

Spatio-Temporal correlations in earthquakes and solar flares: self-organized criticality beyond the sand pile model

Outline

Seismic-occurrence has been recognized as the paradigmatic example of self-organized criticality, since from the origin of the theory

Catastrophic events «must not be considered as anomalous, since could be attributed to abnormal circumstances, but they are intrinsic to the dynamics, the same dynamics that produces small, ordinary events»

"I do not intend to chronicle the suffering it has inflicted on men, nor to provide a list of cities razed to the ground or inhabitants buried under rubble.....

Such a narrative would be moving and, perhaps, touching the heart, might even have an uplifting effect. I, however, entrust this kind of narrative to more experienced hands.

I will describe here only the work of nature, the surprising natural circumstances that accompanied the terrible event and their causes."



I. Kant, *Allgemein Naturgeschichte und Theorie des Himmels*, 1755

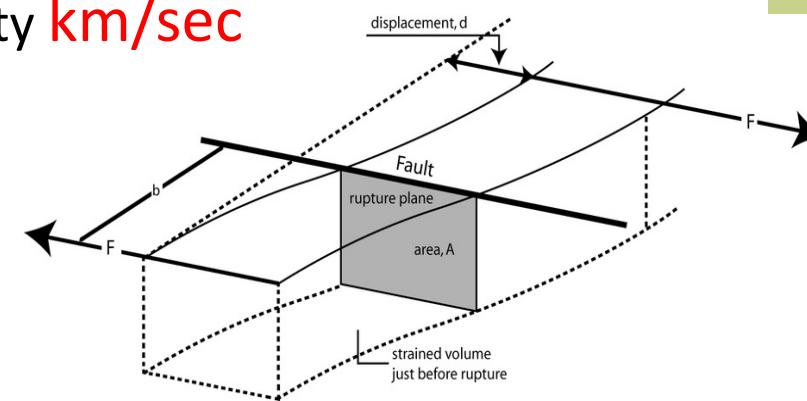
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Catastrophic events «must not be considered as anomalous, since could be attributed to abnormal circumstances, but they are intrinsic to the dynamics, the same dynamics that produces small, ordinary events»

- ★ Separation of time scales
- ★ Gutenberg-Richter law
- ★ Burridge-Knopoff- Spring-block model

{ Strain accumulation rate **cm/year**
Strain propagation velocity **km/sec**

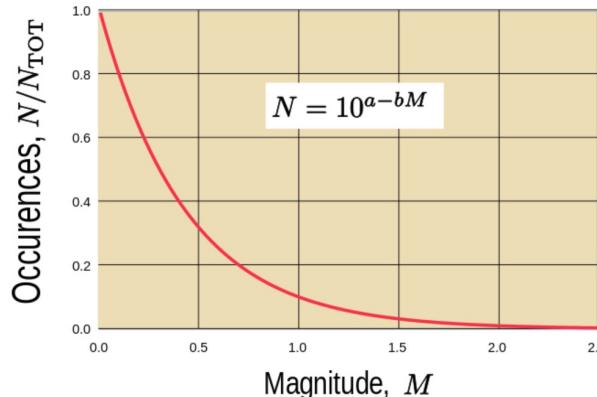


Outline

- ★ Gutenberg-Richter law
- ★ Burridge-Knopoff- Spring-block model
- ★ Non trivial spatio temporal patterns in seismic occurrence
- Similarity and differences with similar patterns in solar flare occurrence
- ★ Epidemic models for seismic and solar flare occurrence
- ★ Generalization of the sand-pile model for realistic temporal correlations

1956

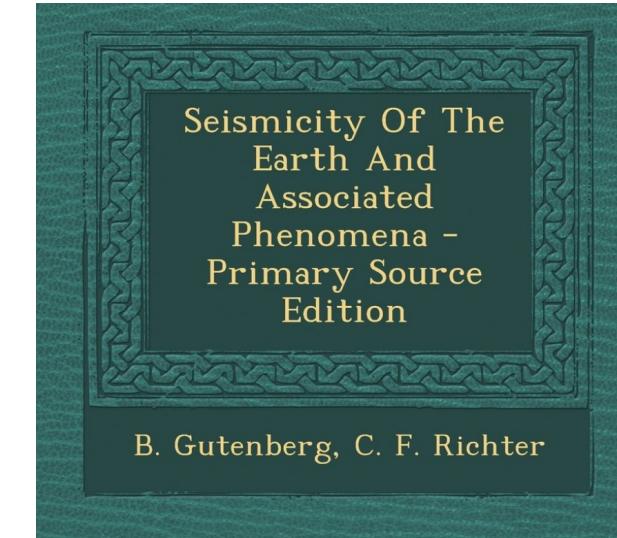
Gutenberg–Richter law



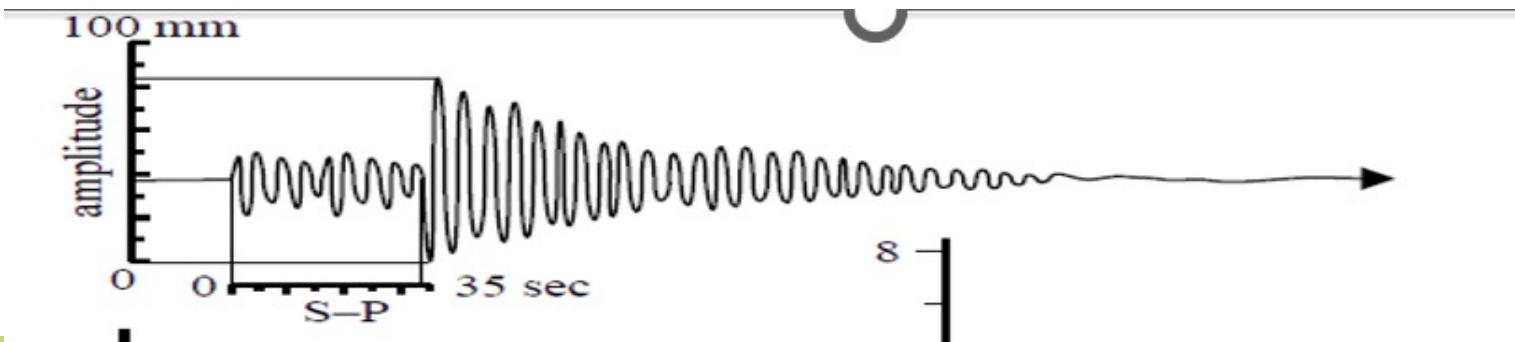
Beno Gutenberg



Charles Francis Richter ...

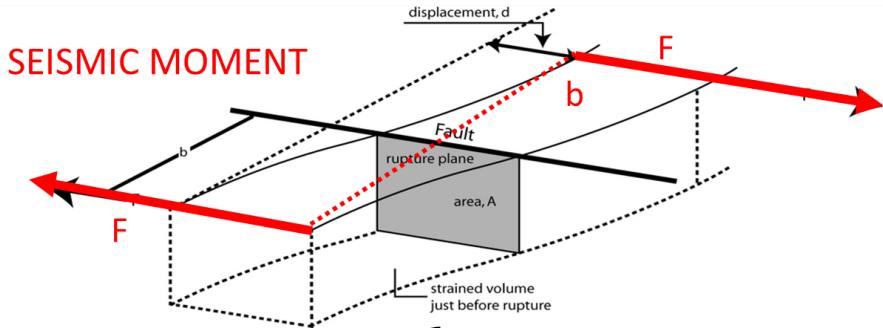


$$P(m) = k 10^{-bm} \quad b \text{ very importat parameter}$$



1979

Moment magnitude



$$M_0 = \mu A d$$

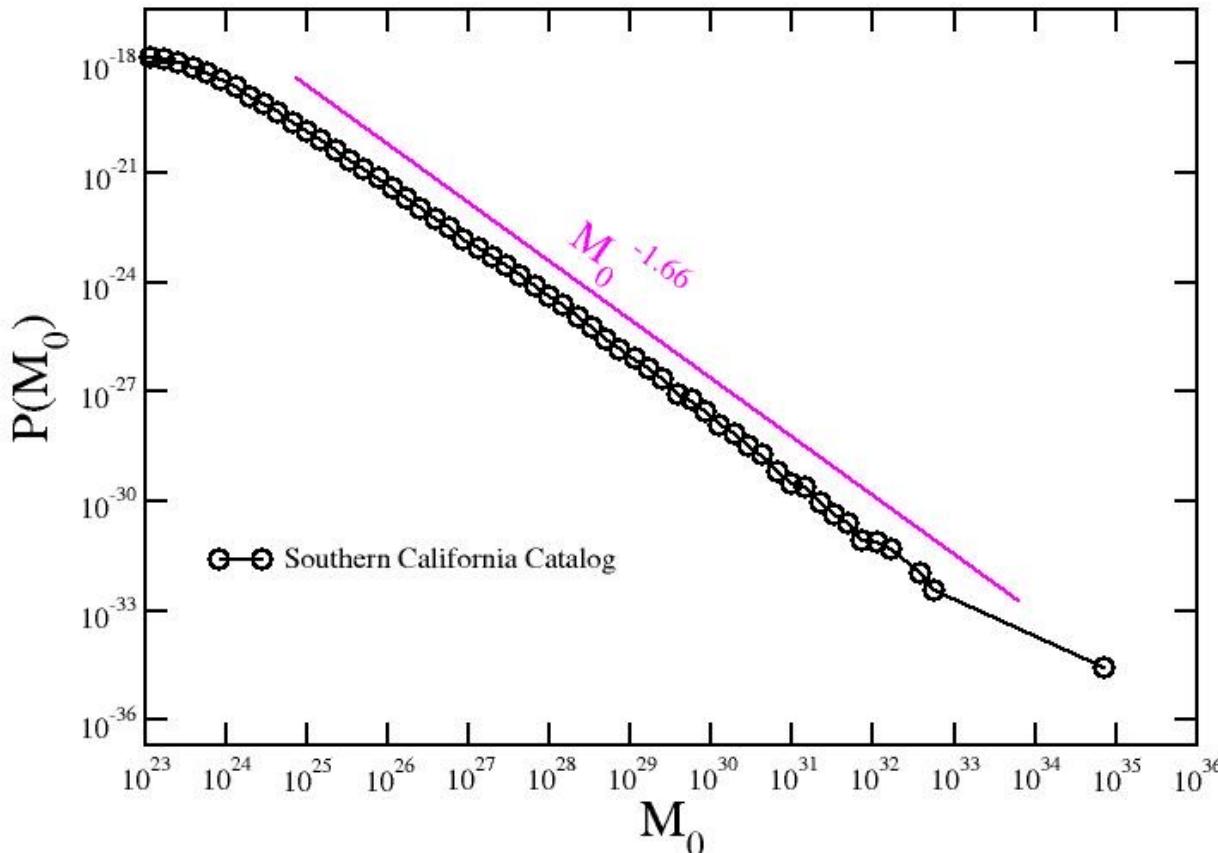
$\left. \begin{array}{l} \mu \text{ Lamè constant (shear modulus)} \\ A \text{ fractured area} \\ d \text{ average displacement} \end{array} \right\}$

$$m = \frac{2}{3} \log_{10}(M_0) - 10.7$$



Hiroo Kanamori |

Power law behavior of size distribution (Gutenberg-Richter law)



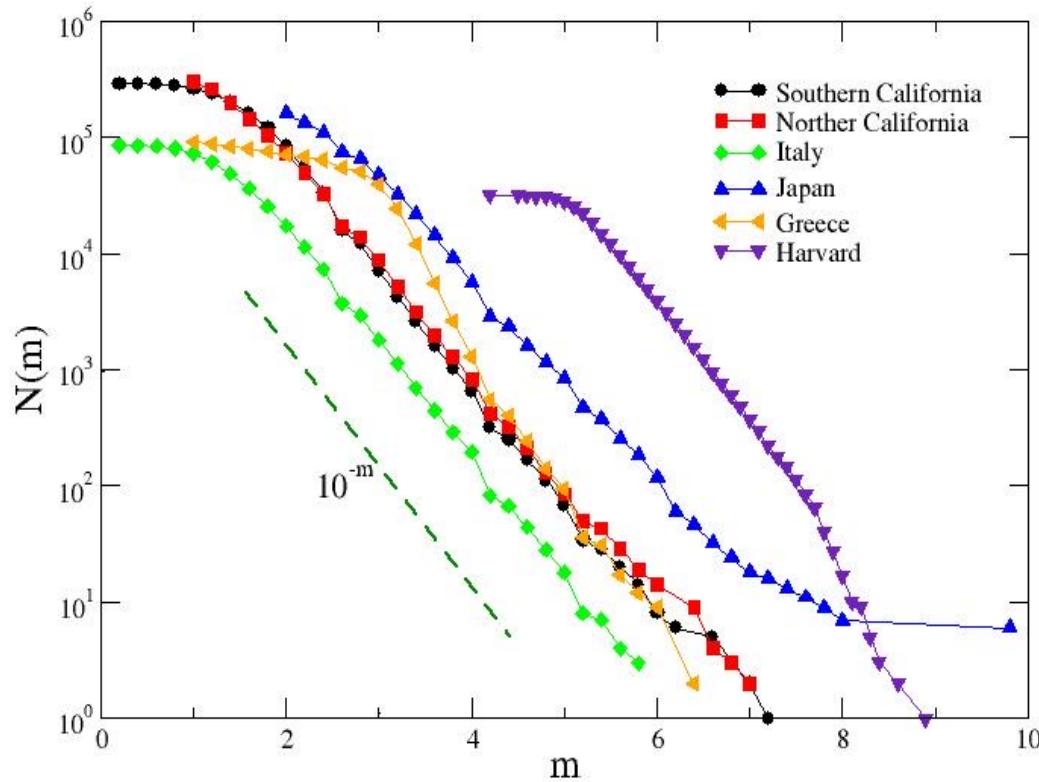
$$P(m) = k 10^{-bm}$$

$$m = \frac{2}{3} \log_{10}(M_0) - 10.7$$

≈

Power law behavior of size distribution (Gutenberg-Richter law)

Universality of the b-value



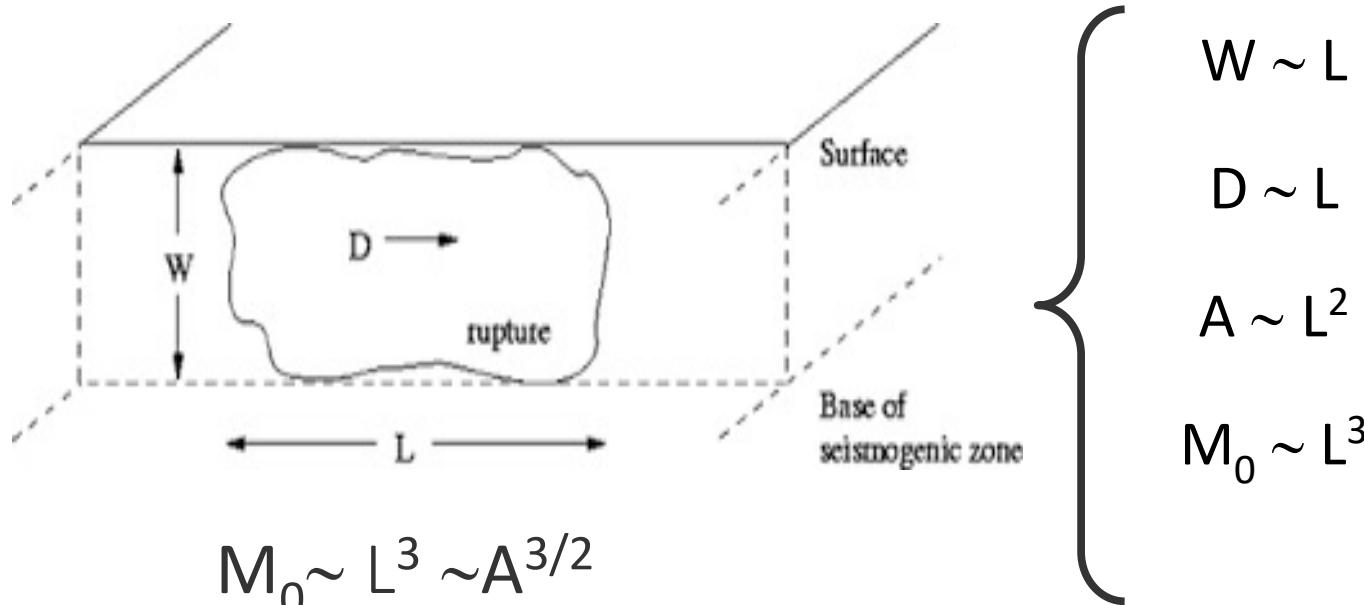
$$P(m) = k10^{-bm}$$

$$m = \frac{2}{3} \log_{10}(M_0) - 10.7$$

≈

SCALING LAWS IN SEISMIC OCCURRENCE

In the geophysical community scaling laws are related to the standard definition of scale invariance:
only one length scale in the process L

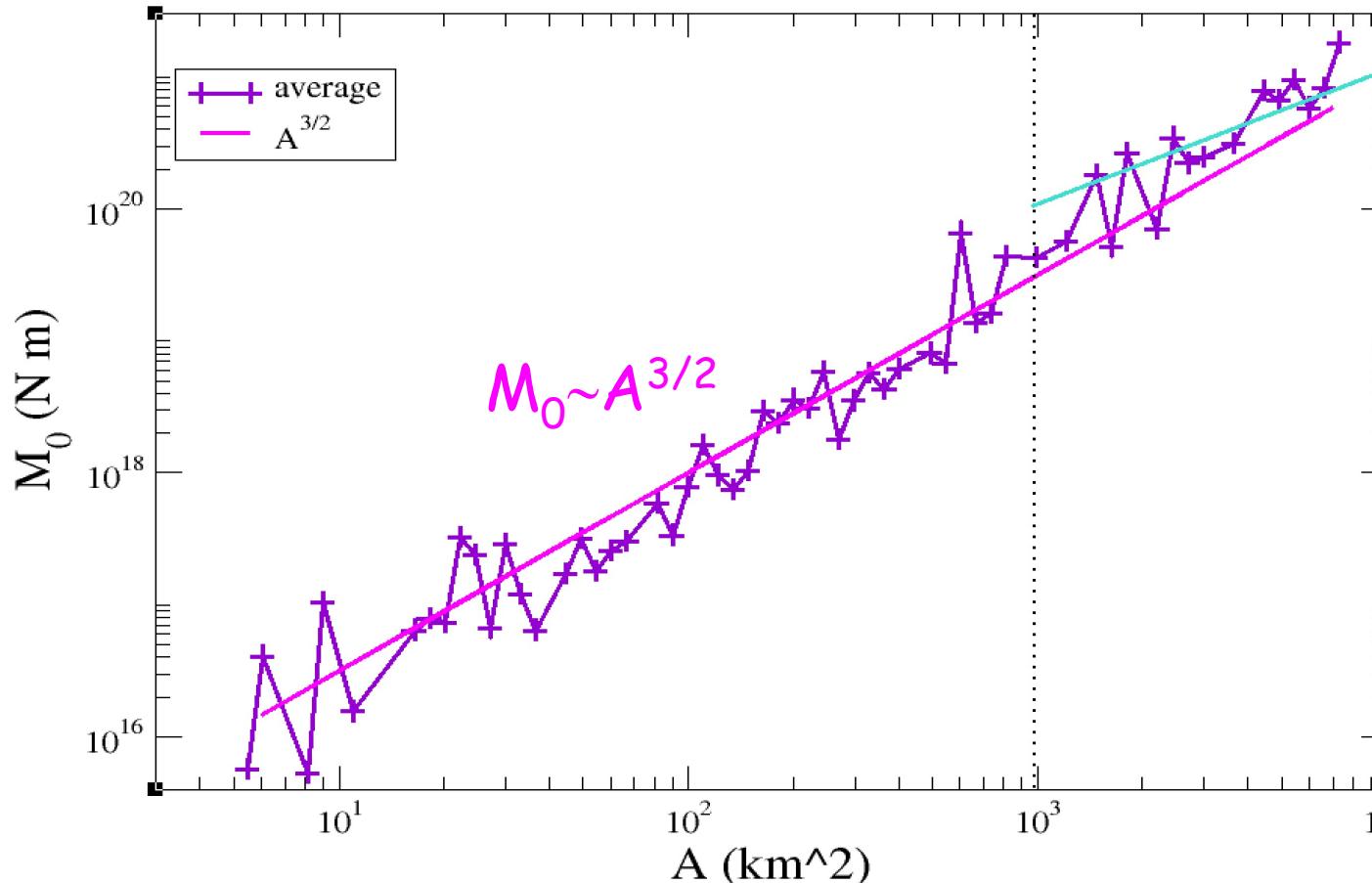


$$m = \frac{2}{3} \log_{10}(M_0) - 10.7$$

$$m \sim \gamma \log(A) \quad \gamma=1$$

SCALING LAWS IN SEISMIC OCCURRENCE

A crossover is observed at large L



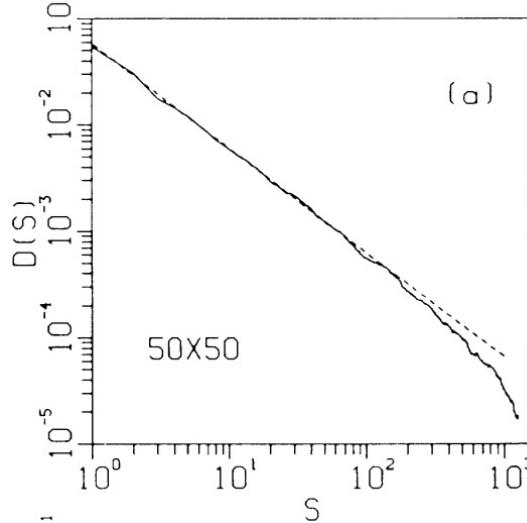
Different scaling for
 $M_0 > 10^{19}$ (m>7.5)

The crossover is usually attributed to
the finiteness of the seismogenic layer:
 W cannot be larger than H so $W \sim L^0$
(SHOLTZ 1982)

1989

SELF-Organized Criticality

Power law in the avalanche size distribution



Bak, Tang & Wiesenfeld 1987

Bak & Tang 1989

JOURNAL OF GEOPHYSICAL RESEARCH, VOL. 94, NO. B11, PAGES 15,635–15,637, NOVEMBER 10, 1989

Earthquakes as a Self-Organized Critical Phenomenon

PER BAK AND CHAO TANG

$$Z(i, j) \rightarrow Z(i, j) - 4$$

$$Z(i \pm 1, j) \rightarrow Z(i \pm 1, j) + 1$$

$$Z(i, j \pm 1) \rightarrow Z(i, j \pm 1) + 1$$

The model is actually very close to the generally accepted “block spring” picture of earthquakes [Burridge and Knopoff, 1967; Mikumo and Miyatake, 1978, 1979]. This is precisely why we believe that our results apply to earthquakes; we do not have to invoke a new and different local mechanism.

1967

Bulletin of the Seismological Society of America. Vol. 57, No. 3, pp. 341-371. June, 1967

MODEL AND THEORETICAL SEISMICITY

BY R. BURRIDGE AND L. KNOPOFF

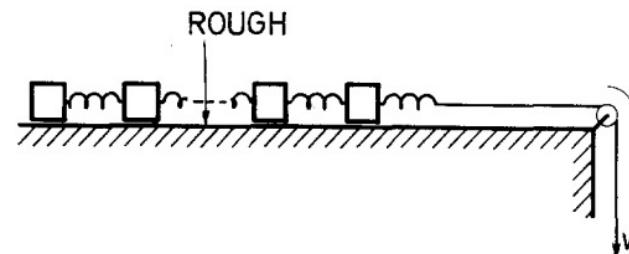
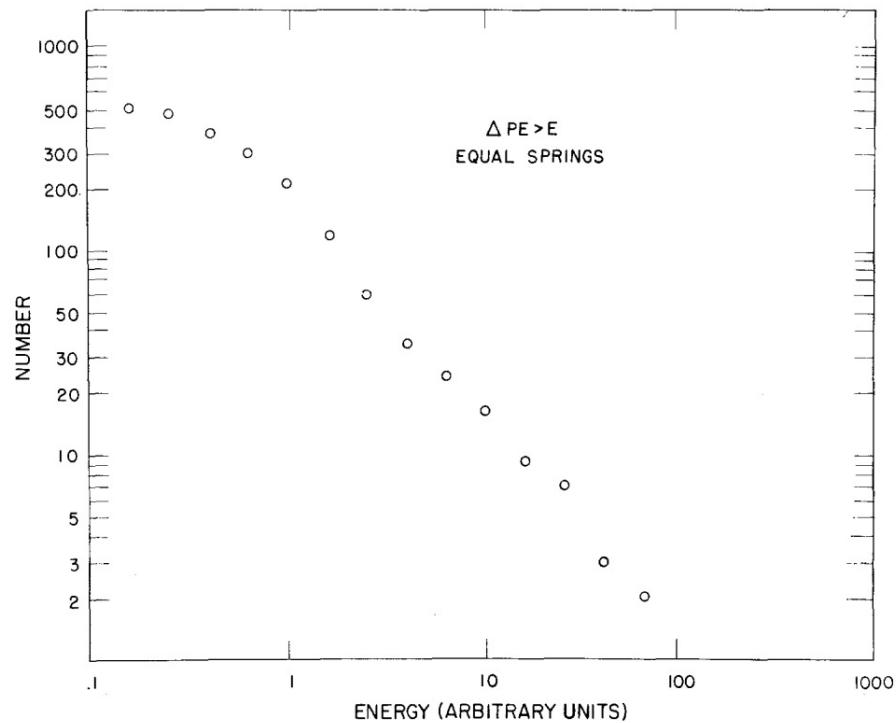


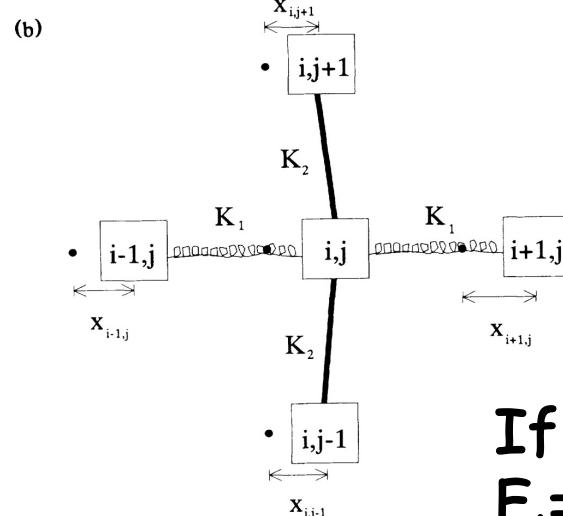
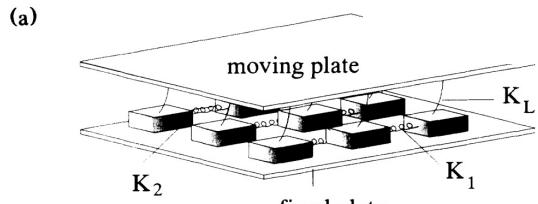
FIG. 3. Schematic diagram of the laboratory model.

Self-Organized Criticality in a Continuous, Nonconservative Cellular Automaton Modeling Earthquakes

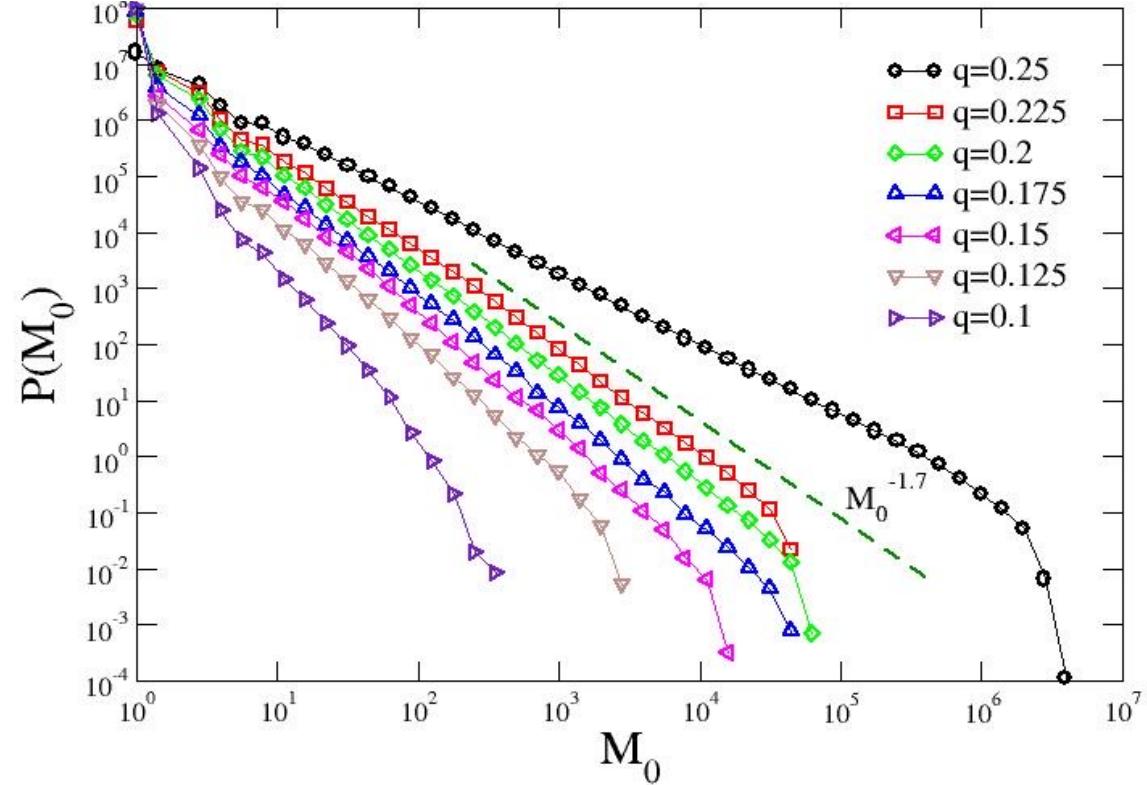
Zeev Olami, Hans Jacob S. Feder, and Kim Christensen ^(a)

Department of Physics, Brookhaven National Laboratory, Upton, New York 11973

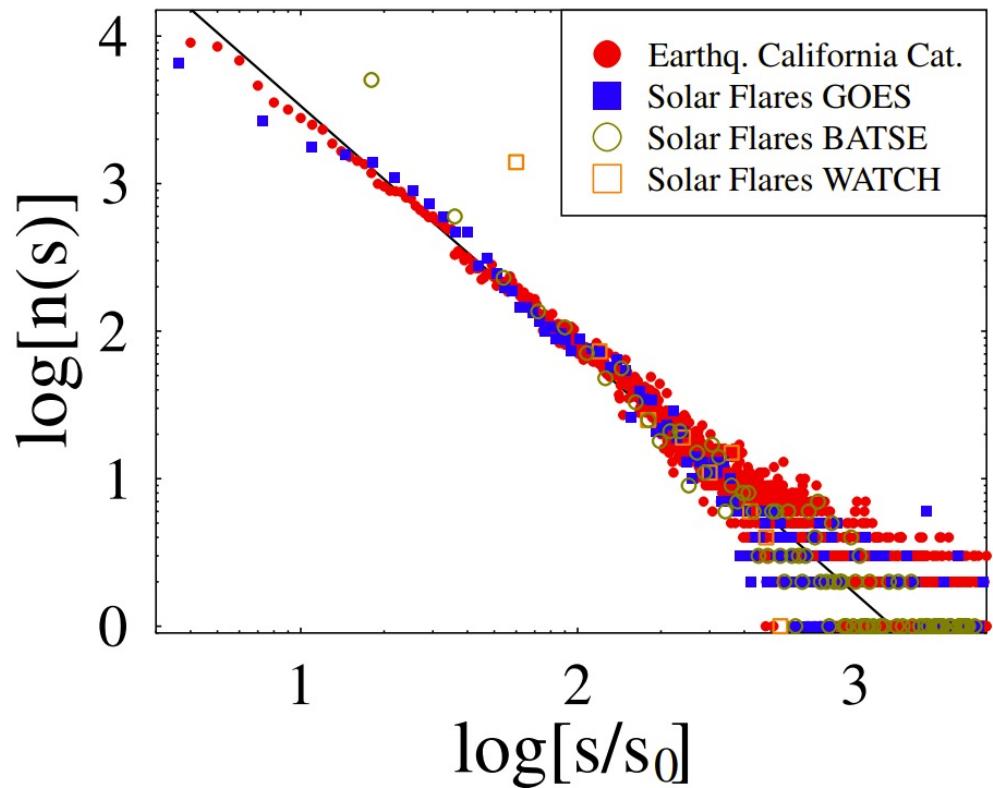
(Received 19 August 1991)



If $F_i > F_{th}$, $F_i = 0$
 $F_j = F_j + qF_i$
with $q = (k_1 + k_0)/4k_1$



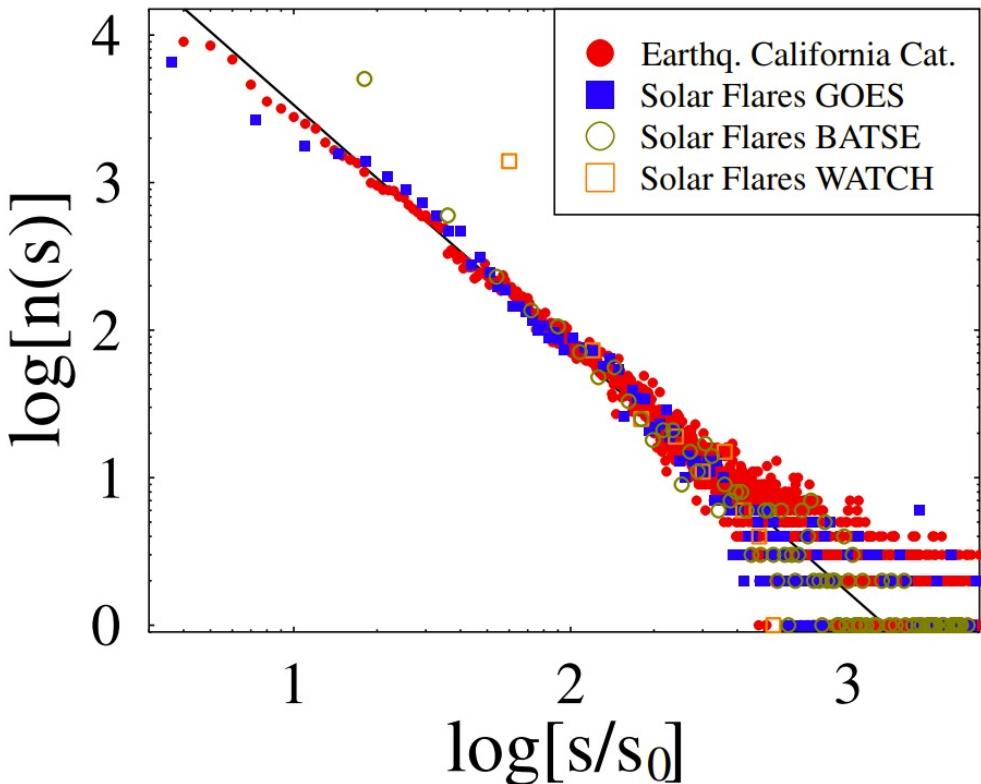
Similar scaling law in Solar Flare occurrence



Similar scaling law in Solar Flare occurrence

debates

11 March 1999



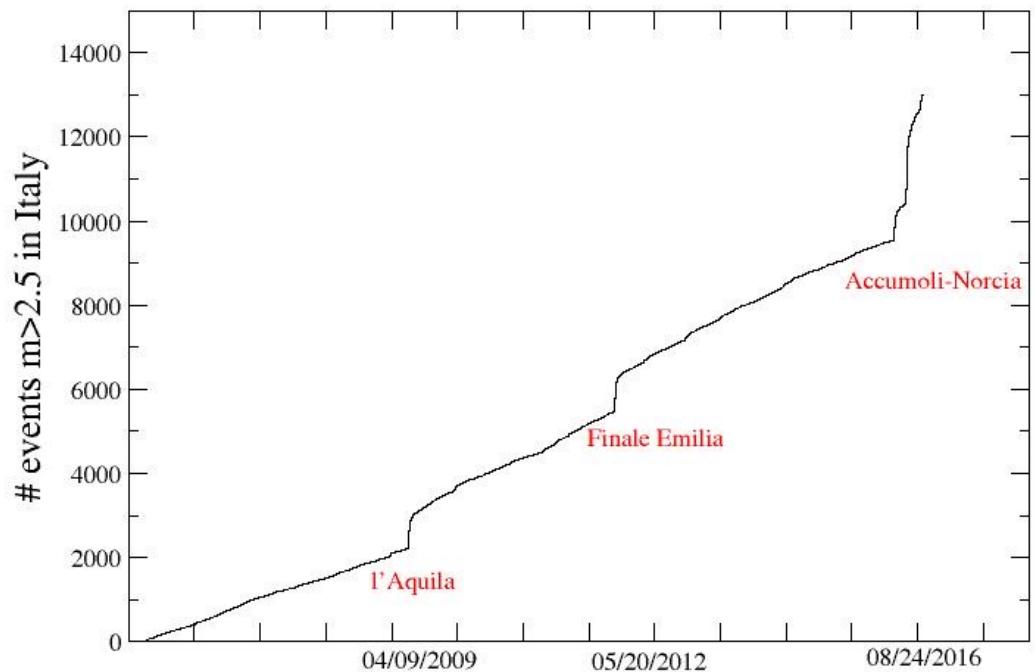
Earthquake prediction is difficult but not impossible

LEON KNOPOFF

Statistics of rare events

The small number of events means that again we need a physics-based theory of the precursory process to amplify the meager data. In the area of physics, another blind alley was followed. The beguiling attractiveness of the illusion of scale-independence of the G-R law suggested that the model of self-organized criticality (SOC), which also yielded scale-independent distributions, might be appropriate. (The logic is evidently faulty: if mammals have four legs, and tables have four legs, it does not follow that tables are mammals, or the reverse.) The model of SOC permits a hierarchical development of large events out of the nonlinear interaction of smaller events, at rates in relation to their sizes, and culminating in the largest event. However, there are several important arguments against the applicability of SOC to the earthquake problem.

Temporal Clustering

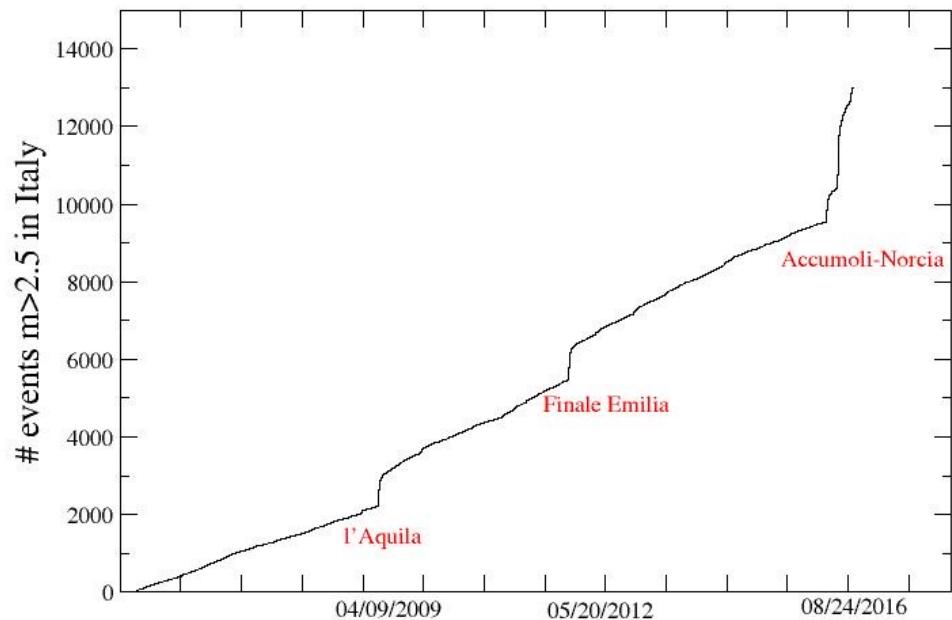


Mainshock: The largest event in a seismic sequence

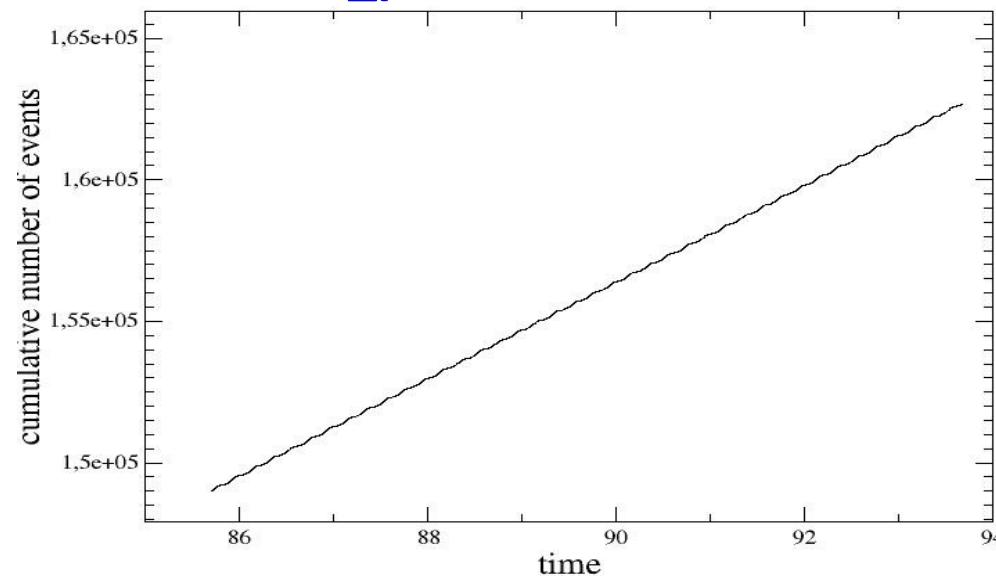
Aftershocks the earthquakes that follow the largest shock of an earthquake sequence. They are typically smaller than the mainshock and occur at few km from the mainshock epicenter

The majority of events in seismic catalogs are aftershocks!

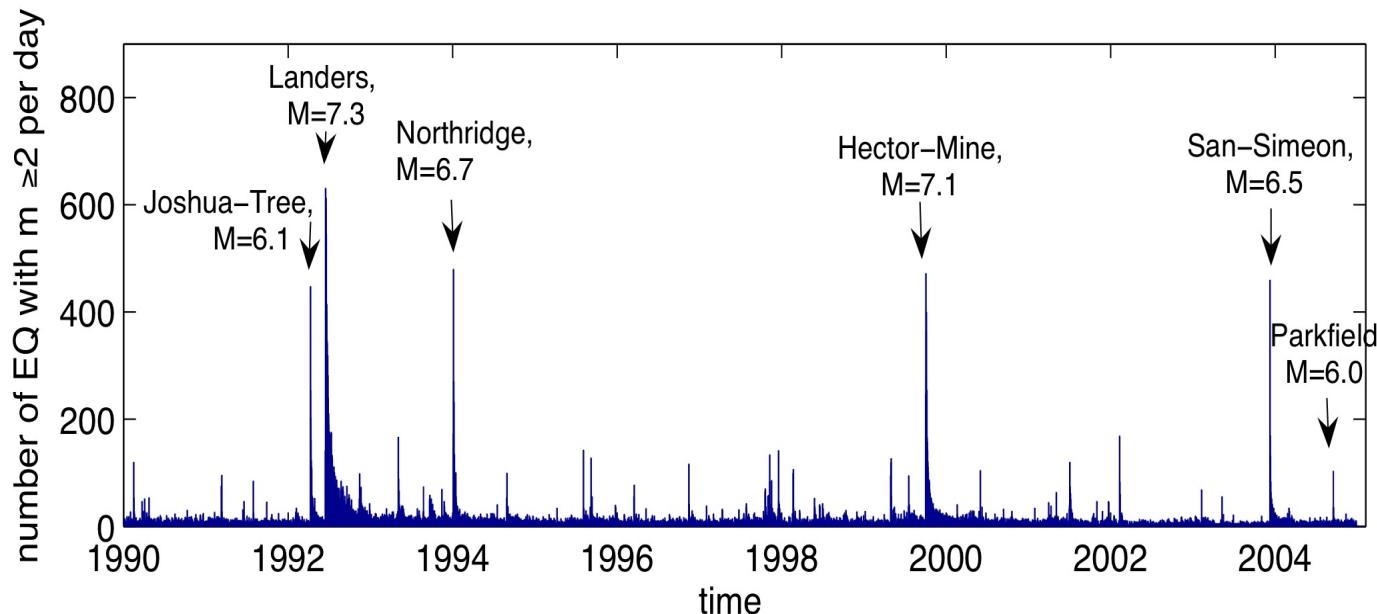
Temporal Clustering



**NO temporal clustering
(NO aftershocks) in the
Sand_pile model**

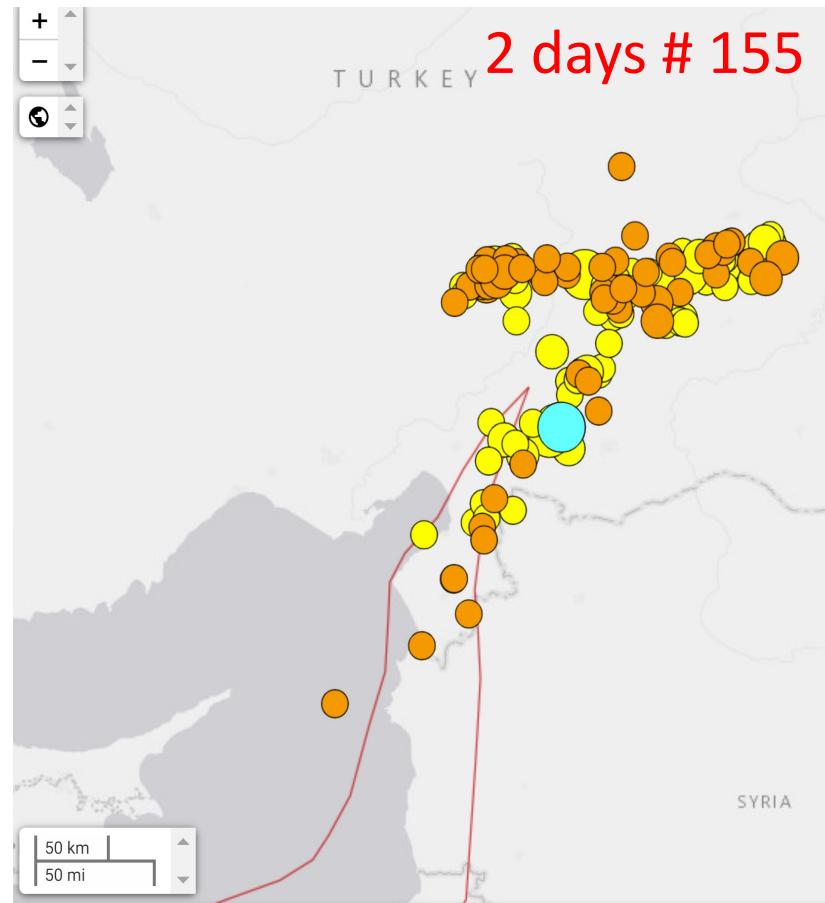
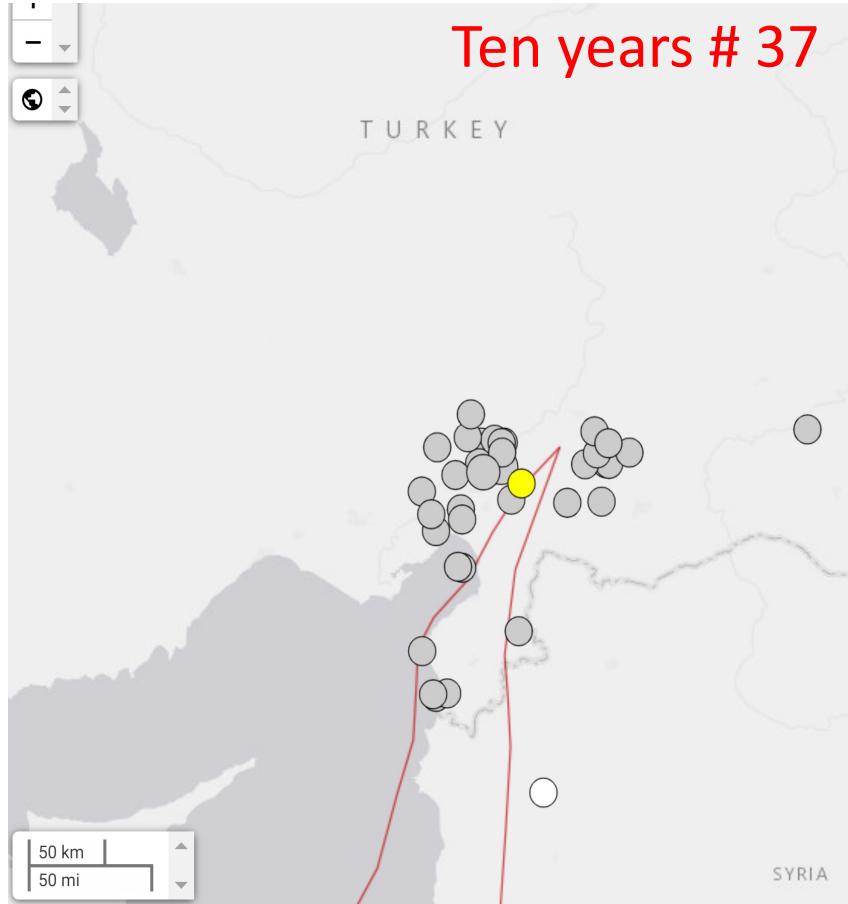


Temporal Clustering



The majority of events in seismic catalogs are aftershocks!

05-02-2023 Earthquake

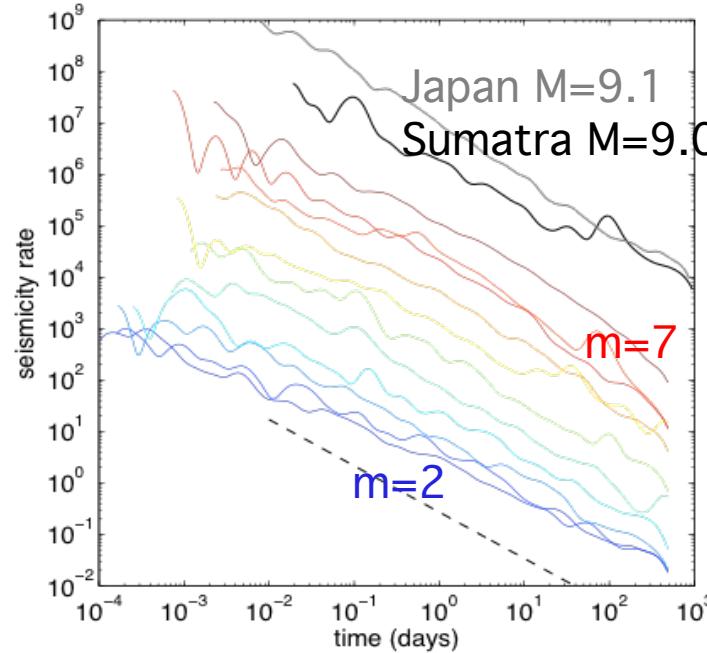
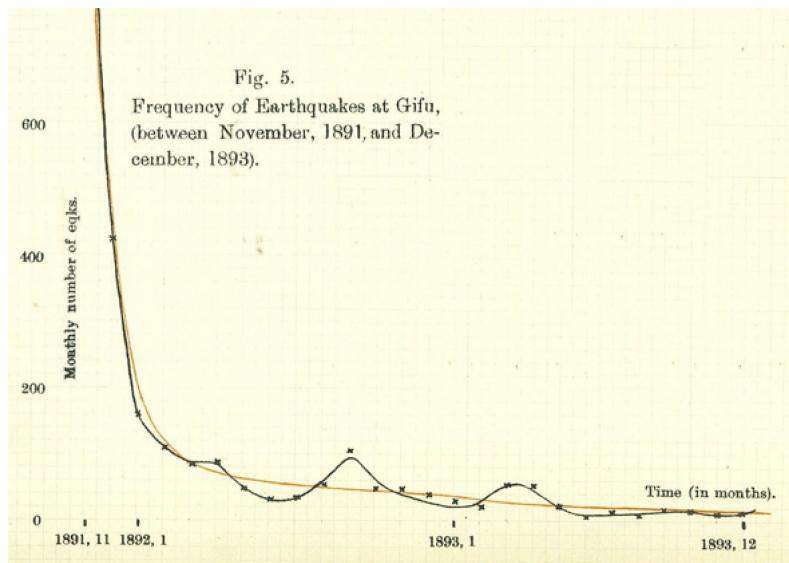


1889

The Omori law

Fusakichi Omori in 1889 published his work on the aftershocks of earthquakes, in which he stated that aftershock frequency decreases by roughly the reciprocal of time after the main shock.

$$n(t) = \frac{k}{t}$$



Fusakichi Omori

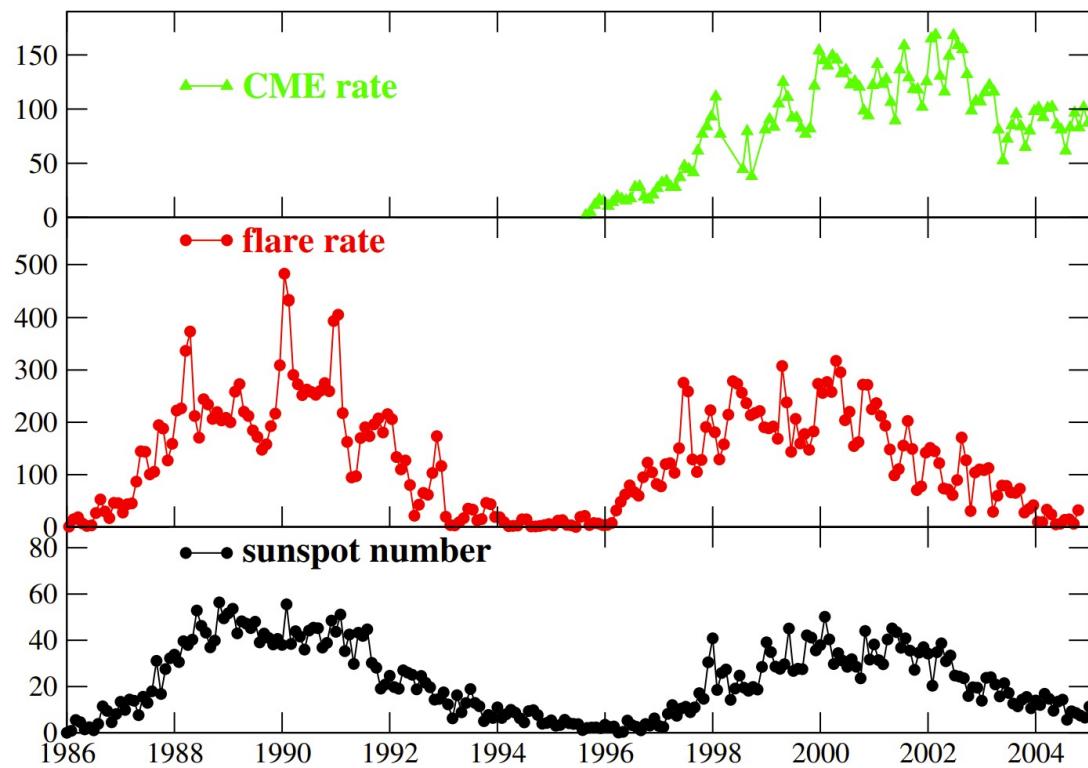
The modified version of Omori's law, now commonly used, was proposed by Utsu in 1961, with typical values of p [0.75:1.5].

$$n(t) = \frac{k}{(t+c)^p}$$

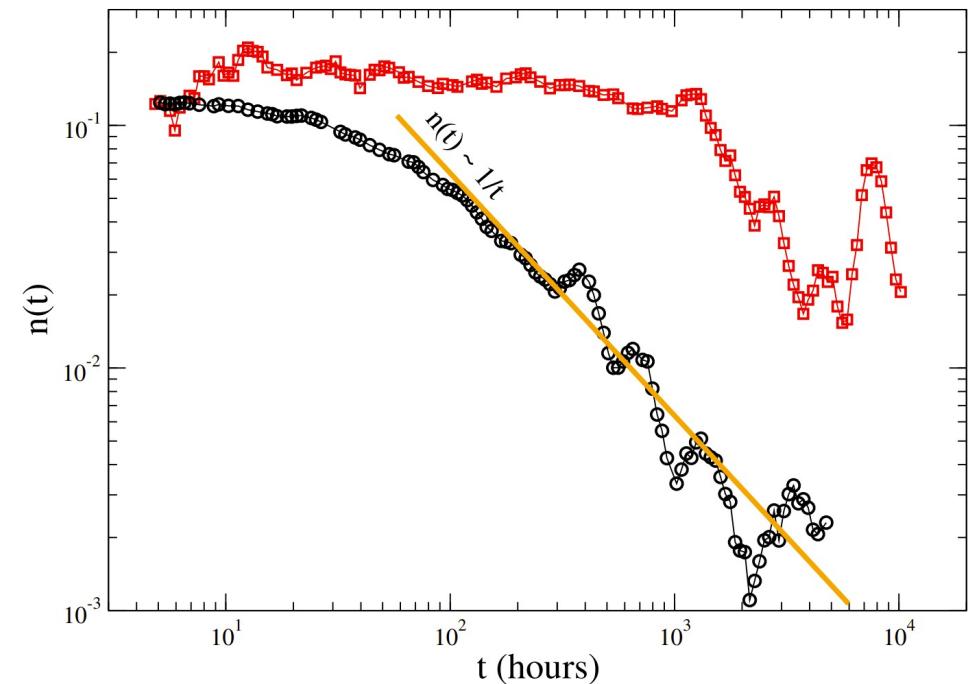
Mainshocks: largest event of a sequence

Aftershocks: its following earthquakes within a given space-time distance

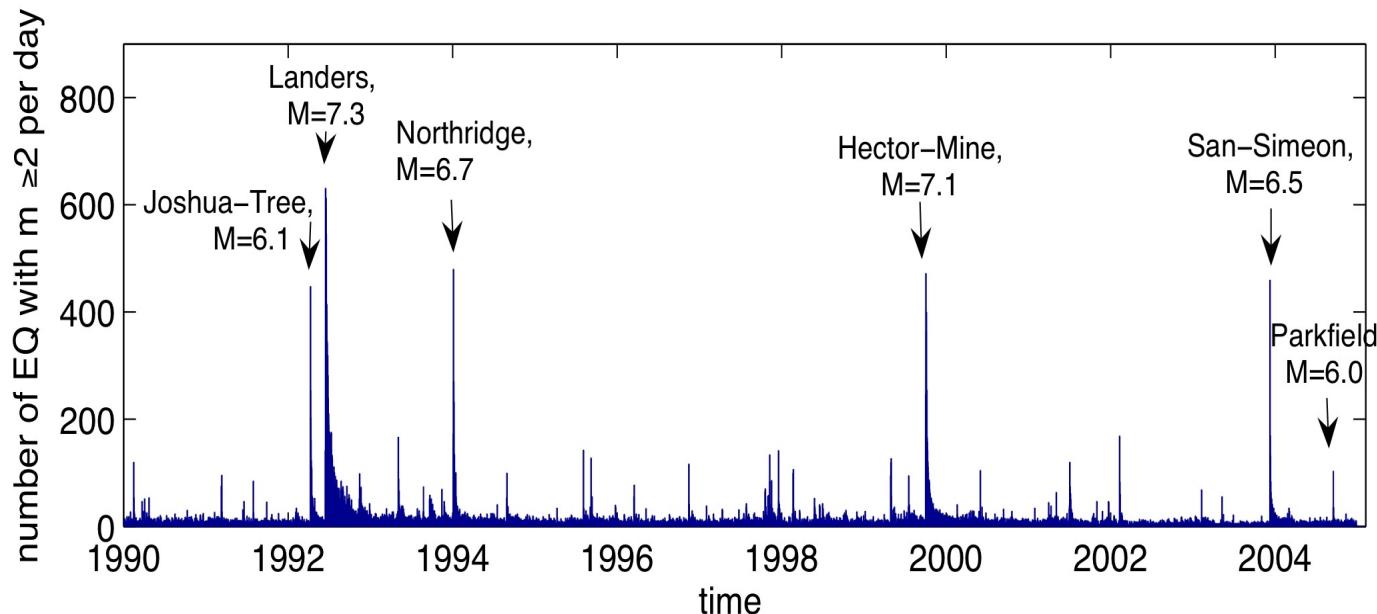
Similar scaling law in Solar Flare occurrence



$$n(t) = \frac{\mathcal{N}(t, e_{th})}{\sum_{i=1}^L \Theta(t - t_{i+1})}$$



Temporal Clustering

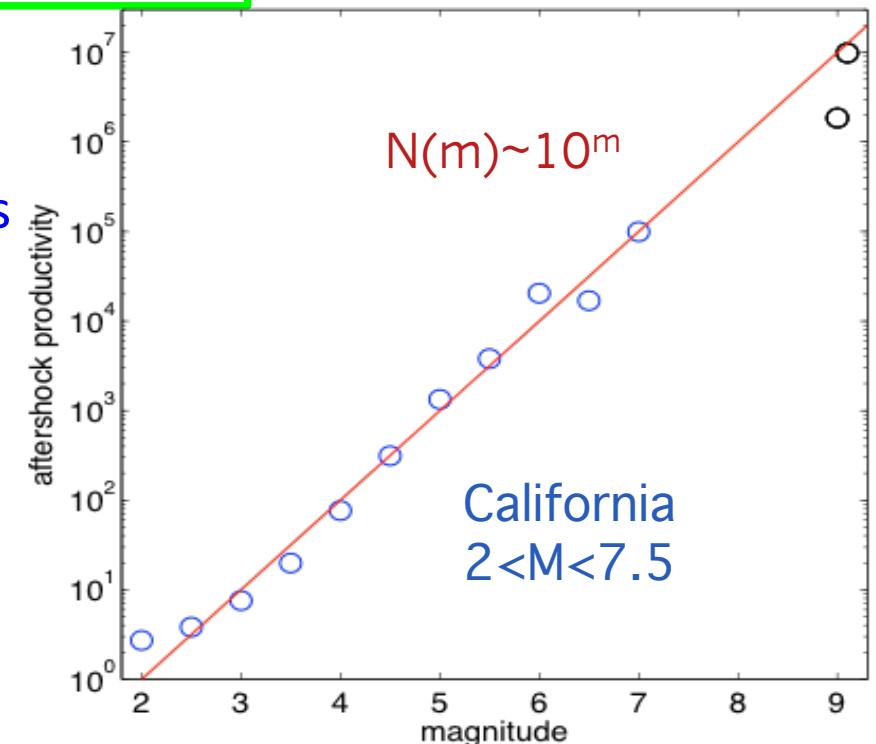


The majority of events in seismic catalogs are aftershocks!

Productivity law

$$N_{\text{aftershocks}} \approx 10^{\alpha m} \quad n \approx 2/3 \log_{10}(M_0) \Rightarrow N_{\text{aftershocks}} \approx M_0^{1+2/3\alpha} \quad \alpha \approx 1$$

In 1970 Utsu observed that the number of aftershocks grows exponentially with the mainshock magnitude



POINT-PROCESS models for seismic forecasting

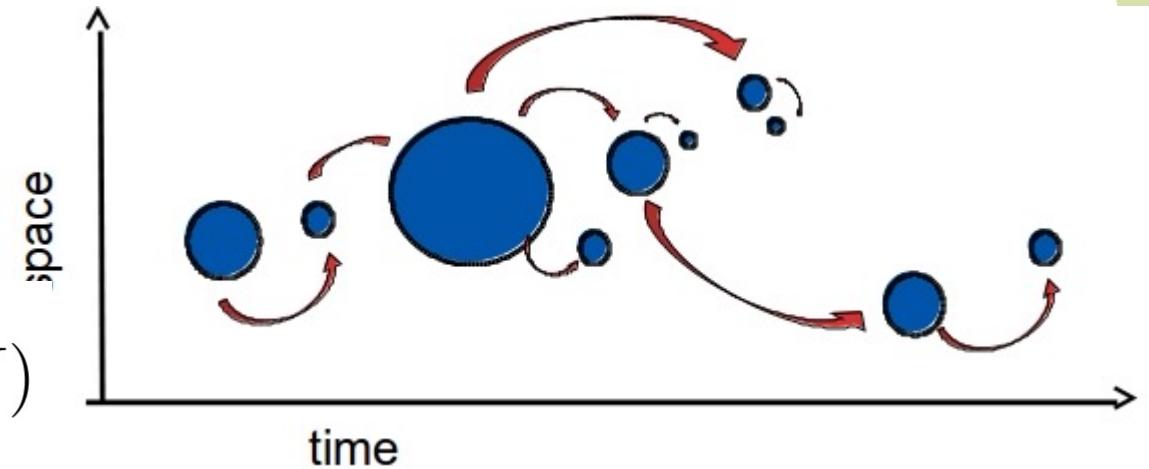
An Earthquakes can trigger more than one daughter earthquake:

EPIDEMIC TYPE AFTERSHOCK SEQUENCE (ETAS) model

Conditional Probability to have an earthquake of magnitude M at time t

$$\Psi(M, t) = \sum_{i:t_i < t} \rho(M, t|M_i, t_i) + \mu P(M)$$

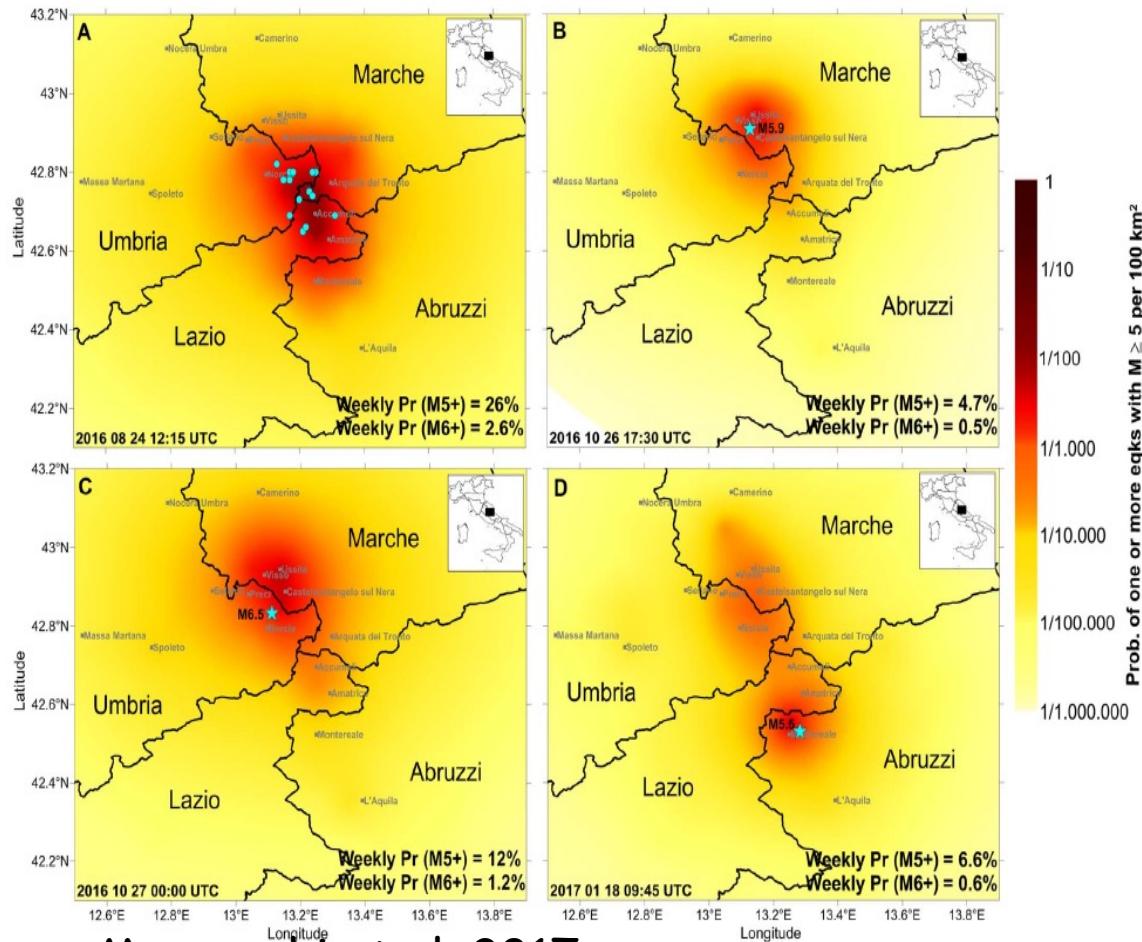
The sum extends over all previous earthquakes



Key approximation:
Factorization between
space, time and magnitude

$$\rho(M, t|M_i, t_i) = K 10^{\alpha M_i} \left[\frac{(t - t_i)}{c} + 1 \right]^{-p} P(M)$$

The ETAS model for the 2016 Amatrice-Norcia sequence



ETAS is usually considered the best available forecasting instrument

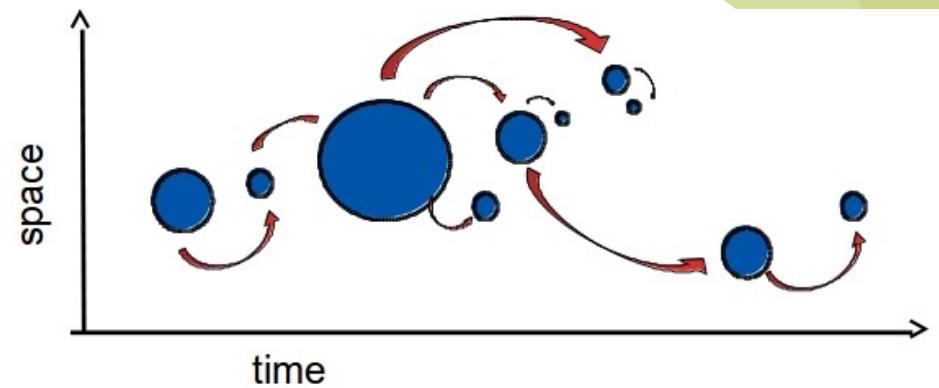
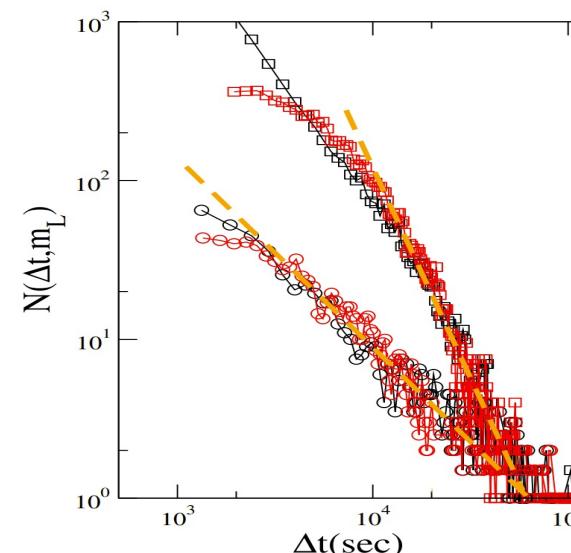
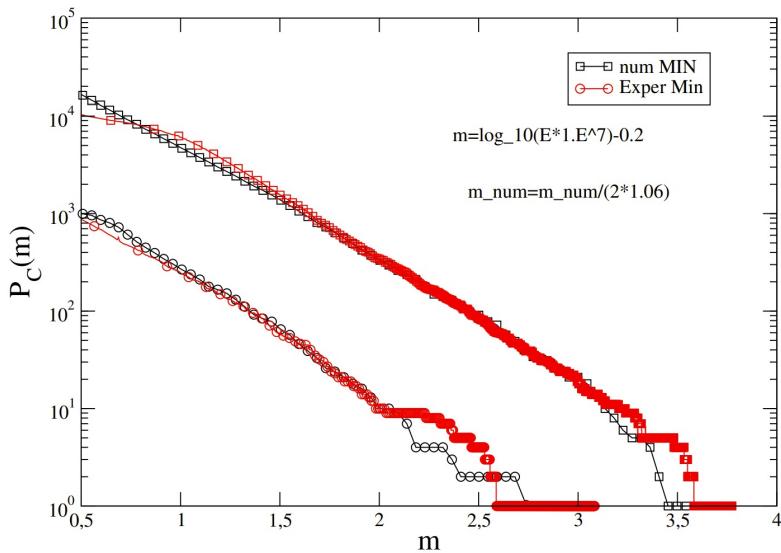
Marzocchi et al. 2017

POINT-PROCESS models for solar flares

$$\Psi(M, t) = \sum_{i:t_i < t} \rho(M, t|M_i, t_i) + \mu P(M)$$

$$\mu(t) = A \cos(2\pi t/T)$$

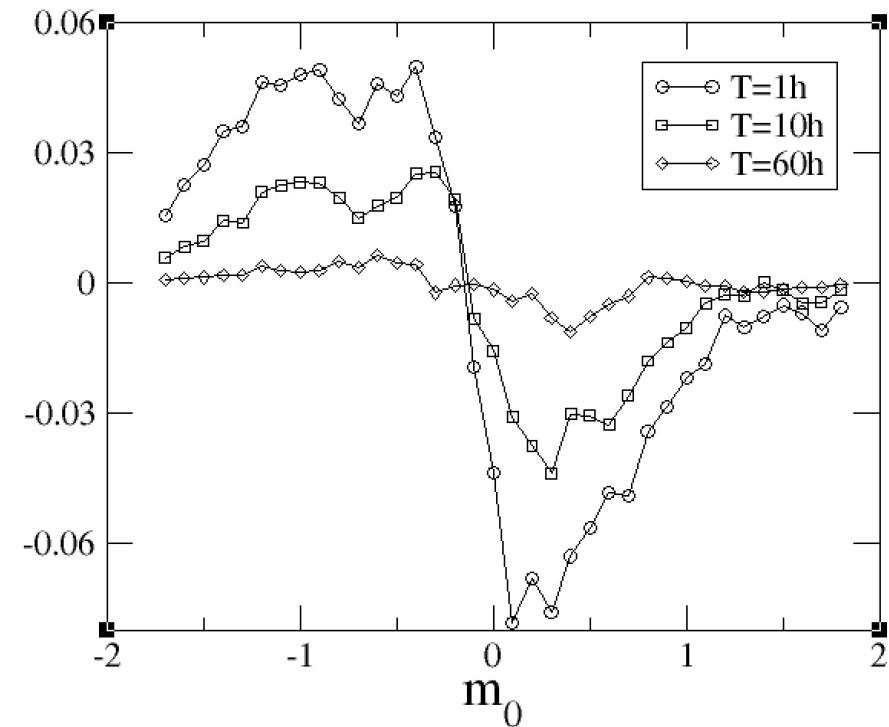
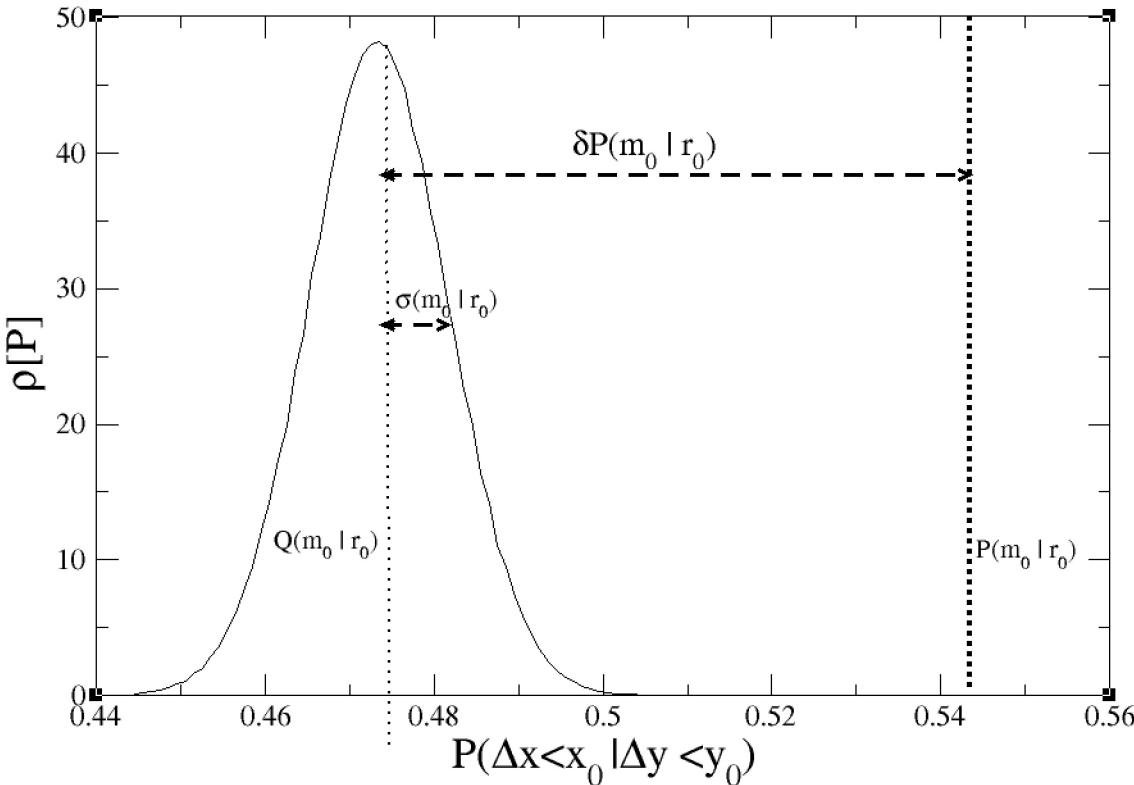
$$\rho(M, t|M_i, t_i) = \frac{k}{t - t_i}$$



Key approximation:
Factorization between
space, time and magnitude

TIME-MAGNITUDE Correlation

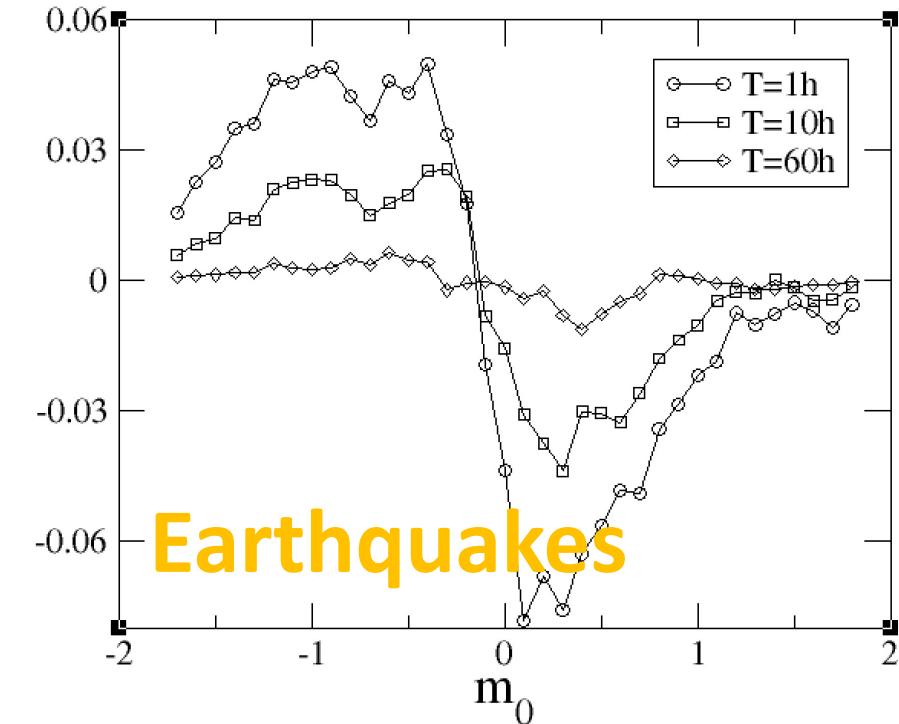
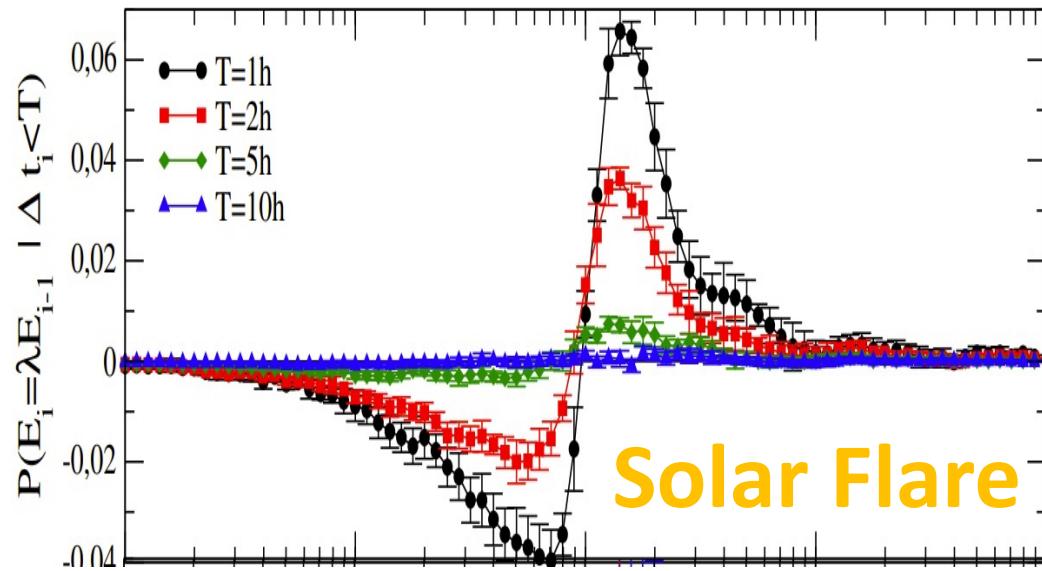
$$dP(M_{i+1} = M_i + m_0 | t_{i+1} - t_i < T) \\ = P(M_{i+1} = M_i + m_0 | t_{i+1} - t_i < T) - P(M_{k^*} = M_{i+1} + m_0 | t_{i+1} - t_i < T)$$



The next earthquake is about -1 smaller than the previous one

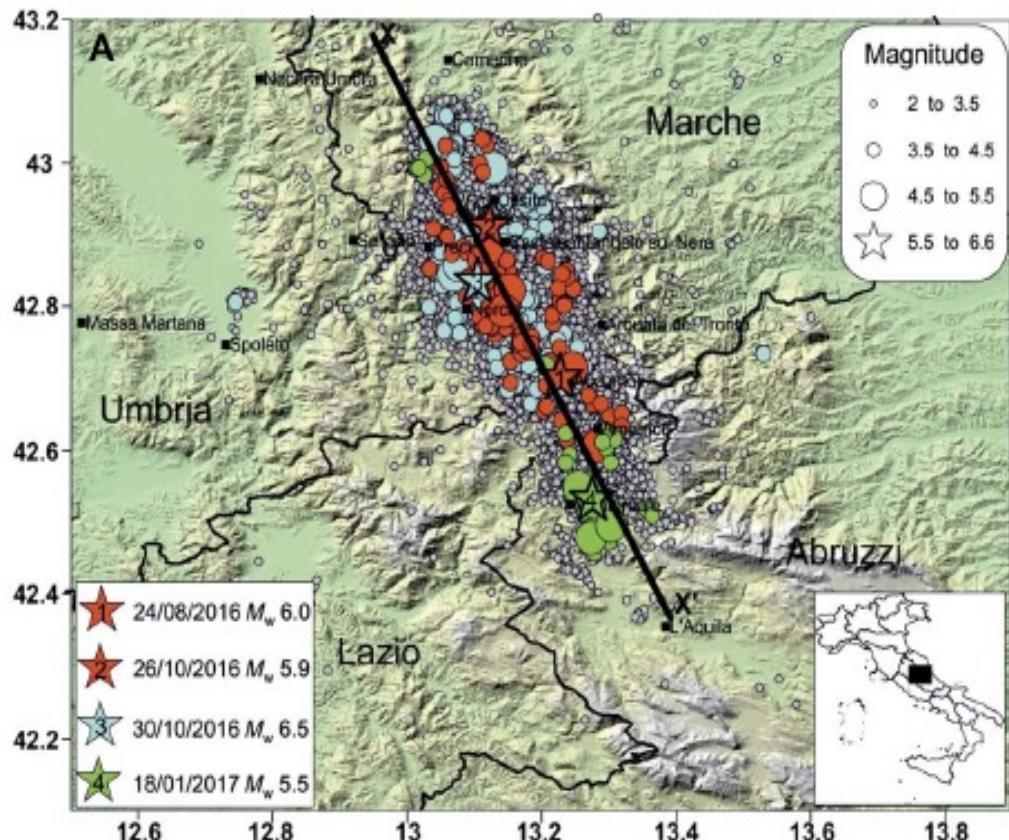
TIME-MAGNITUDE Correlation

$$dP(M_{i+1} = M_i + m_0 | t_{i+1} - t_i < T) \\ = P(M_{i+1} = M_i + m_0 | t_{i+1} - t_i < T) - P(M_{k^*} = M_{i+1} + m_0 | t_{i+1} - t_i < T)$$

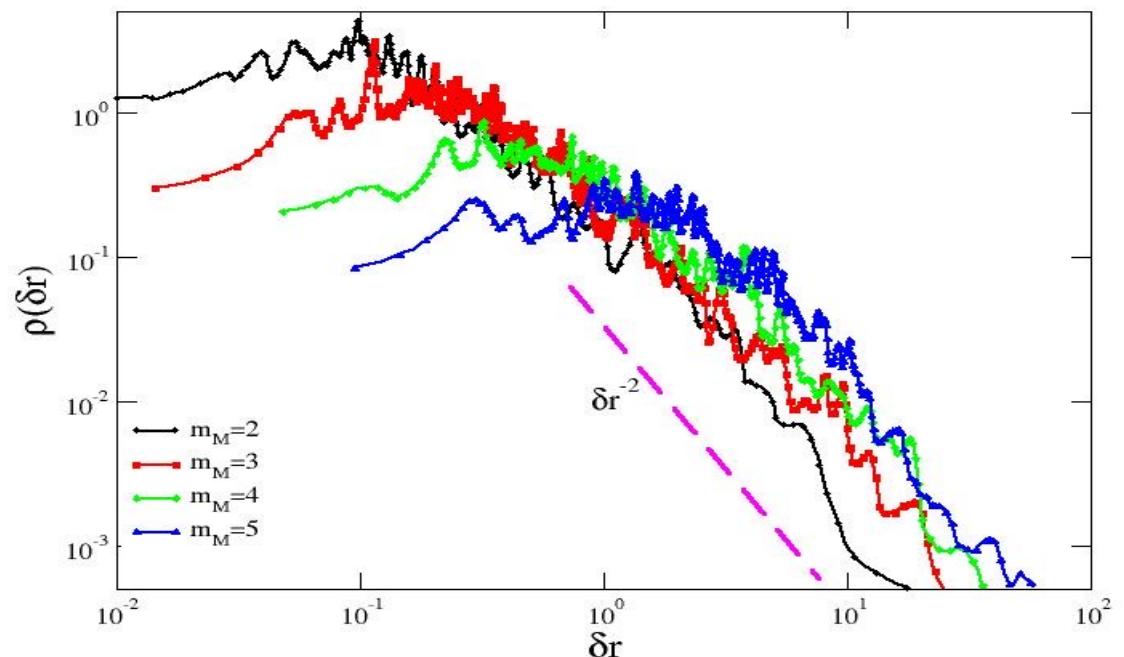


The next earthquake is about -1 smaller than the previous one
The next flare is on average larger than the previous one

SPATIAL clustering



$$n(\delta r) = \frac{k}{(\delta r)^{\gamma}}$$



MISSING INGREDIENTS

debates

11 March 1999

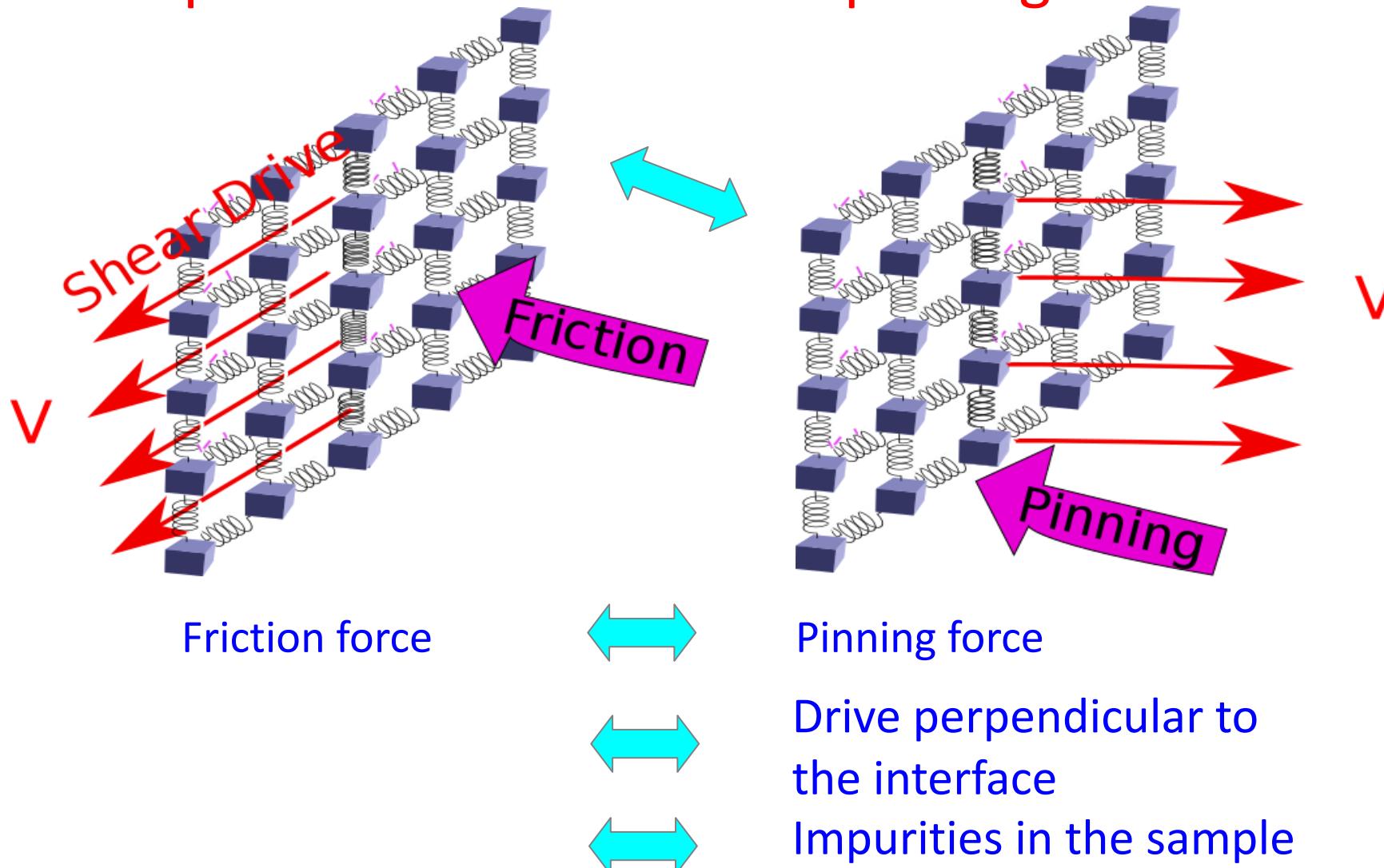
Earthquake prediction is difficult but not impossible

LEON KNOPOFF

their sizes, and culminating in the largest event. However, there are several important arguments against the applicability of SOC to the earthquake problem.

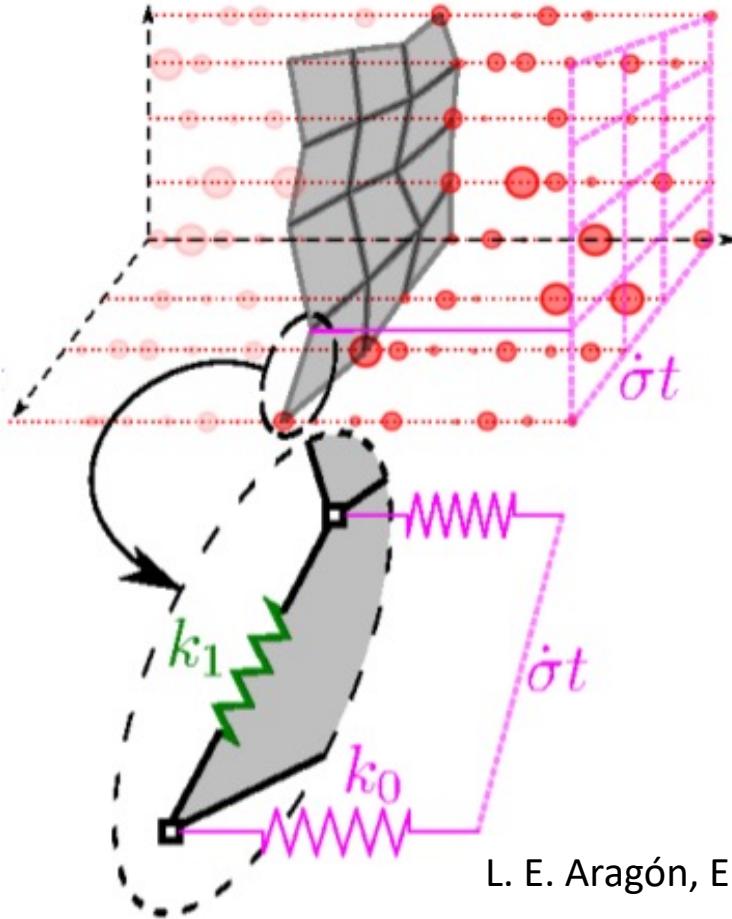
1. Faults and fault systems are inhomogeneous: we have already noted the presence of several scale sizes.

The Map on elastic interface depinning



The Map on elastic interface depinning

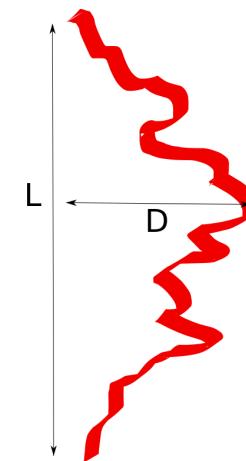
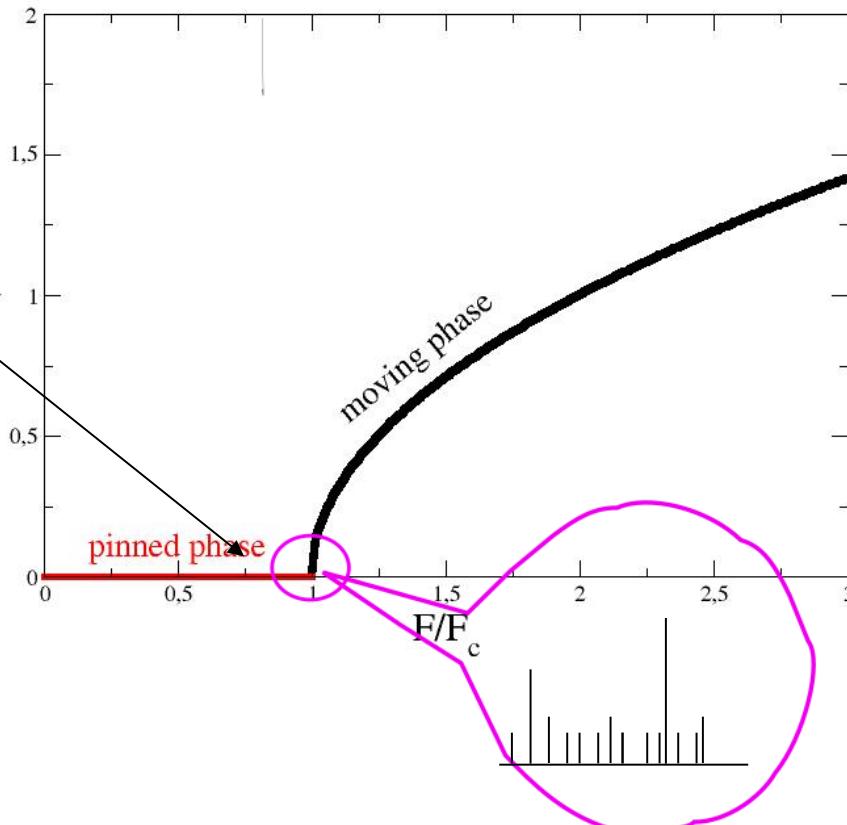
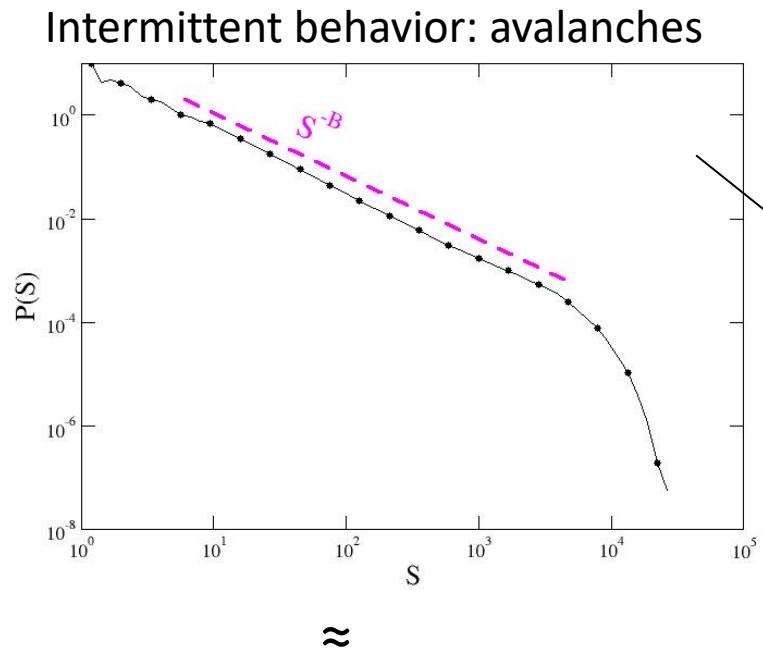
quenched Edwards-Wilkinson (qEW) model



Red dots represent pinning centers
heterogeneous in space and in depth

L. E. Aragón, E. A. Jagla, and A. Rosso, Phys. Rev. E (2012)

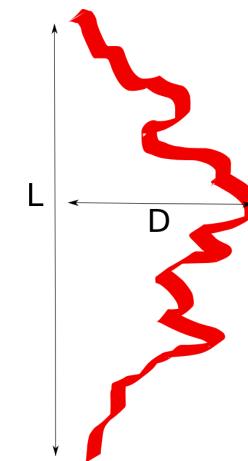
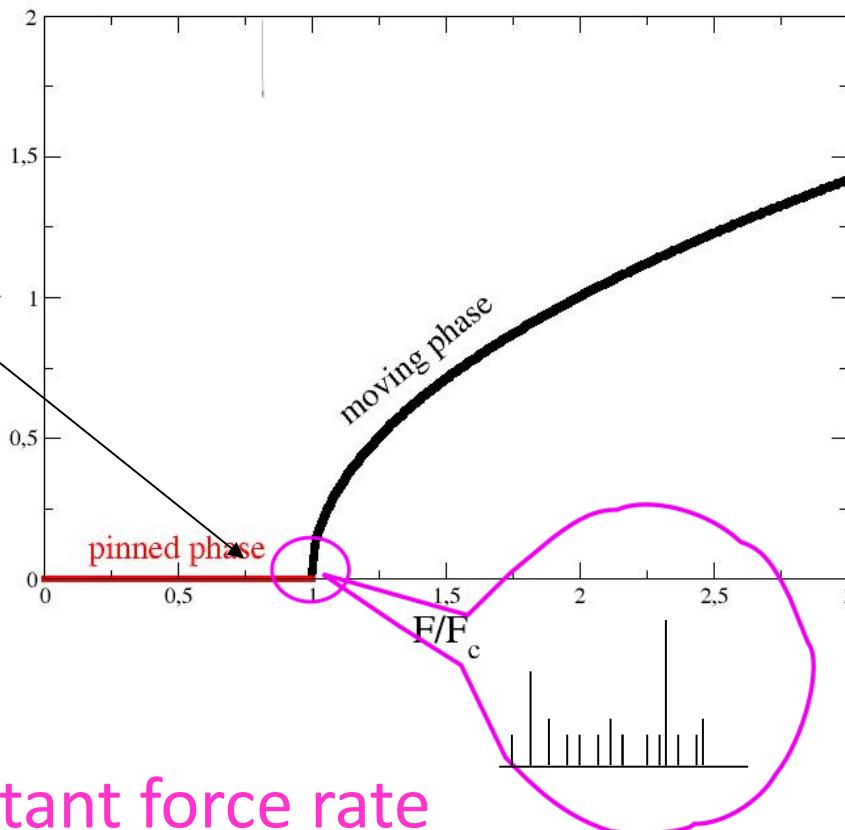
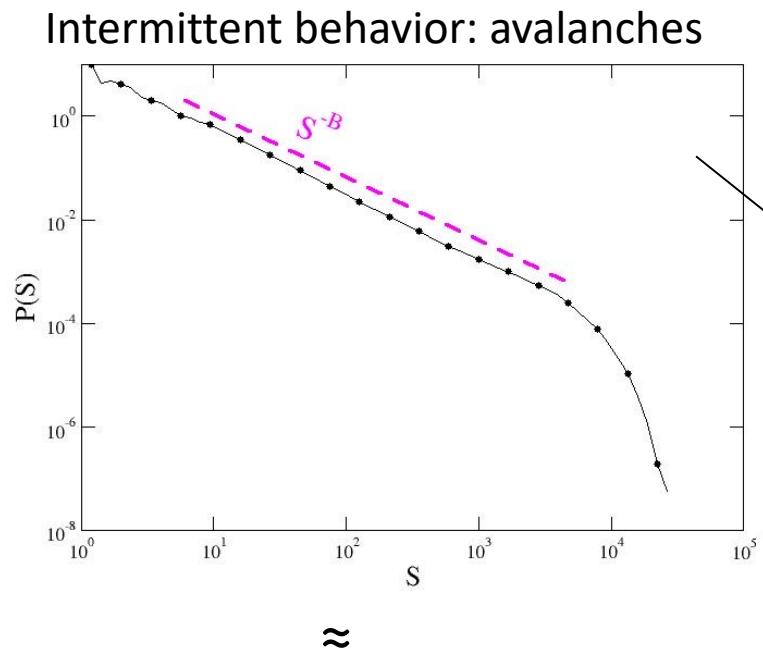
The depinning transition under a costant force



$$B=2-2/(D+\xi)$$

$$M_0 \sim L^{D+\xi}$$

The depinning transition under a costant force



$$B=2-2/(D+\xi)$$

$$M_0 \sim L^{D+\xi}$$

SOC: no costant force but costant force rate

$$\nu \rightarrow 0 \Rightarrow F \rightarrow F_C$$

The Map on elastic interface depinning

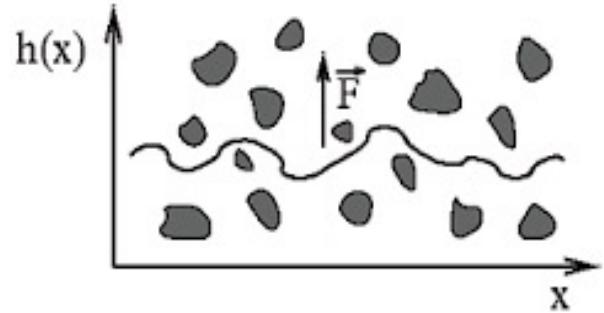
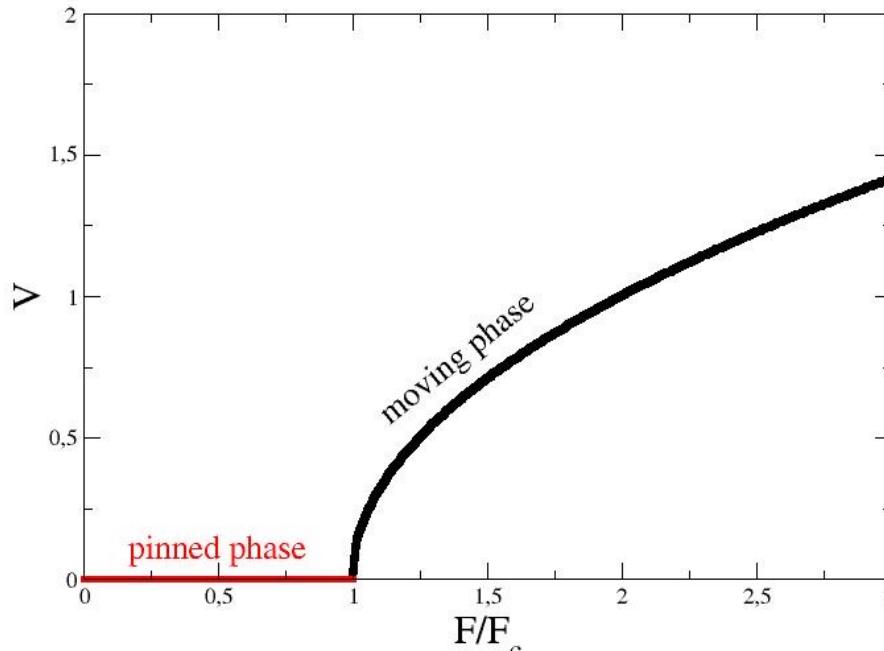


Figure 4. Schematic representation of a driven interface in a disordered media.



Different exponents of those of earthquake occurrence!!

Non uniform force $F=K(V_0 t - x(t))$, when $V_0 \sim 0$ $F \sim F_c$

Always close to a critical point

$$B = 2 - 2/(D + \xi)$$

$$M_0 \sim L^{D+\xi}$$

$$\xi = 0.75$$

$$B = 1.27 < 1.7$$

$$M_0 \sim L^{2.75} < L^3$$

The Map on elastic interface depinning

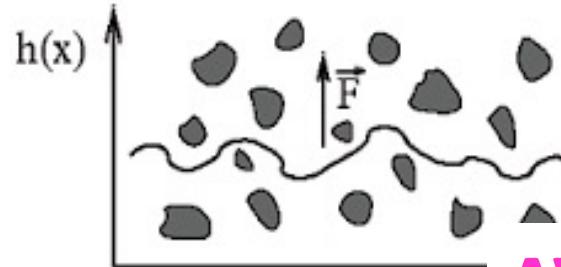
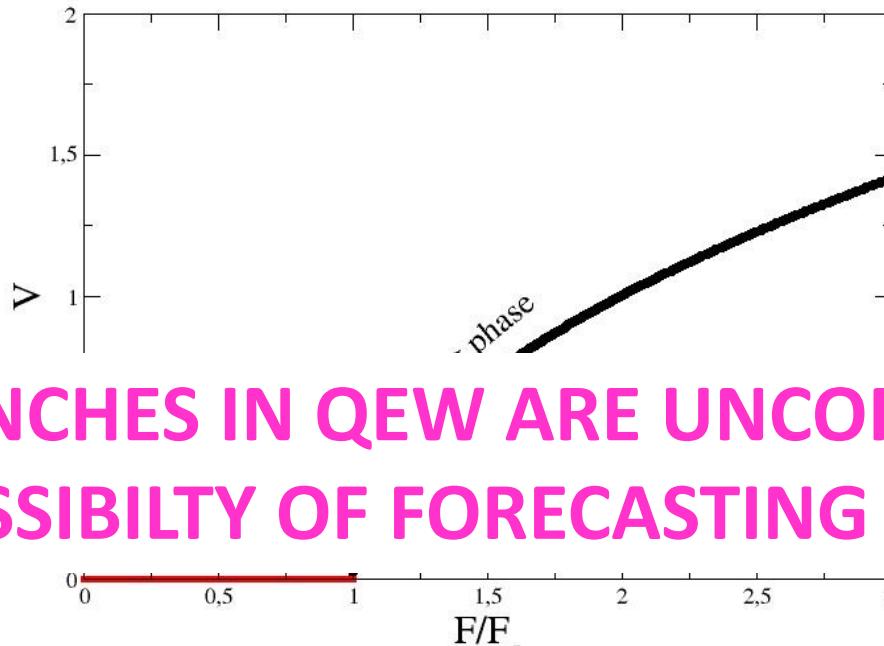


Figure 4. Schematic representation of a driven media.



AVALANCHES IN QEW ARE UNCORRELATED
NO POSSIBILITY OF FORECASTING

Different exponents of
earthquake
!!

Non uniform force $F=K(V_0 t - x(t))$, when $V_0 \sim 0$ $F \sim F_c$

Always close to a critical point

$$B=2-2/(D+\xi)$$

$$M_0 \sim L^{D+\xi}$$

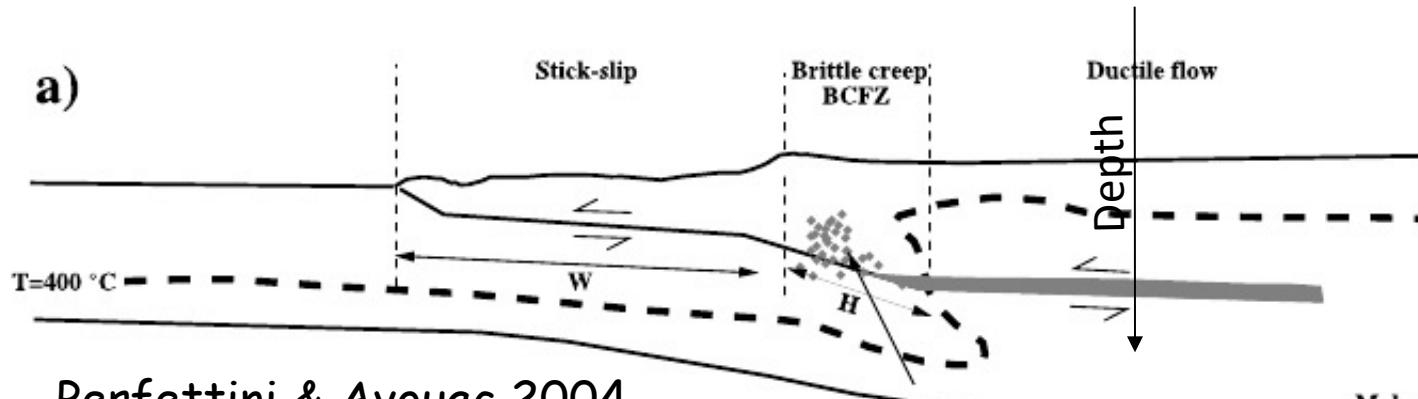
$$\xi=0.75$$

$$B=1.27 < 1.7$$

$$M_0 \sim L^{2.75} < L^3$$

The origin of aftershocks: Afterslip

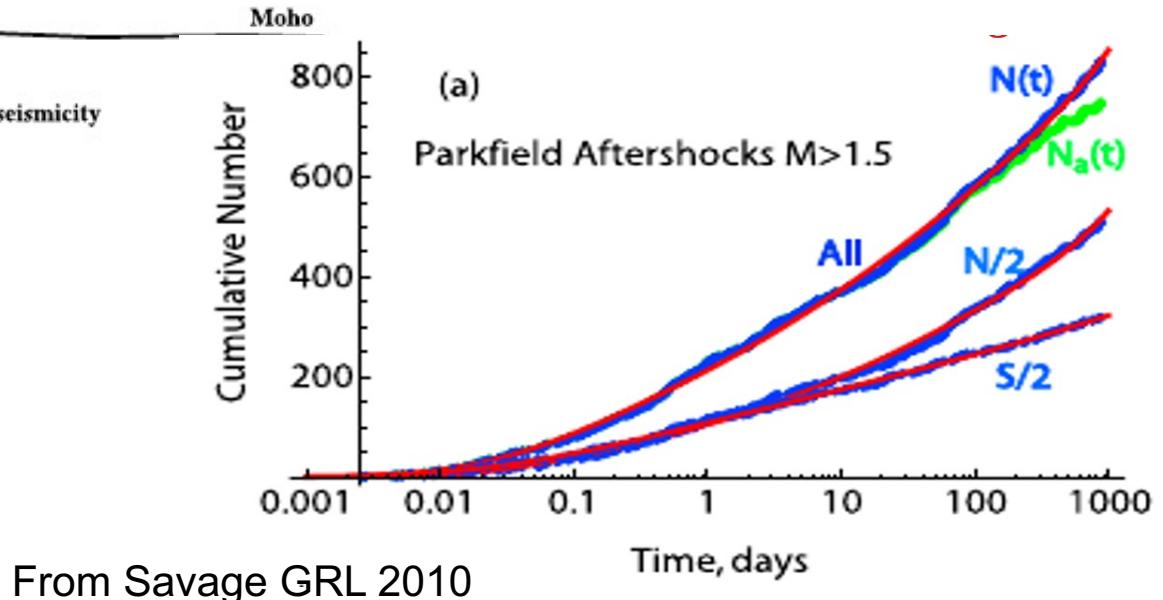
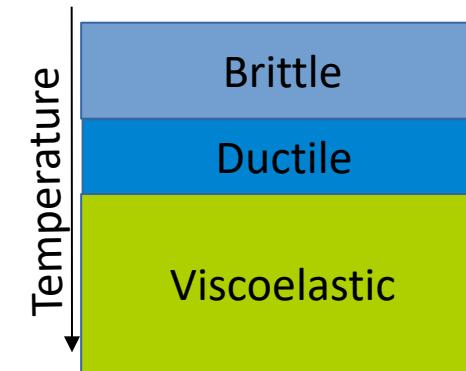
Coupling with the ductile layer



$$u(t) = u_0 + \rho_0 \log \left[1 + e^{\frac{\Delta\tau}{k\rho_0}} \left(\frac{t}{t_R} \right) \right]$$

$$n(t) \sim \frac{1}{t} \Rightarrow N(t) \sim \log(t) \sim u(t)$$

2010

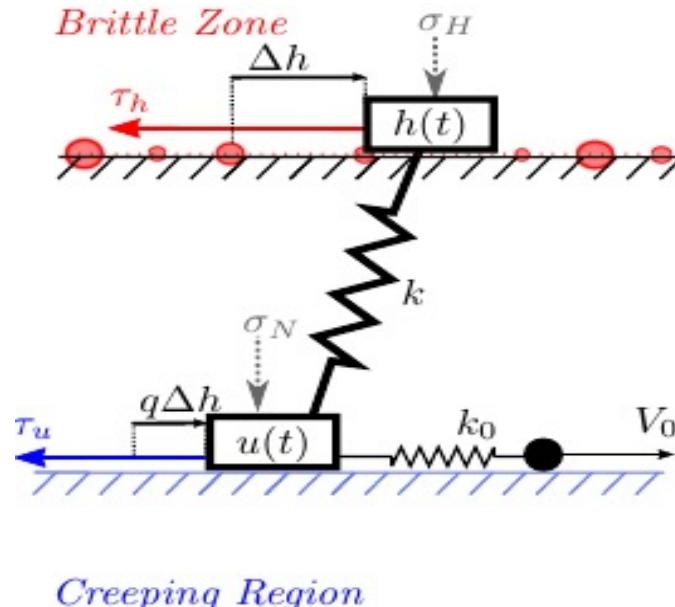


From Savage GRL 2010

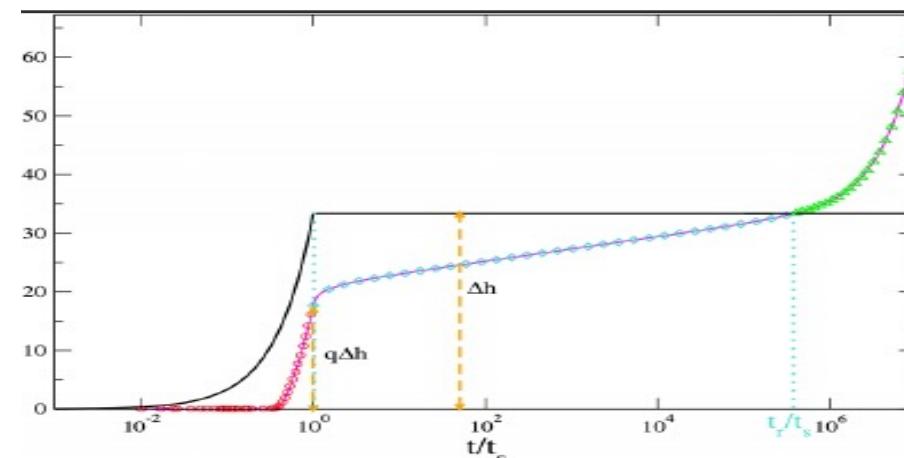
The 2 Layer qEW model

Generalization of the two block model

Two block model (Lippiello et al. 2018)



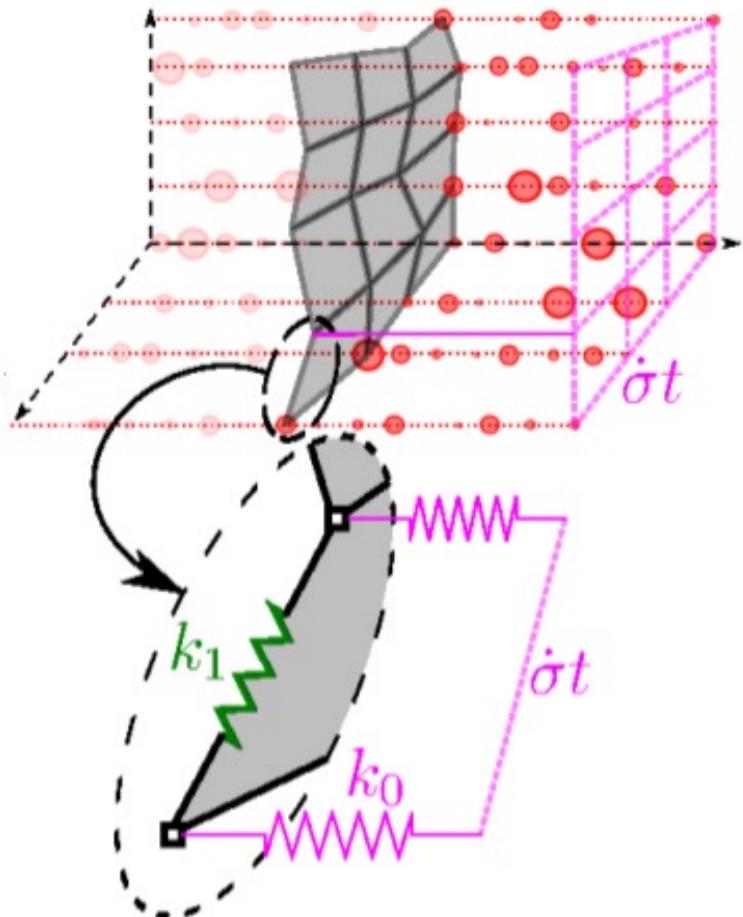
The slip of $h(t)$ induces the coseismic Slip of the $u(t)$ which subsequent relaxes logarithmically because of velocity strengthening friction



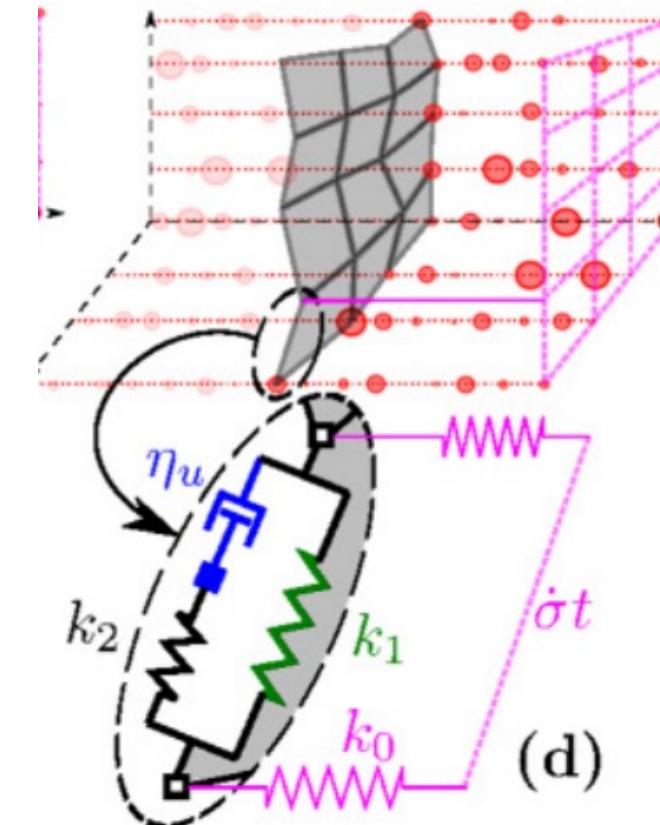
Separation of time scales

- { Strain accumulation rate **cm/year**
- Post seismic defomation velocity **cm/day**
- Strain propagation velocity **km/sec**

qEW model



2L qEW model



The Viscoelastic or 2L qEW model

2L qEW model

All the dynamics can be described in terms of the local stress value

$$F_i = (1-\theta)F_i^{(\text{fast})} + \theta F_i^{(\text{slow})}(t)$$

$F_i^{(\text{fast})}(t)$ is the elastic force caused by usual elastic interaction

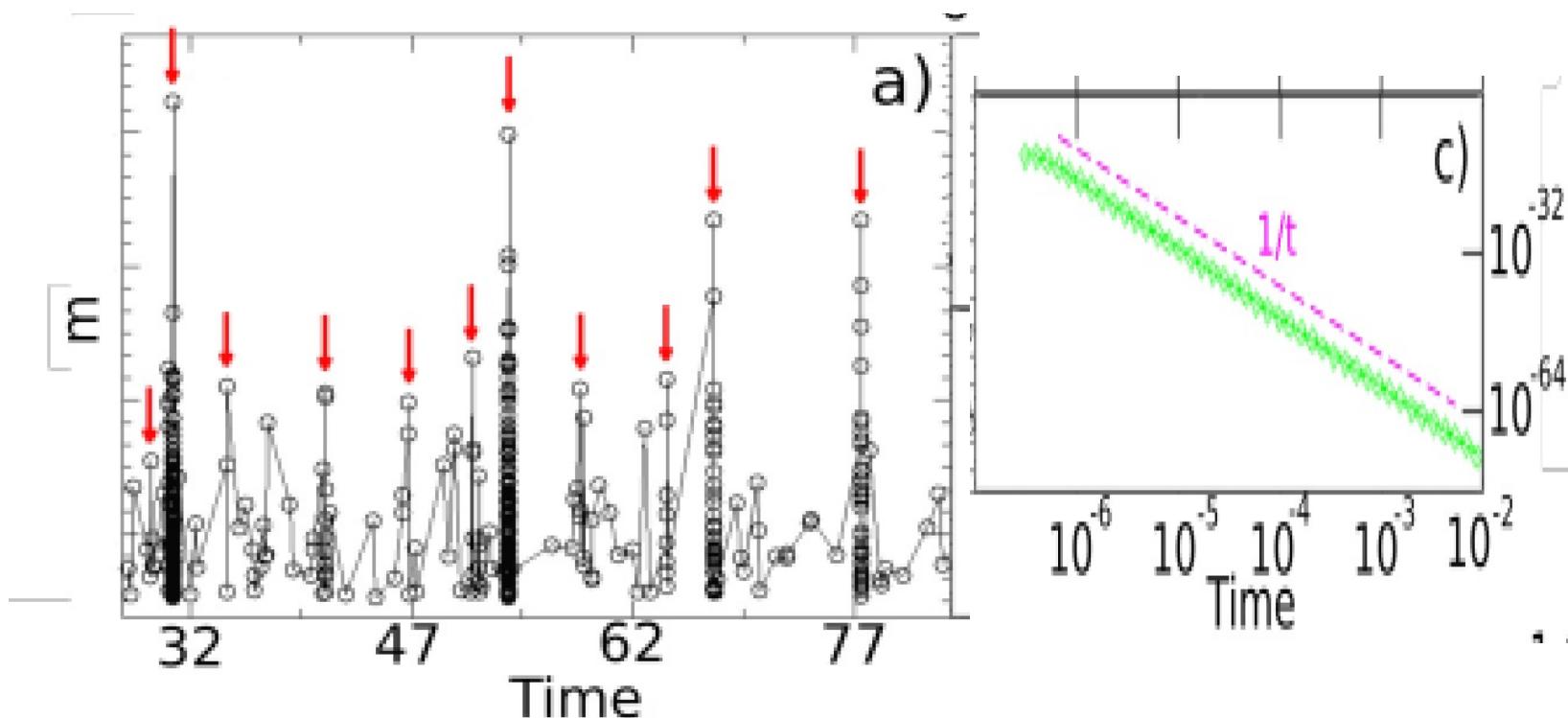
$F_i^{(\text{slow})}(t)$ is the «visco-elastic» force (intraplate force) it decays with a given law $g(t)$

Infinite time separation results are independent of $g(t)$

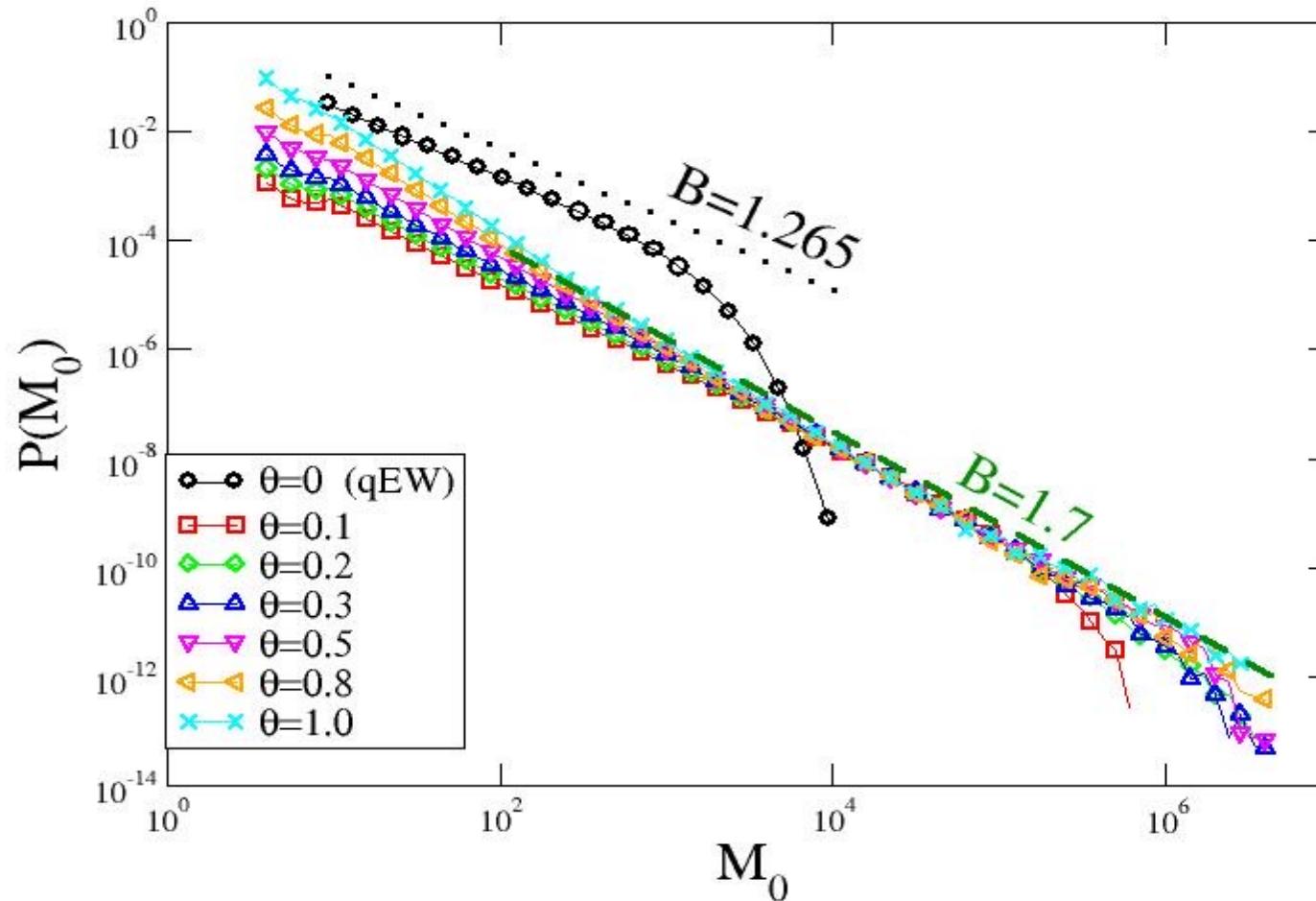
Only one parameter θ (coupling between the two layers)

For $\theta=0$ we have the qEW model

Timing of events in the 2LqEW model

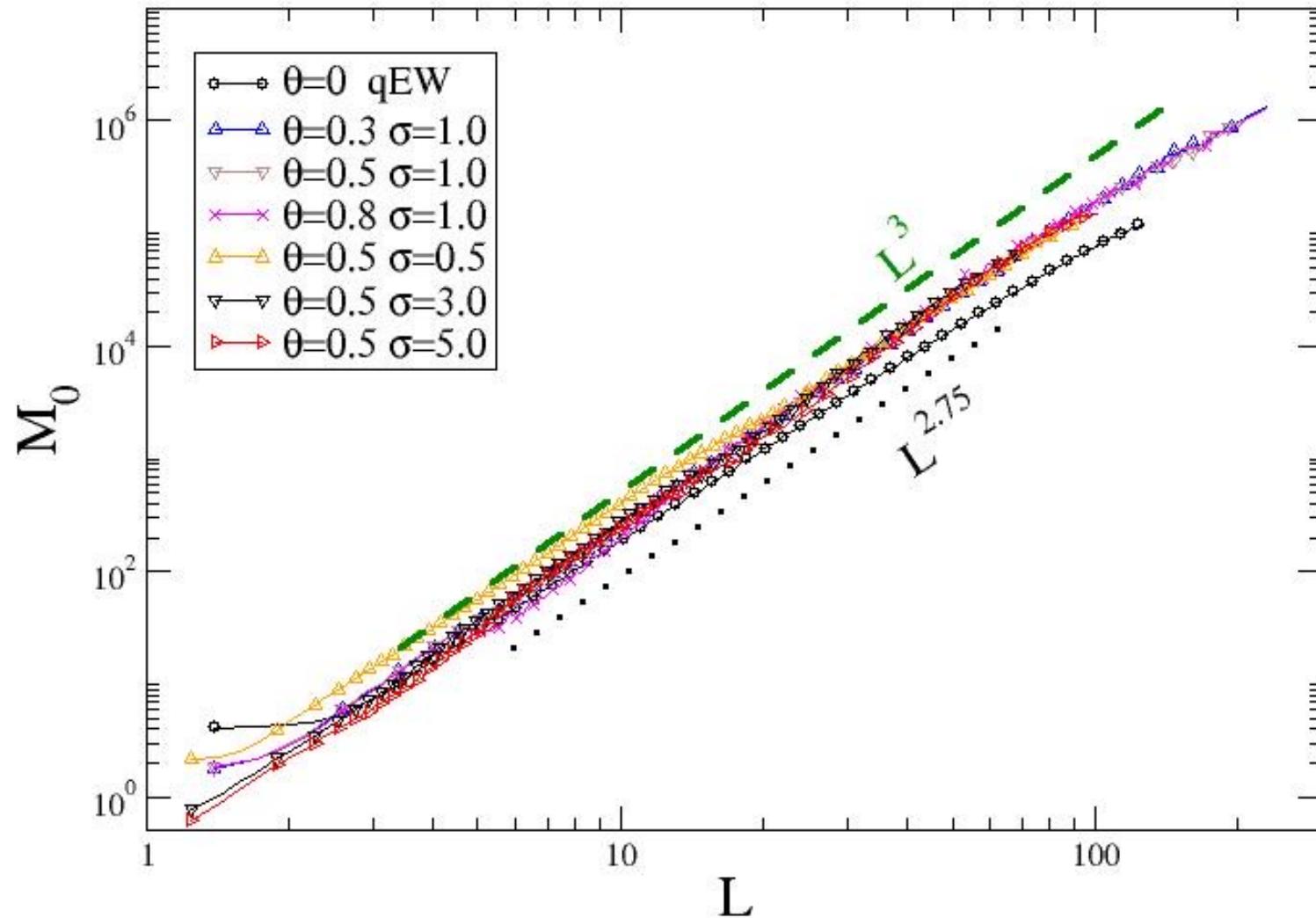


Size distribution in the 2LqEW model

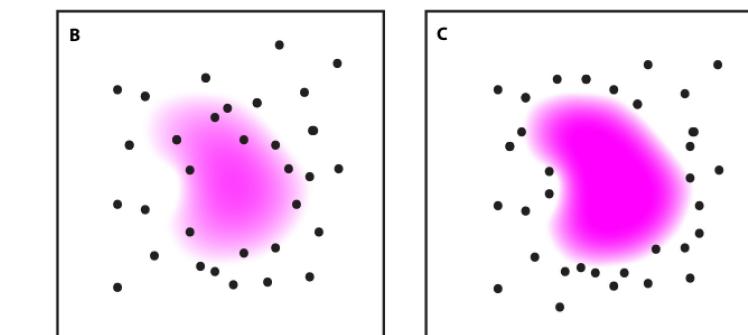
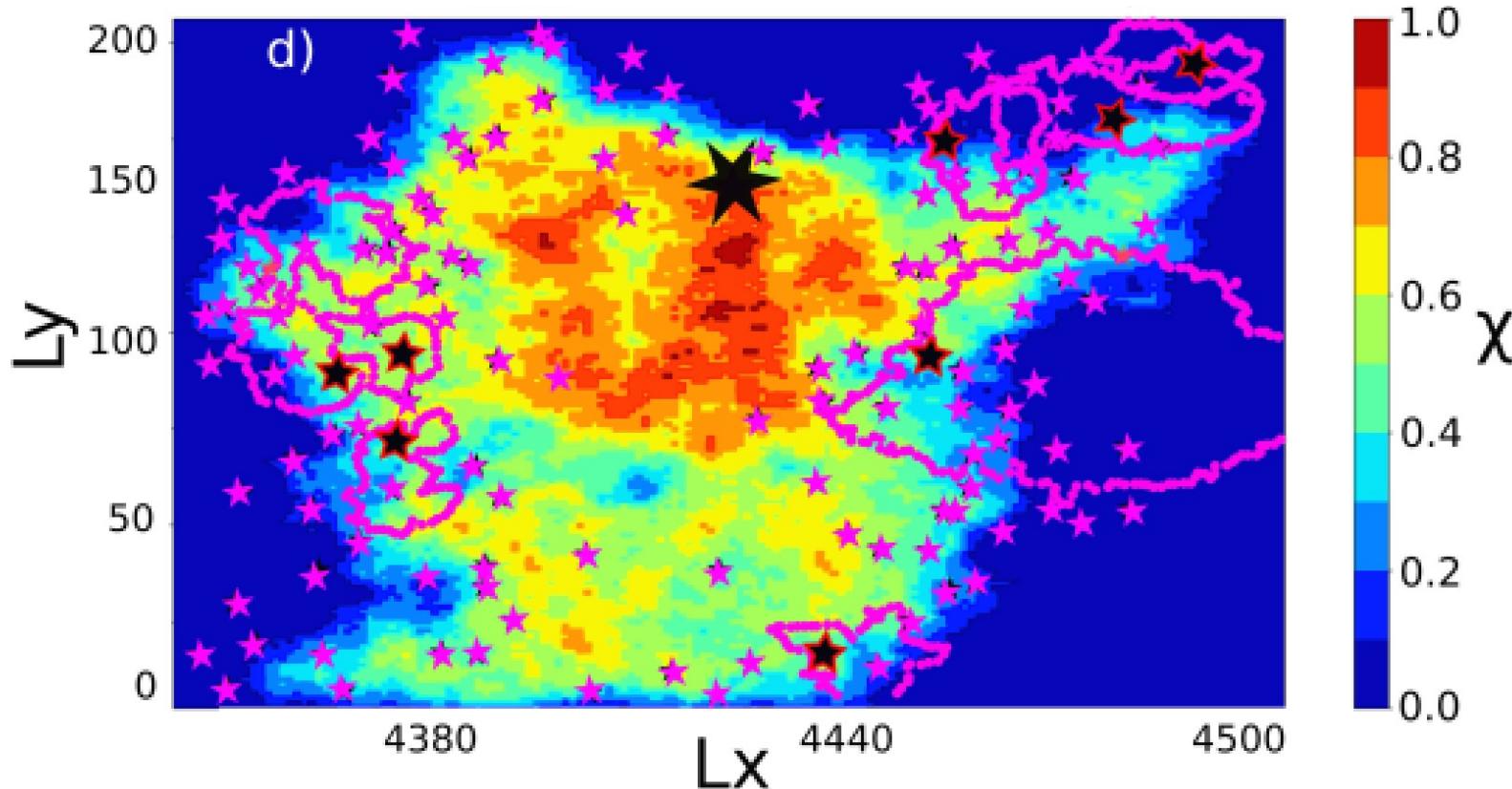


$B=1.7$ indepedent of θ

Scaling in the 2LqEW model



Space organization of events in the 2LqEW model



Final Summary

- The sand-pile model does not capture the temporal organization of earthquakes and also solar flares
- The introduction of an intermediate time scale within the model for depinning transitions leads to a new universality class sharing the same exponents of real seismic catalogs
- The new universality class indicates that large earthquakes are NOT intrinsically unpredictable

Final Summary

"A prince who, moved by a noble heart allows himself to be induced by these misfortunes that touch mankind to remove **the misery of war** from those who are already threatened on all sides by grave misfortunes, is a beneficent instrument working in the benevolent hands of God and a gift that He bestows on the peoples of the earth whose worth they will never know how to estimate in its greatness"

I. Kant, *Allgemein Naturgeschichte und Theorie des Himmels*", 1755

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