Inferring primordial magnetic properties from the resulting gravitational waves

Axel Brandenburg (Nordita)

Collaborators: Emma Clarke, Yutong He, Tina Kahniashvili, Alberto Roper Pol, Jennifer Schober, Ramkishor Sharma, ...

 $\left(\partial_t^2 + 3H\partial_t - c^2\nabla^2\right)h_{ij}(\boldsymbol{x},t) = \frac{16\pi G}{c^2}T_{ij}^{\mathrm{TT}}(\boldsymbol{x},t)$

Relation between spectra:

$$\mathrm{Sp}(\dot{\mathbf{h}}) pprox k^2 \mathrm{Sp}(\mathbf{h}) pprox k^{-2} \mathrm{Sp}(\mathbf{T})$$

GW slope by k^2 steeper Peak at twice magnetic peak



Evolution since electroweak era

Roper

Pol+20



- Power law decay: slope -(β +1)/2
- Hosking+Schekochihin: β =3/2
- Earlier termination (reconnection?)

- Magnetic energy decays
- GW energy does not!

GW energy depends quadratically on energy input & scale



$$\mathcal{E}_{\rm GW}^{\rm sat} \approx (q \mathcal{E}_{\rm M}^{\rm max}/k_{\rm peak})^2$$

Acoustic turbulence more efficient (q~30) Vortical turbulence less efficient (q<5) Helical MHD turbulence least efficient

- Large-scale fields \rightarrow more GW energy
- Generation at electroweak era: need strong fields
- Generation during inflation & reheating

GW energy depends quadratically on energy input & scale



- Regime II: is more resistive \rightarrow unrealistic, but large GW energies
- Regime I: \rightarrow realistic, but small scales & less GW energy

GW energy depends quadratically on energy input & scale



Regime II



$$\tilde{h}(k, t) = \frac{6\tilde{T}_0(k)}{4\gamma_0^2 + k^2} \bigg[e^{2\gamma_0 \tau} - \cos k\tau - \frac{2\gamma_0}{k} \sin k\tau \bigg]_{\tau = t-1}$$

Kinematic phase: Peak determined By growth rate



GW energy depends quintically on limiting CME speed v_{λ}



$$\frac{\partial \boldsymbol{B}}{\partial t} = \boldsymbol{\nabla} \times [\boldsymbol{u} \times \boldsymbol{B} + \eta(\mu_5 \boldsymbol{B} - \boldsymbol{J})], \quad \boldsymbol{J} = \boldsymbol{\nabla} \times \boldsymbol{B}$$

$$\frac{D\mu_5}{Dt} = -\lambda \ \eta(\mu_5 \boldsymbol{B} - \boldsymbol{J}) \cdot \boldsymbol{B} + D_5 \nabla^2 \mu_5 - \Gamma_{\rm f} \mu_5$$

$$v_\lambda=\mu_{50}/\lambda^{1/2}, \qquad v_\mu=\mu_{50}\eta.$$

$$\eta k_1 < v_\mu < v_\lambda$$
 (regime I),
 $\eta k_1 < v_\lambda < v_\mu$ (regime II),

- For realistic parameters \rightarrow very weak GW energy
- Need larger length scales

Magnetic energy also weak

$$\langle \boldsymbol{B}^2 \rangle \, \xi_{\mathrm{M}} = \epsilon \, (k_{\mathrm{B}} T_0)^3 (\hbar c)^{-2},$$

Brandenburg+17

Inflationary magnetogenesis

- Early Universe Turbulence

 Source of gravitational waves
 Information from young universe
- Magnetogenesis
 - Inflation/reheating
 No particles yet, no conductivity
 Coupling with electromagn field
 $f^2 F_{\mu\nu} F^{\mu\nu}$
 - \circ Breaking of conformal invariance \circ Quantum fluct → field stretched

$$\tilde{\mathbf{A}}'' + \left(\mathbf{k}^2 - \frac{f''}{f}\right)\tilde{\mathbf{A}} = 0,$$
$$\tilde{h}''_{+/\times} + \left(\mathbf{k}^2 - \frac{a''}{a}\right)\tilde{h}_{+/\times} = \frac{6}{a}\tilde{T}_{+/\times}$$

CODE (Pencil Code Collaboration et al. 2021), where

$$\mathbf{A}' = \mathbf{u} \times \mathbf{B} + \sigma^{-1} \nabla^2 \mathbf{A} \tag{8}$$

is solved in real space, which includes the induction effect from the velocity and the finite conductivity. It is solved together with (Brandenburg et al. 1996)

$$\mathbf{u}' = -\mathbf{u} \cdot \nabla \mathbf{u} - \frac{1}{4} \nabla \ln \rho + \frac{3}{4\rho} \mathbf{J} \times \mathbf{B} + \boldsymbol{\mathcal{F}}_{\nu} + \boldsymbol{\mathcal{F}}, \quad (9)$$

$$(\ln \rho)' = -\frac{4}{3} \left(\nabla \cdot \mathbf{u} + \mathbf{u} \cdot \nabla \ln \rho \right) + \mathcal{H}, \qquad (10)$$

where $\mathcal{F} = (\nabla \cdot \mathbf{u} + \mathbf{u} \cdot \nabla \ln \rho) \mathbf{u}/3 - [\mathbf{u} \cdot (\mathbf{J} \times \mathbf{B}) + \mathbf{J}^2/\sigma] \mathbf{u}/\rho$, and $\mathcal{H} = [\mathbf{u} \cdot (\mathbf{J} \times \mathbf{B}) + \mathbf{J}^2/\sigma]\rho$ are higher

Software and Data Availability. The source code used for the simulations of this study, the <u>PENCIL CODE</u> (Pencil Code Collaboration et al. 2021), is freely available on https://github.com/pencil-code/. The doi of the code is 10.5281/zenodo. 2315093 (Brandenburg 2018). The simulation setup and the corresponding data are freely available from doi:10.5281/zenodo.4448211, see also http://www.nordita.org/~brandenb/projects/GWfromCME/ for easier access.

Brandenburg & Sharma 2106:03857

Inflationary growth & magnetic decay



Inflationary growth: electric, magnetic, and GW grow



Lorentz force drives smaller scales: surprisingly weak



No small scales in GW field









0.04

0.02

0.00

-0.02

0.04

 $\eta = 16.00$











Circular polarization in chiral inflationary magnetogenesis

$$\iota f^2 F_{\mu\nu} * F^{\mu\nu}$$

$$f(a) = a^{-\beta}, \quad \text{where } a = (\eta + 1)^2/4$$

$$\tilde{A}_{\pm}'' + \left(k^2 \pm 2\iota k \frac{f'}{f} - \frac{f''}{f}\right) \tilde{A}_{\pm} = 0,$$

$$\frac{f'}{f} = -\frac{2\beta}{\eta+1}, \quad \frac{f''}{f} = \frac{2\beta(2\beta+1)}{(\eta+1)^2}$$

- Step I: spectra peaked
- Step II: Inverse cascade
- GW: circularly polarized

$$\mathcal{P}(k) = \int 2 \operatorname{Im} \tilde{h}_{+} \tilde{h}_{\times}^{*} k^{2} \mathrm{d}\Omega_{k} / \int \left(|\tilde{h}_{+}|^{2} + \tilde{h}_{\times}|^{2} \right) k^{2} \mathrm{d}\Omega_{k}$$



Conclusions

- Sp(dh/dt) ~ k^{-2} Sp(T) ~ k^{-2} Sp(B)
- B decays, GW do not \rightarrow messenger
- CME: interesting, but weak B and GWs
- Inflation & reheating: large scale

 Small scales during radiation era
 But surprisingly inefficient
 - $\circ f^{8/3}$ spectrum characteristic of turbulence
- NANOGrav result $\leftarrow \rightarrow$ low energy scale

