Simulations of relativistic jets and recollimation in extreme blazars



Agnese Costa, Fabrizio Tavecchio, (INAF OA Brera -Merate) Gianluigi Bodo (INAF OA Torino)

PASTO 7.09.2022

Simulations of relativistic jets and recollimation in extreme blazars

- ♦ Introduction: TeV blazars and simulations with PLUTO
- ♦ 2D relativistic MHD simulations with lagrangian particles for TeV blazars
- \Rightarrow 2D \rightarrow 3D simulations and turbulence

Blazars with high energy peaked, hard spectrum.

Model: standard leptonic emission via synchrotron and SSC from electrons accelerated through DSA, but the slope is very hard!



Costamante et al. 2018 doi: 10.1093/mnras/sty857

Blazars with high energy peaked, hard spectrum.

Model: standard leptonic emission via synchrotron and SSC from electrons accelerated through DSA, but the slope is very hard!

What's the acceleration process?

 A good guess: due to a series of recollimation and reflection shocks (Zech A., Lemoine M., 2021, doi:10.1051/0004-6361/202141062)



Blazars with high energy peaked, hard spectrum.

Model: standard leptonic emission via synchrotron and SSC from electrons accelerated through DSA, but the slope is very hard!

What's the acceleration process?

 A good guess: due to a series of recollimation and reflection shocks (Zech A., Lemoine M., 2021, doi:10.1051/0004-6361/202141062)

BUT can it happen and work? (my research)

2. Other phenomenology? (the current part of my work)



Blazars with high energy peaked, hard spectrum.

Model: standard leptonic emission via synchrotron and SSC from electrons accelerated through DSA, but the slope is very hard!

What's the acceleration process?

 A good guess: due to a series of recollimation and reflection shocks (Zech A., Lemoine M., 2021, doi:10.1051/0004-6361/202141062)

BUT can it happen and work? (my research)

- 2. Other phenomenologies? (the current part of my work)
- 3. Would these work as acceleration mechanisms? (hear more from A. Sciaccaluga next)



New fit from Tavecchio et al. 2022.

Plasma + Lagrangian particles simulations with PLUTO

Plasma:

Conservation equations evolved with shock-capturing finite volume (or finite difference) methods

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot \left(\rho \boldsymbol{v}\right) &= 0\\ \frac{\partial \boldsymbol{q}}{\partial t} + \nabla \cdot \left[\boldsymbol{q}\boldsymbol{v} - \boldsymbol{B}\boldsymbol{B} + \left(\boldsymbol{p} + \frac{\boldsymbol{B}^2}{2}\right)\right]^T &= -\rho \nabla \Phi + \rho \boldsymbol{g}\\ \frac{\partial \boldsymbol{B}}{\partial t} + \nabla \times (c\boldsymbol{E}) &= 0\\ \frac{\partial \left(\frac{\rho \boldsymbol{v}^2}{2} + \rho \boldsymbol{e} + \frac{\boldsymbol{B}^2}{2} + \rho \Phi\right)}{\partial t} + \nabla \cdot \left[\left(\frac{\rho \boldsymbol{v}^2}{2} + \rho \boldsymbol{e} + \boldsymbol{p} + \rho \Phi\right)\boldsymbol{v} + c\boldsymbol{E} \times \boldsymbol{B}\right] &= \boldsymbol{q} \cdot \boldsymbol{g} \end{aligned}$$

ideal MHD for example

Plasma + Lagrangian particles simulations with PLUTO

Plasma:

Conservation equations evolved with shock-capturing finite volume (or finite difference) methods

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot \left(\rho \boldsymbol{v}\right) &= 0\\ \frac{\partial \boldsymbol{q}}{\partial t} + \nabla \cdot \left[\boldsymbol{q}\boldsymbol{v} - \boldsymbol{B}\boldsymbol{B} + \left(\boldsymbol{p} + \frac{\boldsymbol{B}^2}{2}\right)\right]^T &= -\rho \nabla \Phi + \rho \boldsymbol{g}\\ \frac{\partial \boldsymbol{B}}{\partial t} + \nabla \times (c\boldsymbol{E}) &= 0\\ \frac{\partial \left(\frac{\rho \boldsymbol{v}^2}{2} + \rho \boldsymbol{e} + \frac{\boldsymbol{B}^2}{2} + \rho \Phi\right)}{\partial t} + \nabla \cdot \left[\left(\frac{\rho \boldsymbol{v}^2}{2} + \rho \boldsymbol{e} + \boldsymbol{p} + \rho \Phi\right)\boldsymbol{v} + c\boldsymbol{E} \times \boldsymbol{B}\right] &= \boldsymbol{q} \cdot \boldsymbol{g} \end{aligned}$$

ideal MHD for example

Non-thermal particles: Lagrangian particle module

- Macroparticles of n real particles characterized by a spectral distribution
- They move following the fluid (fluid quantities at particle position is found via standard interpolation methods)
- No feedback on the fluid
- Energy distribution follows the CR transport equation (non diffusive)

Plasma + Lagrangian particles simulations with PLUTO

Plasma:

Conservation equations evolved with shock-capturing finite volume (or finite difference) methods

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot \left(\rho v\right) &= 0\\ \frac{\partial q}{\partial t} + \nabla \cdot \left[qv - BB + \left(p + \frac{B^2}{2}\right)\right]^T &= -\rho \nabla \Phi + \rho g\\ \frac{\partial B}{\partial t} + \nabla \times (cE) &= 0\\ \frac{\partial \left(\frac{\rho v^2}{2} + \rho e + \frac{B^2}{2} + \rho \Phi\right)}{\partial t} + \nabla \cdot \left[\left(\frac{\rho v^2}{2} + \rho e + p + \rho \Phi\right)v + cE \times B\right] &= q \cdot g \end{aligned}$$

ideal MHD for example

Non-thermal particles: Lagrangian particle module

- Macroparticles of n real particles characterized by a spectral distribution
- They move following the fluid (fluid quantities at particle position is found via standard interpolation methods)
- No feedback on the fluid
- Energy distribution follows the CR transport equation (non diffusive)
- After crossing shocks particles are accelerated and the distribution is updated
- Non thermal emission via synchrotron and IC.
- More can be implemented

Setup compatible with TeV blazars

- Relativistic
- Low magnetized and $B = 10^{-3}G$

$$p_e = p_{0e} \left(\frac{z}{z_0}\right)^{-1.8}$$
 and $p_{0j} \ll p_{0e}$ (underpressured)

Setup compatible with TeV blazars

- Relativistic
- Low magnetized and $B = 10^{-3}G$

$$p_e = p_{0e} \left(\frac{z}{z_0}\right)^{-1.8}$$
 and $p_{0j} \ll p_{0e}$ (underpressured)



Setup compatible with TeV blazars

- Relativistic
- Low magnetized and $B = 10^{-3}G$
- $p_e = p_{0e} \left(\frac{z}{z_0}\right)^{-1.8}$ and $p_{0j} \ll p_{0e}$ (underpressured)











Diffusive Shock Acceleration

- $N_e(E) \propto E^{-2}$ for strong shocks
- efficient Fermi I acceleration mechanism:

$$\frac{\langle \Delta E \rangle}{E} \propto \frac{\mathrm{v}}{\mathrm{c}}$$

good for low magnetized jets (otherwise magnetic reconnection too)

In recollimation and reflection shocks the shock surface is curved and the angle between the upstream velocity and the surface is variable

• the strength of the shock increases towards the recollimation point.

acceleration sites in recollimation-reflection shocks



Setup compatible with TeV blazars

- Relativistic
- Low magnetized and $B = 10^{-3}G$
- $p_e = p_{0e} \left(\frac{z}{z_0}\right)^{-1.8}$ and $p_{0j} \ll p_{0e}$ (underpressured)

Non thermal particle injection

- Non thermal density and energy distribution are initialized to be small and mildly relativistic (we expect to have a population of pre-accelerated particles)
- Specifics have little impact on the results.



Results for different particles: 1, 2, 4, 5

- The particles follow the plasma, so they undergo shocks of varying strength depending on the direction of the shock surface, and of the plasma velocity, at their position.
- Particles are accelerated the most near the axis, near the recollimation point.
- The reflection shock is stronger at the axis as well.



Results for different particles:

- Particle 1 is shocked near the axis and soon is reshocked
- Particle 2 experiences two, well distinct, shocks
- Particle 4 experiences two, well distinct, shocks
- Particle 5 seems to be faintly shocked at the beginning and then is accelerated from the increasing background velocity.



Energy distribution for macroparticle 2:

• energized from shocks



Energy distribution for macroparticle 2:

- energized from shocks
- but the reflection shock has no big impact on further acceleration.



Energy distribution for macroparticle 2:

- energized from shocks
- but the reflection shock has no big impact on further acceleration.



First result:

Subsequent shocks might not be able to accelerate further the electrons!



From 2D to 3D: turbulence in low magnetized jets

1. A series of reconfinement and recollimation shocks might not be able to accelerate enough higher energy particles to produce the hard slope of the TeV blazars.



From 2D to 3D: turbulence in low magnetized jets

- 1. A series of reconfinement and recollimation shocks might not be able to accelerate enough higher energy particles to produce the hard slope of the TeV blazars.
- 2. Instabilities that cannot be seen in 2D simulations might be relevant in reality:
 - Centrifugal instability caused by the recollimation shock \rightarrow **turbulence in low magnetized jets** ($\sigma = \frac{B^2}{4\pi\omega} \le 10^{-4}$)

Matsumoto, Komissarov, Gourgouliatos doi:10.1093/mnras/stab828



What now?

3D RHD simulation starting from the **2D stationary solution** to check on turbulence developing with the PLUTO code:



2D stationary solution

user: Agnese Sun Sep 4 11:00:17 2022

What now?

3D RHD simulation starting from the **2D stationary solution** to check on turbulence developing with the PLUTO code:

work in progress!



-0.5

0.0

X-Axis

-1.5

-1.0

DB: data.0026.vtk

Cycle: 26

Pseudocolor

user: Agnese Sun Sep 4 10:57:09 2022

1.0

1.5

0.5

slice of the 3D simulation

What now?

3D RHD simulation starting from the **2D stationary solution** to check on turbulence developing with the PLUTO code:

work in progress!



-0.5

0.0

X-Axis

-1.5

-1.0

DB: data.0026.vtk

Cycle: 26

Pseudocolor

user: Agnese Sun Sep 4 10:57:09 2022

1.0

1.5

0.5

slice of the 3D simulation

What's next?

- 3D RHD to 3D Relativistic MHD to check on turbulence developing
- Particle acceleration (and sync+IC emission) in the 3D final setup

The end!

The end!

Thank you for your kind attention

Recollimation shocks in AGN jets

Relativistic AGN jets expand and cool adiabatically through the environment

$$p_j = p_{0j} \left(\frac{\sqrt{z^2 + r^2}}{z_0} \right)^{-2}$$

while the environment pressure follows a general power law decaying with distance from the central engine

$$p_e = p_{0e} \left(\frac{z}{z_0}\right)^-$$

When there is a pressure unbalance in favour of the environment the jet undergoes a recollimation process that is supersonic and waves/shocks form

a

jet (black) and environment (red) pressure -8 E $\alpha = 0.625$ -9 -10 (d)60 -12 -13 -14 2.0 2.5 3.0 3.5 4.5 5.0 4.0 log(z)

 $\alpha < 2$

Dimensions:

$$z_0 = 0.1 \, pc$$

 $r_0 = 0.1 \, z_0 = 0.01 \, pc$

Plasma parameters:

$$\gamma = 30$$

$$L_{j} = 10^{44} erg$$

$$B_{0} = 3 \cdot 10^{-4}G \text{ poloidal}$$

$$p_{a,0} = 1 \ dyn/cm^{2} \ and \ p_{a} = p_{a,0} \left(\frac{z}{z_{0}}\right)^{-1.8}$$

$$p_{j,0} = 10^{-7}p_{a,0}$$

$$n_{a} = 10^{9} cm^{-3}$$



Particle parameters:

Initialization is meant to provide a minimal Lorentz factor for shock acceleration, because PLUTO itself still doesn't require any minimal energy (updated to be done):

Power law with:

- $\gamma_{nth,0} = (10^2, 5 \cdot 10^3)$
- $n_{nth,0} = 10^{-10} cm^{-3}$

Graph info

- Unit of measure of the energy is erg
- Acceleration up to $\gamma = 10^8$



Results for different particles:

- Particle 1 crosses the strongest shocks
- Particle 2 experiences weaker shocks.
- Particle 4 experiences even weaker shocks.
- Particle 5 does not capture the shock: the injected particle distribution is only updated.





Results for different particles:

- Particle 1 crosses the strongest shocks
- Particle 2 experiences weaker shocks.
- Particle 4 experiences even weaker shocks.
- Particle 5 does not capture the shock: the injected particle distribution is only updated.



BACKUP: references

Vaidya B. et al., 2018 for the details of the implementation of diffusive shock acceleration in PLUTO and all its sources.

Zech A., Lemoine M., 2021 for the emission model of TeV blazars

Komissarov S., Gourgouliatos K., 2017 for 3D relativistic HD simultations of turbulence in recollimated jets

Matsumoto et al., 2021, for 3D relativistic MHD simultations of turbulence in recollimated jets