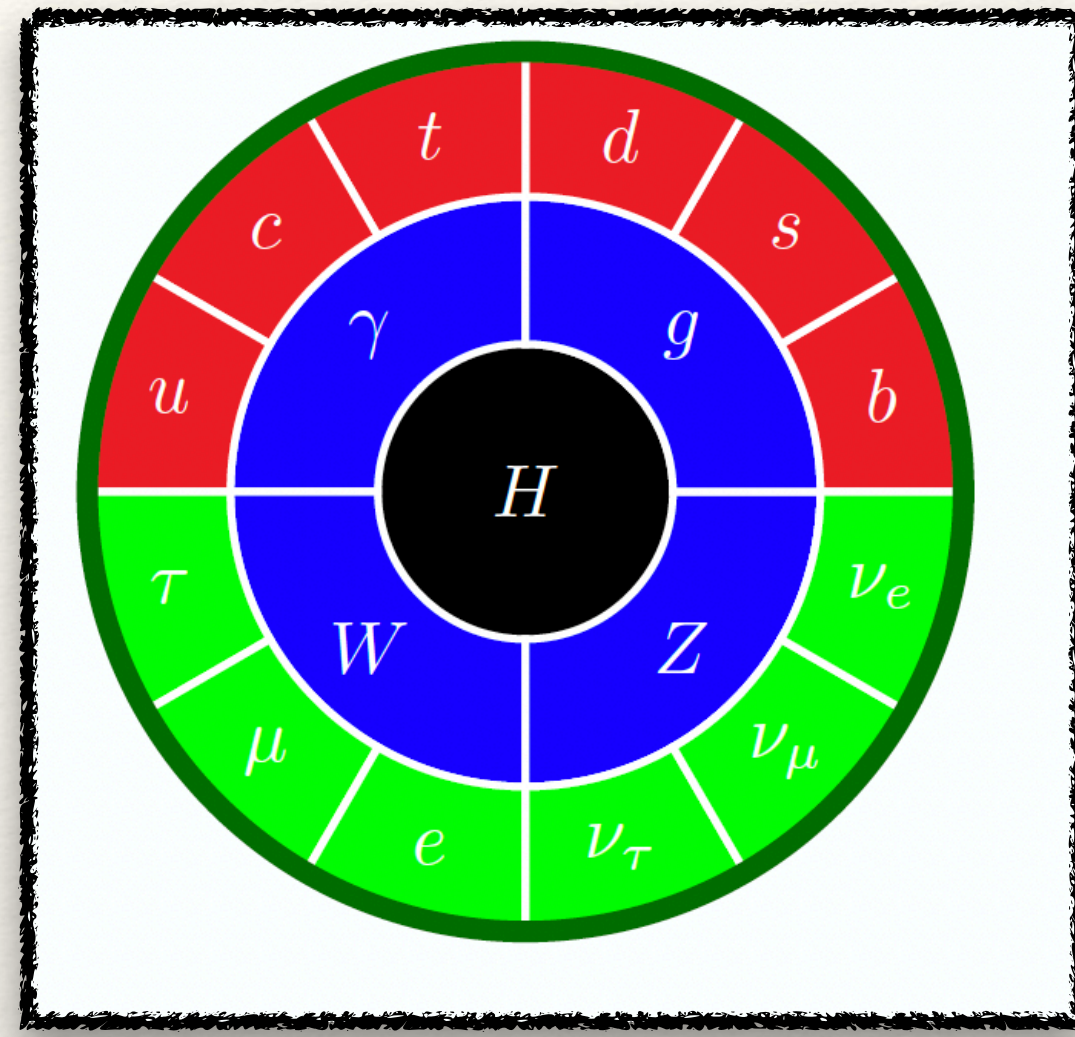
Portal



Dark Sector



Standard Model



Portal

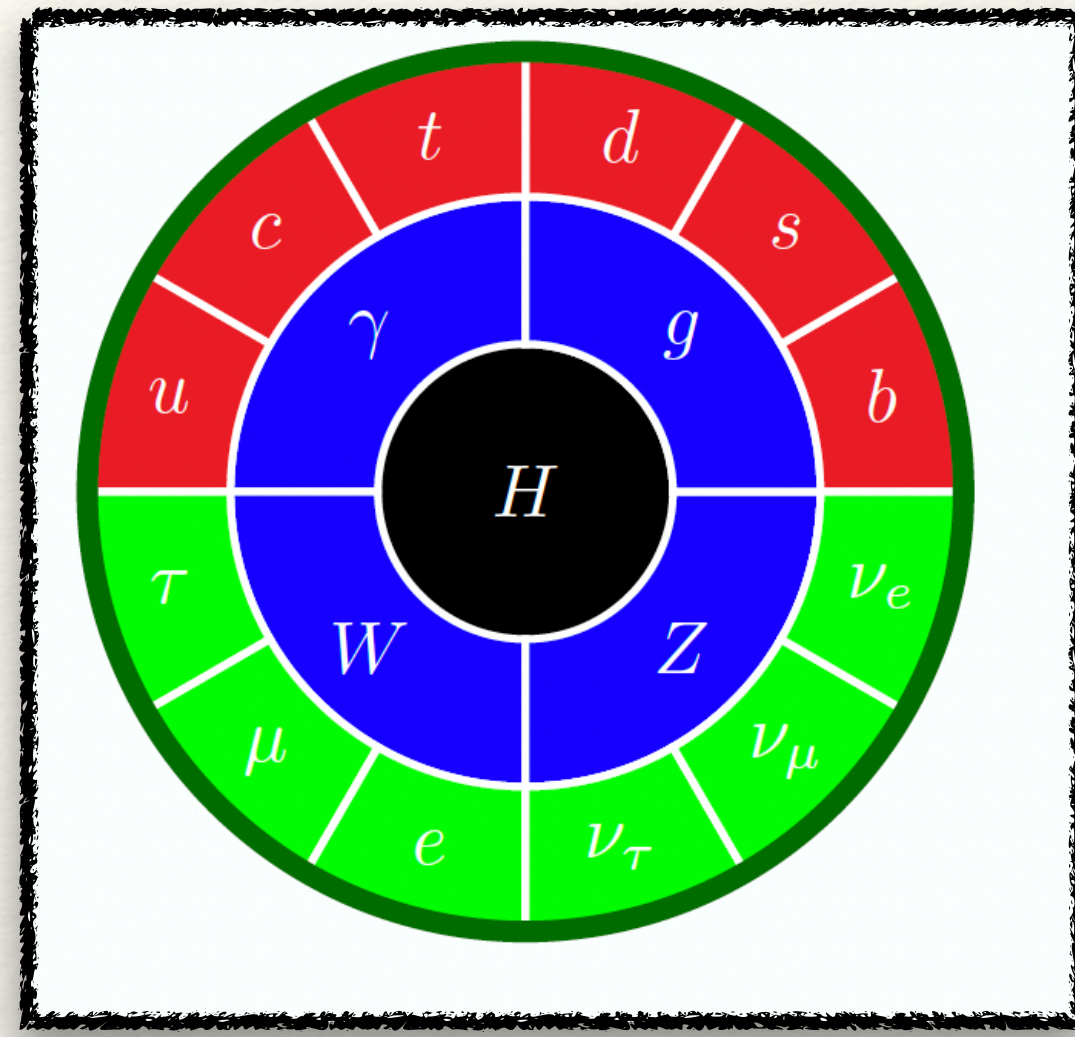


Dark Sector

(Some) Canonical portals:

- **Scalar**
Higgs portal $\lambda H^2 S^2 + \mu H^2 S$
- **Fermion**
Neutrino portal $y(HL)N$
- **Vector**
Kinetic mixing portal $\epsilon F^{\mu\nu} F'_{\mu\nu}$

Standard Model



Portal



Dark Sector

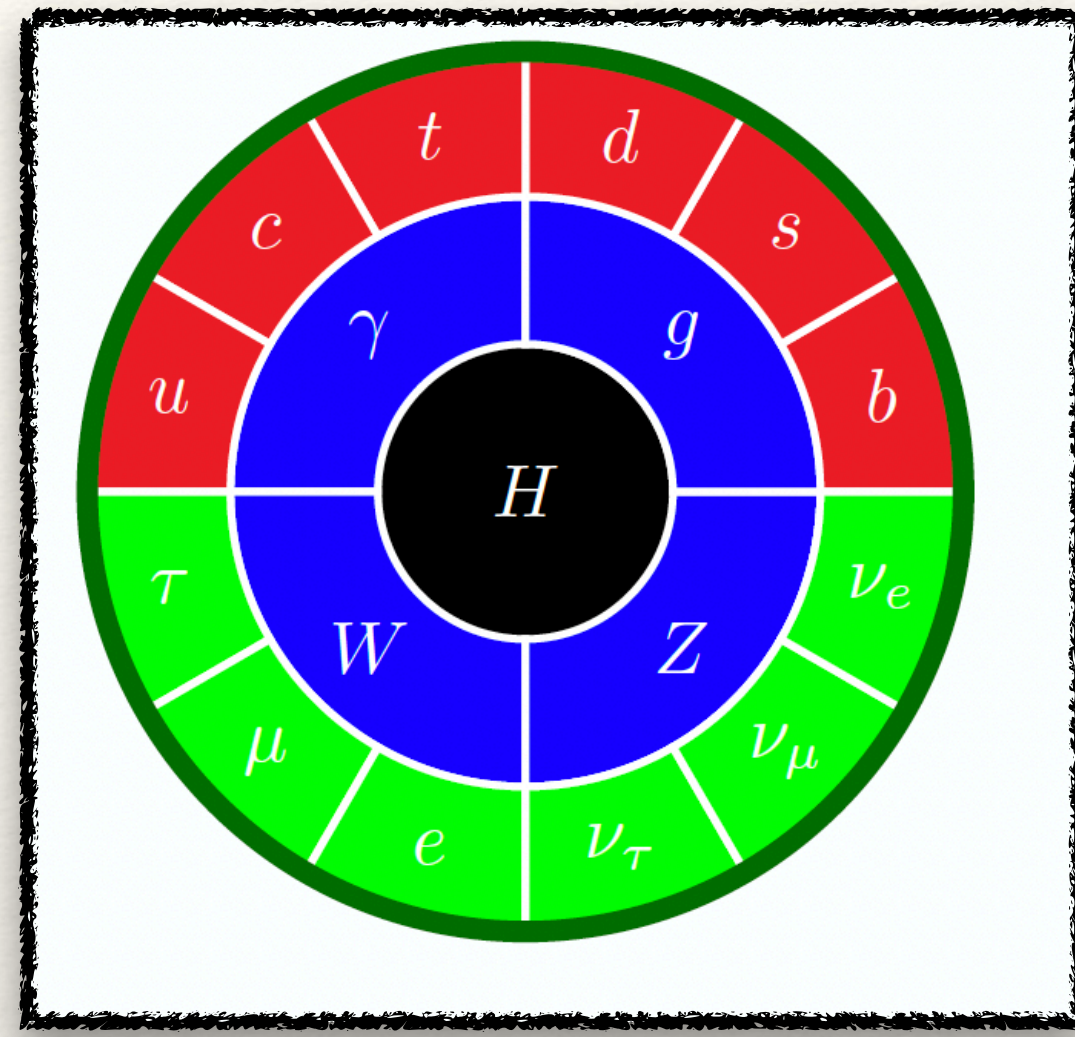
(Some) Canonical portals:

- **Scalar**
Higgs portal $\lambda H^2 S^2 + \mu H^2 S$

- **Fermion**
Neutrino portal $y(HL)N$

- **Vector**
Kinetic mixing portal $\epsilon F^{\mu\nu} F'_{\mu\nu}$

Standard Model



Portal



Dark Sector

(Some) Canonical portals:

- **Scalar**
Higgs portal $\lambda H^2 S^2 + \mu H^2 S$

- **Fermion**
Neutrino portal $y(HL)N$

- **Vector**
Kinetic mixing portal $\epsilon F^{\mu\nu} F'_{\mu\nu}$

Plus axion (like) dark matter $a F' \tilde{F}'$

Light DM a , decaying to two dark photons via and ALP coupling:

$$\mathcal{L} = \frac{1}{2}(\partial_\mu a)^2 - \frac{m_a^2}{2}a^2 + \frac{a}{4f_a}F'_{\mu\nu}\tilde{F}'^{\mu\nu} + \mathcal{L}_{AA'}$$

Light DM a , decaying to two dark photons via and ALP coupling:

$$\mathcal{L} = \frac{1}{2}(\partial_\mu a)^2 - \frac{m_a^2}{2}a^2 + \frac{a}{4f_a}F'_{\mu\nu}\tilde{F}'^{\mu\nu} + \mathcal{L}_{AA'}$$

Vector portal

Dark photon mixes with EM via “familiar” kinetic mixing

$$\mathcal{L}_{AA'} = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}(F'_{\mu\nu})^2 - \frac{\epsilon}{2}F_{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2(A'_\mu)^2 .$$

Light DM a , decaying to two dark photons via and ALP coupling:

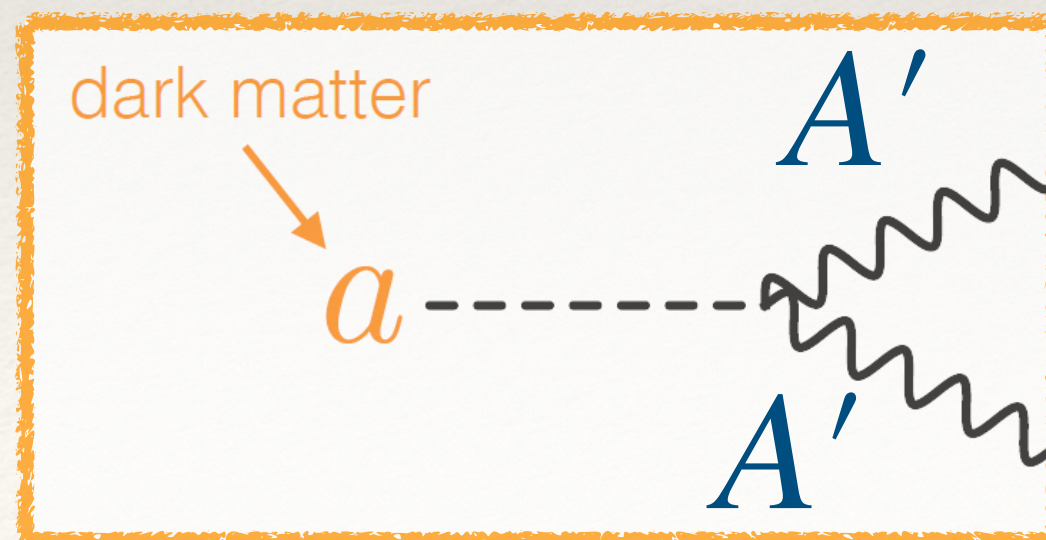
$$\mathcal{L} = \frac{1}{2}(\partial_\mu a)^2 - \frac{m_a^2}{2}a^2 + \frac{a}{4f_a}F'_{\mu\nu}\tilde{F}'^{\mu\nu} + \mathcal{L}_{AA'}$$

Vector portal

Dark photon mixes with EM via “familiar” kinetic mixing

$$\mathcal{L}_{AA'} = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}(F'_{\mu\nu})^2 - \frac{\epsilon}{2}F_{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2(A'_\mu)^2 .$$

Dark matter decays into dark photons



Light DM a , decaying to two dark photons via and ALP coupling:

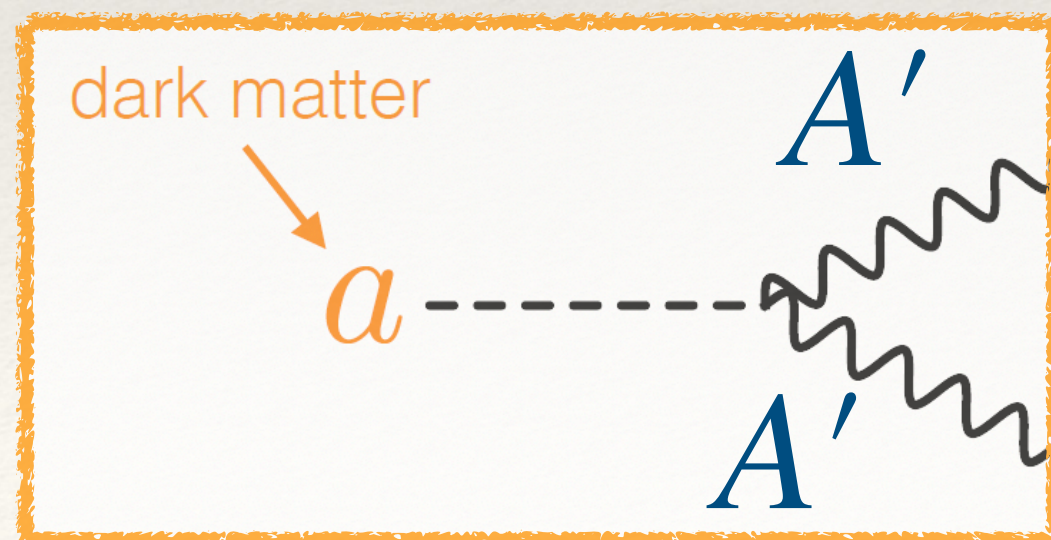
$$\mathcal{L} = \frac{1}{2}(\partial_\mu a)^2 - \frac{m_a^2}{2}a^2 + \frac{a}{4f_a}F'_{\mu\nu}\tilde{F}'^{\mu\nu} + \mathcal{L}_{AA'}$$

Vector portal

Dark photon mixes with EM via “familiar” kinetic mixing

$$\mathcal{L}_{AA'} = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}(F'_{\mu\nu})^2 - \frac{\epsilon}{2}F_{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2(A'_\mu)^2 .$$

Dark matter decays into dark photons



Dark photons resonantly convert into photons



Light DM a , decaying to two dark photons via and ALP coupling:

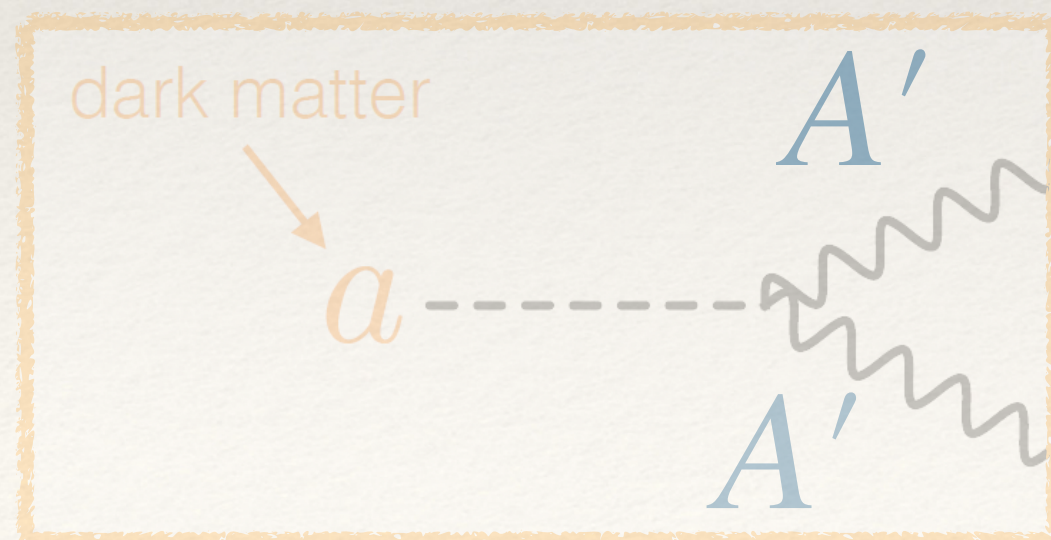
$$\mathcal{L} = \frac{1}{2}(\partial_\mu a)^2 - \frac{m_a^2}{2}a^2 + \frac{a}{4f_a}F'_{\mu\nu}\tilde{F}'^{\mu\nu} + \mathcal{L}_{AA'}$$

Vector portal

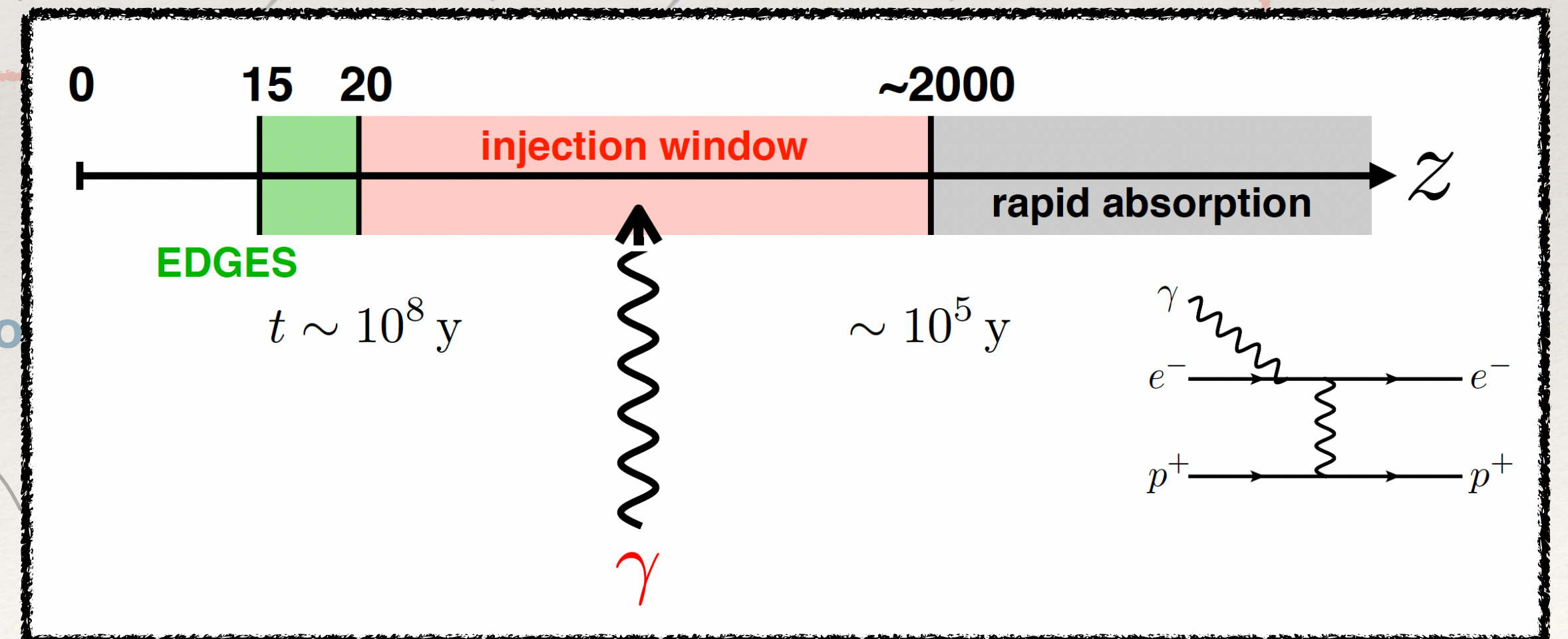
Dark photon mixes with EM via “familiar” kinetic mixing

$$\mathcal{L}_{AA'} = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}(F'_{\mu\nu})^2 - \frac{\epsilon}{2}F_{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2(A'_\mu)^2.$$

Dark matter decays into dark photons

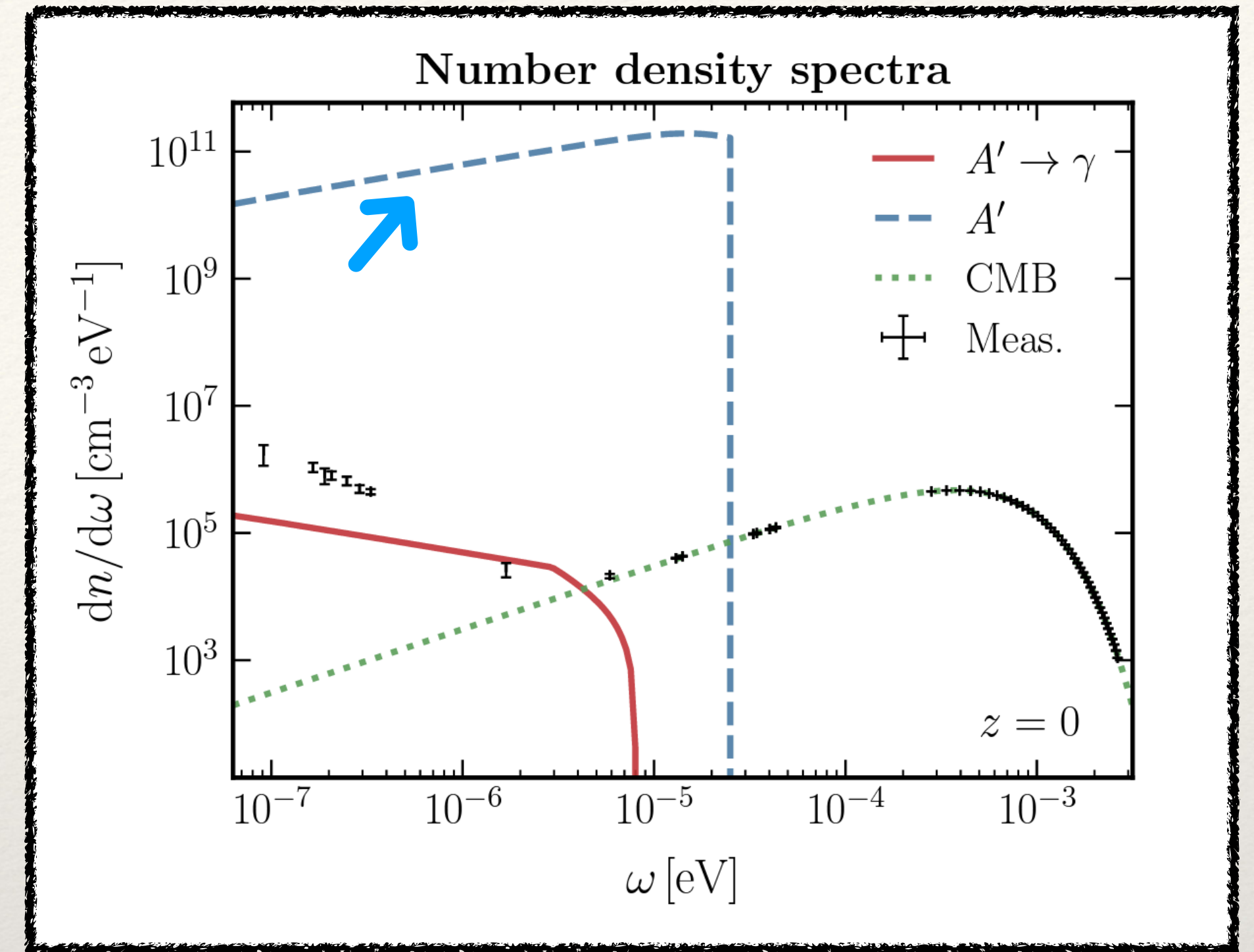


Dark photons resonantly co



First, axion decays (that's easy)

$$\frac{dn_{A'}}{d\omega} = \frac{2\rho_{\text{DM}}(z_{\text{dec}})(1+z)^3}{\tau_a H(z_{\text{dec}}) m_a \omega (1+z_{\text{dec}})^3} \Theta\left(\frac{m_a}{2} - \omega\right)$$

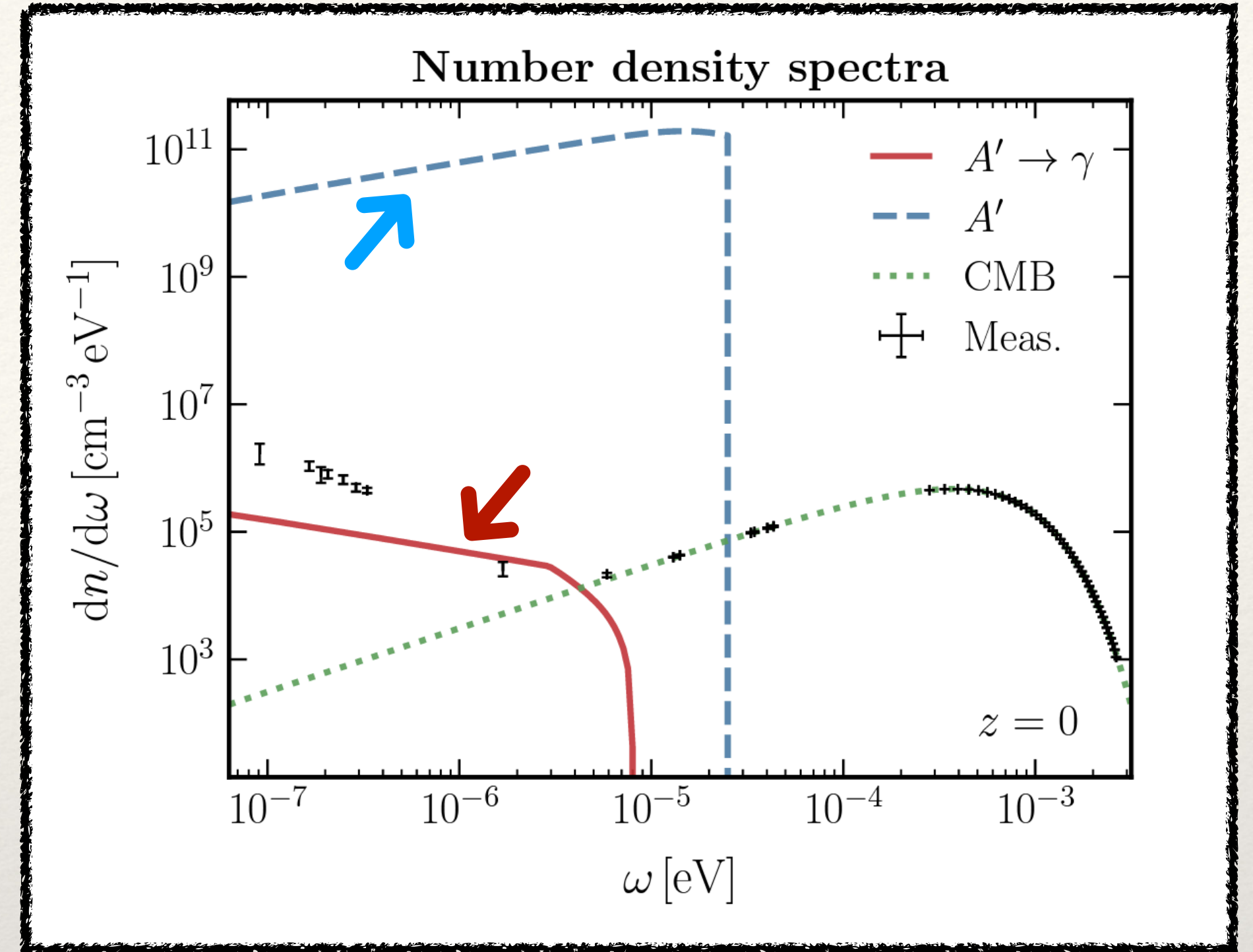


AC, H. Liu, S. Mishra-Sharma, M. Pospelov, J. T. Ruderman, A. Urbano,
Phys.Rev.Lett. 127 (2021) 1, 011102

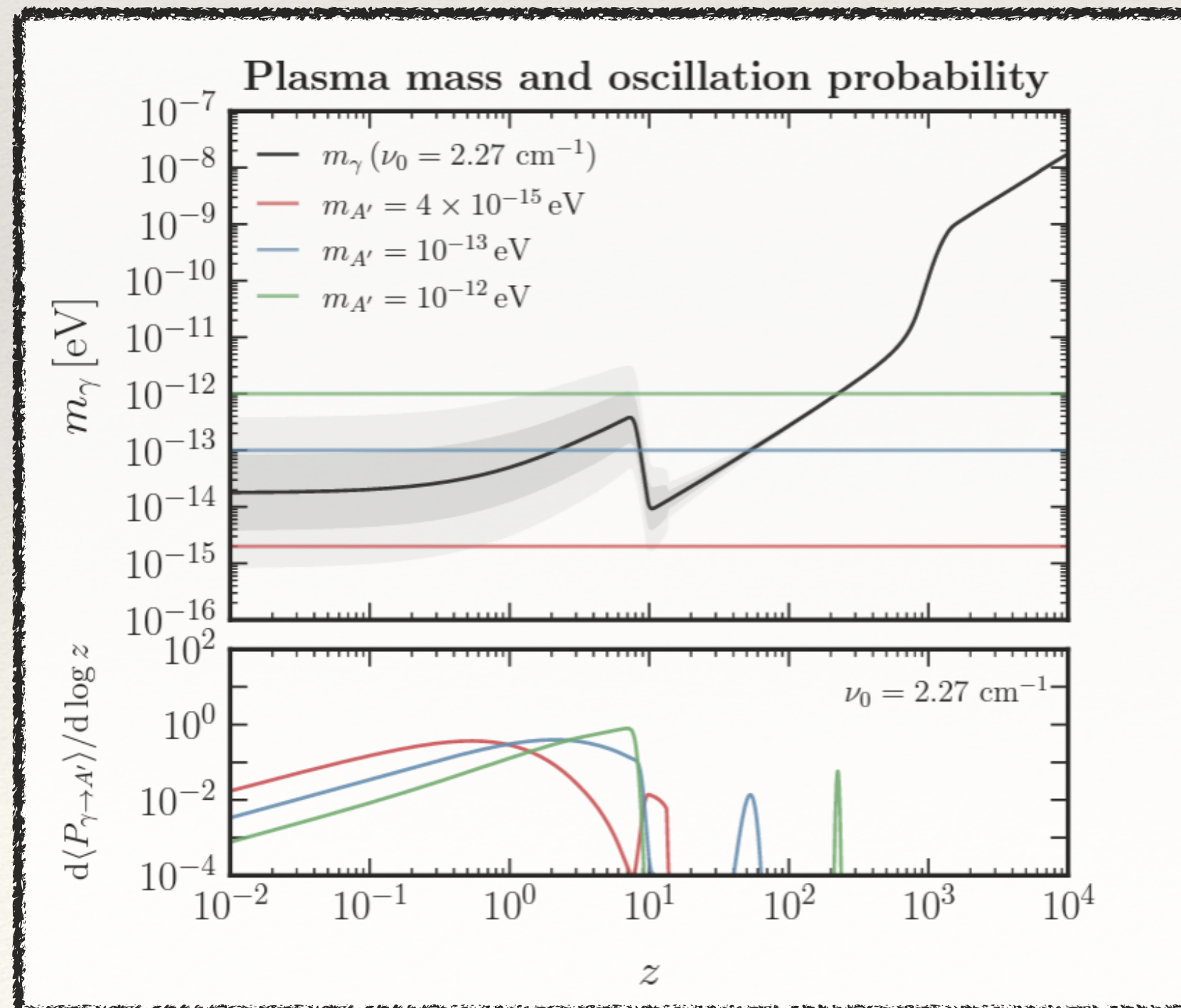
First, axion decays (that's easy)

$$\frac{dn_{A'}}{d\omega} = \frac{2\rho_{\text{DM}}(z_{\text{dec}})(1+z)^3}{\tau_a H(z_{\text{dec}}) m_a \omega (1+z_{\text{dec}})^3} \Theta\left(\frac{m_a}{2} - \omega\right)$$

Second, resonant conversion of photons to dark photons



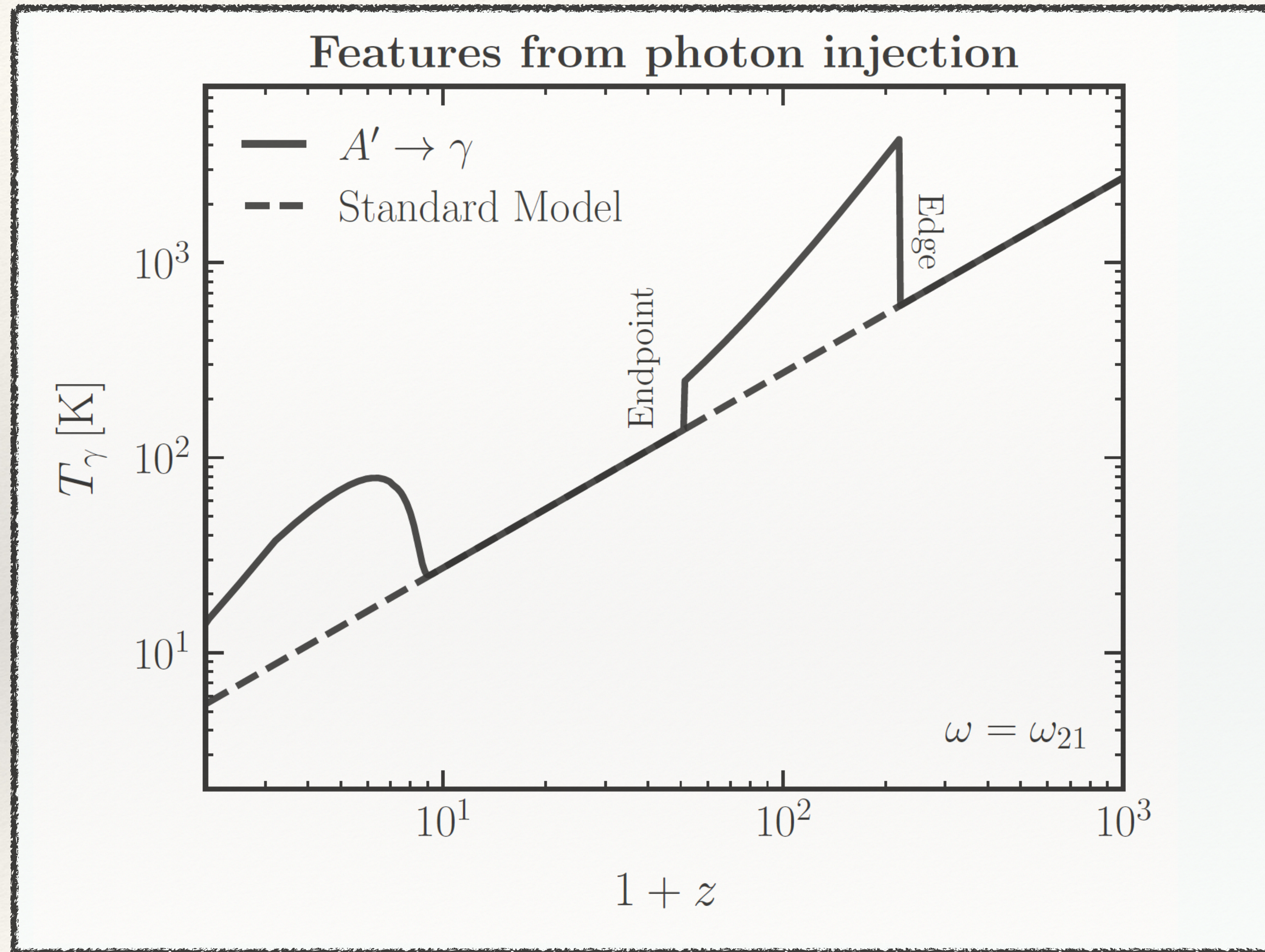
AC, H. Liu, S. Mishra-Sharma, M. Pospelov, J. T. Ruderman, A. Urbano, *Phys.Rev.Lett.* 127 (2021) 1, 011102



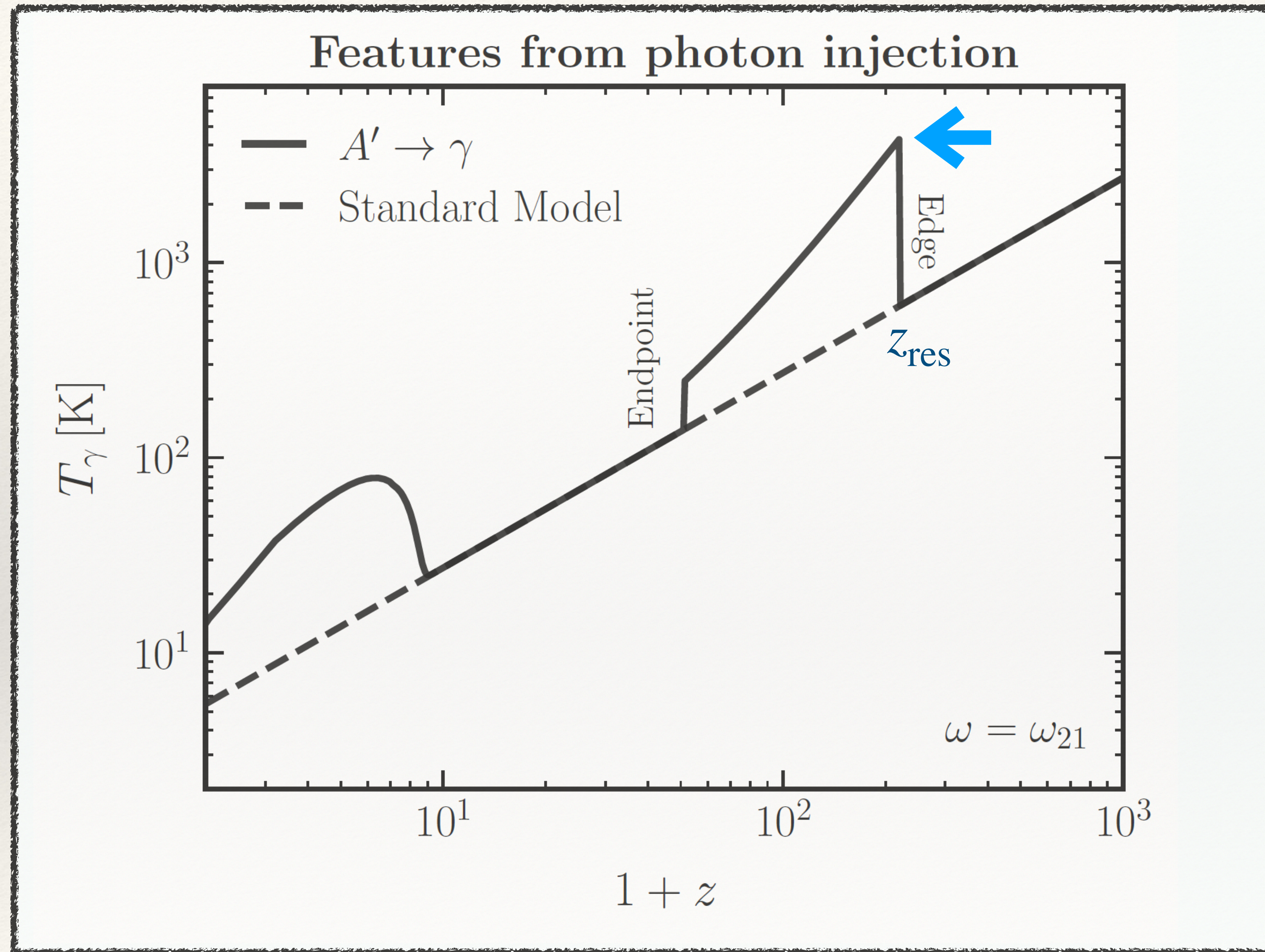
$$\frac{d\langle P_{\gamma \rightarrow A'} \rangle}{dz} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(t)} \left| \frac{dt}{dz} \right| \times \int dm_\gamma^2 f(m_\gamma^2; t) \delta_D(m_\gamma^2 - m_{A'}^2) m_\gamma^2$$

AC, H. Liu, S. Mishra-Sharma, J. T. Ruderman, *Phys.Rev.Lett.* 125 (2020) 22, 221303
AC, H. Liu, S. Mishra-Sharma, J. T. Ruderman, *Phys.Rev.D* 102 (2020) 10, 103533

Some nice spectral features from this type of models

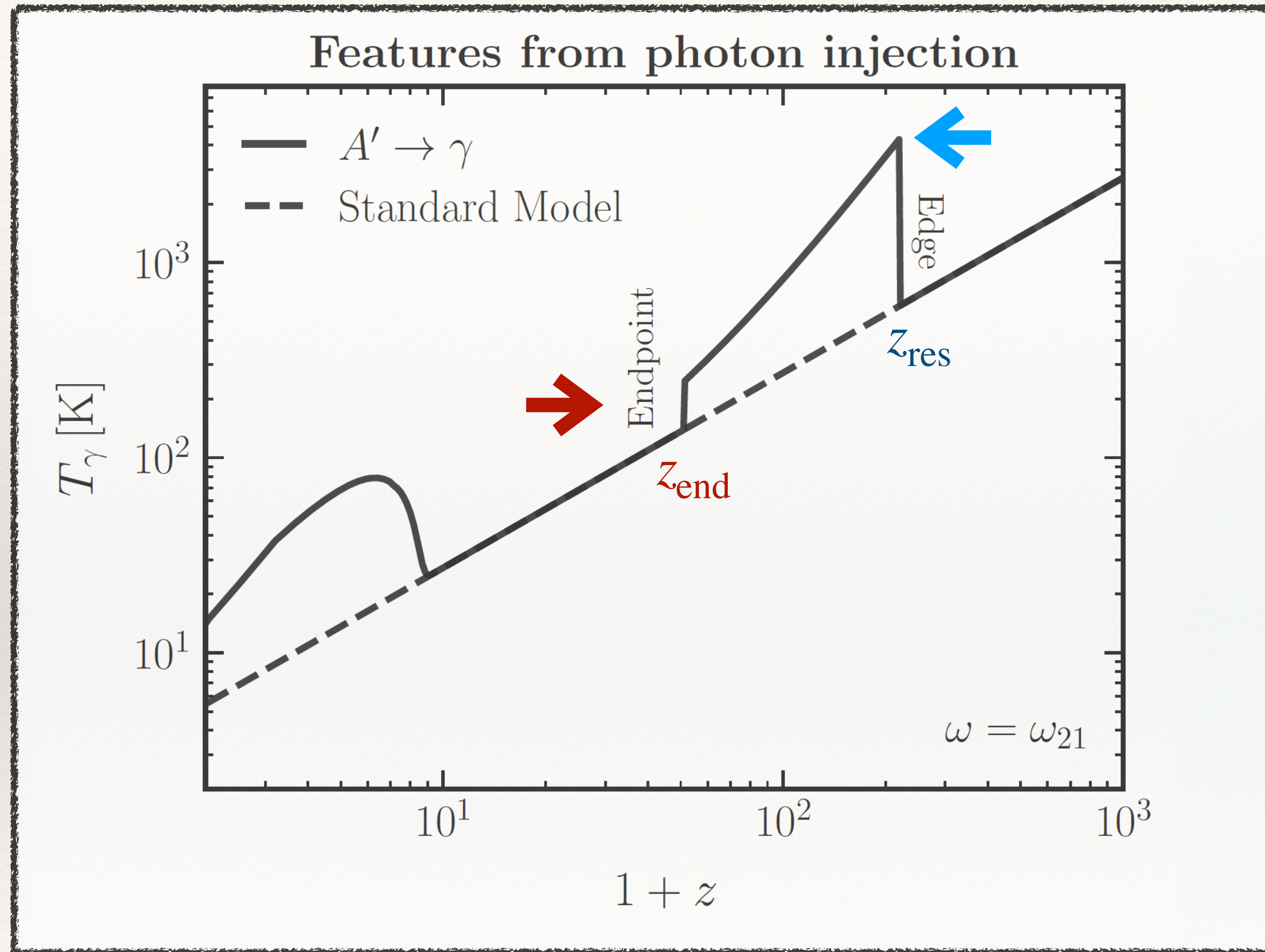


Some nice spectral features from this type of models



$$m_\gamma(z_{\text{res}}) \sim m_{A'}$$

Some nice spectral features from this type of models

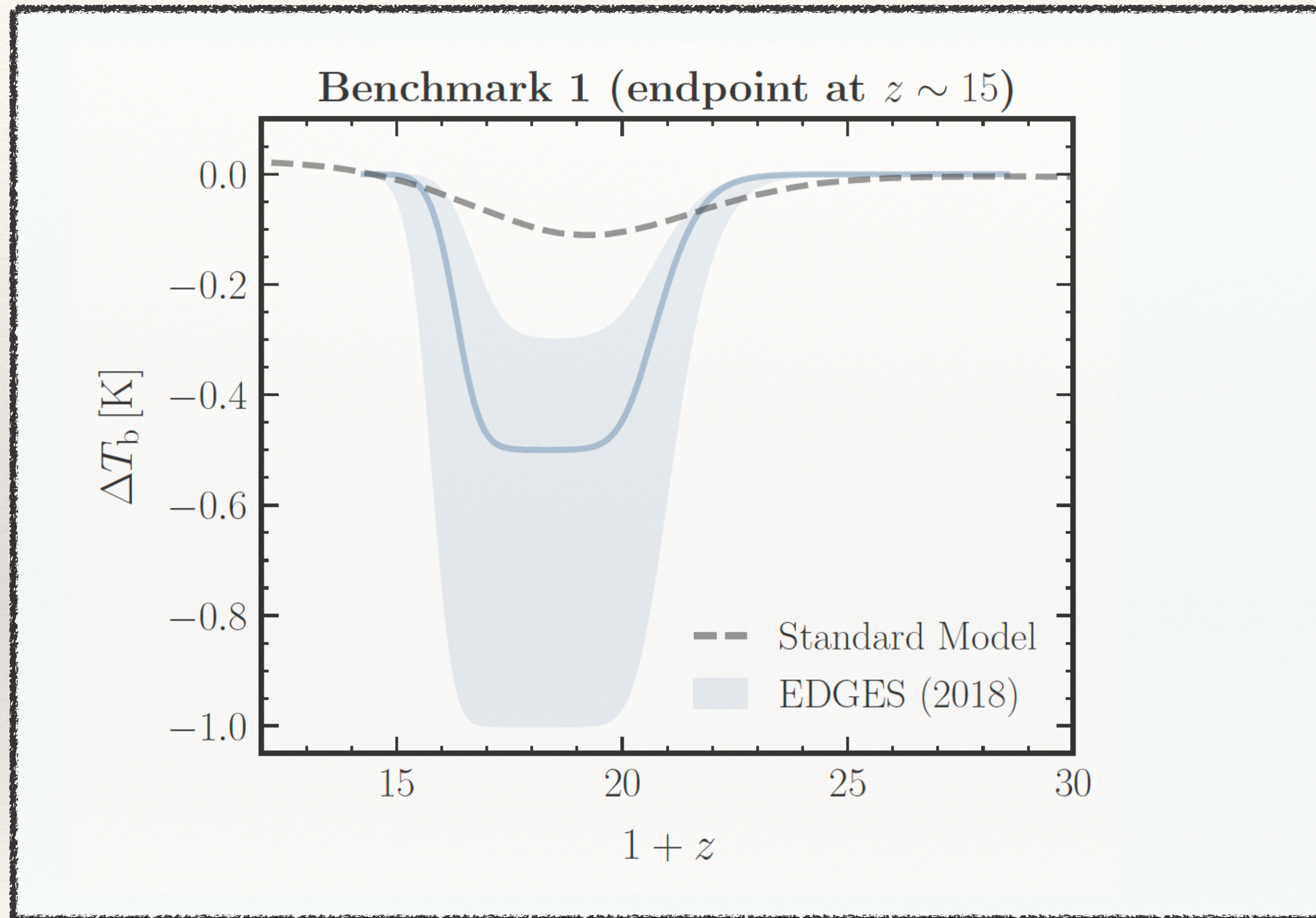


$$m_\gamma(z_{\text{res}}) \sim m_{A'}$$

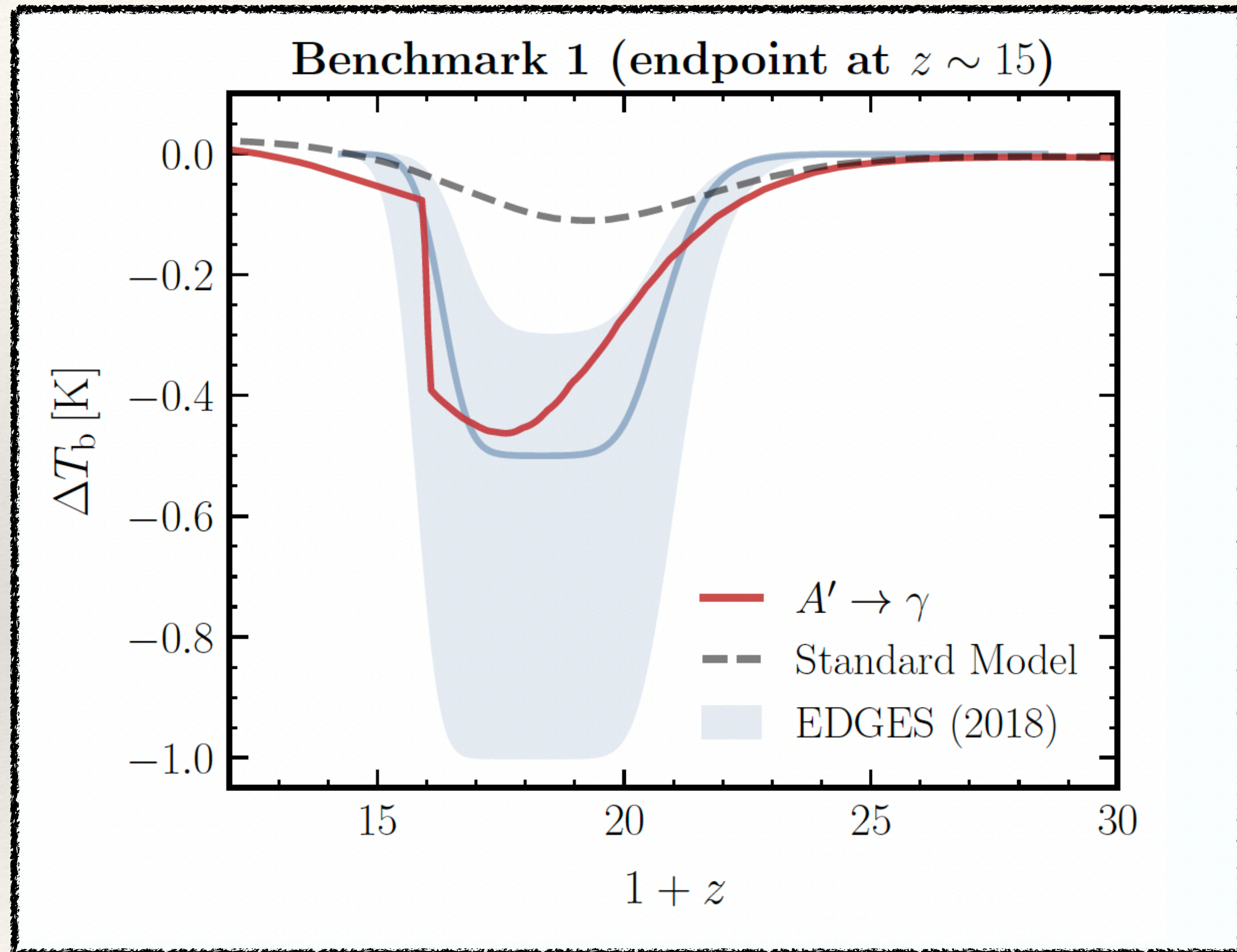
Kinematic endpoint

$$\frac{m_a}{2} \left(\frac{1 + z_{\text{end}}}{1 + z_{\text{res}}} \right) \sim \omega_{21}$$

Effect on the brightness temperature



Effect on the brightness temperature

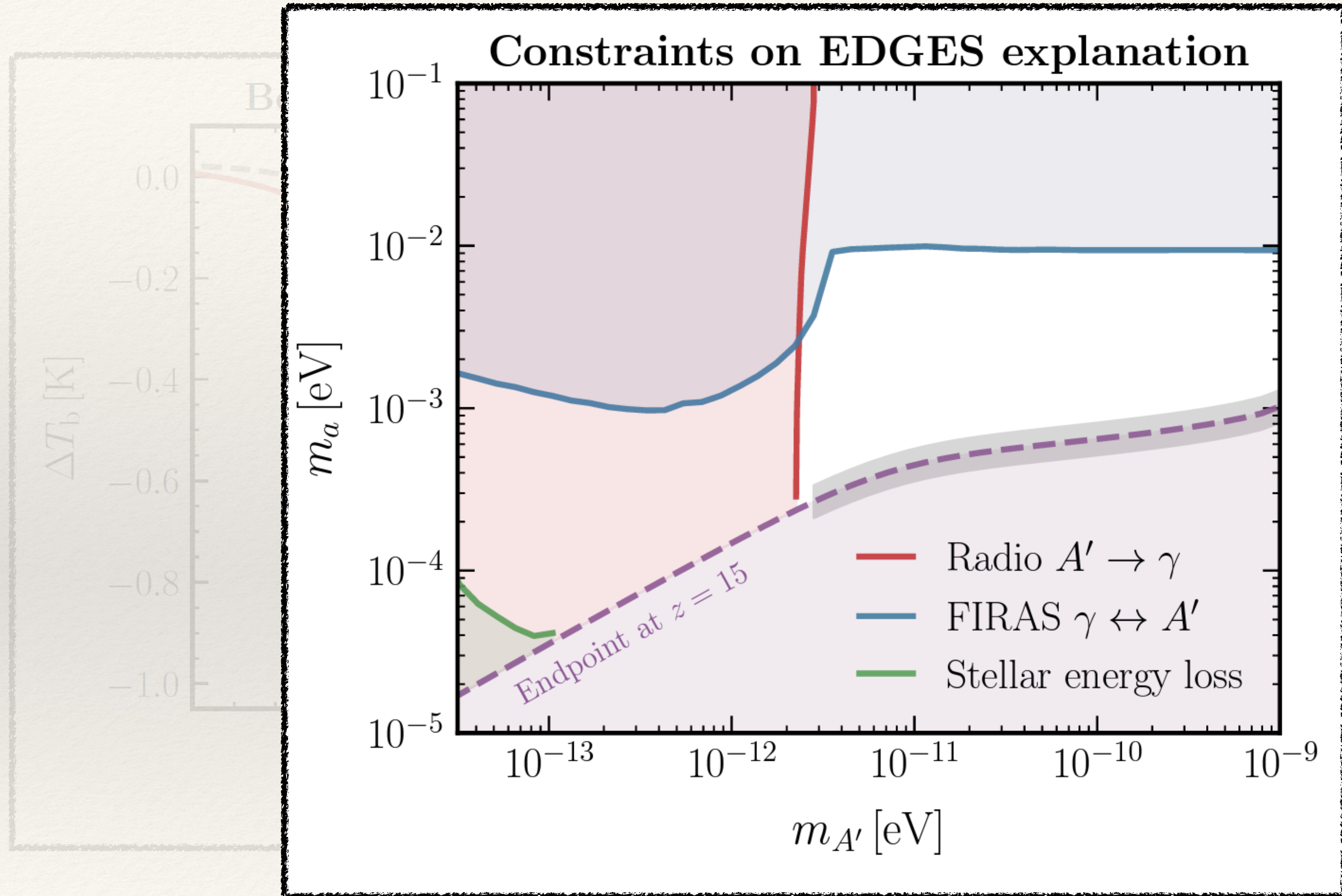


$$m_{A'} = 10^{-11} \text{ eV}$$

$$m_a = 5 \times 10^{-4} \text{ eV}$$

$$\epsilon = 5 \times 10^{-8}$$

Effect on the brightness temperature



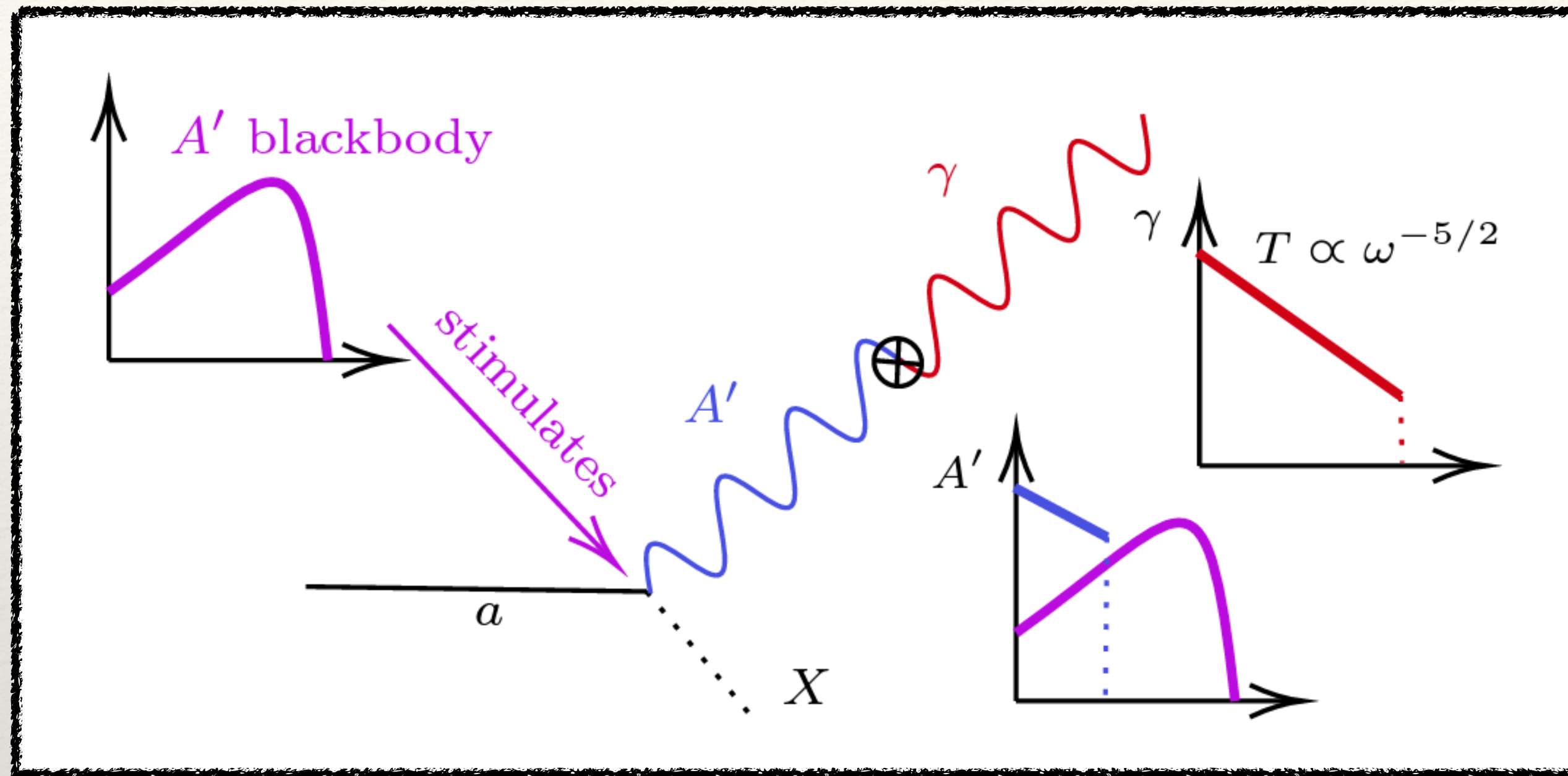
$= 10^{-11}$ eV

$= 5 \times 10^{-4}$ eV

5×10^{-8}

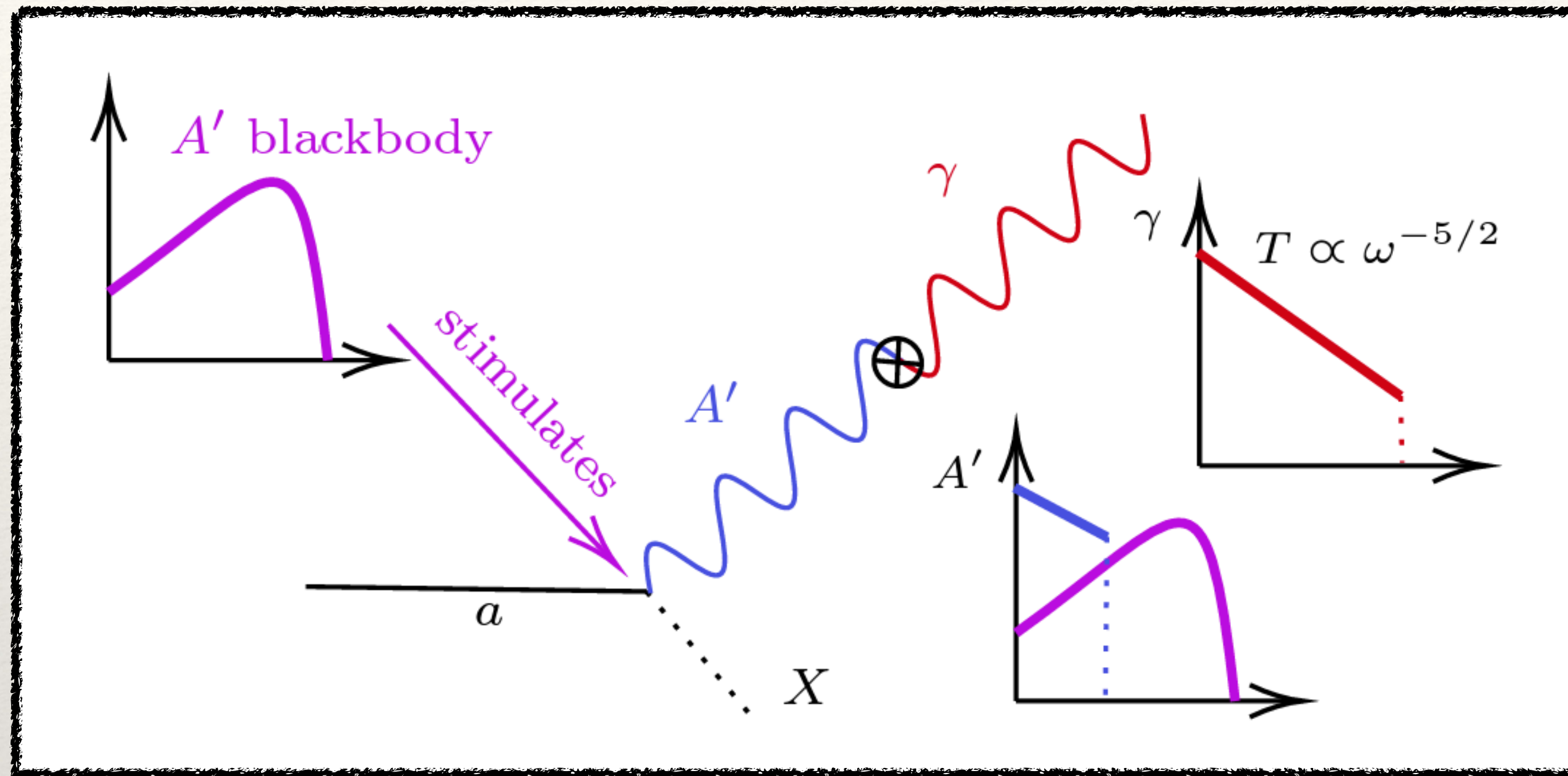
A lot of constraints

But the model — despite of EDGES — is interesting, for example it can be enlarged to explain ARCADE excess



Assume there is a background of dark photons

But the model — despite of EDGES — is interesting, for example it can be enlarged to explain ARCADE excess



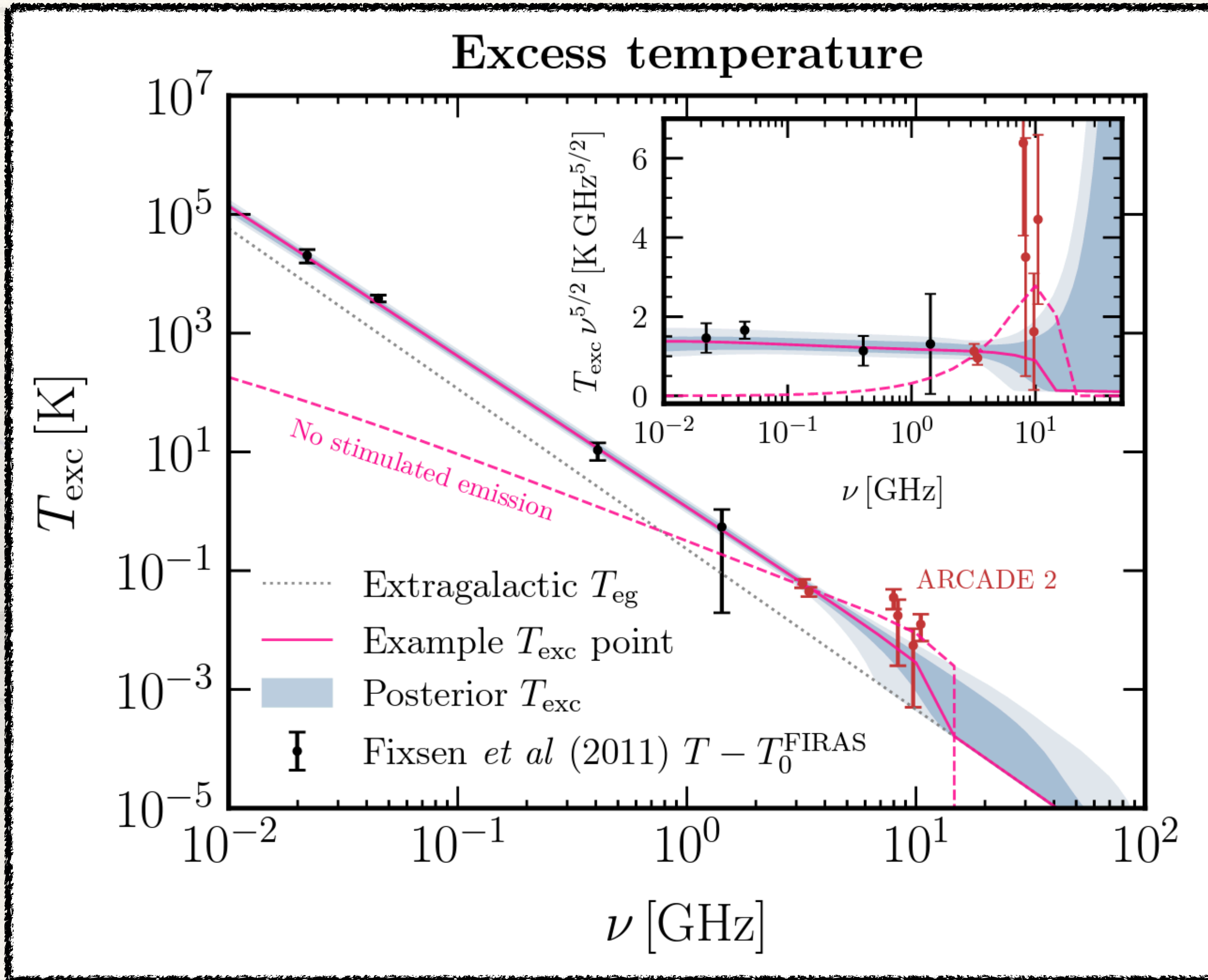
Assume there is a background of dark photons

$$\frac{dn_\gamma}{dx}(x, z) = \frac{\rho_a(z)}{m_a} \frac{\alpha}{x} \underbrace{\frac{1}{\tau(z_*)}}_{\propto x^{-1}} \underbrace{\frac{1}{H(z_*)}}_{\propto x^{3/2}} \underbrace{\int_z^{z_*} dz' \frac{d\langle P_{A' \rightarrow \gamma} \rangle}{dz'}}_{\propto x^{-1}}$$

$$\tau(z) = \tau_{\text{vac}} [1 + n f_{A'}^{\text{BB}}(z)]^{-1}$$

Stimulated decay

But the model — despite of EDGES — is interesting, for example it can be enlarged to explain ARCADE excess



is a background
k photons

$$\frac{dn_\gamma}{dx}(x, z) =$$

$$\tau(z) = \tau_{\text{vac}} [1 + n f_{A'}^{\text{BB}}(z)]^{-1}$$

Stimulated decay

To be done..

1) Other models, e.g axion-photon oscillations in magnetic fields

To be done..

1) Other models, e.g axion-photon oscillations in magnetic fields

2) Better treatment of astrophysical uncertainties

To be done..

- 1) Other models, e.g axion-photon oscillations in magnetic fields**
- 2) Better treatment of astrophysical uncertainties**
- 3) Study CMB and 21 cm anisotropy**

To be done..

1) Other models, e.g axion-photon oscillations in magnetic fields

2) Better treatment of astrophysical uncertainties

3) Study CMB and 21 cm anisotropy

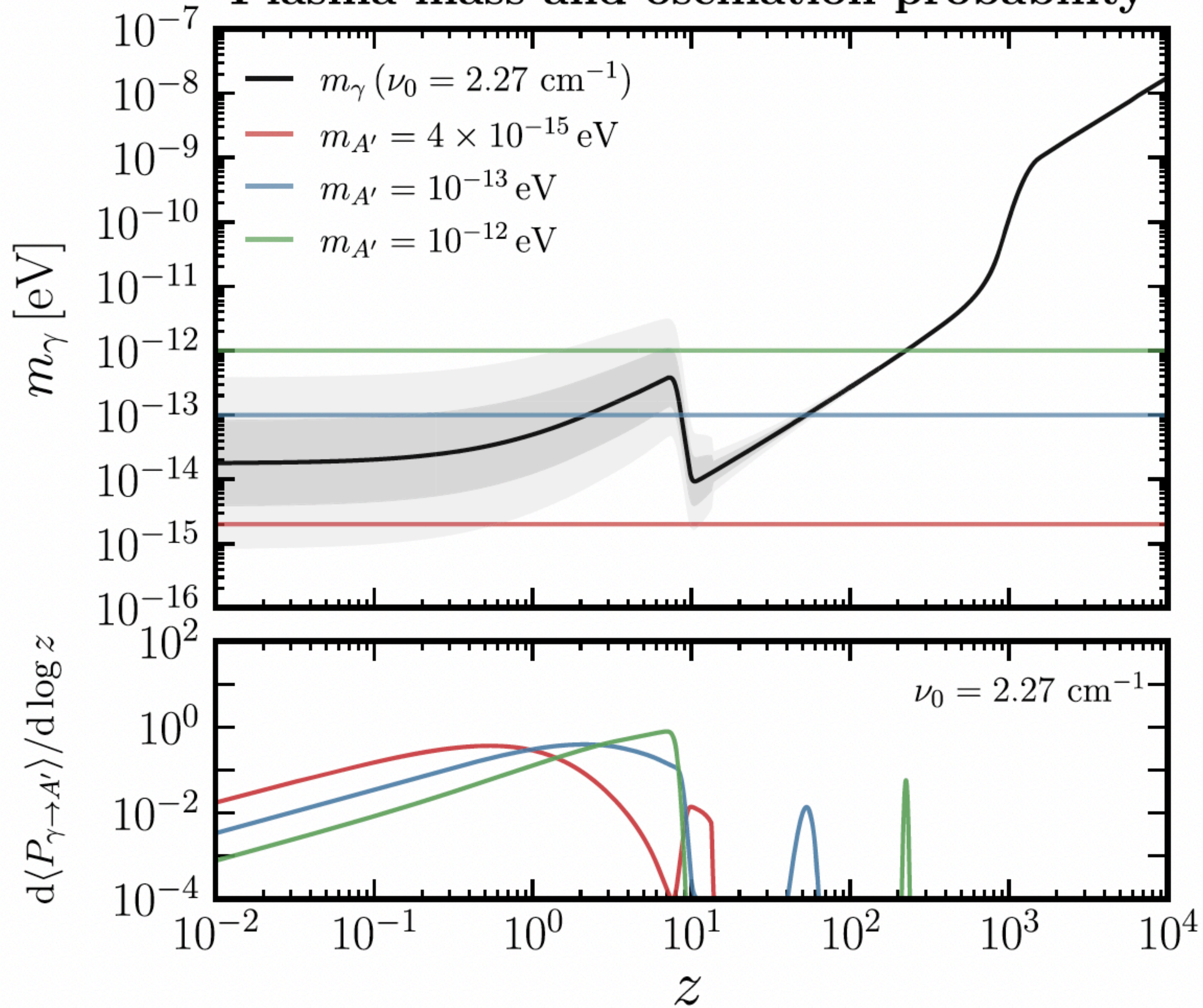
Thanks for the attention!



Back-up

Back-up

Plasma mass and oscillation probability



Dark Photon Mass Kinetic mixing

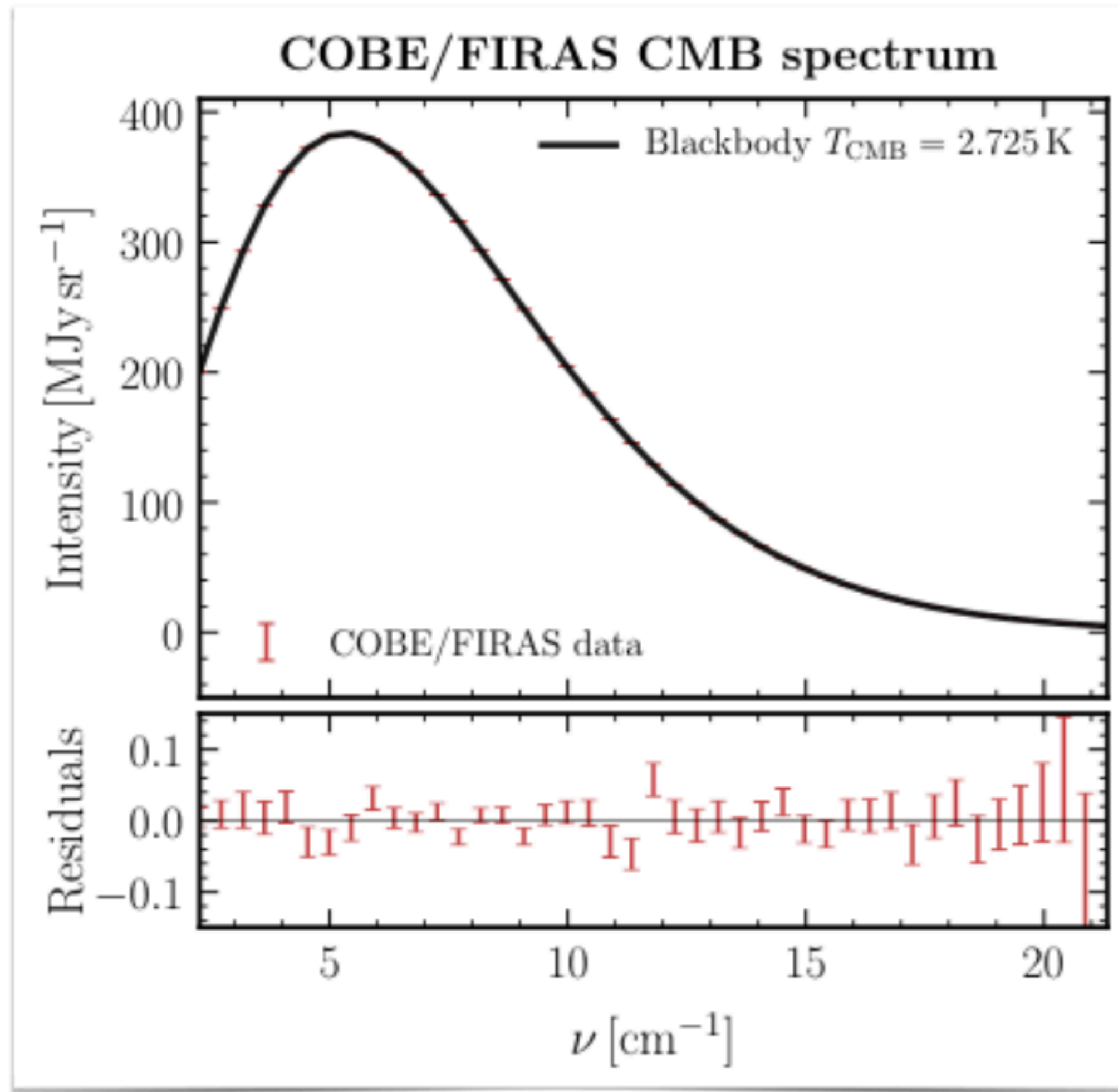
$$\frac{d\langle P_{\gamma \rightarrow A'} \rangle}{dz} = \frac{\pi m_{A'}^2 \epsilon^2}{\omega(t)} \left| \frac{dt}{dz} \right|$$

$$\times \int dm_\gamma^2 f(m_\gamma^2; t) \delta_D(m_\gamma^2 - m_{A'}^2) m_\gamma^2,$$

Photon Mass PDF Photon Mass

This will then affect CMB

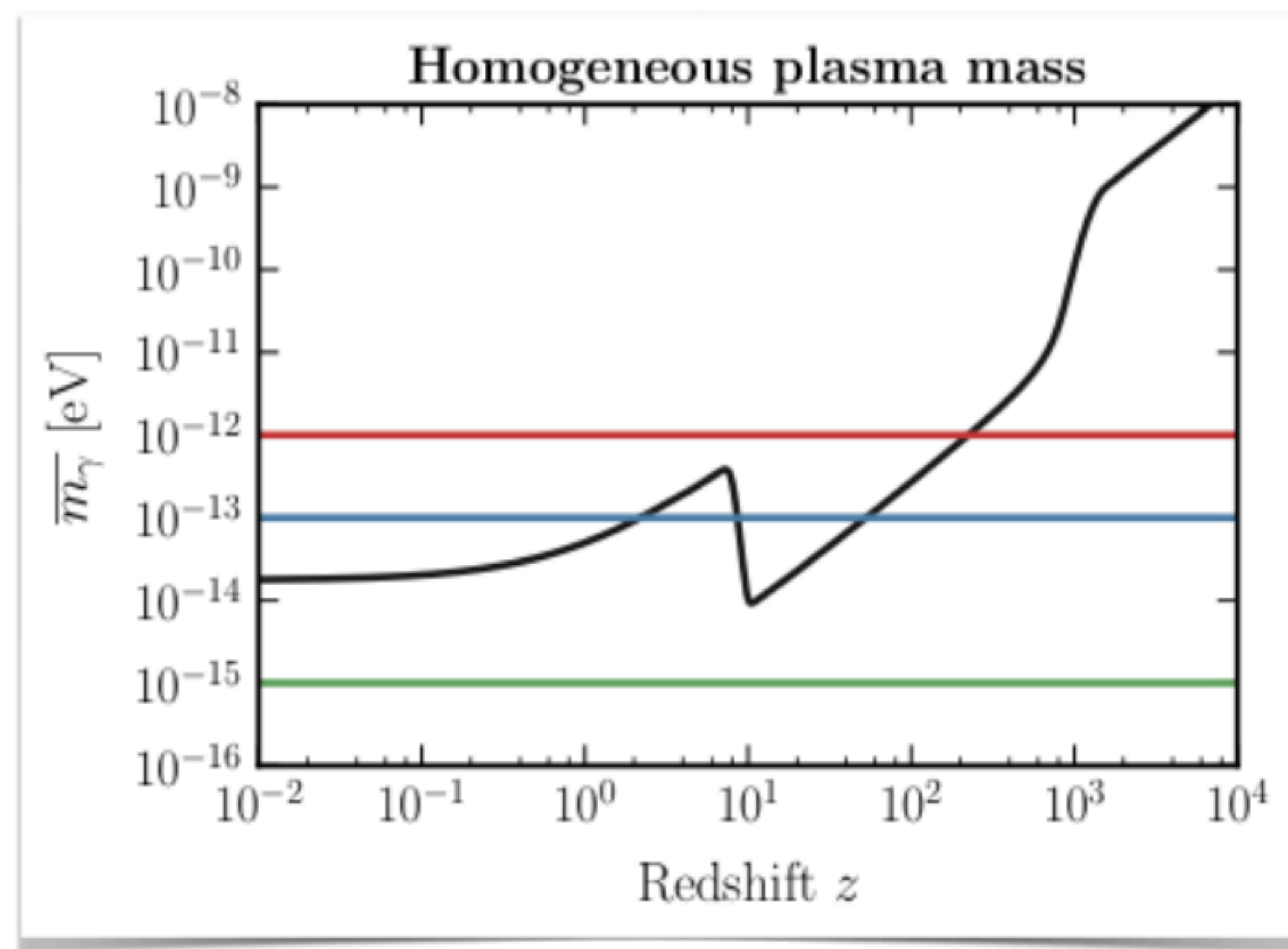
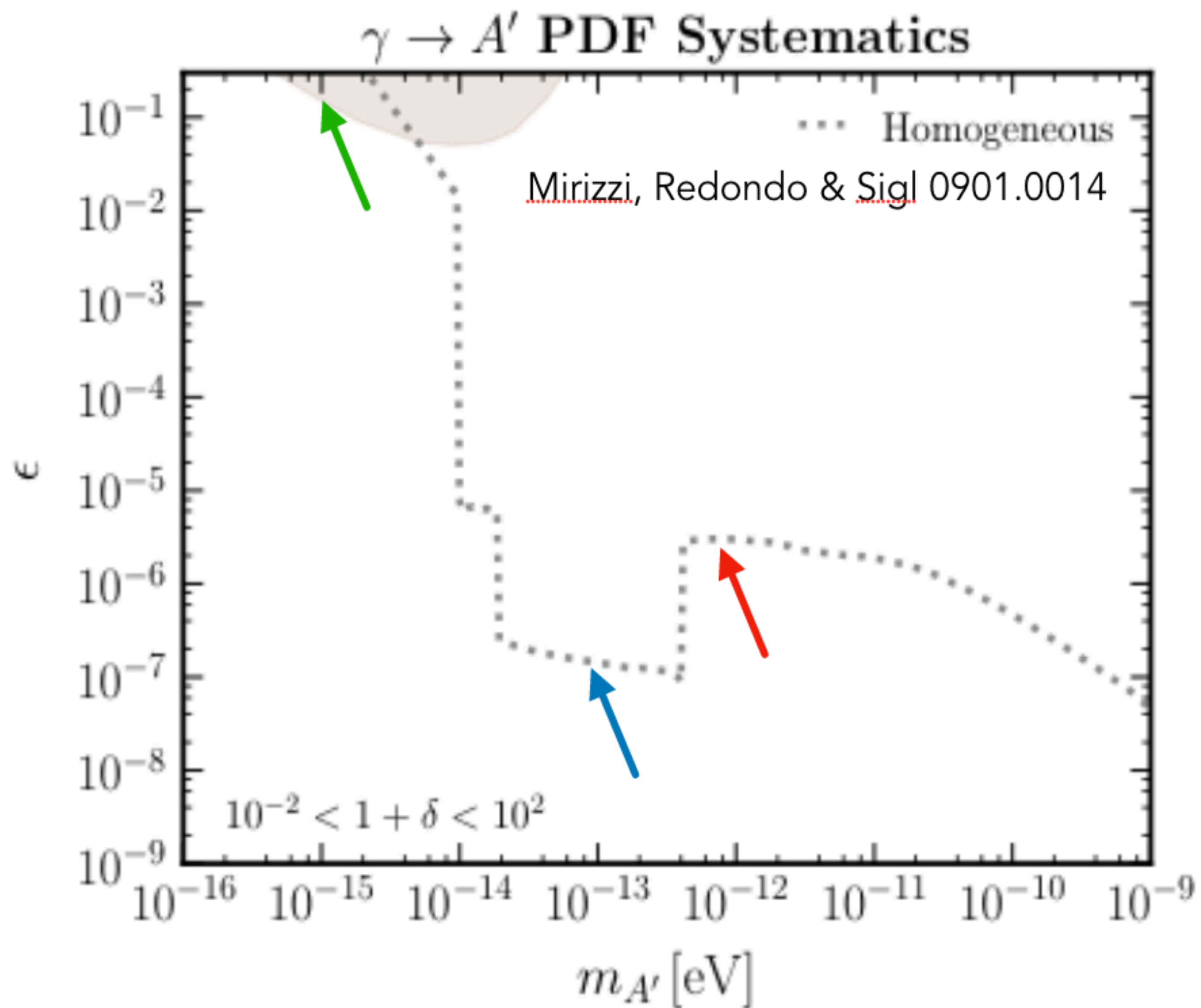
Fixsen+ astro-ph/9605054



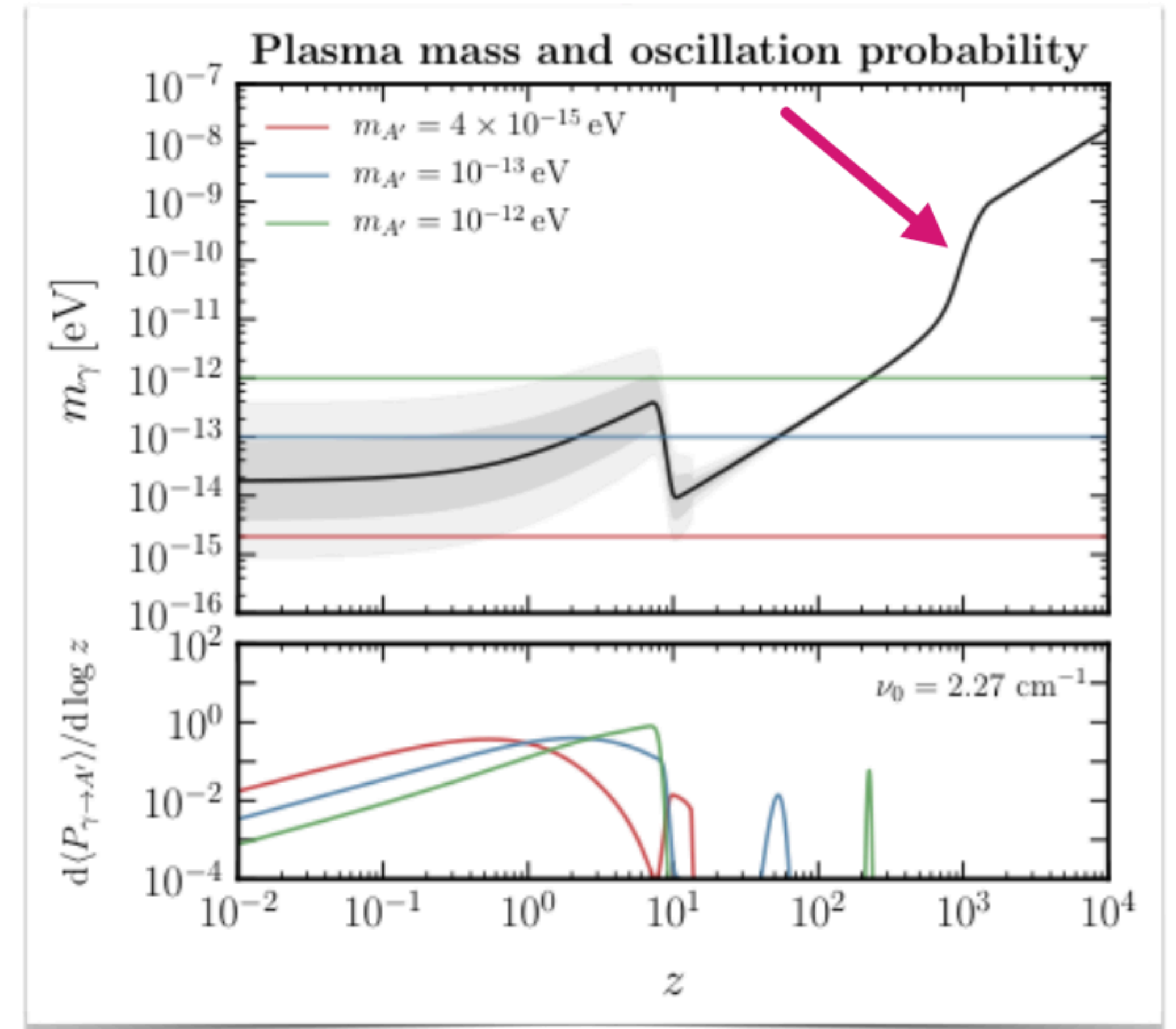
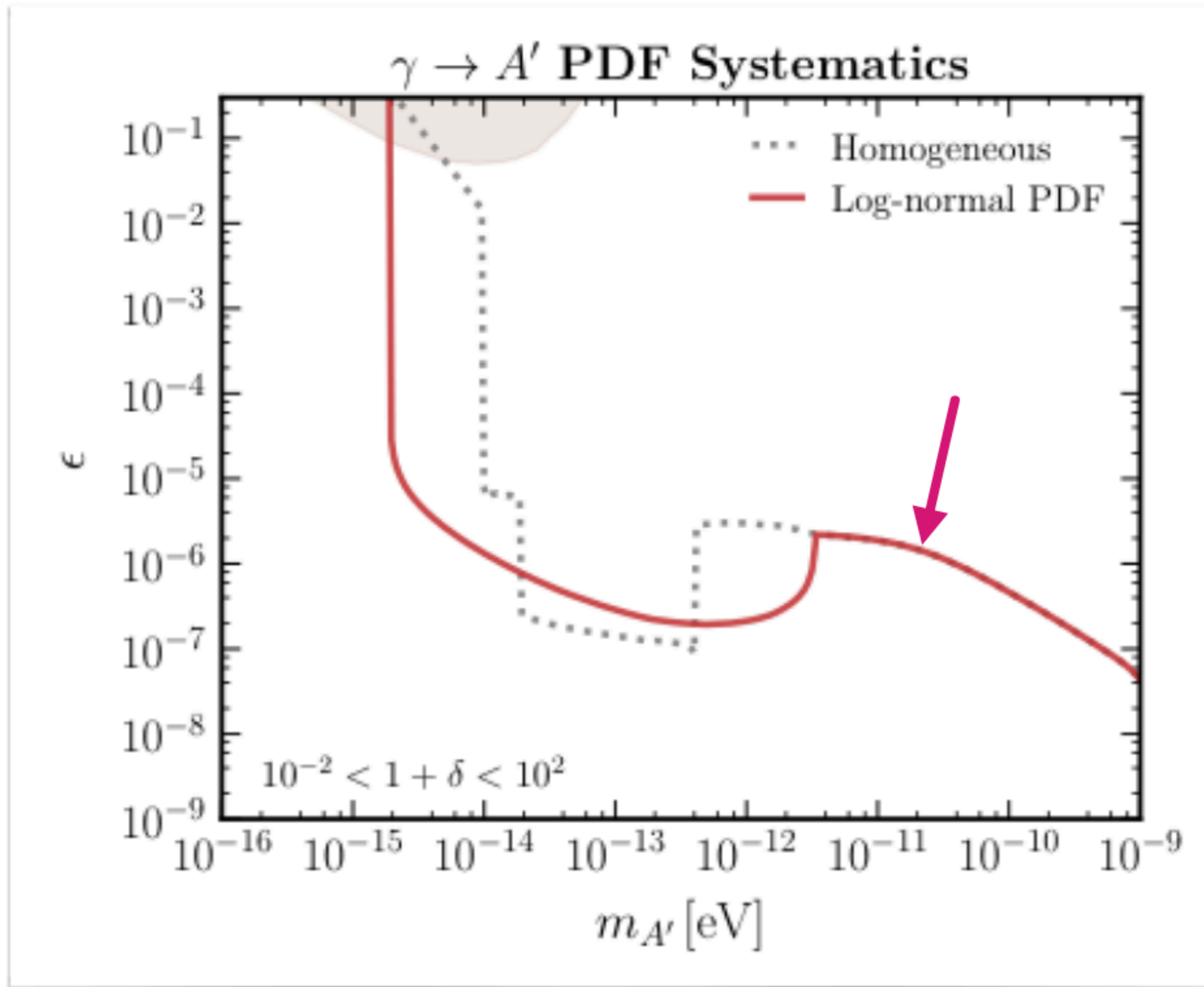
The CMB is very close to a **perfect blackbody**.

Spectral distortions due to disappearing photons are **highly constrained**.

CMB Constraints

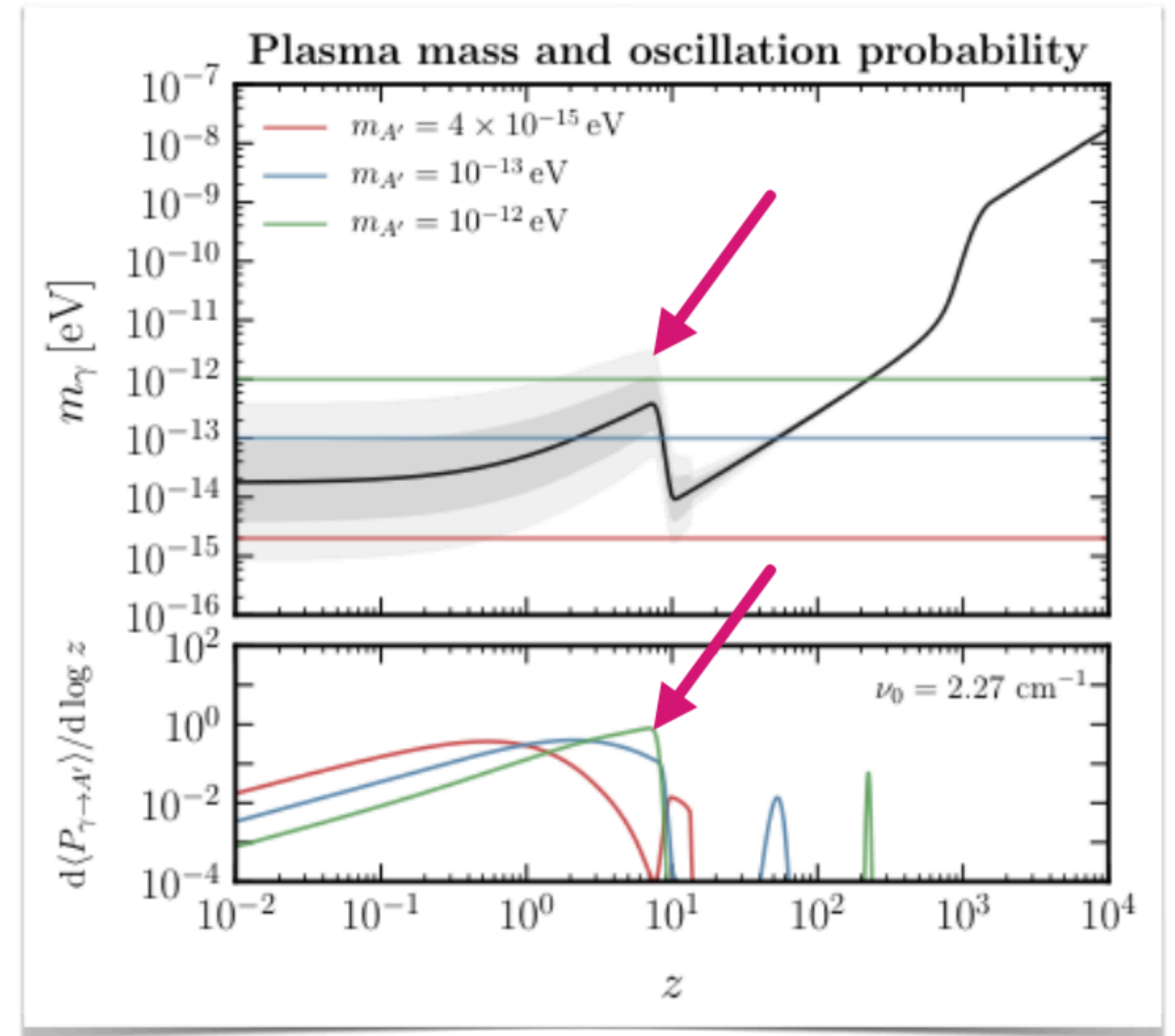
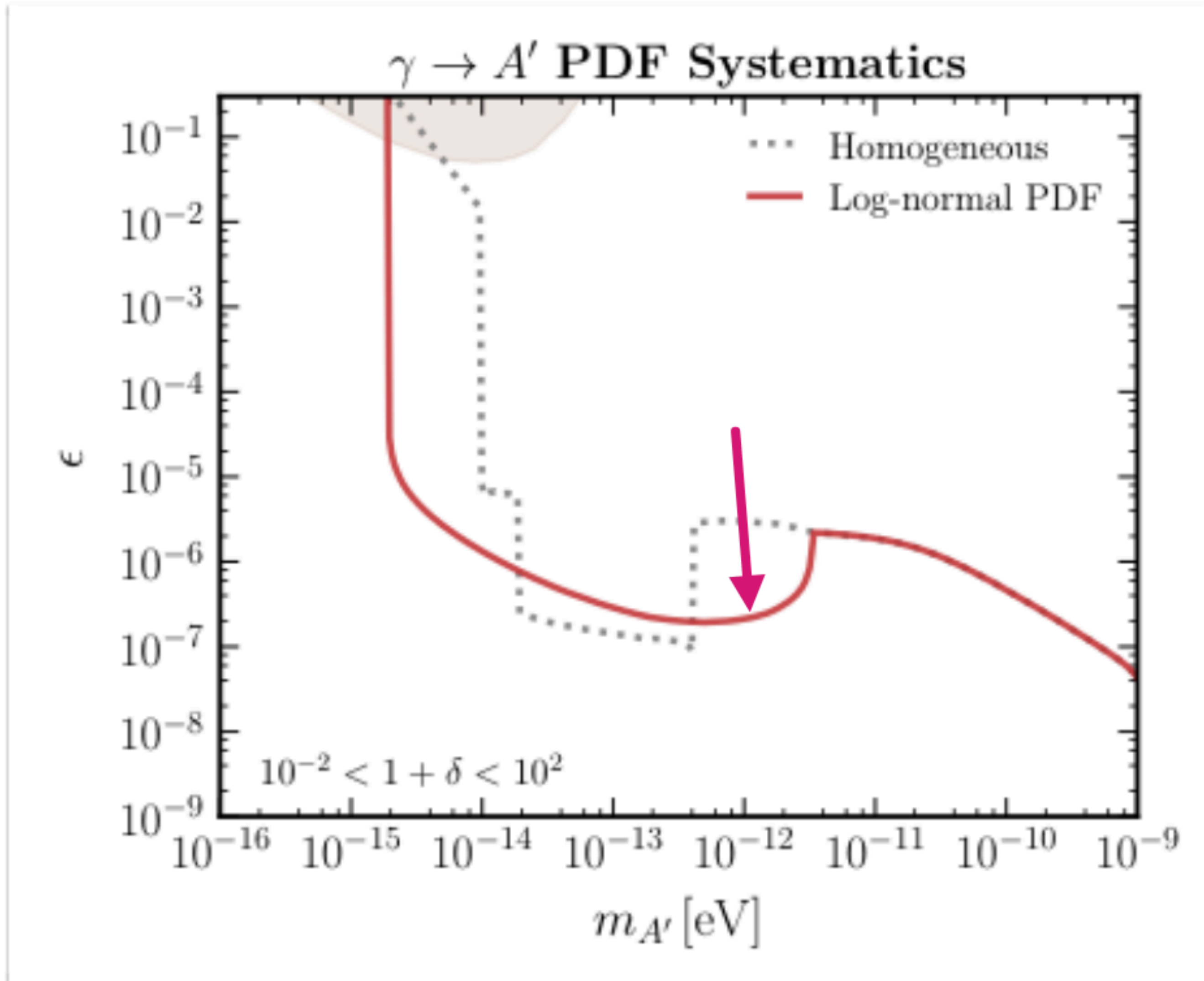


CMB Constraints



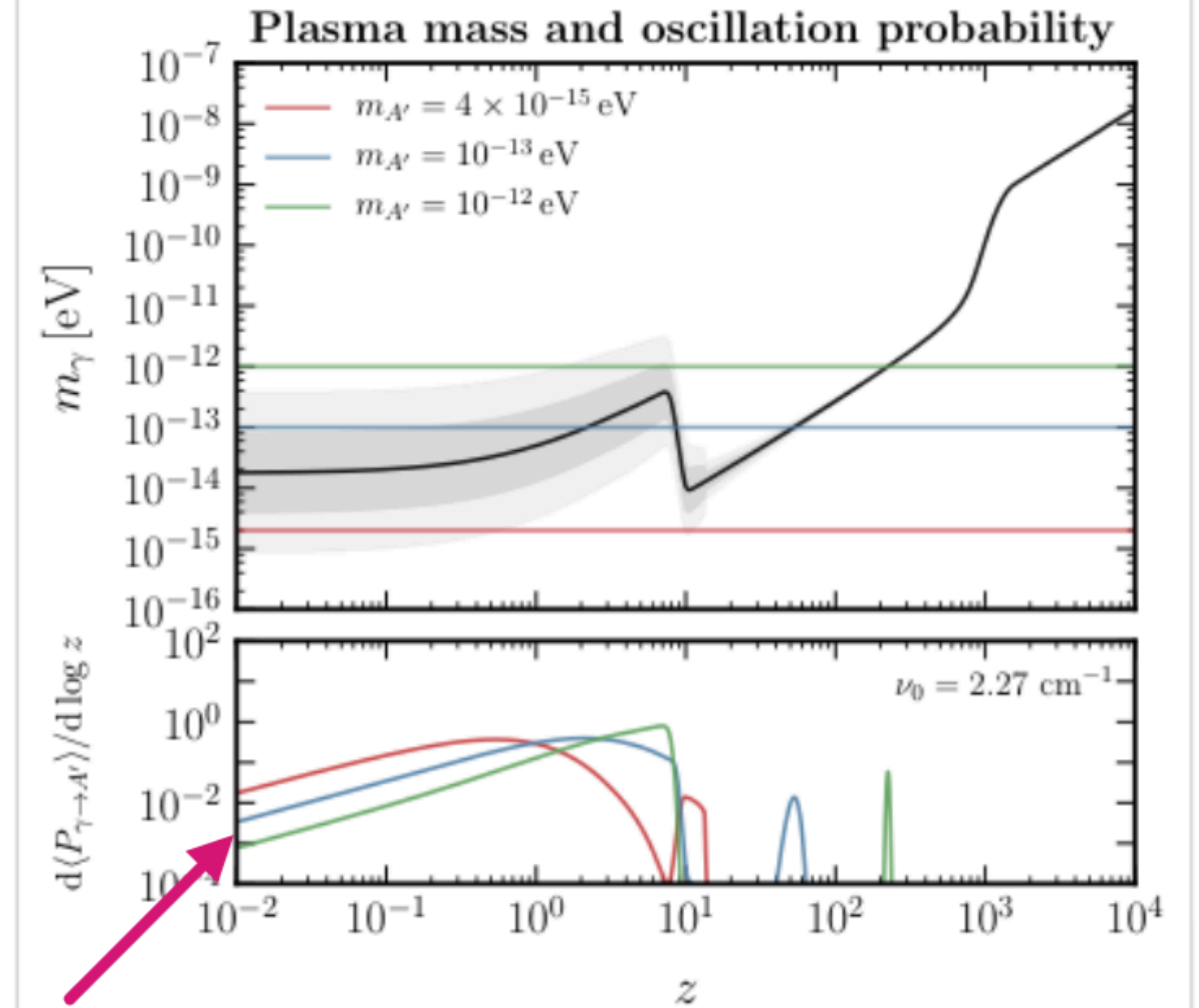
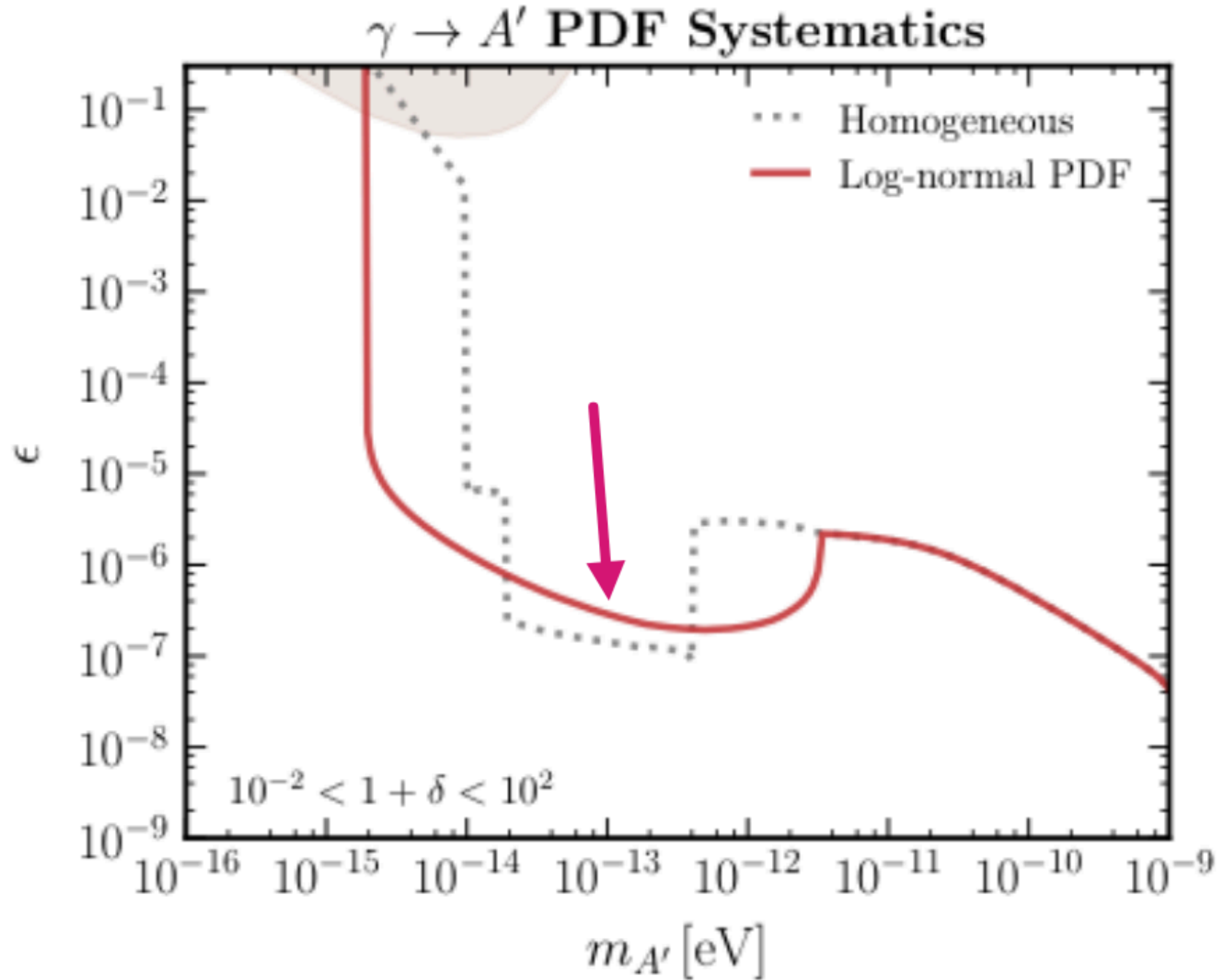
Small fluctuations at high redshift:
similar to the homogeneous case

CMB Constraints



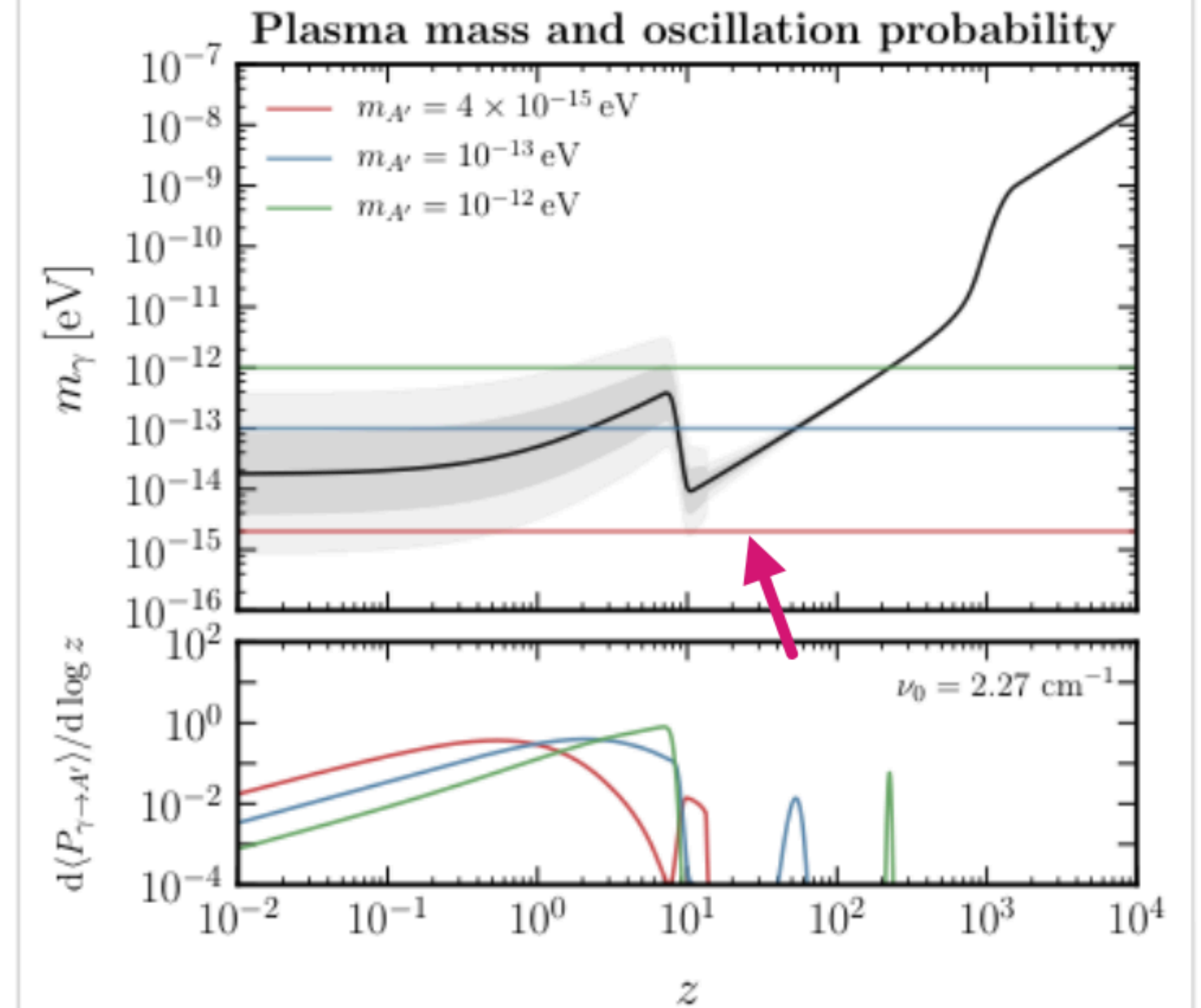
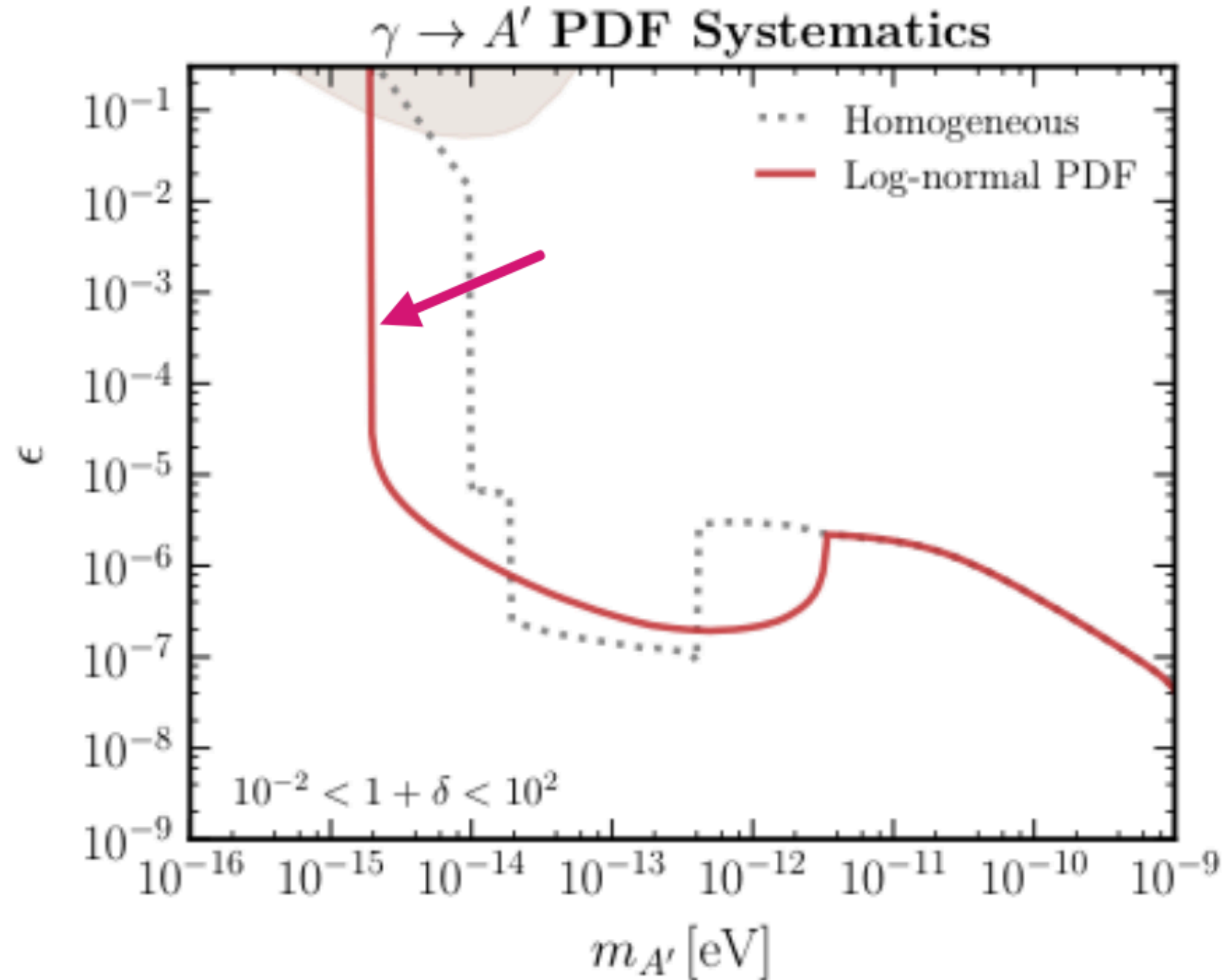
Later conversions available

CMB Constraints



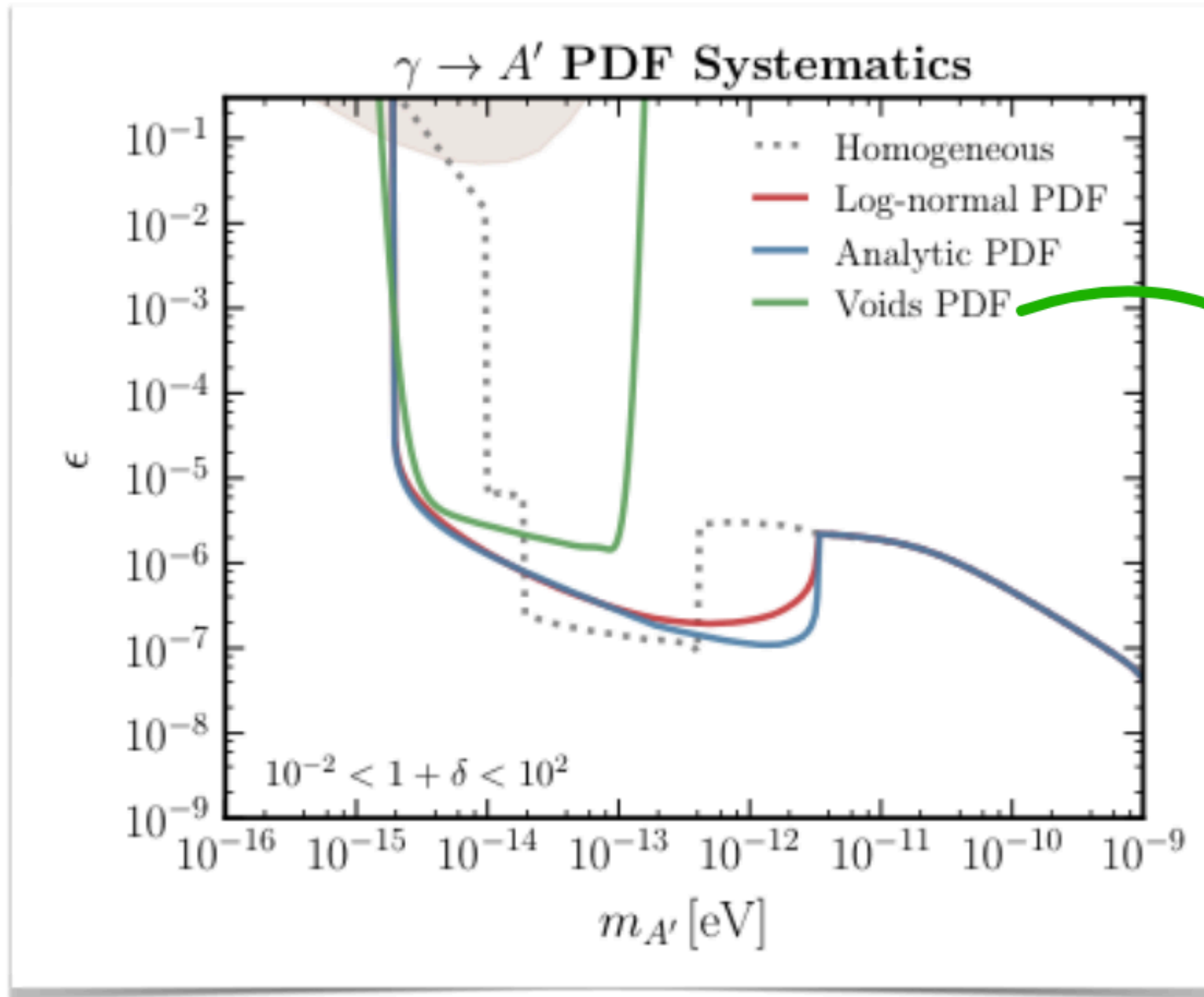
Constraints weakened here, as some conversion probability is in the future

CMB Constraints



Limits extend to lower masses because of under-fluctuations

Check some systematics



check for
underdensities