

THEORY GRANT

# Simulations of Sun-to-Earth Propagation of Solar Eruptions (SSEPSE)

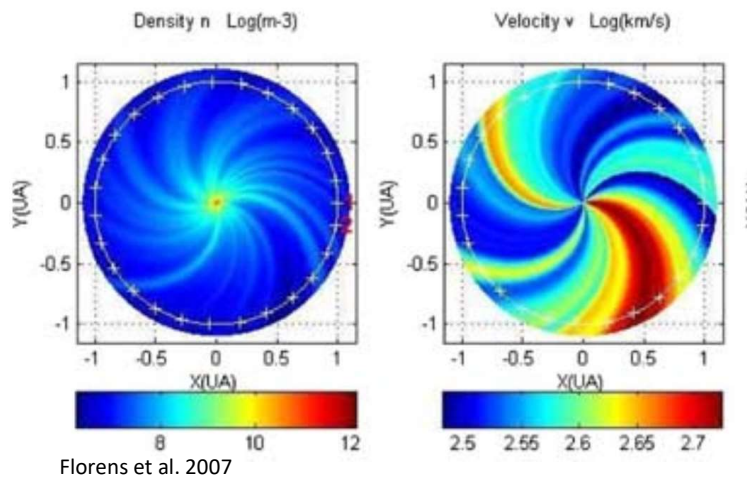
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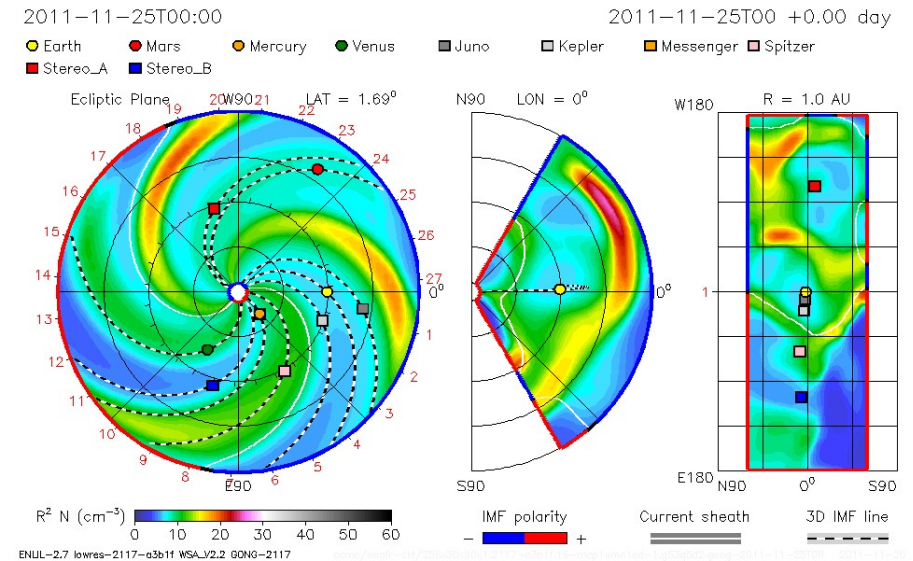
# Analytical Models



Assuming *stationarity* in the solar wind conditions it is possible to **track backwards** the plasma trajectories by following the **Parker Spiral** starting from 1 AU in-situ data.

Schatten, K. H., N. F. Ness, and J. M. Wilcox (1968) *Influence of a Solar Active Region on the Interplanetary Magnetic Field*, Sol.Phys.,5(2), 240–256  
 Schatten, K. H. (1971) *Current sheet magnetic model for the solar corona*, Cosmic Electrodynamics,2, 232–245  
 Florens et al. (2007) *Data-driven solar wind model and prediction of type II bursts*, Geophys. Res. Lett. 34  
 Tasnim et al. (2018) *A generalised equatorial model for the accelerating solar wind*, JGR (SP) 123

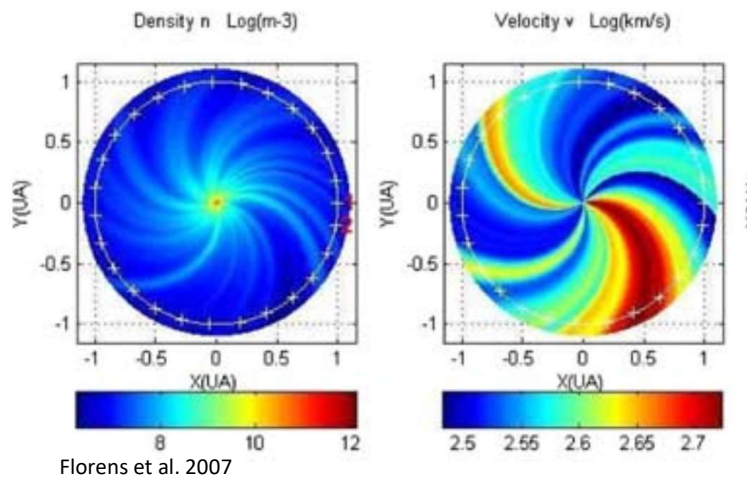
# Numerical Models



Starting from 2D **photospheric magnetograms**, different techniques reconstruct the 3D magnetic fields in the inner and intermediate corona, using then numerical and/or empirical methods **to build the i.b.** of a MHD simulation.

Mikić et al. (1999): *Magnetohydrodynamic modeling of the global solar corona*, Phys. Of Plasmas, vol 6, issue 5  
 Pizzo et al. (2011): *Wang-Sheeley-Arge-ENLIL Cone Model Transitions to Operations*, Space Weather, Volume 9, S03004  
 Pomoell & Poedts (2018): *EUHFORIA: European heliospheric forecasting information asset*, JSWSC 8, A35

# Analytical Models



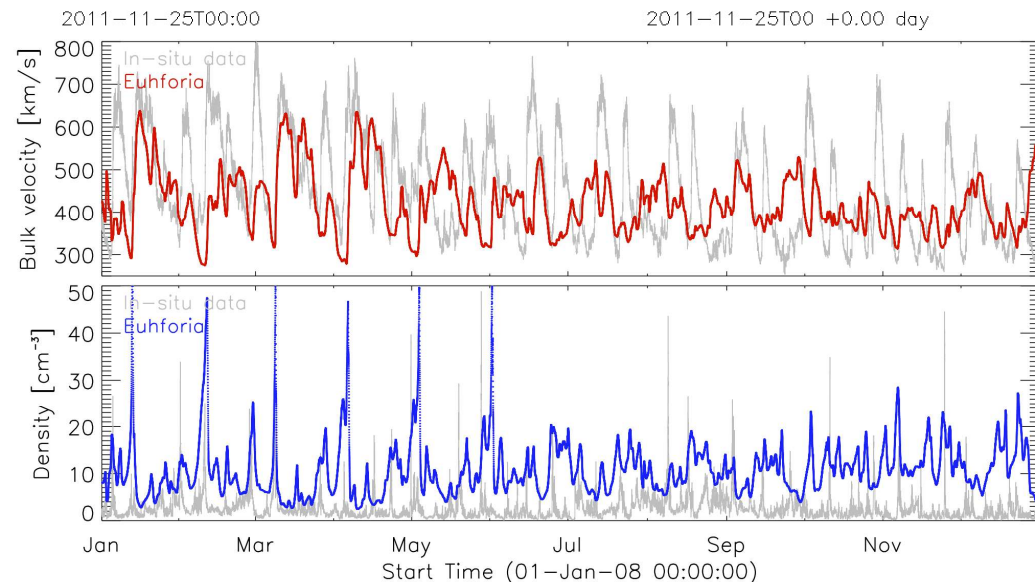
- Reconstructions **complicated** by stream interaction regions, magnetic clouds, shocks and other **transient phenomena**
- **Not time-dependent**

Pizzo (1981): *On the application of numerical models to the inverse mapping of solar wind flow structures*, JGR 86(A8)

Jian et al. (2016): *Validation for global solar wind prediction using Ulysses comparison: Multiple coronal and heliospheric models installed at the Community Coordinated Modeling Center*, SW 14, 8

Riley et al. (2018): *Forecasting the Arrival Time of Coronal Mass Ejections: Analysis of the CCMC CME Scoreboard*, SW 16, 9

# Numerical Models



- Agreement between the **reconstructed plasma ambient condition** and the available **in-situ measurements** still **suboptimal**
- **Forecasting capability** of solar disturbances still **limited**: standard deviation of predicted CMEs **arrival times can exceed 20 h**

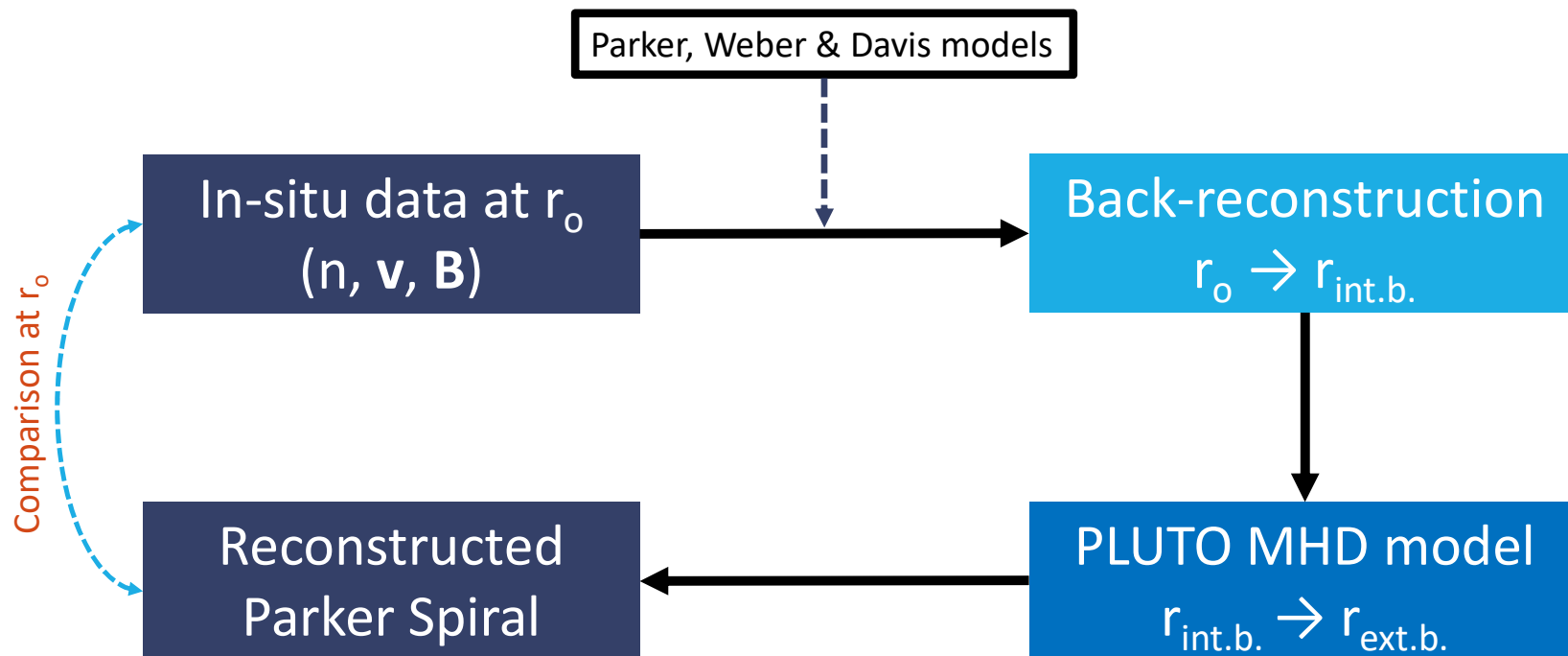
Hinterreiter et al. (2019): *Assessing the Performance of EUHFORIA Modeling the Background Solar Wind*, Sol. Phys. 294:170

Riley et al. (2018): *Forecasting the Arrival Time of Coronal Mass Ejections: Analysis of the CCMC CME Scoreboard*, SW 16, 9

# RIMAP

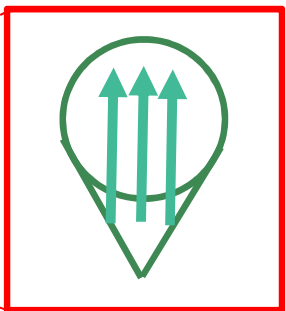
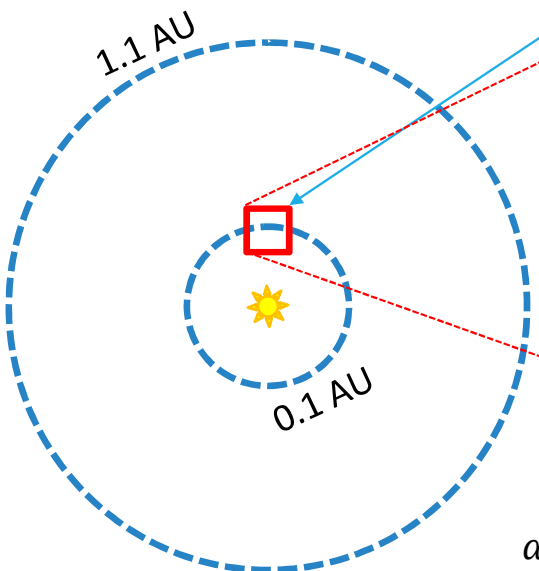
## Reverse In-situ and MHD Approach

Combining the best of both approaches to offer **an alternative method** to the reconstruction of the Parker Spiral



# Modeling of ICMEs in RIMAP: cone model

We insert in the inner radial boundary of RIMAP an **ice-cream cone model** (Zhao et al. 2002, Xie et al. 2004, Gopalswamy et al. 2009).

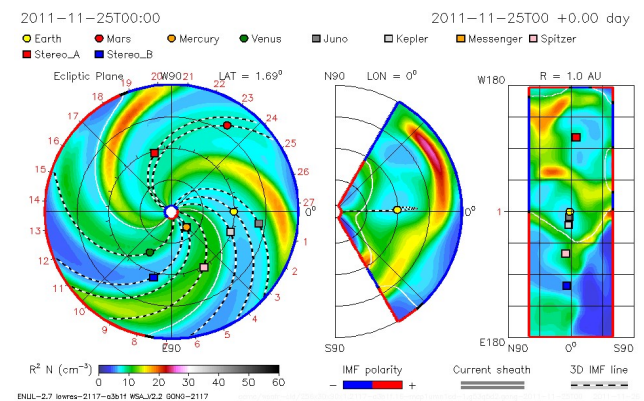


Cone models assume that close to the Sun, **in the early phases\* of the ICME propagation, its angular width and velocity remain constant.** Thus the perturbation can be described as a homogeneous plasma cloud, isotropic in expansion and with entirely radial bulk velocity.

Due to this simple geometry, cone models are particularly **convenient for routine applications** in space-weather forecasting (es. Scolini et al. 2018).

$$\alpha(t) = \frac{\omega}{2} \sin\left(\frac{\pi(t - t_0)v_0}{2 R_b \sin\left(\frac{\omega}{2}\right)}\right)$$

if  $(\varphi - \varphi_0)^2 \leq \alpha^2(t)$ , the background **solar wind parameters are replaced with the perturbation ones.**

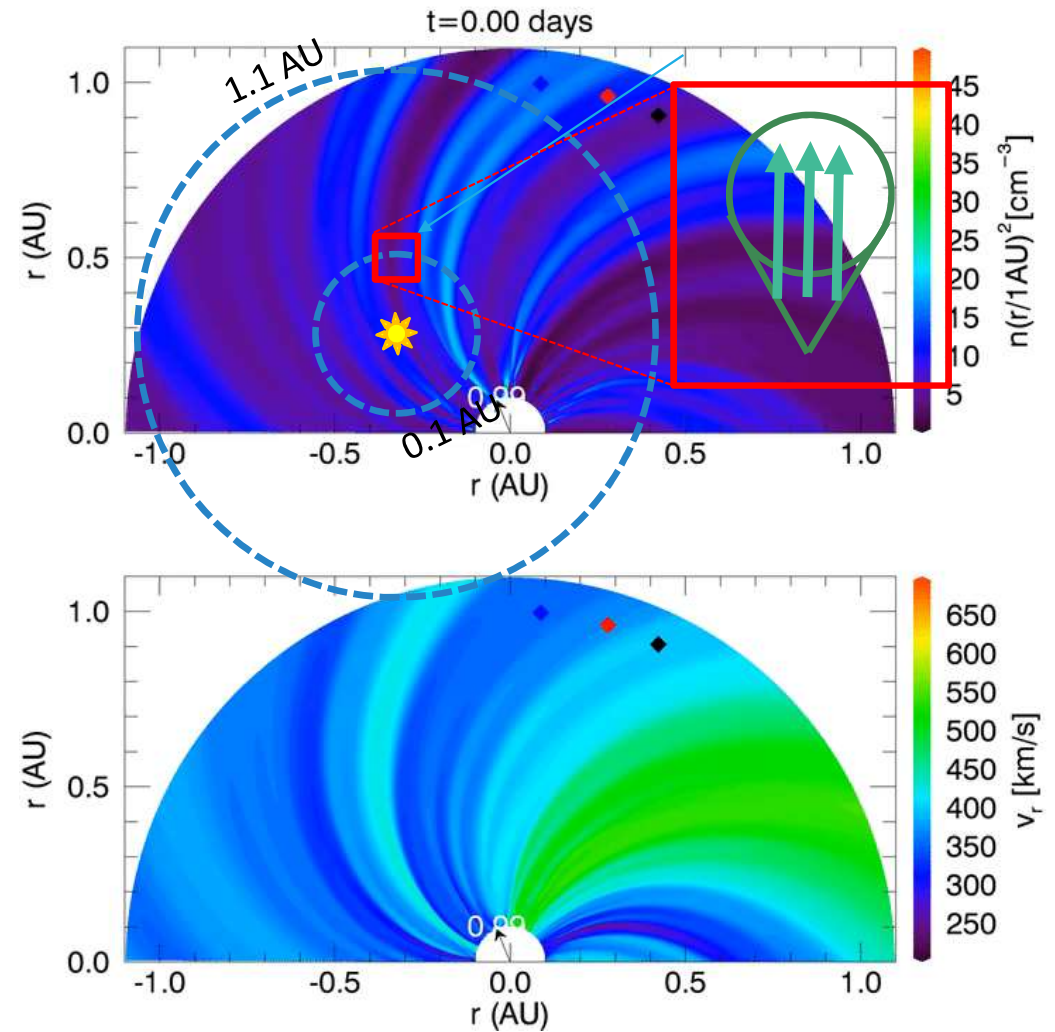


Unmagnetized ICME propagation with ENLIL

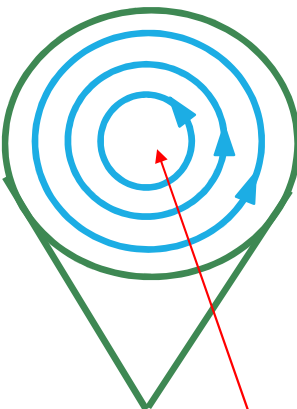


# Modeling of ICMEs in RIMAP

- **Cone model** (*e.g.* Zhao et al. 2002, Gopalswamy et al. 2009) in the RIMAP inner radial boundary;
- Density and plasma temperature inside the ICME **homogeneous** ( $n_0 = 600 \text{ cm}^{-3}$  and  $T_0 = 5 \cdot 10^5 \text{ K}$ );
- An artificial, passive scalar is added to the ICME model  $\rightarrow$  **flow tracer  $T_r$** , obeys a simple advection equation and does not interact with the physical quantities of the model;
- ICME initial velocity  $v_0 = 800 \text{ km/s}$ ;  
Parker spiral reconstructed by RIMAP using WIND data of March 2009;  
Super-imposed, in white, the iso-lines of  $T_{rg}$ ;
- The **ICME remain compact**, even in the absence of internal flux-rope, traveling across the highly structured configuration.



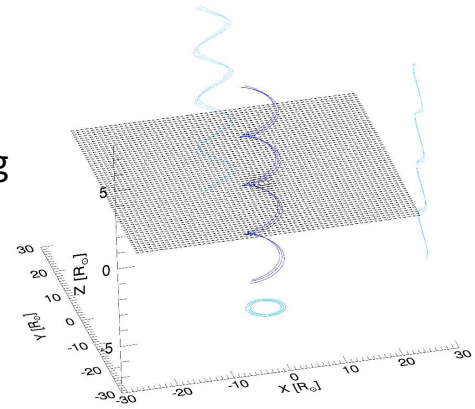
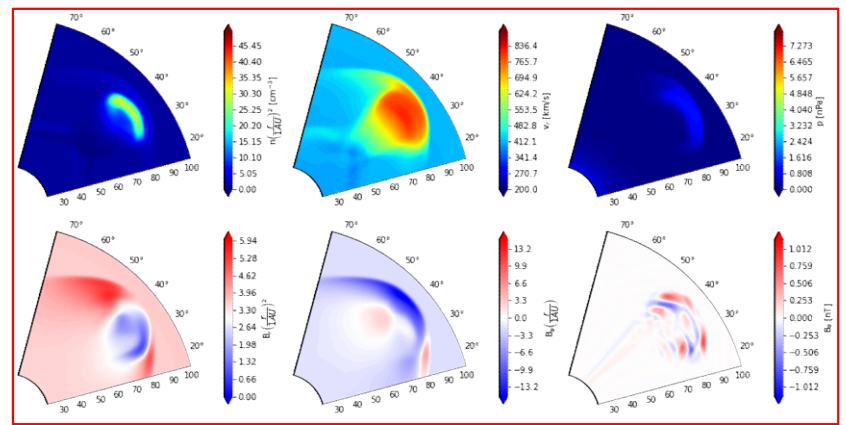
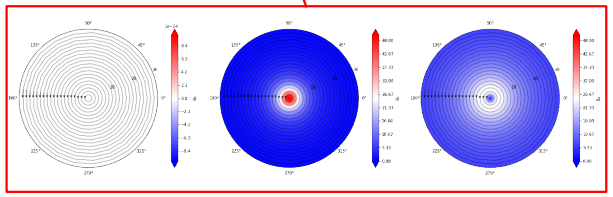
# Modeling of ICMEs in RIMAP: work in progress



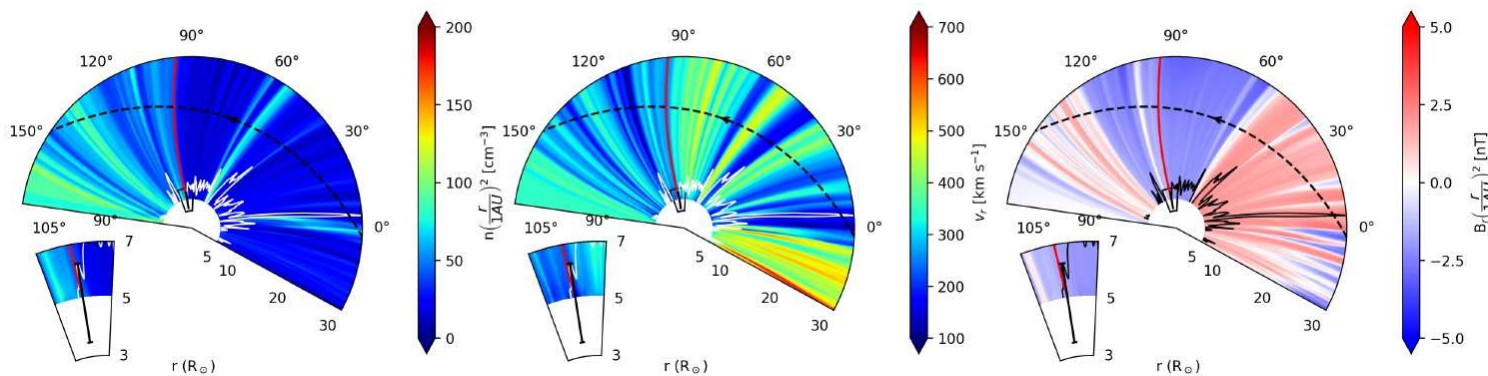
- **1/3** of ICMEs at 1AU shows clear **magnetic cloud** signatures (e.g. Cane and Richardson 2003), but this fraction is probably larger (e.g. Jian et al. 2006, Kilpua et al. 2011);
- Magnetic clouds can be described by cylindrical symmetric flux ropes with **force-free B**. One of the simplest solution is the Gold & Hoyle 1960 one, generalized by Low & Berger 2003:

$$B_R = 0, B_z = \frac{B_0}{1 + k^2 R^2}, B_\phi = kRB_z$$

- We are inscribing different flux-ropes into the cone-modelled ICME, studying its propagation across the RIMAP-reconstructed Parker spiral.



# Boundary conditions from observations



- **Pre-CME background plasma conditions** will be reconstructed with RIMAP combining Solar Orbiter/ Metis observations and/or Parker Solar Probe data (e.g. [Biondo et al. 2022](#));

- **CME kinematic and thermodynamic conditions** will be reconstructed combining Solar Orbiter/ Metis observations and SOHO/ LASCO-C2 & C3 observations; in particular, the plasma thermodynamic conditions inside the CME will be derived combining VL and UV (H I Lyman- $\alpha$ ) images acquired by Metis (e.g. [Bemporad 2022](#)), as well as the **plasma conditions across the shock front**.

